

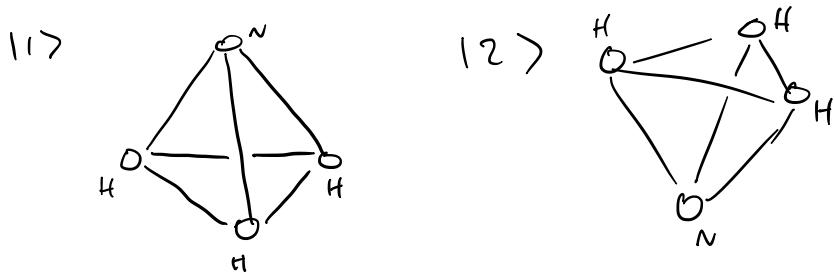
QM Lecture 8



EXAMPLE AMMONIA MOLECULE & AMMONIA MASER

* previous, \hat{H} was nearly diag, but small perturbation ($B_1 \neq 0$) allowed transitions between the would be \hat{H} eigenstates.

Ammonia is another famous example.



* WANT: QM DESCRIPTION OF SYSTEM.

* N SAME DIST. FROM 3H-PLANE

* SYMMETRY $\Rightarrow \langle 1 | \hat{H} | 1 \rangle \approx \langle 2 | \hat{H} | 2 \rangle =: E_0$, WHATEVER \hat{H} IS.

* FACT: SOMETIMES NH_3 "FLIPS" ORIENTATION

(Q: what does this imply about \hat{H} ?)

$$\Rightarrow \langle 1 | \hat{H} | 2 \rangle \neq 0 \quad \langle 2 | \hat{H} | 1 \rangle \neq 0$$

* $H^+ = H \Rightarrow \langle 1 | \hat{H} | 2 \rangle = \langle 2 | \hat{H} | 1 \rangle^*$ TAKE $= -A \in \mathbb{R}$
(choice of sign for later convenience)

$$\hat{H} \xrightarrow{1,2 \text{ BASIS}} \begin{pmatrix} \langle 1 | \hat{H} | 1 \rangle & \langle 1 | \hat{H} | 2 \rangle \\ \langle 2 | \hat{H} | 1 \rangle & \langle 2 | \hat{H} | 2 \rangle \end{pmatrix} = \begin{pmatrix} E_0 & -A \\ -A & E_0 \end{pmatrix}$$

Q: WHAT ARE ENERGY E-VACS?

$$\hat{H}|4\rangle = E|4\rangle$$

$$\begin{pmatrix} E_0 & -A \\ -A & E_0 \end{pmatrix} \begin{pmatrix} \langle 1|4\rangle \\ \langle 2|4\rangle \end{pmatrix} = E \begin{pmatrix} \langle 1|4\rangle \\ \langle 2|4\rangle \end{pmatrix}$$

$$\det \begin{pmatrix} E_0 - E & -A \\ -A & E_0 - E \end{pmatrix} = 0 \Rightarrow (E_0 - E)^2 - A^2 = 0 \Rightarrow E_0 - E = \pm A \Rightarrow \boxed{E = E_0 \pm A}$$

* DEFINE $|I\rangle$, $\hat{H}|I\rangle = (E_0 - A)|I\rangle$

$$\begin{pmatrix} E_0 - E_0 + A & -A \\ -A & E_0 - E_0 + A \end{pmatrix} \begin{pmatrix} \langle 1|I\rangle \\ \langle 2|I\rangle \end{pmatrix} = 0 \Rightarrow A \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \langle 1|I\rangle \\ \langle 2|I\rangle \end{pmatrix} = 0 \Rightarrow \langle 1|I\rangle = \langle 2|I\rangle$$

$$\Rightarrow \boxed{|I\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle)}$$

* SIMILARLY: DEF. $|II\rangle$ $\hat{H}|II\rangle = (E + A)|II\rangle$

$$\boxed{|II\rangle = \frac{1}{\sqrt{2}}(|1\rangle - |2\rangle)}$$

* NB ① EXISTENCE OF TRANSITION $\Leftrightarrow A \neq 0 \Leftrightarrow$ ENERGIES SPLIT

② EN. E-STATES ARE LIN. COMB. OF $|1\rangle, |2\rangle$

③ $|I\rangle$ SYMM. UNDER $|1\rangle \rightarrow |2\rangle$
 $|2\rangle \rightarrow |1\rangle$

$|II\rangle$ ANTI SYMM. ($I \in |II\rangle \rightarrow |-II\rangle$)

④ EXPERIMENT $E_0 \sim \text{eV}$ $A \sim 10^{-4} \text{ eV}$

⑤ $E_{II} - E_I = 2A \Rightarrow A \in \mathbb{R}$ (if not ifid our assumption)

TIME EVOLUTION

RECALL: $|I\rangle$ NOT \hat{f} E-STATE. \Rightarrow TIME EVOLVES

$$|\psi(0)\rangle = |I\rangle = \frac{1}{\sqrt{2}}(|I\rangle + |II\rangle)$$

$$\begin{aligned} |\psi(t)\rangle &= e^{-i\hat{H}t/\hbar} \frac{1}{\sqrt{2}}(|I\rangle + |II\rangle) = \frac{1}{\sqrt{2}}\left(e^{-i(E_0-A)t/\hbar}|I\rangle + e^{-i(E_0+A)t/\hbar}|II\rangle\right) \\ &= e^{-i(E_0-A)t/\hbar}\left(\frac{1}{\sqrt{2}}|I\rangle + e^{-2iAt/\hbar}\frac{1}{\sqrt{2}}|II\rangle\right) \end{aligned}$$

* TIME DEP. REL STATE \Rightarrow NOT STATIONARY STATE.

* IF $A \rightarrow 0$, BECOMES STATIONARY STATE

(Q: why is this expected?)

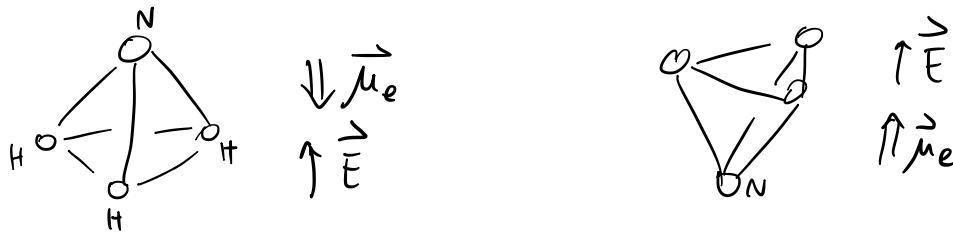
NOTE $|\psi(0)\rangle = |I\rangle$, BUT $|\psi(\frac{\hbar\pi}{2A})\rangle = (\text{phase})\left(\frac{1}{\sqrt{2}}|I\rangle - \frac{1}{\sqrt{2}}|II\rangle\right)$

$$= (\text{phase})|12\rangle$$

\Rightarrow AMMONIA FLIPPING BETWEEN $|I\rangle, |2\rangle$!

(empirically, the freq. ≈ 24 GHz)

PERTURBATION: STATIC EXTERNAL ELECTRIC FIELD



$$U = -\vec{\mu}_e \cdot \vec{E}$$

* SYMMETRY BROKEN BY $|\vec{E}| \neq 0$!

$$\Rightarrow \langle 1 | \hat{H} | 1 \rangle \neq \langle 2 | \hat{H} | 2 \rangle$$

* IE DIAG. ENTRIES CHANGE!

$$\begin{pmatrix} 1 & 2 \\ 1 & 2 \end{pmatrix} \xrightarrow{\text{BS}} \begin{pmatrix} E_0 + \mu_e |\vec{E}| & -A \\ -A & E_0 - \mu_e |\vec{E}| \end{pmatrix} \quad (\text{Q: why these signs?})$$

\Rightarrow EN. E-VACS

$$E = E_0 \pm \sqrt{(\mu_e |\vec{E}|)^2 + A^2}$$

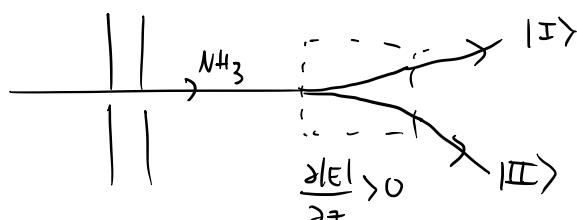
* WE TURNED ON E-FIELD. CAN MAKE $|\vec{E}|$ SMALL

$$\Rightarrow E = E_0 \pm A \left[1 + \frac{(\mu_e |\vec{E}|)^2}{2A} \right] = (E_0 \pm A) \pm \frac{1}{2} \frac{(\mu_e |\vec{E}|)^2}{A}$$

\Rightarrow ENERGIES SPLIT FURTHER

* NOTE: CAN USE \vec{E} -FIELD TO SPLIT

$$F_z \cong -\frac{\partial}{\partial z} \left[\pm \frac{(\mu_e |\vec{E}|)^2}{2A} \right] \quad (\text{sim. to SG dev})$$



INDUCED TRANSITIONS

(ammonia maser: induce transitions with a time dependent electric field)

PUT \hat{H} IN $|I\rangle, |II\rangle$ BASIS

$$(\hat{H})_{I\bar{I}} = S^+ (\hat{H})_{I\bar{I}} S \quad (\hat{H})_{I\bar{I}} = \begin{pmatrix} E_0 + \mu_e |\vec{E}| & -A \\ -A & E_0 - \mu_e |\vec{E}| \end{pmatrix}$$

$$S = \begin{pmatrix} \langle 1|I\rangle & \langle 1|II\rangle \\ \langle 2|I\rangle & \langle 2|II\rangle \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$\Rightarrow \hat{H} \xrightarrow{I\bar{I}} \left(\frac{1}{\sqrt{2}} \right)^2 \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} E_0 + \mu_e |\vec{E}| & -A \\ -A & E_0 - \mu_e |\vec{E}| \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$= \begin{pmatrix} E_0 - A & \mu_e |\vec{E}| \\ \mu_e |\vec{E}| & E_0 + A \end{pmatrix}$$

INTRODUCE $\vec{E} = \vec{E}(t) = \vec{E}_0 \cos \omega t$

$$(\hat{H})_{I\bar{I}} = \begin{pmatrix} E_0 - A & \mu_e |\vec{E}_0| \cos \omega t \\ \mu_e |\vec{E}_0| \cos \omega t & E_0 + A \end{pmatrix}$$

* THIS \hat{H} HAS SAME FORM AS $\hat{H}_{\text{MAG RES.}}$ (so just use that soln)

DICTIDONARY

$$\begin{aligned} \frac{\hbar \omega_0}{2} &\longleftrightarrow E_0 - A \\ -\frac{\hbar \omega_0}{2} &\longleftrightarrow E_0 + A \\ \pm \frac{\hbar \omega_1}{2} &\longleftrightarrow \mu_e |\vec{E}_0| \end{aligned}$$

* $\omega \approx \omega_0 = \frac{\hbar(E_0 - A)}{2\mu_e}$ RESONANCE

$$|\psi(t)\rangle = \begin{pmatrix} \cos\left(\frac{\omega_1 t}{4}\right) e^{-i(E_0 - A)t/\hbar} \\ -i \sin\left(\frac{\omega_1 t}{4}\right) e^{i(E_0 + A)t/\hbar} \end{pmatrix} \implies$$

$$|\langle I | \psi(t) \rangle|^2 \approx \cos^2 \frac{\mu_e |\vec{E}_0| t}{2\hbar}$$

$$|\langle II | \psi(t) \rangle|^2 \approx \sin^2 \frac{\mu_e |\vec{E}_0| t}{2\hbar}$$

BUILD AMMONIA MASER (Microwave Amplification by Stimulated Emission of Radiation)

- (1) AMMONIA BEAM THROUGH CONST \vec{E} , FILTER |II> (higher energy \uparrow)
- (2) |II> THROUGH RES. TUNED \vec{E} -FIELD
- (3) WAIT ONE PERIOD
- (4) $\Rightarrow |II\rangle \xrightarrow{\text{trans}} |I\rangle$, ENERGY RELEASE

(energy split $\simeq 2A + O(|\vec{E}|^2)$ cov.)

$$\begin{aligned} (\text{microwaves}) \quad \lambda &\sim 1 \text{ cm} \\ \Rightarrow A &\sim 10^{-4} \text{ eV} \end{aligned}$$

ENERGY UNCERTAINTY

|I> NOT EN. EIG-STATE.

\Rightarrow WHAT IS EN. UNC.?

$$\begin{aligned} (\Delta H)_1 &= (\langle H^2 \rangle - \langle H \rangle^2)^{\frac{1}{2}} = \left\{ \frac{1}{2}(\epsilon_0 - A)^2 + \frac{1}{2}(\epsilon_0 + A)^2 \right\} - \left[\frac{1}{2}(\epsilon_0 - A) + \frac{1}{2}(\epsilon_0 + A) \right]^2 \left\{ \frac{1}{2} \right\} \\ &= A \end{aligned}$$

(makes sense based on level splitting)