

A Dynamical Point of View on Quantum Optimization

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“Quantum Hamiltonian Descent”, arXiv:2303.01471

“A quantum-classical performance separation in nonconvex optimization”, manuscript in preparation



JOINT CENTER FOR
QUANTUM INFORMATION
AND COMPUTER SCIENCE

Quantum computing 101

Classical (digital) Computer

- Unit of information: bit ($b \in \{0,1\}$)
- Classical information \rightarrow bitstrings (e.g., 10001001)
- Computation = manipulation of finite-length bitstrings
- Readout/Return: bitstrings

Quantum Computer

- Unit of information: quantum bit, or **qubit**
 $|\psi\rangle = a|0\rangle + b|1\rangle, a, b \in \mathbb{C}, |a|^2 + |b|^2 = 1$
- Quantum information \rightarrow quantum state (i.e., superposition of $|bitstring\rangle$)

$$|GHZ\rangle = \frac{1}{\sqrt{2}} (|000\rangle + |111\rangle)$$

- Computation = unitary (i.e., preserving 2-norm) operations on quantum states
- Readout: quantum measurement \rightarrow a sample of bitstrings

Measure GHZ state:
000: 50%, 111: 50%, others: 0%

Digital Quantum Computers

Use quantum gates & quantum circuits, not ready in 5-10 years.
Must do error correction.

Analog Quantum Computers

Solving problems by emulating a real quantum system. Ready & easy to scale. No error correction.

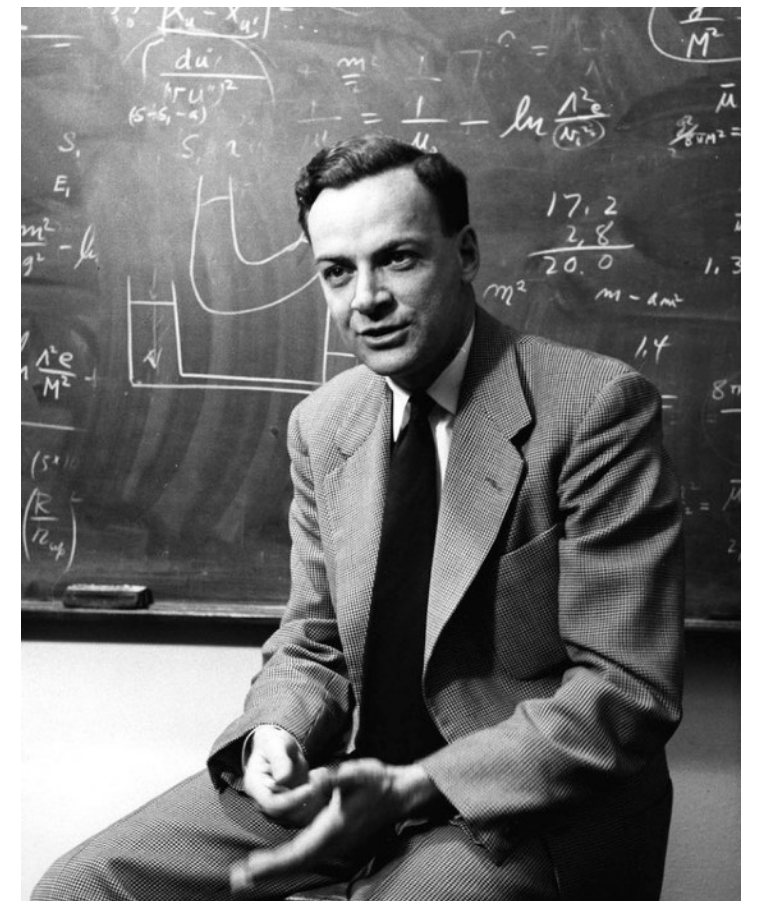
Hamiltonian simulation

- **Quantum simulation** (or **Hamiltonian simulation**): a prominent application of quantum computers. **Exponential** quantum speedup expected.

$$i\frac{\partial}{\partial t}\psi = \hat{H}(t)\psi$$

$\hat{H}(t)$: quantum Hamiltonian (Hermitian/self-adjoint)

- **Feynman** (“*Simulating Physics with Computers*”, 1982): to simulate quantum systems, we would need to build quantum computers.
- **Applications**: quantum chemistry, quantum field theory, condensed matter physics, numerical optimization, etc.



Overview

- Part I: Quantum Hamiltonian Descent (QHD)
- Part II: A Quantum-Classical Performance Separation
- Part III: Implementation on Quantum Computers

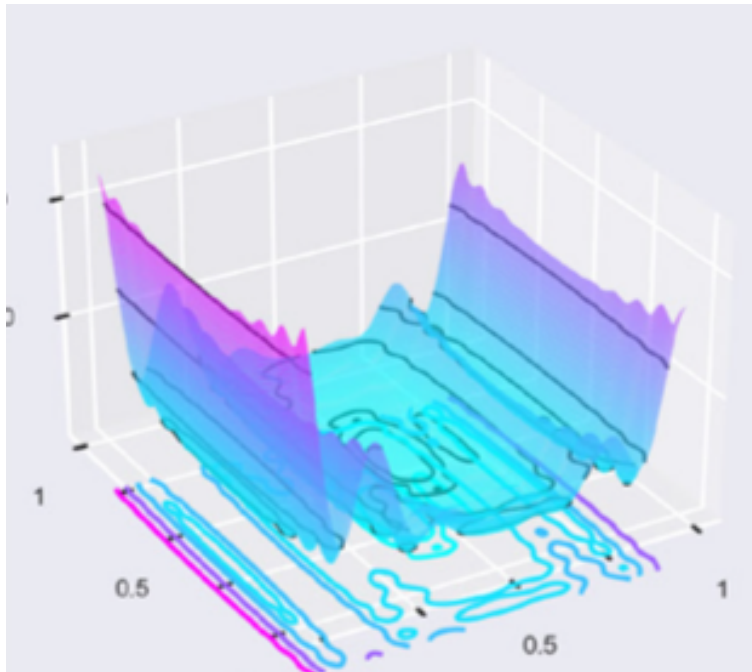
Part I:

Quantum Hamiltonian Descent

Problem formulation

Continuous optimization

$$\min_{x \in \mathbb{R}^d} f(x)$$



- Important in practice: machine learning, operations research, scientific computing, etc.
- A **challenging problem** for quantum: different nature in (quantum) algorithm design, requires new mathematical tools to prove convergence.
- **Opportunity:** new *primitives* of quantum speedups!

Classical algorithms: Gradient Descent (GD)

- Standard GD: $x_{k+1} = x_k - s \nabla f(x_k)$.
- Nesterov's accelerated GD: $x_{k+1} = y_k - s \nabla f(y_k)$, $y_{k+1} = x_{k+1} + \frac{k}{k+3} (x_{k+1} - x_k)$.

A Lagrangian formulation of accelerated methods

Su, Boyd, & Candes
NeurIPS 2014

$$\begin{aligned}x_k &= y_{k-1} - s \nabla f(y_{k-1}), \\ y_k &= x_k + \frac{k-1}{k+2}(x_k - x_{k-1}).\end{aligned}$$

Nesterov's accelerated GD



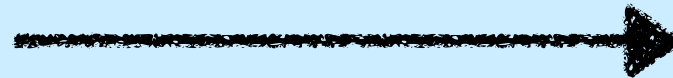
$$\ddot{X} + \frac{3}{t}\dot{X} + \nabla f(X) = 0,$$

ODE model

Wibisono, Wilson, & Jordan
PNAS Nov 2016, 113 (47)

$$\begin{aligned}x_k &= y_{k-1} - s \nabla f(y_{k-1}), \\ y_k &= x_k + \frac{k-1}{k+2}(x_k - x_{k-1}).\end{aligned}$$

Nesterov's accelerated GD



$$\mathcal{L}(t, \dot{X}, X) = t^3 \left(\frac{1}{2} |\dot{X}|^2 - f(X) \right)$$

Lagrangian formulation

$$H(t, X, P) = \frac{1}{2t^3} |P|^2 + t^3 f(X)$$

Hamiltonian formulation

- Accelerated GD can be modeled by [accelerated gradient flows](#).
- GD algorithms can be generated by discretizing continuous-time dynamics.

Accelerated Hamiltonian flows: classical v.s. quantum

We can make the *classical* dynamics *quantum*!!

**Classical
Hamiltonian
Systems**

*Path integral formulation of
Quantum Mechanics*

Due to Feynman

**Quantum
Hamiltonian
Evolution**

$$H(t, X, P) = \frac{1}{2t^3} |P|^2 + t^3 f(X)$$



$$\hat{H}(t) = \frac{1}{t^3} \left(-\frac{1}{2} \Delta \right) + t^3 f(x)$$

Bregman Lagrangian

$$\mathcal{L}(X, \dot{X}, t) = e^{-\varphi_t} \left(\frac{1}{2} |\dot{X}|^2 \right) - e^{\chi_t} f(X)$$

Quantum propagator

$$K(b, t_b; a, t_a) \propto \sum_{\text{paths from a to b}} e^{\frac{i}{\hbar} S[X(t)]}$$

Schrodinger equation (QHD)

$$i\hbar \frac{\partial \Psi}{\partial t} = \left[e^{\varphi_t} \left(-\frac{\hbar^2}{2} \nabla^2 \right) + e^{\chi_t} f(x) \right] \Psi$$

Infinitesimal expansion

$$\Psi(x, t + \epsilon) \propto \int \exp \left\{ \frac{i}{\hbar} \epsilon \mathcal{L} \left(\frac{x+y}{2}, \frac{x-y}{\epsilon}, t \right) \right\} \Psi(y, t) dy$$

Quantum Hamiltonian Descent

$$H(t, X, P) = e^{\varphi_t} \left(\frac{1}{2} |P|^2 \right) + e^{\chi_t} f(X)$$

Accelerated gradient flows

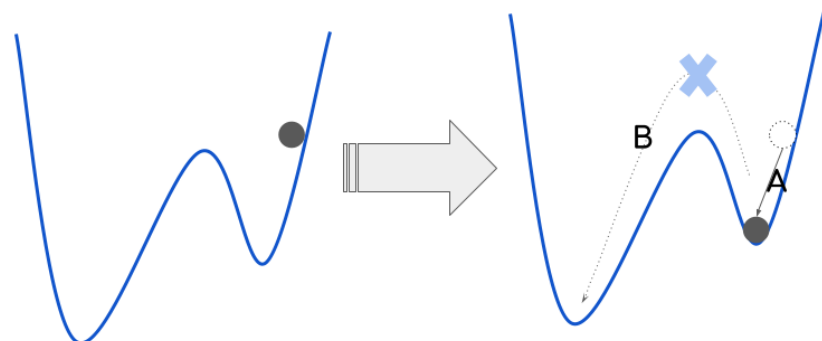


$$\hat{H}(t) = e^{\varphi_t} \left(-\frac{1}{2} \nabla^2 \right) + e^{\chi_t} f(x)$$

Quantum Hamiltonian Descent (QHD)

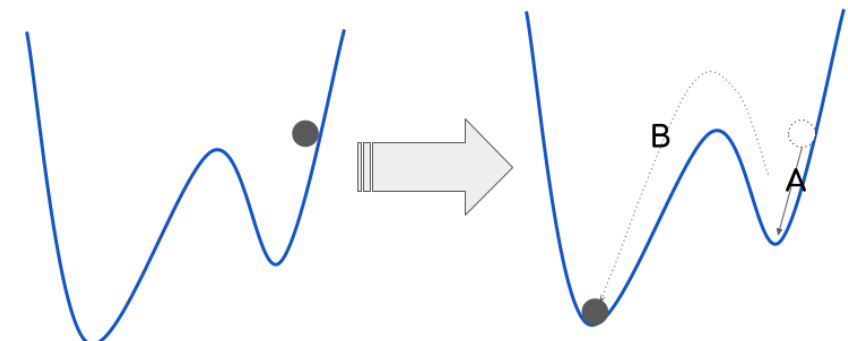
- This is a quantum dynamical system with *damping* ($e^{\varphi_t - \chi_t} \rightarrow 0$).
- The energy damping allows the quantum system to *converge* to a low-energy configuration (i.e., minimizing the objective f).
- **Intuition:** “Path integral” of classical GD \rightarrow does it help **nonconvex** optimization?

Classical Gradient methods



Optimization paths = A
B is prohibited by classical mechanics.

Quantum Hamiltonian Descent (QHD)



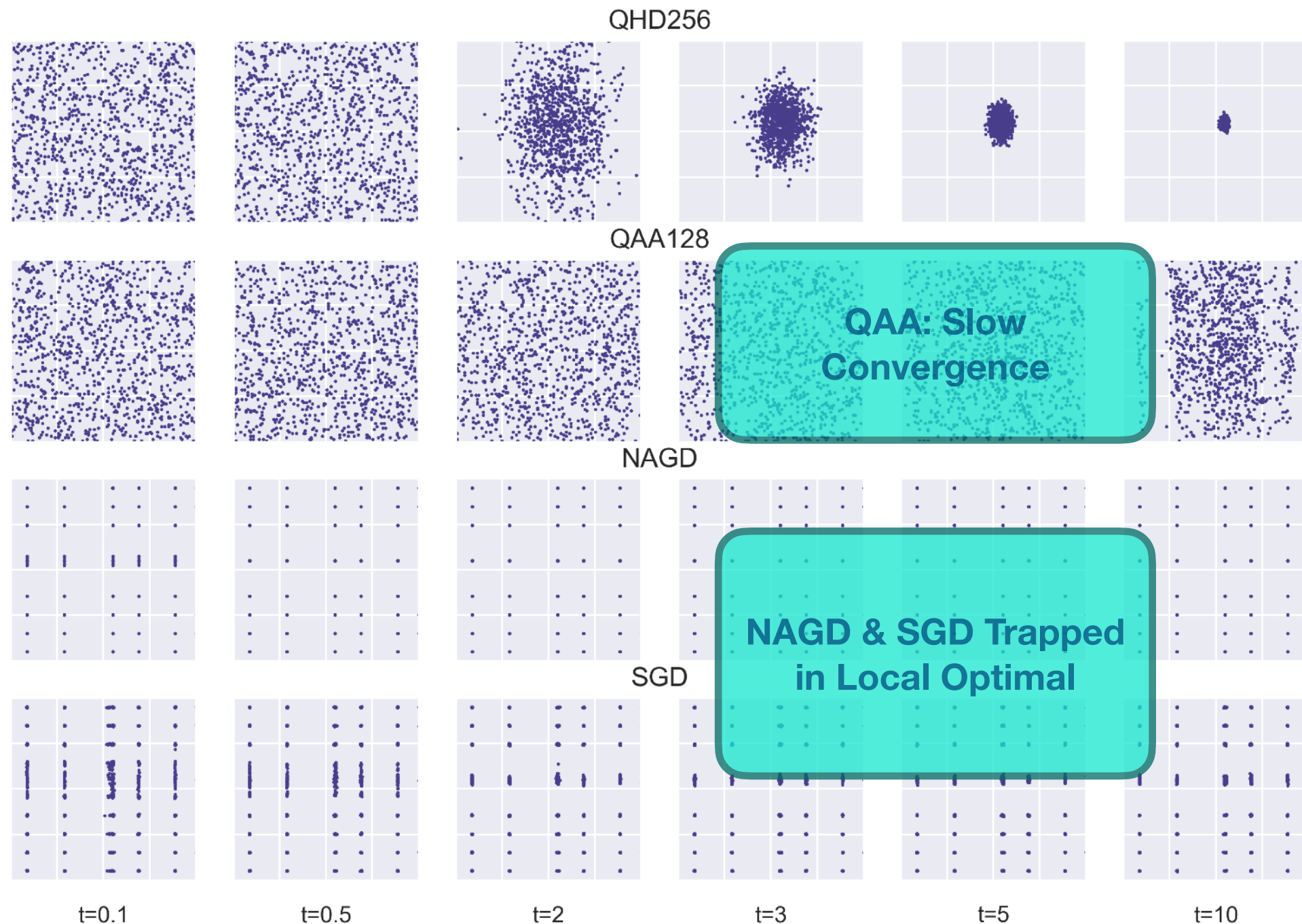
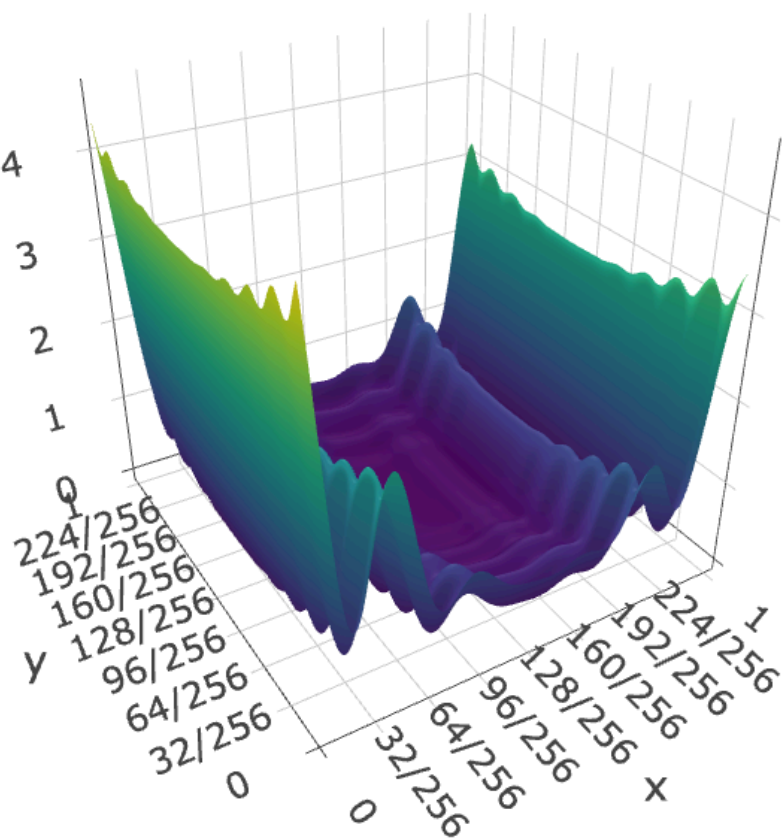
Optimization paths = $\int_{\text{all possible paths}}$
A, B both contribute to quantum optimization paths.

Quantum Hamiltonian Descent - numerical example

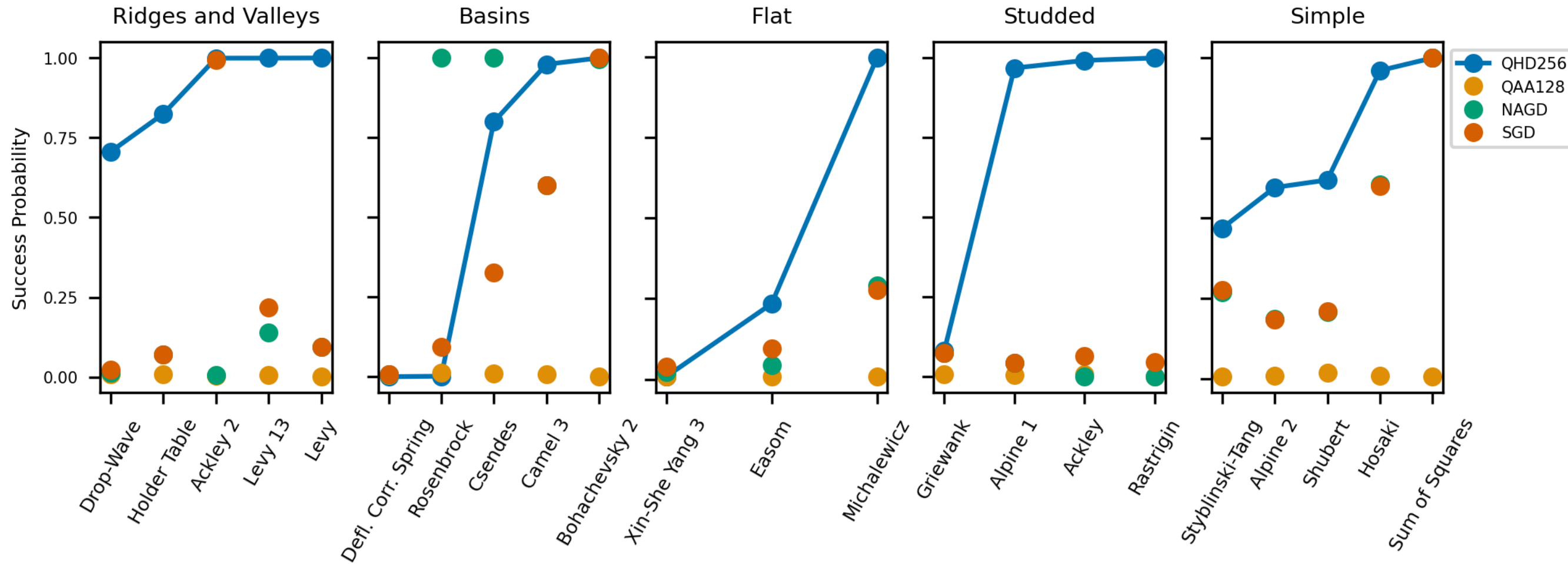
QAA: quantum adiabatic algorithm. **NAGD:** Nesterov's accelerated gradient descent.

SGD: stochastic gradient descent.

**QHD showcase:
the Levy function.**



Behavior of QHD on 2D benchmark



Measure in Success Probability for some given time:

the probability of generating a solution within radius r of the optimal solution from random initial points.

QAA: use standard *linear interpolation* in the adiabatic method (i.e., $H_T(t) = (1 - t/T)H_0 + (t/T)H_1$) and simulate $H_T(t)$ for $T=10$ over a **128-by-128** grid on $[0,1] \times [0,1]$.

QHD: simulate QHD Hamiltonian with Nesterov's AGD parameters (i.e., $H(t) = -\frac{1}{2t^3}\nabla^2 + t^3f(x)$) from $t = 0$ to 10 over a **256-by-256** grid on $[0,1] \times [0,1]$.

Convergence guarantee: convex optimization

$$H(t, P, X) = e^{\alpha_t} \left(-\frac{e^{-\gamma_t}}{2} P^2 + e^{\gamma_t + \beta_t} f(X) \right) \longrightarrow \frac{d}{dt} (e^{-\alpha_t} \dot{X}_t + X_t) = e^{\alpha_t + \beta_t} \nabla f(X_t)$$

Theorem [Wibisono, Wilson, Jordan (2016)]

Assume $f(x)$ is continuously differentiable and convex, and $\dot{\beta}_t \leq e^{\alpha_t}$, $\dot{\gamma}_t = e^{\alpha_t}$ (aka, *ideal scaling condition*). Then, the solution $X(t)$ satisfies

$$f(X_t) - f(x^*) \leq O(e^{-\beta_t}).$$

$$\hat{H}(t) = e^{\alpha_t} \left(-\frac{e^{-\gamma_t}}{2} \Delta + e^{\gamma_t + \beta_t} f(x) \right) \longrightarrow \frac{\partial}{\partial t} \Psi(t, x) = \hat{H}(t) \Psi(t, x)$$

Theorem [Leng, Hickman, Li, Wu (2023)]

Assume $f(x)$ is continuously differentiable and convex, and $\dot{\beta}_t \leq e^{\alpha_t}$, $\dot{\gamma}_t = e^{\alpha_t}$ (aka, *ideal scaling condition*). Define $\mathbb{E}[f]_{\sim \Psi(t)} = \langle \Psi(t) | f | \Psi(t) \rangle = \int f | \Psi(t) |^2 dx$, then we have

$$\mathbb{E}[f]_{\sim \Psi(t)} - f(x^*) \leq O(e^{-\beta_t}).$$

A Lyapunov function approach

Theorem [Leng, Hickman, Li, Wu (2023)]

Assume $f(x)$ is continuously differentiable and convex, and $\dot{\beta}_t \leq e^{\alpha_t}$, $\dot{\gamma}_t = e^{\alpha_t}$ (aka, ideal scaling condition). Define $\mathbb{E}[f]_{\sim \Psi(t)} = \langle \Psi(t) | f | \Psi(t) \rangle = \int f | \Psi(t) |^2 dx$, then we have

$$\mathbb{E}[f]_{\sim \Psi(t)} - f(x^*) \leq O(e^{-\beta_t}).$$

- We construct a Lyapunov function ($\langle O \rangle_t = \langle \Psi_t | O | \Psi_t \rangle$, $\hat{p} = -i \nabla$):

$$\mathcal{W}(t) = \langle \hat{J}^2 / 2 \rangle_t + e^{\beta_t} \langle f \rangle_t$$

$$\hat{J} := e^{-\gamma_t} \hat{p} + \hat{x}.$$

- We can prove this Lyapunov function is **non-increasing in t**.
- Therefore, we have $e^{\beta_t} \mathbb{E}[f]_{\sim \Psi_t} \leq \mathcal{W}(t) \leq \mathcal{W}(0)$, $\mathbb{E}[f]_{\sim \Psi_t} \leq \mathcal{W}(0) e^{-\beta_t} \leq O(e^{-\beta_t})$.

- A sanity check: we have the same converge rate as classical [WWJ16].
- Nesterov's GD: $e^{\alpha_t} = 2/t$, $e^{\beta_t} = e^{\gamma_t} = t^2$, implies the convergence rate $O(t^{-2})$.
- QHD is more stable & robust with larger time discretization steps.
- Question: can we achieve $O(e^{-\sqrt{\beta_t}})$ convergence rate for β -strongly convex f ?

Convergence guarantee: nonconvex optimization

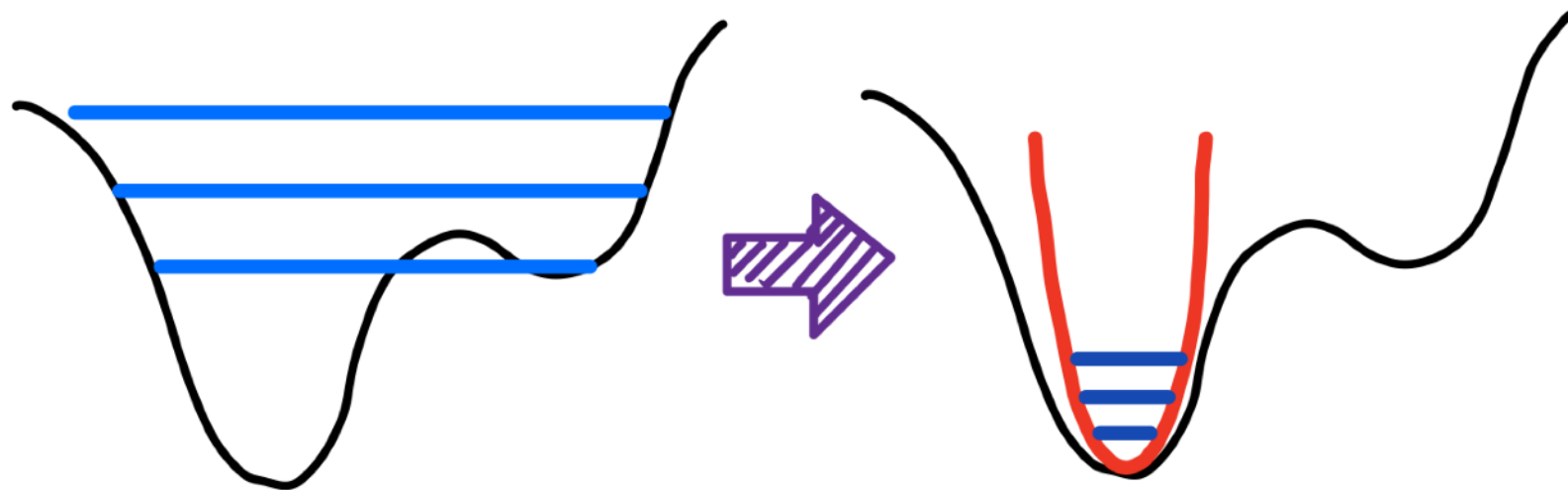
$$\hat{H}(t) = \left(-\frac{e^{\varphi_t}}{2} \Delta + e^{\chi_t} f(x) \right) \longrightarrow \frac{\partial}{\partial t} \Psi(t, x) = \hat{H}(t) \Psi(t, x)$$

Theorem (informal) [Leng, Hickman, Li, Wu (2023)]

Suppose $f(x)$ be smooth, unbounded at infinity, and has a unique non-degenerate global minimum x^* . Let the initial wave Ψ_0 be in the low-energy subspace of $H(0)$ and $e^{\varphi_t - \chi_t} \rightarrow 0$, $|\dot{\varphi}_t|, |\dot{\chi}_t| \ll 1$ (i.e., $H(t)$ is slow-varying), then

$$\lim_{t \rightarrow \infty} \mathbb{E}[f]_{\sim \Psi_t} = f(x^*).$$

- Slow-varying $H(t)$ \rightarrow the wave function $\Psi(t)$ stays in the low-energy subspace for all t .
- The low-energy subspace of $H(t)$ will *migrate* to the global minimum of f as $t \rightarrow \infty$.



Quantum Hamiltonian Descent: complexity analysis

1. Prepare an initial state $|\psi_0\rangle$.
2. Simulating the Schrodinger equation: $i\partial_t\psi = \hat{H}(t)\psi$, $\psi(0) = \psi_0$.
3. At time $t = T$, *measure* the final state $|\Psi_T\rangle$ (i.e., to sample from the corresponding distribution $|\Psi(T)|^2$).
4. Ideally, the measurement results will cluster around the global minimizer of f .

$$\hat{H}(t) = e^{\varphi_t} \left(-\frac{1}{2} \nabla^2 \right) + e^{\chi_t} f(x)$$

Runtime of QHD

QHD is comparable to GD in terms of simplicity and resource cost:

- Running QHD = time-dependent Hamiltonian simulation [Childs, Leng, Li, Liu & Zhang, 2022]
- The runtime of QHD is $\tilde{O}(dT)$ (d : dimension of f , T : total evolution time).
- Resource cost (on quantum computer) is **comparable** to that of classical GD.
- Resource cost (on classical computer) is **exponential** in d .

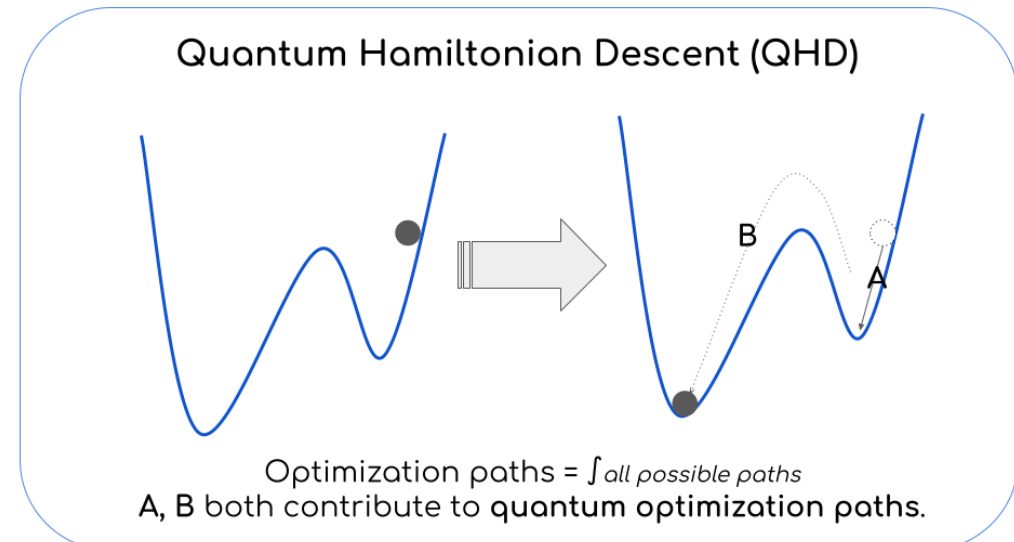
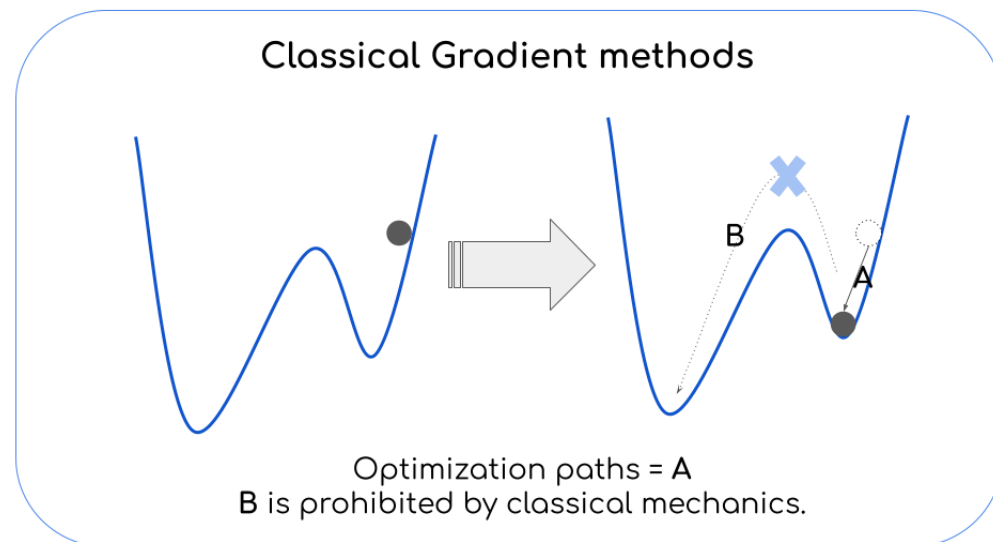
Therefore, *QHD is a quantum-upgraded version of GD*. We will demonstrate how powerful QHD could be by itself. One could also build on top of QHD like what we've done with GD.

Part II:

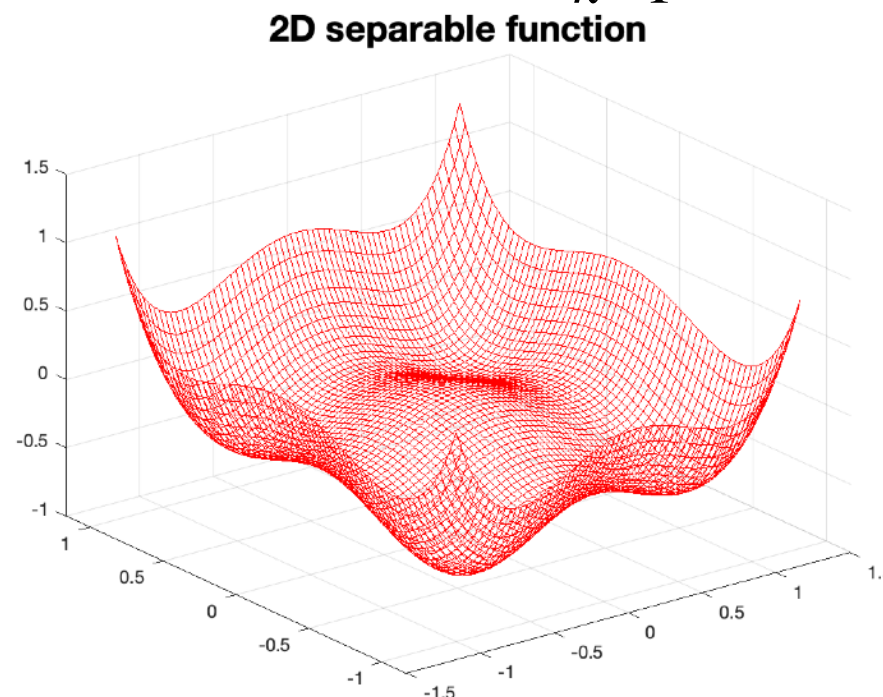
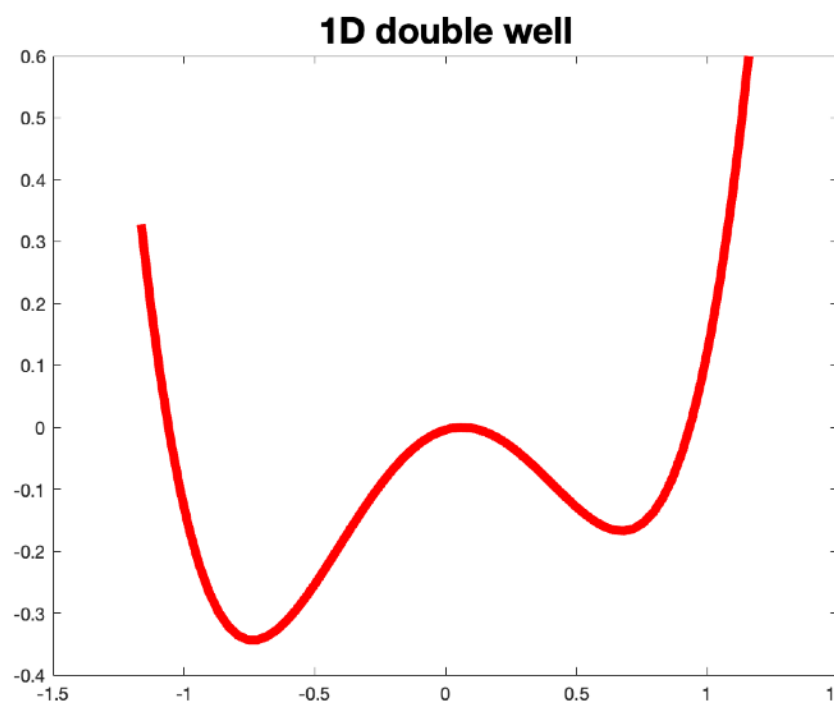
A Quantum-Classical Performance Separation

Quantum Hamiltonian Descent — practical setting?

1-dim nonconvex model problem: double-well potential — $f(x)$



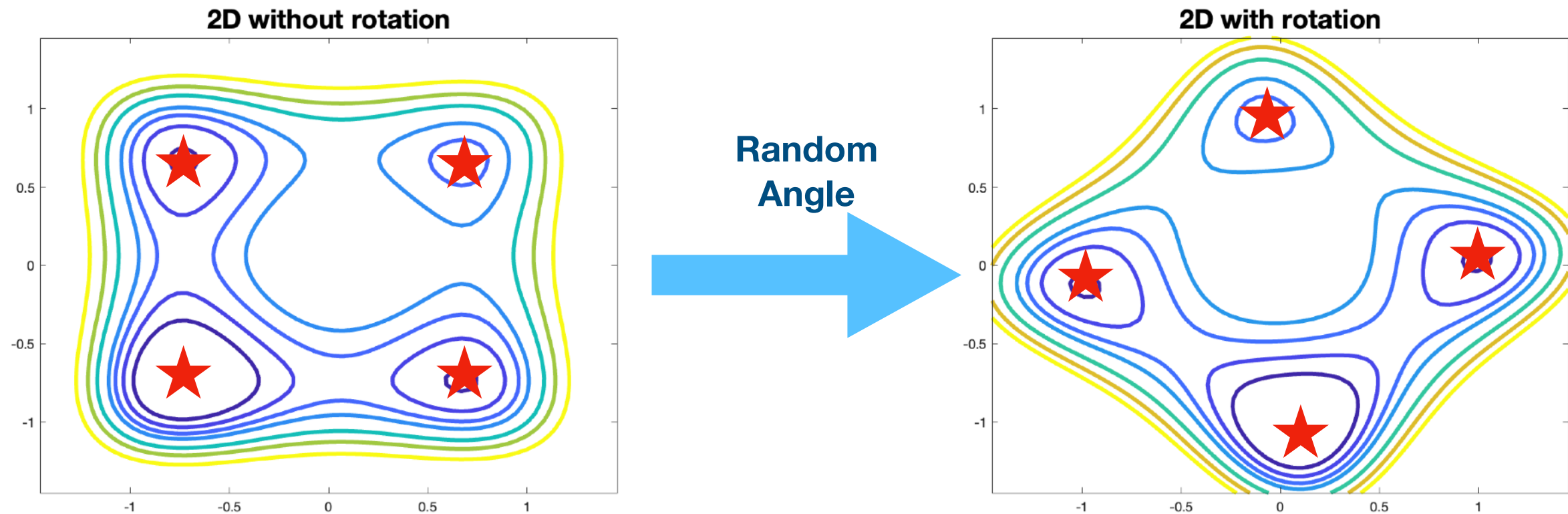
d-dim objective function: $F(x_1, \dots, x_d) = \sum_{k=1}^d f(x_k),$



- $F(x)$ has 2^d local minima! (Only **one unique** global minimum).
- $F(x)$ is **separable**: not difficult for classical algorithms if the closed-form formula is given.

Construction of the optimization instances

Our instances = d-dim separable functions + random rotation



$$F(x_1, \dots, x_d) = \sum_{k=1}^d f(x_k) \mapsto F_U(x) = F(Ux)$$

- $F_U(x)$ still has 2^d local minima, with a **unique** global minimum.
- $F(x)$ is **non-separable**: difficult to recover the rotation even with the closed-form formula!

QHD: a polynomial-time quantum algorithm

Given an optimization problem $f(x): \mathbb{R}^d \rightarrow \mathbb{R}$ with a unique global minimizer x^* .

x is a **δ -approximate** solution if $\|x - x^*\| < \delta$.

Theorem (Informal) [Leng, Zheng, Wu (2023)]

Let $f(x): \mathbb{R} \rightarrow \mathbb{R}$ be a double-well potential function. Define $F_U(x) = F(Ux)$ where

$F(x) = \sum_{k=1}^d f(x_k)$ and U is an arbitrary orthogonal matrix. For any small $\delta > 0$, QHD can

produce a δ -approximate solution with probability at least $2/3$ using

- $\tilde{\mathcal{O}}(d^3/\delta^2)$ quantum queries to F_U , and
- $\tilde{\mathcal{O}}(d^4/\delta^2)$ additional 1- and 2-qubit gates.

Manuscript in preparation.

- The QHD Hamiltonian (more precisely, the Laplacian operator) is rotationally invariant.
- The ground state of the QHD Hamiltonian is the *vehicle* of quantum optimization.
- We use an adiabatic theorem for unbounded Hamiltonian.

A quantum-classical performance separation

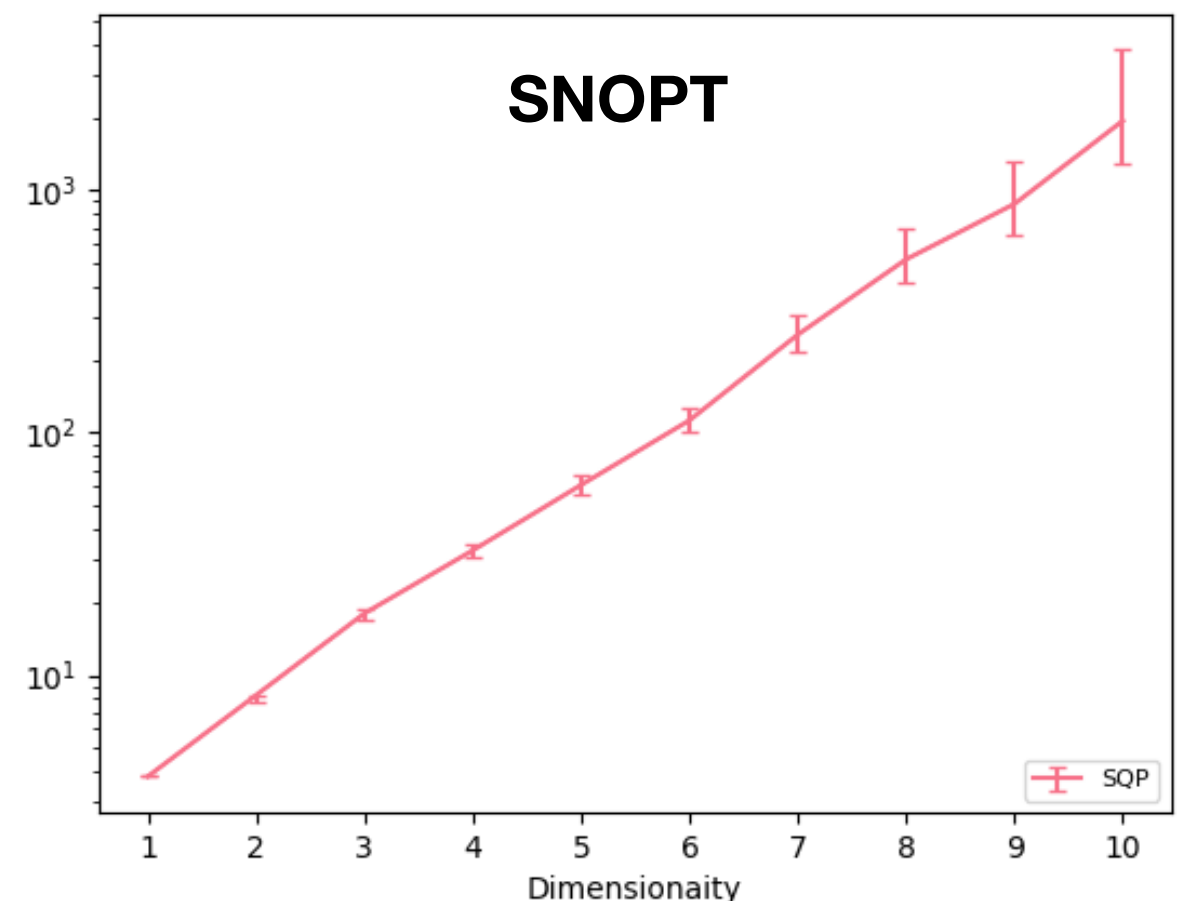
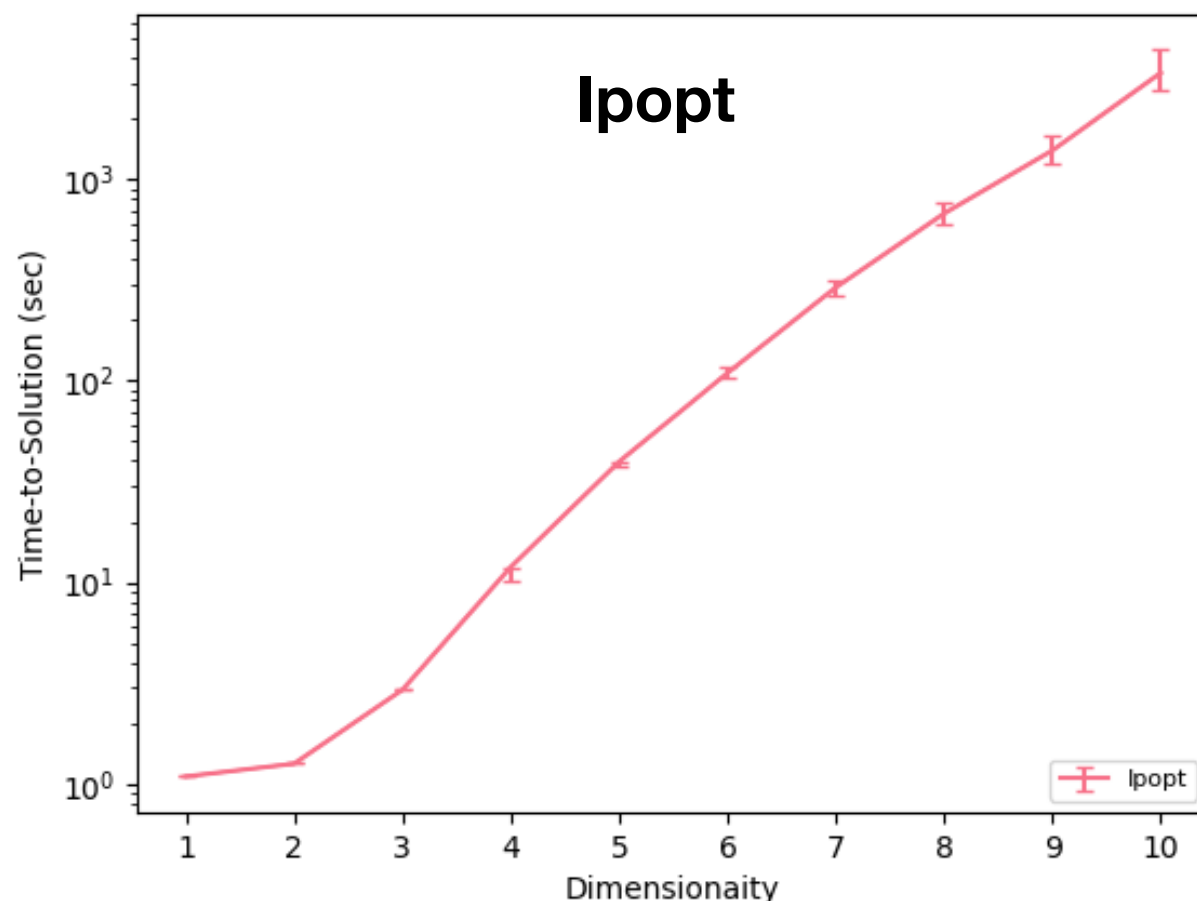
Time-To-Solution (TTS)

$$\text{TTS} = t_f \left\lceil \frac{\ln(1 - 0.99)}{\ln(1 - P_g)} \right\rceil$$

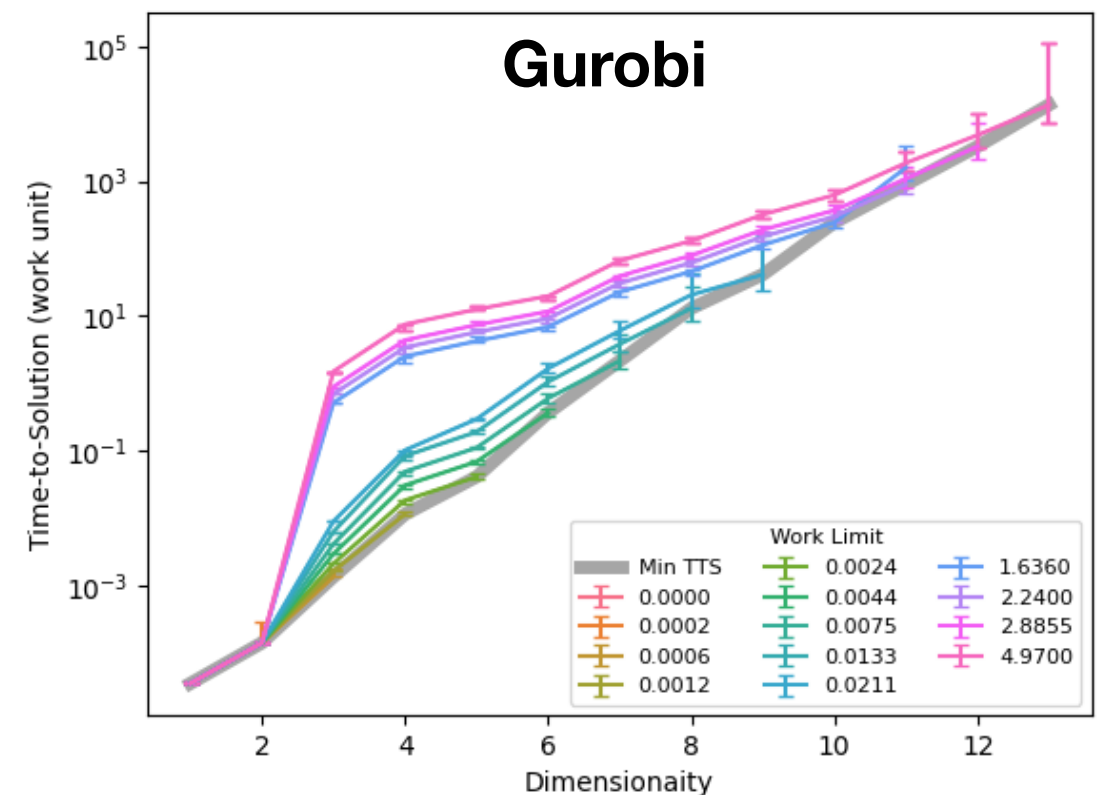
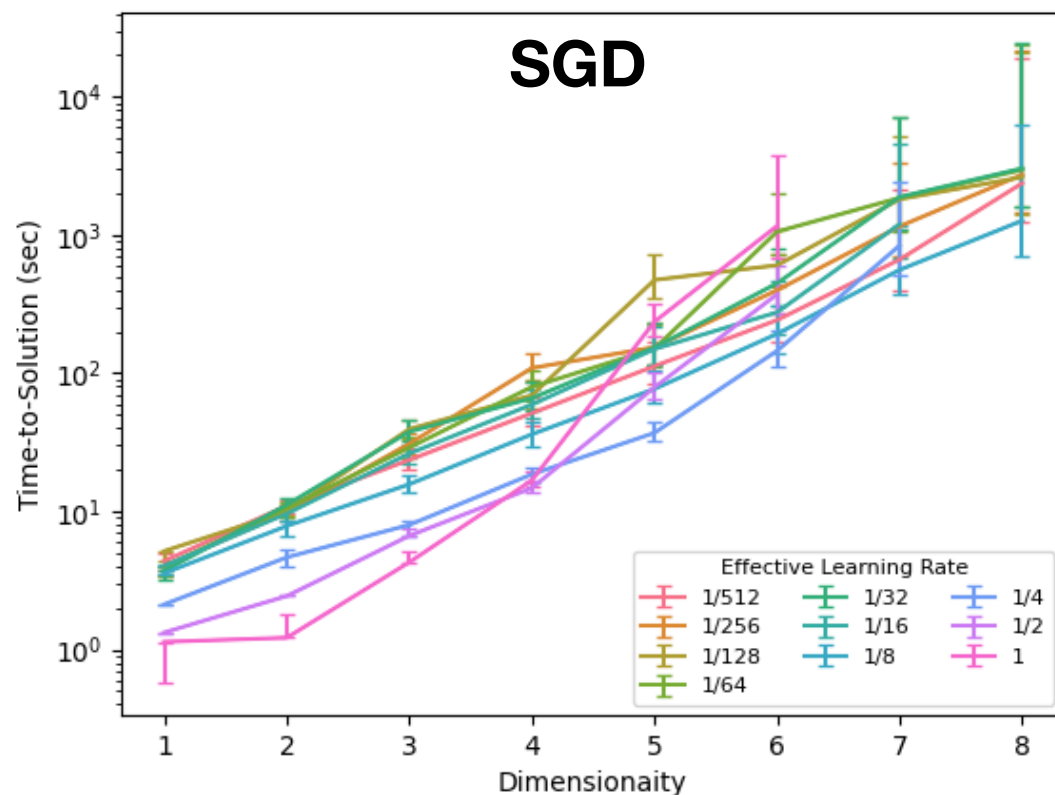
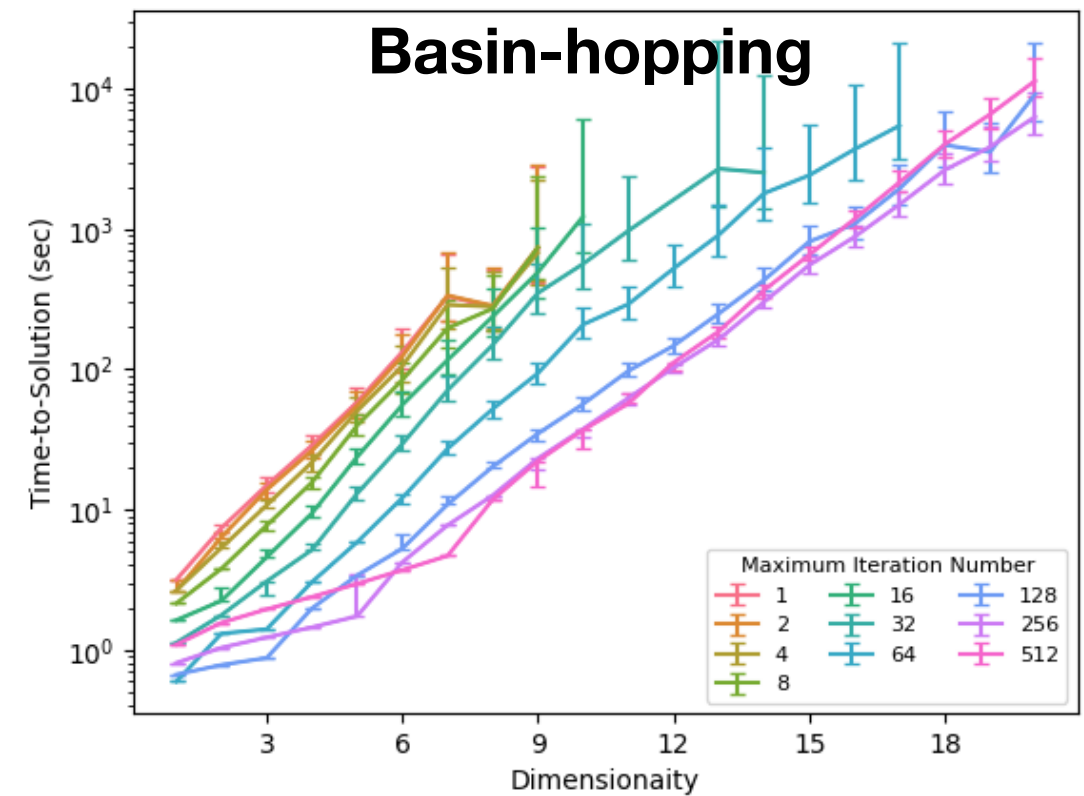
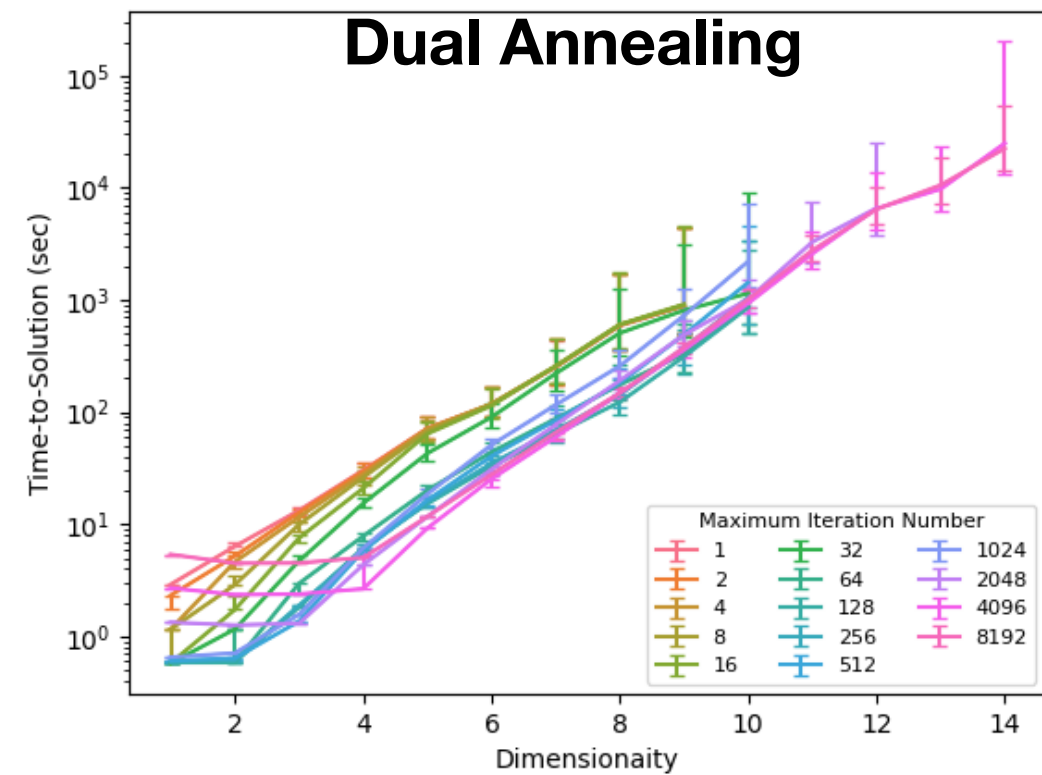
- t_f - algorithm/solver running time (wall clock time)
- p_g - success probability per run

QHD TTS - $\mathcal{O}(d^4)$ (given fixed δ)

Classical TTS - numerical results suggest **super-polynomial** scaling in d



A quantum-classical performance separation



Part III:

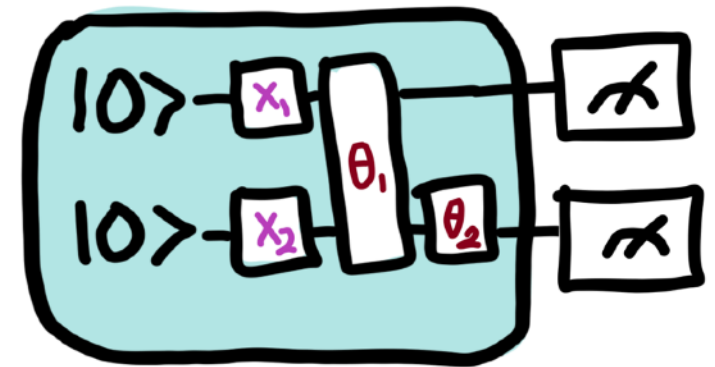
Implementation on Quantum Computers

QHD with digital quantum computers

$$H(t) = e^{\varphi(t)}\left(-\frac{1}{2}\nabla^2\right) + e^{\chi(t)}f(x), t \geq 0$$

Digital Quantum Implementation

- The algorithm is effectively a **time-dependent Hamiltonian simulation** in the real space (we know poly-time algorithms).
- Requires **hundreds of millions of gates!**
- **Conclusion:** the digital implementation is **far from being feasible** in near term!



Corresponding **T-gate** Count with Digital Quantum Computing (before fault-tolerance)

Dimensions	3-qubit format	16-qubit format	32-qubit format
50	5.49e+8	7.8386e+9	2.672e+10
60	6.588e+8	9.4063e+9	3.2064e+10
75	8.235e+8	1.1758e+10	4.008e+10

QHD with analog quantum computers

Analog Quantum Implementation

- **Analog simulation:** problem solving by *emulating real quantum systems*.
- Abstraction: **Quantum Ising Machine (QIM)**, e.g., D-Wave, QuEra, etc.)

$$H(t) = -\frac{A(t)}{2} \left(\sum_j \sigma_x^{(j)} \right) + \frac{B(t)}{2} \left(\sum_j h_j \sigma_z^{(j)} + \sum_{j>k} J_{j,k} \sigma_z^{(j)} \sigma_z^{(k)} \right)$$

- **Programmability:** coefficients $h_j, J_{j,k}$ and functions $A(t), B(t)$.



- QHD is formulated as a quantum Hamiltonian evolution → suitable for analog implementation!
- Patter mismatch: $\text{QHD } H(t) = e^{\varphi_t(-\Delta/2)} + e^{\chi_t f}$
- We develop a new technique named **Hamiltonian embedding**: mapping our target Hamiltonian (QHD) to a “diagonal block” of the machine Hamiltonian (QIM).
- If $H = H_0 \oplus H_1$, we have $e^{-iHt} = e^{-iH_0t} \oplus e^{-iH_1t}$.
- We implement QHD on today’s analog quantum computers (D-Wave’s **advantage_system6.1**).
- **Advantage:** resource-efficiency → first large-scale empirical study for nonlinear optimization using quantum computers.

Hamiltonian embedding of QHD

$$i \frac{\partial}{\partial t} \Psi(t, x) = \left[\underbrace{-\frac{e^{\varphi_t}}{2} \nabla^2}_{\text{Kinetic part}} + \underbrace{e^{\chi_t} f(x)}_{\text{Potential part}} \right] \Psi(t, x),$$

Finite difference 

$$i \frac{d}{dt} |\phi_t\rangle = \left[\underbrace{-\frac{e^{\varphi_t}}{2} \hat{A}}_{\text{Kinetic part}} + \underbrace{e^{\chi_t} \hat{F}}_{\text{Potential part}} \right] |\phi_t\rangle$$

- **Kinetic operator:** $\hat{L} = \frac{1}{(1/r)^2} \begin{bmatrix} -2 & 1 & & \\ 1 & -2 & 1 & \\ \dots & \dots & \dots & \dots \\ & 1 & -2 & 1 \\ & & 1 & -2 \end{bmatrix} \rightarrow \hat{A} = \frac{1}{(1/r)^2} \begin{bmatrix} 0 & 1 & & \\ 1 & 0 & 1 & \\ \dots & \dots & \dots & \dots \\ & 1 & 0 & 1 \\ & & 1 & 0 \end{bmatrix}$
Tri-diagonal
- **Potential operator:** $\hat{F} = \begin{bmatrix} f(a_0) & & & \\ & f(a_1) & & \\ \dots & \dots & \dots & \dots \\ & & f(a_{r-1}) & \\ & & & f(a_r) \end{bmatrix}$
Diagonal

Hamming states — an orthonormal basis

$$|H_j\rangle = \frac{1}{\sqrt{C_j}} \sum_{|b|=j} |b\rangle \quad \text{where } C_j = \binom{N}{j} \text{ and } 0 \leq j \leq N$$

$$|H_0\rangle = |0000\rangle, |H_1\rangle = \frac{1}{2}(|1000\rangle + |0100\rangle + |0010\rangle + |0001\rangle), \text{ etc.}$$

Lemma (informal version). [Leng, Hickman, Li, Wu (2023)]

Given n qubits, the subspace \mathcal{S} spanned by all $(n + 1)$ Hamming states is an invariant subspace of the QIM Hamiltonian. The projection of the QIM Hamiltonian into the subspace \mathcal{S} approximates the discretized QHD Hamiltonian.

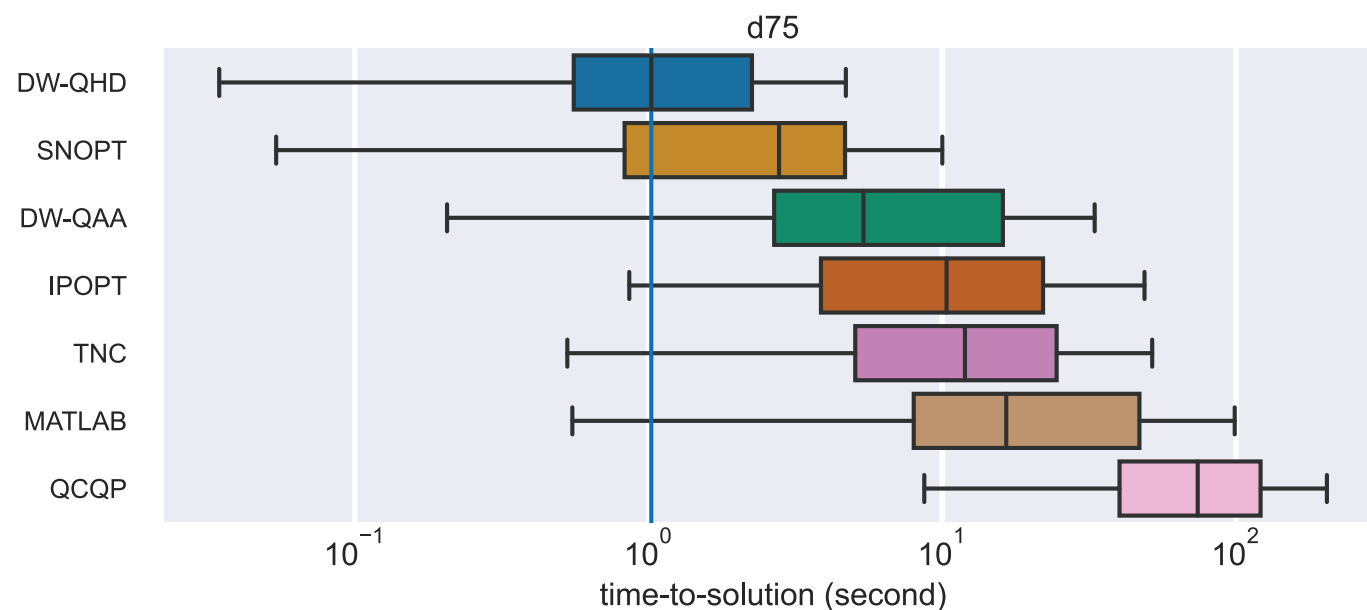
Large-scale empirical study on real quantum computers

We identify a class of *non-trivial* and *self-interesting* optimization problems that can be mapped to QIMs — **Quadratic Programming (QP) with box constraints**.

$$\begin{array}{ll} \text{minimize} & f(\mathbf{x}) = \frac{1}{2} \mathbf{x}^\top \mathbf{Q} \mathbf{x} + \mathbf{b}^\top \mathbf{x} \\ \text{subject to} & \mathbf{0} \preceq \mathbf{x} \preceq \mathbf{1}, \end{array}$$

QP — NP-hard with indefinite Q

(For implementation details, see our paper [arXiv:2303.01471](https://arxiv.org/abs/2303.01471) Appendix F)



Time-To-Solution (TTS)

The lower, the better!!

$$\text{TTS} = t_f \left[\frac{\ln(1 - 0.99)}{\ln(1 - P_g)} \right]$$

t_f - quantum anneal time + post-processing
or classical runtime (**wall-clock time**)

P_g - success probability per run

- **DW-QHD** is better than the rest, including **DW-QAA**, and **classical GD, interior points**, and **some local search heuristic**.
- Assuming DW-QHD is no worse than the ideal QHD, this provides **a very strong empirical evidence supporting QHD**.
- QHD does not beat SOTA classical solvers (e.g., Gurobi, CPLEX) in D75. However, such branch-and-bound solvers are **not scalable!**

Summary & Future Work

- **QHD** is an upgraded version of **classical GD and variants**.
- QHD leverages the continuous structure of the problem and converges faster than QAA.
- QHD has different solution path compared to classical GD.
- **QHD** can be used as *subroutines* for more complicated algorithms like branch-and-bound.

All data & codes are available online!

- Source code (Github): <https://github.com/jiaqileng/quantum-hamiltonian-descent>
- Raw data (Box): <https://umd.app.box.com/s/vq747fvjnt8qrkboxprexhoh44n0q9m0i>
- Website: <https://jiaqileng.github.io/quantum-hamiltonian-descent/>



arXiv:2303.01471

Thank You!



QHD Website