

# Hamiltonian Embedding

## & A Paradigm Shift in Quantum Algorithm Design

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# I. Input models of matrices

# Input models of a matrix $A$

## Sparse-input oracles

$$O_r: |i\rangle|k\rangle \rightarrow |i\rangle|r_{ik}\rangle, \quad O_c: |\ell\rangle|j\rangle \rightarrow |c_{\ell j}\rangle|j\rangle$$
$$O_A: |i\rangle|j\rangle|0\rangle^{\otimes b} \rightarrow |i\rangle|j\rangle|a_{ij}\rangle$$

- $A$  sparse, not necessarily Hermitian
- Applications: quantum simulation, quantum linear system solvers, Carleman linearization, etc.
- **Limitations:** no guarantee on the actual cost.

## Block-encoding

$$U = \begin{bmatrix} A/\alpha & \\ & \cdot \end{bmatrix}, \quad A = \alpha(\langle 0| \otimes I)U(|0\rangle \otimes I)$$

- $A$  not necessarily sparse, not necessarily Hermitian
  - Applications: quantum signal processing, quantum singular value transformation, etc.
  - **Limitations:** no guarantee on the actual cost. Still far from NISQ era.
- 
- Dalzell et. al. “End-to-end resource analysis for quantum interior point methods and portfolio optimization”. [\[arXiv: 2211.12489\]](https://arxiv.org/abs/2211.12489). Portfolio optimization with **100** shares (easy for a personal laptop): T-count  $\approx 2 \times 10^9$  in a **single** block-encoding, and **repeat**  $\sim 10^7$  times.
  - Liu et. al. “An efficient block encoding quantum circuits for a pair Hamiltonian”. **10** fermions: T-count  $\approx 3 \times 10^3$ , 2-qubit entanglement gates  $\sim 6 \times 10^3$  (can be further improved!)

# Input models of a matrix $A$ (Con't)

## Pauli basis decomposition (Pauli strings)

$$A = \sum_s a_s P_s, \quad a_s = \frac{1}{2^n} \text{Tr}[A P_s]$$

$\{P_s\}$  = the set of Pauli basis (e.g., I, X, Y, Z)  
 $P_s$  is  $k$ -local if involving  $k$  non-identity Paulis

- **Limitations:**  $A$  must be Hermitian. The coefficients  $a_s$  could require exponential time to compute.
- **Advantages:** Easy to simulate on real devices if matching the hardware native *instruction set*.

## Analog Quantum Computers

Example: Rydberg atoms (QuEra)

$$H = \Omega \sum_j X_j + \Delta \sum_j Z_j + \sum_{j < k} V_{jk} Z_j \otimes Z_k, \quad V_{jk} = C_6 / |r_j - r_k|^6$$

It is efficient to run spin-glass or TFI models on Rydberg atoms.

## Digital Quantum Computers

Example: Trapped ions (IonQ)

$$\text{MS}(\phi_0, \phi_1) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & -ie^{-i(\phi_0+\phi_1)} \\ 0 & 1 & -ie^{-i(\phi_0-\phi_1)} & 0 \\ 0 & -ie^{-i(\phi_0-\phi_1)} & 1 & 0 \\ -ie^{-i(\phi_0+\phi_1)} & 0 & 0 & 1 \end{bmatrix}$$

$$\text{MS}(\phi_0, \phi_1) = e^{-i\frac{\pi}{4}[\cos(\phi_0)X+\sin(\phi_0)Y] \otimes [\cos(\phi_1)X+\sin(\phi_1)Y]}$$

# Sparse matrices $\rightarrow$ Pauli strings?

- Finding the Pauli basis decomposition of a matrix is **HARD!**

$$H = \begin{bmatrix} L & 0 \\ 0 & 0 \end{bmatrix}, L = \begin{bmatrix} 1 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 1 \end{bmatrix}$$

```
from qiskit.quantum_info import SparsePauliOp
H = np.zeros((8, 8))
H[:5, :5] = L_chain
SparsePauliOp.from_operator(H)
```

```
SparsePauliOp(['III', 'IIX', 'IXX', 'IYY', 'XXX', 'XYY', 'YXY', 'YYX', 'ZII', 'ZIX', 'ZIZ', 'ZXX', 'ZYY', 'ZZI', 'ZZZ'],
               coeffs=[ 1. +0.j, -0.5 +0.j, -0.25+0.j, -0.25+0.j, -0.25+0.j,  0.25+0.j,
                        -0.25+0.j, -0.25+0.j,  0.75+0.j, -0.5 +0.j, -0.25+0.j, -0.25+0.j,
                        -0.25+0.j, -0.25+0.j, -0.25+0.j])
```

## Target Hamiltonian: $H$

Can we find a Hamiltonian  $H'$  (as a sum of Pauli strings), such that the information of  $e^{-iHt}$  can be extracted/recovered from  $e^{-iH't}$ ?

# II. Hamiltonian Embedding & Quantum Simulation

# Mathematical formulation

Suppose that  $H'$  is a  $q$ -qubit operator, let  $\mathcal{S}$  be an  $n$ -dim subspace of  $\mathbb{C}^{2^q}$ . We write

$$H' = \begin{bmatrix} H & R \\ R^\dagger & B \end{bmatrix},$$

with  $H = P_{\mathcal{S}} H' P_{\mathcal{S}}$ ,  $R = P_{\mathcal{S}} H' P_{\mathcal{S}^\perp}$ ,  $B = P_{\mathcal{S}^\perp} H' P_{\mathcal{S}^\perp}$ .

If  $R = 0$ , we have

$$e^{-iH't} = e^{-iHt} \oplus e^{-iBt} \implies \left( e^{-iH't} \right) \Big|_{\mathcal{S}} = e^{-iHt}$$

To simulate  $H$ , we only need to look at the **projection** of the *full evolution* in the subspace  $\mathcal{S}$ .

# Mathematical formulation

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If  $R \neq 0$ , in general

$$\left( e^{-iH't} \right) \Big|_{\mathcal{S}} \neq e^{-iHt}$$

There will be **leakage** from the subspace  $\mathcal{S}$  to  $\mathcal{S}^\perp$ .

A natural idea: we want to penalize the leakage!



# Shriefer-Wolff transformation

Suppose that  $H'$  is a  $q$ -qubit operator, let  $\mathcal{S}$  be an  $n$ -dim subspace of  $\mathbb{C}^{2^q}$ . We write

$$H' = \begin{bmatrix} H & R \\ R^\dagger & B \end{bmatrix},$$

$$\text{with } H = P_{\mathcal{S}} H' P_{\mathcal{S}}, \quad R = P_{\mathcal{S}} H' P_{\mathcal{S}^\perp}, \quad B = P_{\mathcal{S}^\perp} H' P_{\mathcal{S}^\perp}.$$

Assume  $\Delta = \lambda_{\min}(B) - \lambda_{\max}(H) > 0$ . Define  $\kappa = \|R\|/\Delta$ .

When  $\kappa < 1/2$ , there exists a **Shriefer-Wolff transformation**  $U$  such that  $U^\dagger H' U$  is block-diagonal, and

$$H_{\text{eff}} = (U^\dagger H' U) \big|_{\mathcal{S}} = H + O\left(\frac{\|R\|^2}{\Delta}\right)$$

**Theorem. (Quantum simulation by Hamiltonian embedding)**

$$\text{For any } t \geq 0, \text{ we have } \left\| (e^{-iH't}) \big|_{\mathcal{S}} - e^{-iHt} \right\| \leq (8\kappa^2 + 4\sqrt{2}\kappa) \|R\| t.$$

# Explicit constructions

$$L = \begin{bmatrix} 1 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 1 \end{bmatrix} = \text{The graph Laplacian of } \begin{array}{c} \bigcirc - \bigcirc - \bigcirc - \bigcirc - \bigcirc \end{array}$$

**Penalty Hamiltonian:**  $H_{\text{pen}} = Z_1 + Z_4 + \sum_{j=1}^3 Z_j Z_{j+1}.$

Ground-energy subspace of  $H_{\text{pen}}$ :  $\mathcal{S} = \text{span}\{ |0101\rangle, |0100\rangle, |0110\rangle, |0010\rangle, |1010\rangle \}$

**Encoding Hamiltonian:**  $Q = \hat{n}_1 + \hat{n}_4 - \sum_{j=1}^4 X_j$ , where  $\hat{n}_j = \frac{1}{2} (I - Z_j).$

Consider  $H' = Q + gH_{\text{pen}}.$

$$H' = \begin{bmatrix} L & R \\ R^\dagger & B \end{bmatrix} + \text{const.}$$

(1) For sufficiently large  $g > 0$ ,  $0 < \kappa < 1/2$ . Recall that  $\kappa = \|R\|/\Delta$ ,  $\Delta = \lambda_{\min}(B) - \lambda_{\max}(L) > 0$ .

( $g > 0$  is called the **penalty** coefficient  $\rightarrow$  penalizing the leakage from  $\mathcal{S}$  to  $\mathcal{S}^\perp$ )

(2) Easy to check:  $Q|_{\mathcal{S}} = L$ , or  $H'|_{\mathcal{S}} = L + \text{const.}$

# Explicit constructions

$$L = \begin{bmatrix} 1 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 2 \end{bmatrix} \quad \hookrightarrow \quad H' = \left( -\frac{1}{2}Z_1 - \frac{1}{2}Z_4 - \sum_{j=1}^4 X_j \right) + g \left( Z_1 + Z_4 + \sum_{j=1}^4 Z_j Z_{j+1} \right)$$

- We call this the **antiferromagnetic embedding** of the graph Laplacian  $L$ .
- “Antiferromagnetic”  $\rightarrow$  the *antiferromagnetic* ordering in the quantum Ising model.
- **Analog quantum simulation:** Rydberg atoms (e.g., QuEra) or spin-glass simulators (e.g., D-Wave).
- **Digital quantum simulation:**  $H_{\text{pen}}$  is **fast-forwardable**  $\rightarrow$  moving to the interaction picture and run **qDRIFT**.



# Rules for building Hamiltonian embeddings

1. (Addition) For  $j = 1, 2$ , suppose  $H_j^{\text{ebd}} = gH_j^{\text{pen}} + Q_j$  is an embedding of  $H_j$  with the same embedding subspace  $\mathcal{S}$ . Then,  $H^{\text{ebd}} = gH^{\text{pen}} + Q_1 + Q_2$  is an embedding of  $H = H_1 + H_2$  with embedding subspace  $\mathcal{S}$ .

2. (Multiplication) Suppose  $H^{\text{ebd}} = gH^{\text{pen}} + Q$  is an embedding of  $H$  with embedding subspace  $\mathcal{S}$ . Then, for any real number  $\alpha$ ,  $H'' = gH^{\text{pen}} + \alpha Q$  is an embedding of  $\alpha H$ .

3. (Composition) For  $j = 1, 2$ , suppose that the  $q_j$ -qubit operator  $H_j^{\text{ebd}} = gH_j^{\text{pen}} + Q_j$  is an embedding of  $H_j$  with the embedding subspace  $\mathcal{S}_j$  being the ground-energy subspace of  $H_j^{\text{pen}}$ . Then,  $H^{\text{ebd}} = gH^{\text{pen}} + Q$  with

$$H^{\text{pen}} = H_1^{\text{pen}} \otimes \mathbb{I} + \mathbb{I} \otimes H_2^{\text{pen}}, \quad Q = Q_1 \otimes \mathbb{I} + \mathbb{I} \otimes Q_2 \quad (\text{A.14})$$

is an embedding of  $H = H_1 \otimes \mathbb{I} + \mathbb{I} \otimes H_2$  with embedding subspace  $\mathcal{S} = \mathcal{S}_1 \otimes \mathcal{S}_2$ .

4. (Tensor product) For  $j = 1, 2$ , suppose that  $H_j^{\text{ebd}} = gH_j^{\text{pen}} + Q_j$  is an embedding of  $H_j$  with the embedding subspace  $\mathcal{S}_j$  being the ground-energy subspace of  $H_j^{\text{pen}}$ . Then,  $H^{\text{ebd}} = gH^{\text{pen}} + Q_1 \otimes Q_2$  with

$$H^{\text{pen}} = H_1^{\text{pen}} \otimes \mathbb{I} + \mathbb{I} \otimes H_2^{\text{pen}} \quad (\text{A.15})$$

is an embedding of  $H = H_1 \otimes H_2$  with embedding subspace  $\mathcal{S} = \mathcal{S}_1 \otimes \mathcal{S}_2$ .

# Hamiltonian embedding of sparse matrices

1. **Band matrix:** non-zero entries are confined to a diagonal band

- Unary embedding, antiferromagnetic embedding
- Embedding Hamiltonian is  $\max(2, d)$ -local,  $d$  = bandwidth

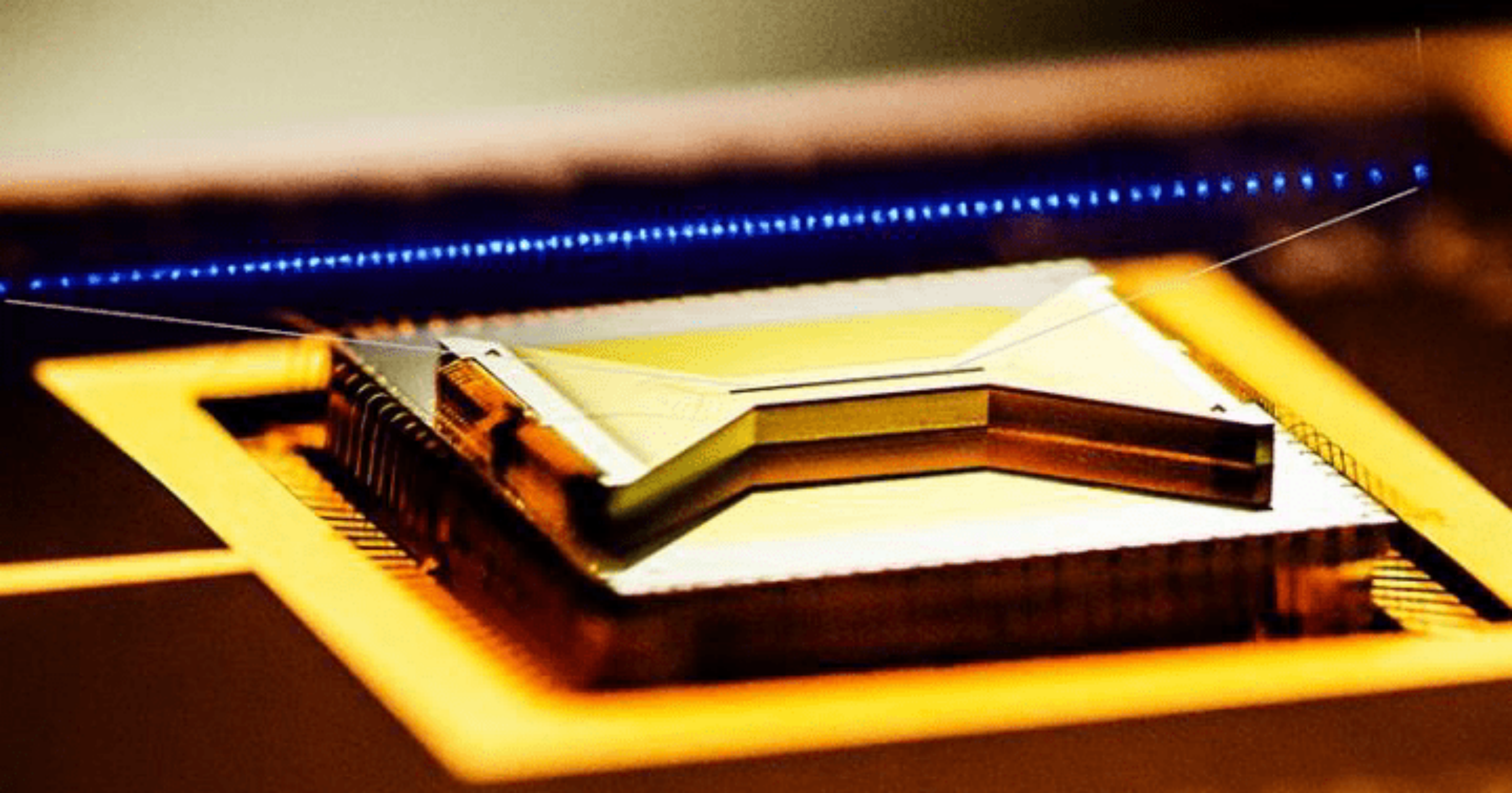
2. **Banded circulant matrix:**

- Circulant unary embedding, circulant antiferromagnetic embedding

3. **Arbitrary sparse matrix:**

- One-hot embedding (with or without penalty Hamiltonian)

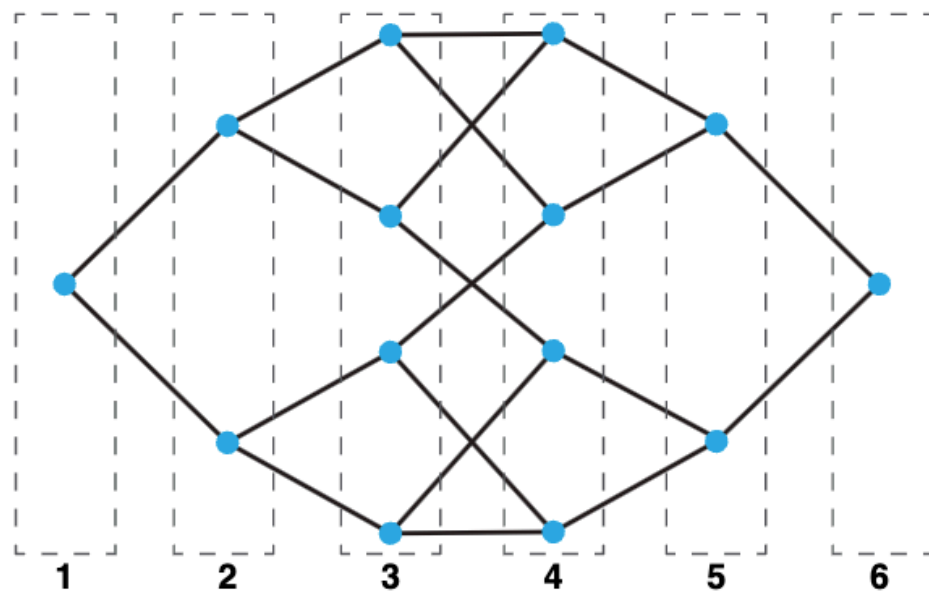




# Real-Machine Experiments

# Traversing the glued trees graph

A 6-layer Glued Trees Graph



Childs et. al. “Exponential algorithmic speedup by quantum walk”. [arXiv:quant-ph/0209131](https://arxiv.org/abs/quant-ph/0209131)

**Task:** to simulate  $e^{-iLt}$ ,  $L$  = graph Laplacian

**Initial state:**  $|\text{node1}\rangle$

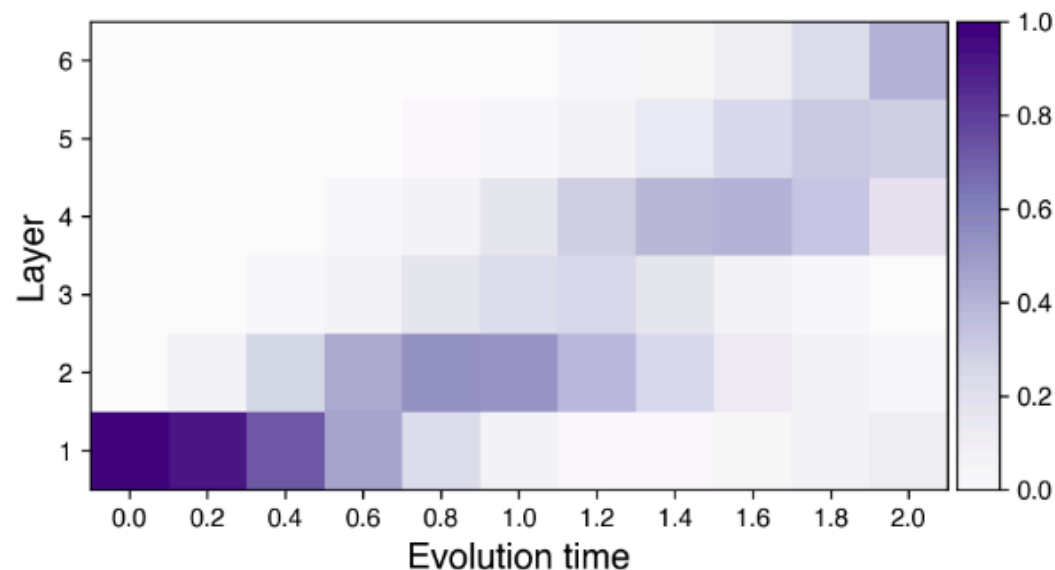
**Embedding scheme:** penalty-free one-hot

**Device:** IonQ Aria-1

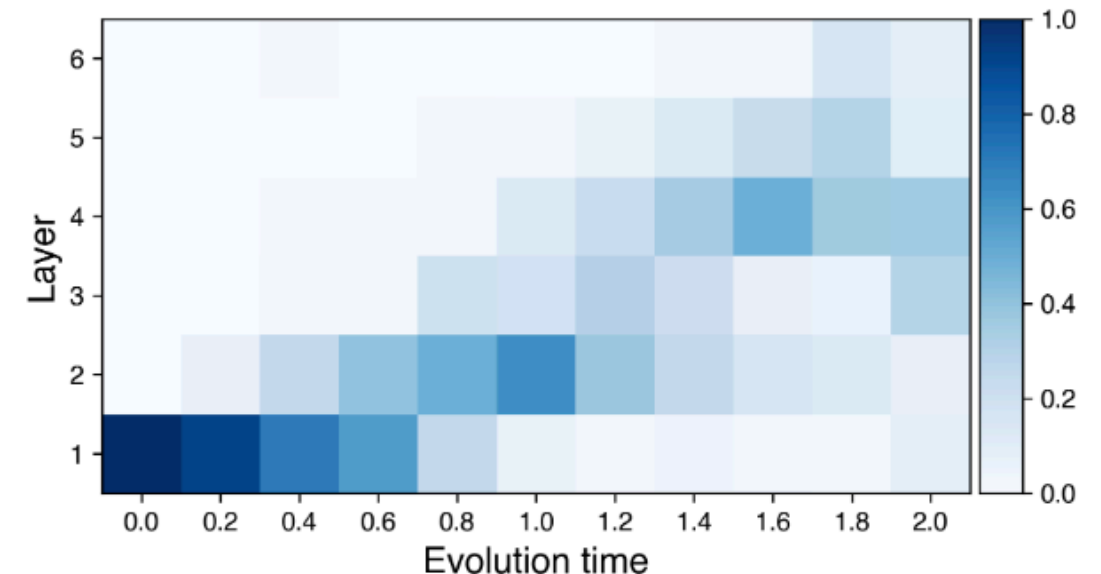
**Resources:** 14 qubits, 1 single-qubit gate, 120 two-qubit gates.

**Budget:** < \$100.

Quantum Walk (numerical simulation)



Quantum Walk (IonQ, one-hot)



Heat map on layers over evolution time



# Simulating real-space quantum dynamics

**Task:** to simulate the following 1D Schrodinger equation (1 Bosonic mode),

$$i\frac{\partial}{\partial t}\Psi(t, x) = \left[ -\frac{1}{2}\frac{\partial^2}{\partial x^2} + \left( \frac{1}{2}ax^2 + bx \right) \right] \Psi(t, x) \quad \text{We choose } a = 2, \quad b = -1/2.$$

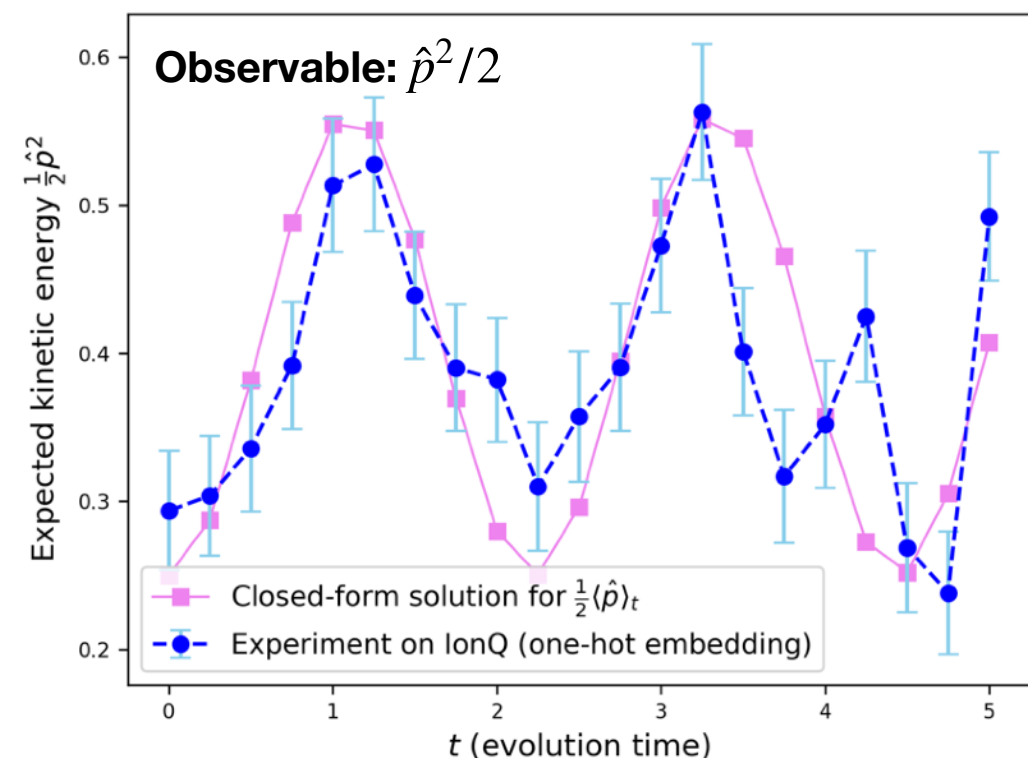
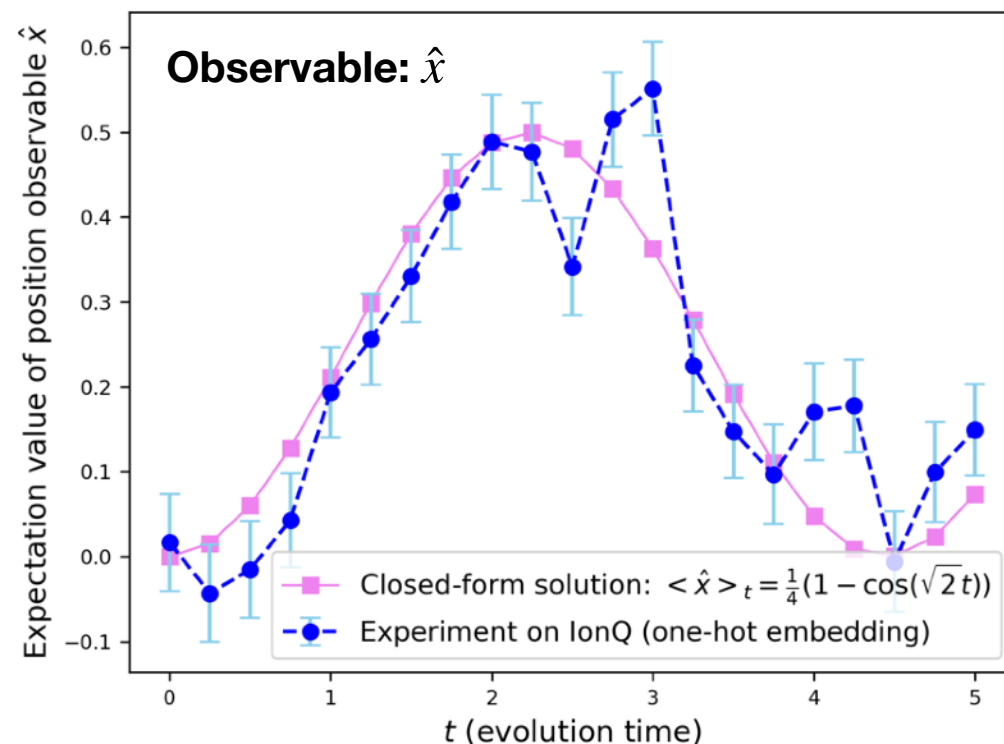
with Gaussian initial state  $\Psi(0, x) = \left( \frac{1}{2\pi\sigma^2} \right)^{1/4} e^{-\frac{x^2}{4\sigma^2}}$

**Embedding scheme:** penalty-free one-hot

**Device:** IonQ Aria-1

**Resources:** 5 qubits, 1 single-qubit gate, 154 two-qubit gates.

**Budget:** < \$1,300.

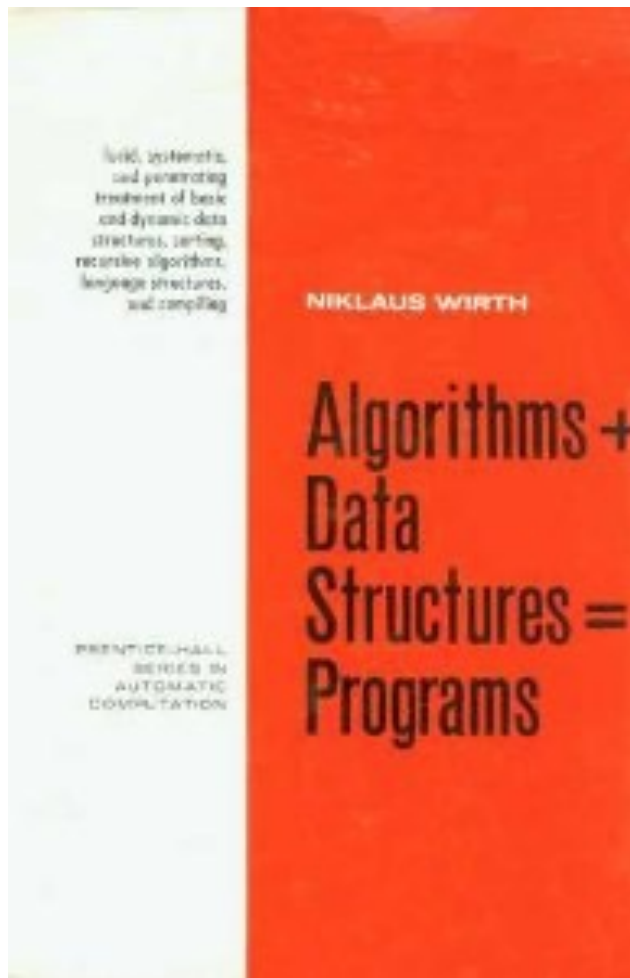




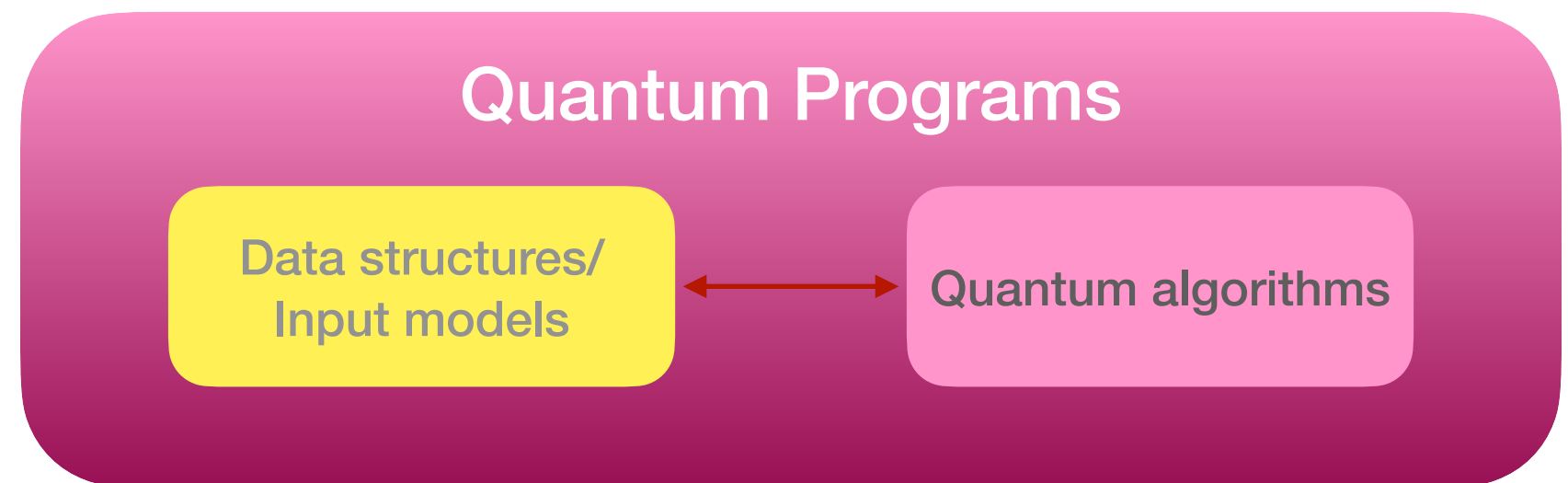
# II. Hamiltonian-oriented Quantum Algorithms

**A New Paradigm in Quantum Algorithm Design**

# Programs = Algorithms + Data Structures



By Niklaus Wirth, 1976.



\* Output/measurement is also important in quantum algorithm design!

- **Algorithms** and **data structures (input models)** are *inherently* related!
- **Classical scientific computing** → “matrix-oriented” algorithms (using **numerical linear algebra** as major subroutines)
- **Block-encodings** → QSP/QSVT
- **Hamiltonian embeddings** → **Hamiltonian-oriented Quantum Algorithms** (using *Hamiltonian simulation* as a major subroutine)

# Hamiltonian-oriented Quantum Algorithms

\* Feasible on NISQ devices and even early fault-tolerant devices

- Quantum adiabatic algorithms (for discrete optimization)
- (Continuous-time) quantum walk
- Phase estimation (Trotterized controlled evolution, [arXiv:1512.06860](#))
- Linear combination of Hamiltonian Simulations (non-unitary dynamics, [arXiv:2303.01029](#))
- Schrodingerization (non-unitary PDEs, [arXiv:2212.13969](#))
- Quantum Hamiltonian Descent (for continuous optimization, [arXiv:2303.01471](#))
- Pulse-level variational quantum algorithms & quantum optimal control ([arXiv:2210.15812](#))
- And many more ...

# SimuQ: SIMUlation language for Quantum

arXiv:2303.02775. Yuxiang Peng, Jacob Young, Pengyu Liu (CMU), and Xiaodi Wu. Accepted by POPL 2024.

