DUNE PDELab Tutorial 01

Conforming FEM for a Nonlinear Poisson Equation

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Motivation

This tutorial extends on tutorial 00 by

- 1) Solving a **nonlinear** stationary PDE
- 2) Using conforming finite element spaces of arbitrary order
- 3) Using different types of (conforming) meshes (simplicial, cubed and mixed)
- 4) Using multiple types of boundary conditions

PDE Problem

We consider the problem

$$-\Delta u + q(u) = f \qquad \text{in } \Omega, \tag{1a}$$

$$u = g$$
 on $\Gamma_D \subseteq \partial \Omega$, (1b)

$$-\nabla u \cdot \nu = j \qquad \text{on } \Gamma_{N} = \partial \Omega \setminus \Gamma_{D}. \tag{1c}$$

- ▶ $q: \mathbb{R} \to \mathbb{R}$ is possibly nonlinear function
- ▶ $f: \Omega \to \mathbb{R}$ the source term
- $\triangleright \nu$ unit outer normal to the domain

Weak Formulation

Find
$$u \in U$$
 s.t.: $r^{NLP}(u, v) = 0 \quad \forall v \in V,$ (2)

with the continuous residual form

$$r^{\mathsf{NLP}}(u,v) = \int_{\Omega} \nabla u \cdot \nabla v + (q(u) - f)v \, dx + \int_{\Gamma_N} jv \, ds$$

and the function spaces

- $V = \{v \in H^1(\Omega) : "v = g" \text{ on } \Gamma_D\} \text{ (affine space)}$
- $V = \{ v \in H^1(\Omega) : "v = 0" \text{ on } \Gamma_D \}$

We assume that a unique solution exists

Algebraic Problem

→ Solve weak formulation in finite-dimensional spaces

$$U_h = \operatorname{span}\{\phi_1, \dots, \phi_n\}, \quad V_h = \operatorname{span}\{\psi_1, \dots, \psi_m\}$$

Expanding the solution $u_h = \sum_{j=1}^n (z)_j \phi_j$ results in an algebraic equation for $z \in \mathbb{R}^n$:

Find
$$u_h \in U_h$$
 s.t.: $r(u_h, v) = 0 \quad \forall v \in V_h$
$$\Leftrightarrow \quad r\left(\sum_{j=1}^n (z)_j \phi_j, \psi_i\right) = 0 \quad \forall i = 1, \dots, m$$

$$\Leftrightarrow \qquad R(z) = 0,$$

where $R: \mathbb{R}^n \to \mathbb{R}^m$ given by $R_i(z) = r_h\left(\sum_{j=1}^n (z)_j \phi_j, \psi_i\right)$ is a nonlinear, vector-valued function.

Solution of Algebraic Problem

Use *iterative* methods to solve R(z) = 0, fixed point iteration:

$$z^{(k+1)} = G(z^{(k)}) = z^{(k)} - \lambda^k W(z^{(k)}) R(z^{(k)}).$$
 (3)

- $\triangleright \lambda^k$ is a damping factor
- $W(z^{(k)})$ is a preconditioner matrix, e.g. in Newton's method one has

$$W(z^{(k)}) = (J(z^{(k)}))^{-1}$$
 where $(J(z^{(k)}))_{i,j} = \frac{\partial R_i}{\partial z_i}(z^{(k)})$

i.e. need to solve
$$J(z^{(k)}) w = R(z^{(k)})$$

The following algorithmic building blocks are required:

- i) residual evaluation R(z),
- ii) Jacobian evaluation J(z) (or an approximation of it),
- iii) matrix-free Jacobian application J(z)w (or an approximation).

Note on Matrix-free Evaluation

Nonlinear case:

$$(J(z)w)_i = \sum_{i=1}^n (J(z))_{i,j}(w)_j = \sum_{i=1}^n \frac{\partial}{\partial z_j} r_h \left(\sum_{l=1}^n (z)_l \phi_l, \psi_i \right) (w)_j.$$

Linear case: $r_h(u, v) = a(u, v) - l(v)$, a BLF, I LF

$$(J(z)w)_{i} = \sum_{j=1}^{n} \frac{\partial}{\partial z_{j}} r_{h} \left(\sum_{l=1}^{n} (z)_{l} \phi_{l}, \psi_{i} \right) (w)_{j}$$

$$= \sum_{j=1}^{n} \frac{\partial}{\partial z_{j}} \left(a_{h} \left(\sum_{l=1}^{n} (z)_{l} \phi_{l}, \psi_{i} \right) - I_{h}(\psi_{i}) \right) (w)_{j}$$

$$= \sum_{j=1}^{n} \frac{\partial}{\partial z_{j}} \left(\sum_{l=1}^{n} (z)_{l} a_{h} (\phi_{l}, \psi_{i}) \right) (w)_{j}$$

$$= \sum_{j=1}^{n} a_{h} (\phi_{j}, \psi_{i}) (w)_{j} = a_{h} \left(\sum_{l=1}^{n} (w)_{j} \phi_{j}, \psi_{i} \right) = (Aw)_{i}$$

Recall Finite Element Mesh Notation

i) Ordered sets of vertices and elements:

$$\mathcal{X}_h = \{x_1, \dots, x_N\}, \quad \mathcal{T}_h = \{T_1, \dots, T_M\}$$

ii) Partitioning of vertex index set $\mathcal{I}_h = \{1, \dots, N\}$ into $\mathcal{I}_h = \mathcal{I}_h^{int} \cup \mathcal{I}_h^{\partial \Omega}$:

$$\mathcal{I}_h^{int} = \{i \in \mathcal{I}_h \, : \, x_i \in \Omega\}, \quad \mathcal{I}_h^{\partial \Omega} = \{i \in \mathcal{I}_h \, : \, x_i \in \partial \Omega\}.$$

iii) For every element $T \in \mathcal{T}_h$ a local-to-global map

$$g_T: \{0,\ldots,n_T-1\} \to \mathcal{I}_h$$

iv) For every element $T \in \mathcal{T}_h$ an element transformation map

$$\mu_T: \hat{T} \to T$$

 μ_T is differentiable with invertible Jacobian and consistent with g_T :

$$\forall i \in \{0,\ldots,n_T-1\} : \mu_T(\hat{x}_i) = x_{g_T(i)}$$

Conforming Finite Element Space

with polynomial degree k in dimension d on mesh \mathcal{T}_h :

$$V_h^{k,d}(\mathcal{T}_h) = \left\{ v \in C^0(\overline{\Omega}) : \forall T \in \mathcal{T}_h : v|_T = \mu_T \circ p_T \land p_T \in \mathbb{P}_T^{k,d} \right\}$$

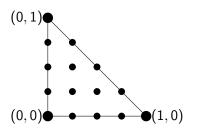
where the multivariate polynomials \mathbb{P} depend on element type:

$$\mathbb{P}_{T}^{k,d} = \begin{cases} \begin{cases} p : p(x_{1}, \dots, x_{d}) = \sum\limits_{0 \leq \|\alpha\|_{1} \leq k} c_{\alpha} x_{1}^{\alpha_{1}} \cdot \dots \cdot x_{d}^{\alpha_{d}} \end{cases} & \hat{T} = \hat{S} \text{ (simplex)}, \\ p : p(x_{1}, \dots, x_{d}) = \sum\limits_{0 \leq \|\alpha\|_{\infty} \leq k} c_{\alpha} x_{1}^{\alpha_{1}} \cdot \dots \cdot x_{d}^{\alpha_{d}} \end{cases} & \hat{T} = \hat{C} \text{ (cube)} \end{cases}$$

The dimension of $\mathbb{P}_T^{k,d}$ is:

$$n_{\hat{C}}^{k,d} = (k+1)^d \text{ (cube)}, \ n_{\hat{S}}^{k,d} = \begin{cases} 1 & k = 0 \lor d = 0\\ \sum_{i=0}^k n_{\hat{S}}^{i,d-1} & \text{else} \end{cases}$$
 (simplex)

Local Lagrange Basis



Lagrange points and poynomials (shape functions) on \hat{T} :

$$L_{\hat{T}} = \left\{ \hat{x}_0^{\hat{T}}, \dots, \hat{x}_{n_{\div}^{k,d}-1}^{\hat{T}} \right\}, \quad P_{\hat{T}} = \left\{ p_0^{\hat{T}}, \dots, p_{n_{\div}^{k,d}-1}^{\hat{T}} \right\}$$

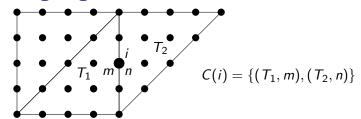
such that

$$p_i^{\hat{T}}(\hat{x}_i^{\hat{T}}) = \delta_{i,i}$$

Extend local to global map:

$$g_{\mathcal{T}}: \{0,\dots,n_{\hat{\mathcal{T}}}^{k,d}-1\} \to \mathcal{I}\left(V_h^{k,d}(\mathcal{T}_h)\right) = \left\{0,\dots,\dim V_h^{k,d}(\mathcal{T}_h)-1\right\}$$

Global Lagrange Basis



Define inversion of the global map:

$$C(i) = \{(T, m) \in \mathcal{T}_h \times \mathbb{N} : g_T(m) = i\}$$

then the global Lagrange basis functions are

$$\phi_i(x) = \left\{ \begin{array}{ll} p_m^{\hat{T}}(\mu_T^{-1}(x)) & x \in T \land (T,m) \in C(i) \\ 0 & \text{else} \end{array} \right., \quad i \in \mathcal{I}\left(V_h^{k,d}(T_h)\right).$$

corresponding to the global Lagrange points

$$\mathcal{X}_h^{k,d} = \left\{ x_i \in \overline{\Omega} : x_i = \mu_T(\hat{x}_m^{\hat{T}}) \land (T, m) \in C(i) \right\}$$

Dirichlet Boundary Conditions

Indices of Lagrange points on the Dirichlet boundary are:

$$\mathcal{I}^{D}\left(V_{h}^{k,d}(\mathcal{T}_{h})\right) = \left\{i \in \mathcal{I}\left(V_{h}^{k,d}(\mathcal{T}_{h})\right) \,:\, x_{i} \in \mathcal{X}_{h}^{k,d} \cap \Gamma_{D}\right\}.$$

Then the test space with zero Dirichlet condition is:

$$V_{h,0}^{k,d}(\mathcal{T}_h) = \left\{ v \in V_h^{k,d}(\mathcal{T}_h) \ : \ v(x_i) = 0 \quad \forall i \in \mathcal{I}^D\left(V_h^{k,d}(\mathcal{T}_h)\right) \right\}$$

For the trial space choose any extension

$$u_{h,g} = \sum_{i \in \mathcal{I}\left(V_h^{k,d}(\mathcal{T}_h)\right)} u_g(x_i)\phi_i \qquad u_g(x_i) = g(x_i) \ \forall i \in \mathcal{I}^D\left(V_h^{k,d}(\mathcal{T}_h)\right)$$

Then

$$U_h^{k,d}(\mathcal{T}_h) = \left\{ u \in V_h^{k,d}(\mathcal{T}_h) : u = u_{h,g} + w \wedge w \in V_{h,0}^{k,d}(\mathcal{T}_h) \right\}$$

General Constraints

Task: Given $U_h = \operatorname{span} \{\phi_j : j \in J_h = \{1, \dots, n\}\}$ construct $\tilde{U}_h \subseteq U_h$

This is how it is done in PDELab:

- 1) Assume $U_h = \operatorname{span}\{\phi_i : i \in J_h\}$
- **2)** Select a subset of indices $\tilde{J}_h \subset J_h$
- 3) Set $\tilde{U}_h=\operatorname{span}\left\{\tilde{\phi}_j:j\in \tilde{J}_h\right\}$, where the new basis functions have the form

$$\tilde{\phi}_j = \phi_j + \sum_{I \in J_h \setminus \tilde{J}_h} (B)_{j,I} \phi_I \quad \forall j \in \tilde{J}_h.$$

Any subspace is thus characterized by $C = (\tilde{J}_h, B)$

Element-wise Computations

Come back to the residual form which is element-wise

$$r^{\mathsf{NLP}}\left(u,v
ight) = \sum_{T \in \mathcal{T}_h} lpha_T^V(u,v) + \sum_{T \in \mathcal{T}_h} \lambda_T^V(v) + \sum_{F \in \mathcal{F}_h^{\partial\Omega}} \lambda_F^B(v)$$

with

$$\alpha_T^V(u, v) = \int_T \nabla u \cdot \nabla v + q(u)v \, dx$$
$$\lambda_T^V(v) = -\int_T fv \, dx$$
$$\lambda_F^B(v) = \int_{F \cap \Gamma_N} jv \, ds$$

 $\mathcal{F}_h^{\partial\Omega}$: intersections of elements with the domain boundary, i.e.

$$F = T_{E}^{-} \cap \partial \Omega$$

with $T_F^- \in \mathcal{T}_h$ the "minus" element associated with F

λ Volume Term

For any $(T, m) \in C(i)$ we obtain

$$\lambda_T^V(\phi_i) = -\int_T f\phi_i \, dx = -\int_{\hat{T}} f(\mu_T(\hat{x})) \rho_m^{\hat{T}}(\hat{x}) |\det J_{\mu_T}(\hat{x})| \, d\hat{x}.$$

 J_{μ_T} is the Jacobian of the element map μ_T

This integral is computed using numerical quadrature

Collect all contributions of T in a small vector:

$$(\mathcal{L}_T^V)_m = -\int_{\hat{T}} f(\mu_T(\hat{x})) p_m^{\hat{T}}(\hat{x}) |\det J_{\mu_T}(\hat{x})| \, d\hat{x}.$$

λ Boundary Term

For $F \in \mathcal{F}_h^{\partial\Omega}$ with $F \subseteq \Gamma_N$ and $(T_F^-, m) \in C(i)$ we obtain

$$\lambda_T^B(\phi_i) = \int_F j v \, ds = \int_{\hat{F}} j(\mu_F(\hat{x})) p_m^{\hat{T}}(\eta_F(\hat{x})) \sqrt{|\text{det}(J_{\mu_F}^T(\hat{x})J_{\mu_F}(\hat{x}))|} \, ds$$

The integration is more involved here because it is over a face:

- ightharpoonup $\mu_F:\hat{F} o F$ maps the reference element of the face to the face
- $\eta_F:\hat{F}\to\hat{T}_F^-$ maps reference element of the face to reference element of T_F^-

Collect all contributions of F in a small vector:

$$(\mathcal{L}_T^B)_m = \int_{\hat{F}} j(\mu_F(\hat{x})) p_m^{\hat{T}}(\eta_F(\hat{x})) \sqrt{|\text{det}J_{\mu_T}^T(\hat{x})J_{\mu_T}(\hat{x})|} \ ds.$$

Numerical quadrature is applied to compute the integral

α Volume Term

For any $(T, m) \in C(i)$ we get

$$\begin{split} \alpha_T^V(u_h,\phi_i) &= \int_T \nabla u \cdot \nabla \phi_i + q(u)\phi_i \, dx, \\ &= \int_T \sum_j (z)_j \left(\nabla \phi_j \cdot \nabla \phi_i \right) + q \left(\sum_j (z)_j \phi_j \right) \phi_i \, dx, \\ &= \int_{\hat{T}} \sum_n (z)_{g_T(n)} (J_{\mu_T}^{-1}(\hat{x}) \hat{\nabla} p_n^{\hat{T}}(\hat{x})) \cdot (J_{\mu_T}^{-1}(\hat{x}) \hat{\nabla} p_m^{\hat{T}}(\hat{x})) \\ &+ q \left(\sum_n (z)_{g_T(n)} p_n^{\hat{T}}(\hat{x}) \right) p_m^{\hat{T}}(\hat{x}) |\det J_{\mu_T}(\hat{x})| \, dx \end{split}$$

and again collect all contributions from T in a small vector:

$$(\mathcal{R}_T^V(R_T z))_m = \alpha_T^V \left(\sum_n (R_T z)_n p_n^{\hat{T}}, p_m^{\hat{T}} \right)$$

Putting It All Together

With these local contributions the evaluation of the algebraic residual is

$$R(z) = \sum_{T \in \mathcal{T}_h} R_T^T \mathcal{R}_T^V(R_T z) + \sum_{T \in \mathcal{T}_h} R_T^T \mathcal{L}_T^V + \sum_{F \in \mathcal{F}_h^{\partial \Omega} \cap \Gamma_N} R_T^T \mathcal{L}_F^B$$

where R_T is the "picking out" matrix of element T

The Jacobian of R is

$$(J(z))_{i,j} = \frac{\partial R_i}{\partial z_j}(z) = \sum_{\substack{(T,m,n): (T,m) \in C(i) \land (T,n) \in C(j)}} \frac{\partial (\mathcal{R}_T^V)_m}{\partial z_n}(R_T z)$$

Note that:

- i) Entries of the Jacobian can be computed element by element.
- ii) The derivative is independent of the λ -terms
- iii) Jacobian entries may be computed by numerical differentiation

Implementation Overview

- 1) tutorialO1.ini holds parameters controlling the execution
- 2) tutorial01.cc includes the necessary C++, DUNE and PDELab header files; contains main function calling driver
- 3) Function driver in driver.hh instantiates the necessary PDELab classes for solving a nonlinear stationary problem and finally solves the problem.
- 4) File nonlinearpoissonfem.hh contains class NonlinearPoissonFEM realizing a PDELab local operator
- 5) File problem.hh contains a "parameter class" which encapsulates the user-definable part of the PDE problem
- **6)** Finally, the tutorial provides some mesh files.