NANOPHYSIQUE INTRODUCTION PHYSIQUE AUX NANOSCIENCES

Ch 5. Carbon Structres

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Lecture 6, 2021-2022

Carbon Structures

- Introduction
 - Histoire
 - Carbone
 - Hybridization
- Tight-binding calculations
 - Principe
 - Example: trans
 - Graphene π-liens
 - Graphene σ-liens
- Structure de carbone nanotubes
- Structure Electronique
- Synthesis
- Transport
- Multi-walled nanotubes

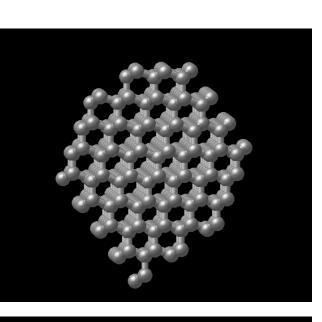
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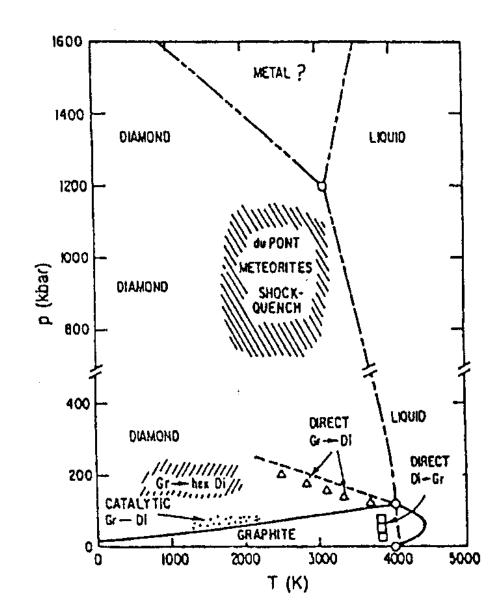
Une petite histoire

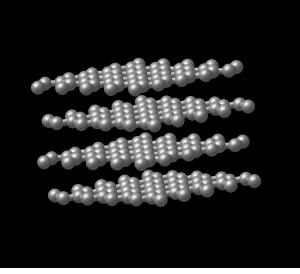
- La première fibre de carbone: Edison, 1879
- Recherche lente jusqu'à ce que les années 1950
- Aéronautique -> recherche de matériaux légers et résistants → "carbon whisker"
- Découverte de fullerenes (Kroto, Smalley 1985)
- Recherche invité à l'échelle du nanomètre fibres
- Spéculation théorique sur la forme et symétrie de carbone nanotubes (Smalley, Dresselhaus, ...)
- Observation de carbone nanotubes (Iijima, 1992 utilisant TEM).

Propriétés du carbone





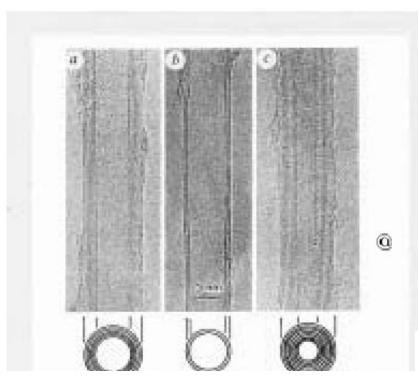




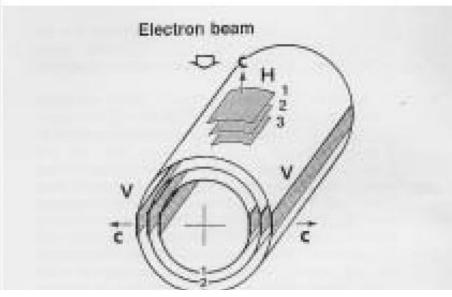
graphite

Discovery of Carbon Nanotubes

Iijima, Nature 354, 56 (1992)



Electron microscope image

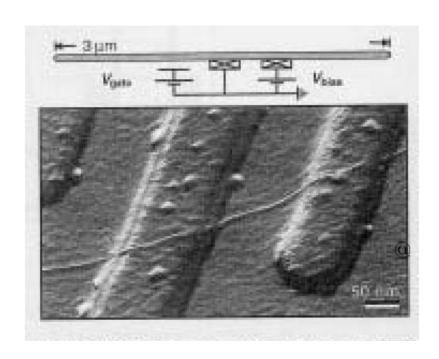


Interpretation of the images

diameter 55 ms, which has the smallest forew diameter (2.2 ms).

Current-voltage characteristics of CNT

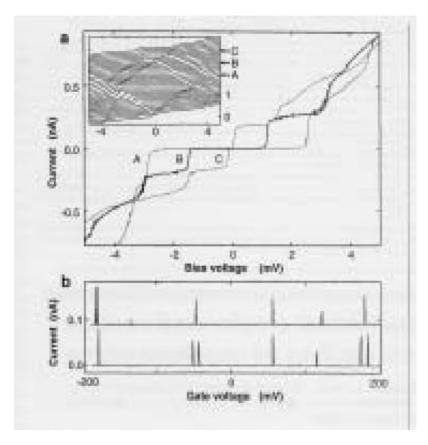
S.J. Tans et al. Nature 386 474 (1997)



Electron microscope image of the system

thin filament: Single-wall CNT

• hills: electrodes



a.Nonlinear conductance (Coulomb staircase)b.Controlling the number of electrons

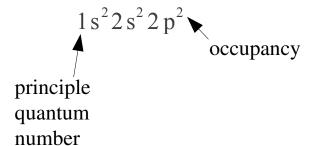
Carbone

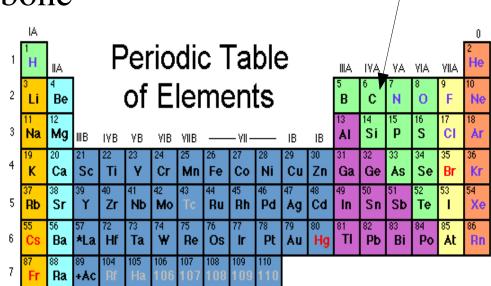
Carbone:

$$Z = 6$$
¹²C, ¹³C--> stable

¹⁴C--> half-life = 5,730 an (carbon dating)

Configuration Electronique:



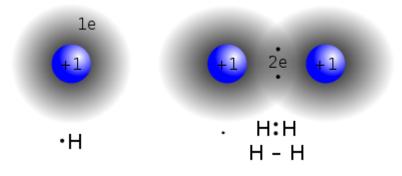


*Lanthanide	58	59	60	61	62	63	64		66	67	68	69	70	71
Series	Ce	Pr	Nd	Pm	Sm	Eu	Gd		Dy	Ho	Er	Tm	Yb	Lu
+ Actinide	90	91	92	93	94	95	96	97	98	99	100	101	102	103
Series	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr



1s 2p

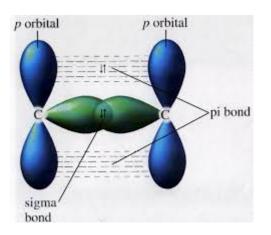
Covalent bonds: sharing electrons

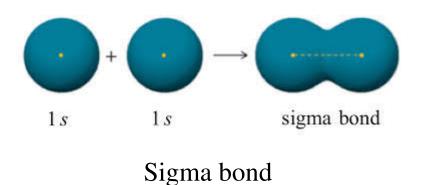


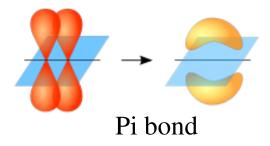
T (kinetic energy) : lower energy by delocalizing electrons

V (potential energy): lower energy by localizing electrons near ions

Types of molecular bonds







T (kinetic energy) : lower energy by delocalizing electrons

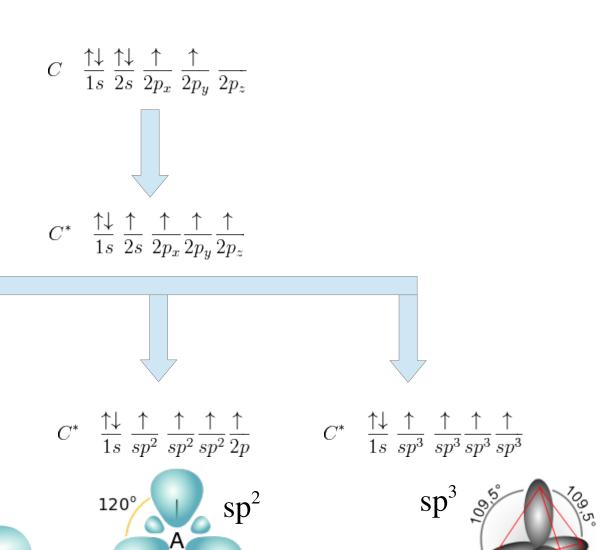
V (potential energy): lower energy by localizing electrons near ions

Hybridization

 $C^* \quad \frac{\uparrow\downarrow}{1s} \; \frac{\uparrow}{sp} \; \frac{\uparrow}{sp} \frac{\uparrow}{p} \frac{\uparrow}{p}$

sp

180°

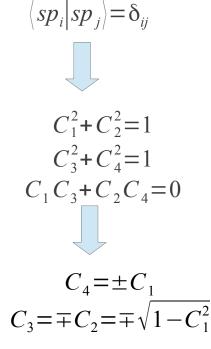


SP Hybridization

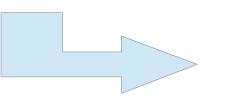
$$|sp_a\rangle = C_1|s\rangle + C_2|p_x\rangle$$

$$|sp_b\rangle = C_3|s\rangle + C_4|p_x\rangle$$

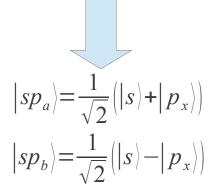
Orthonormality:



R. Saito, G. Dresselhaus & M. S. Dresselhaus, *Physical Properties of Carbon Nanotubes* (Imperial College Press, London, 1998)



$$\begin{split} E_{a} &= C_{1}^{2} E_{s} + C_{2}^{2} E_{p} \\ &= C_{1}^{2} \underbrace{(E_{s} - E_{p})}_{<0} + E_{p} \\ &= \underbrace{E_{b} = C_{3}^{2} E_{s} + C_{4}^{2} E_{p}}_{<0} \\ &= E_{s} + C_{1}^{2} (E_{p} - E_{s}) \\ &= E_{a} + (1 - 2 C_{1}^{2}) (E_{s} - E_{p}) \end{split}$$



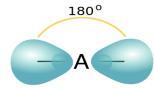
Orbitals and hybridizations for C

Structure électronique d'un atome de carbone = $1s^2 2s^2 2p^2$ coeur = $1s^2$ 4 électrons de valence = $2s^2 2p^2$

Hybridation sp: $2s+2p_x$, $2p_y$, $2p_z$

Example : acétylène: HCCH liaison triple: 2 lien σ + 2 liens π

$$|sp_{a,b}\rangle = \frac{1}{\sqrt{2}}(|2s\rangle \pm |2p_x\rangle)$$

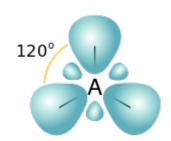


Hybridation sp²: $2s + 2p_x + 2p_y$, $2p_z$

Example : polyacétylène: $(HCCH)_n$ liaison double: 3 liens $\sigma + 1$ lien π

$$|sp_a^2\rangle = \frac{1}{\sqrt{3}} |2s\rangle - \sqrt{\frac{2}{3}} |2p_x\rangle$$

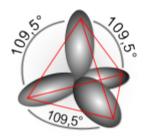
$$|sp_{b,c}^2\rangle = \pm \frac{1}{\sqrt{3}} |2s\rangle \pm \frac{1}{\sqrt{2}} |2p_x\rangle + \frac{1}{\sqrt{6}} |2p_z\rangle$$



Hybridation sp³: $2s + 2p_x + 2p_y + 2p_z$

Example : méthane: CH₄ liaison simple: 4 liens σ

$$|sp_a^3\rangle = \frac{1}{2}(|2s\rangle \pm |2p_x\rangle \pm |2p_y\rangle \pm |2p_z\rangle)$$



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1. **Théorème de Bloch:** Donné d'un potentiel qui est périodique sur un réseau de Bravais, U($\mathbf{r}+\mathbf{R}$)=U(\mathbf{r}), la solution de l'équation Schrodinger d'un électron est $\psi(\mathbf{r})$ avec $\psi(\mathbf{r}+\mathbf{R})=e^{i\mathbf{k}\cdot\mathbf{R}}\psi(\mathbf{r})$ \forall \mathbf{R} dans le reseaux .

Preuve: Soit T_R etre l'opérateur de translation. Puis,

$$T_R H \psi(r) = H(r+R) \psi(r+R) = H(r) \psi(r+R) = H(r) T_R \psi(r) \rightarrow [H, T_R] = 0$$

Alors, il ya vecteurs propres simultanés: $H \psi = E \psi$ $T_R \psi = c(R) \psi$

C'est evident que $T_{R_1}T_{R_2} = T_{R_2}T_{R_1} = T_{R_1+R_2}$

donc
$$c(\mathbf{R}_1)c(\mathbf{R}_2)=c(\mathbf{R}_1+\mathbf{R}_2)$$

$$c(\mathbf{R}_{n}) = c(n_{1}\mathbf{a}_{1} + n_{2}\mathbf{a}_{2} + n_{3}\mathbf{a}_{3}) = c(\mathbf{a}_{1})^{n_{1}}c(\mathbf{a}_{2})^{n_{2}}c(\mathbf{a}_{3})^{n_{3}} = e^{2\pi i(n_{1}x_{1} + n_{2}x_{2} + n_{3}x_{3})}, \qquad x_{j} = \frac{\ln c(\mathbf{a}_{j})}{2\pi i}$$

$$= e^{i(x_{1}\mathbf{b}_{1} + x_{2}\mathbf{b}_{2} + x_{3}\mathbf{b}_{3}) \cdot \mathbf{R}_{n}}, \qquad \mathbf{a}_{i} \cdot \mathbf{b}_{j} = 2\pi \delta_{ij}$$

Conditions à la limite periodique

$$\psi(\mathbf{r} + N_i \mathbf{a}_i) = \psi(\mathbf{r}), \quad i = 1,2,3$$

 $\Rightarrow x_i = \frac{m_i}{N_i}, \quad m_i \text{ nombre entier } \leq N_i$

2. Tight-binding fonctions de base

 $\Phi_{jk}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}_n \in B}^{N} e^{i\mathbf{k} \cdot \mathbf{R}_n} \phi_j(\mathbf{r} - \mathbf{R}_n), \text{ o } \dot{\mathbf{u}} \phi_j \text{ sont les fonctions atomique est } \mathbf{k} \in \tilde{B}$

Vérification:
$$\Phi_{jk}(\mathbf{r}+\mathbf{R}_m) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}_n \in B}^{N} e^{i\mathbf{k}\cdot\mathbf{R}_n} \Phi_j(\mathbf{r}+\mathbf{R}_m-\mathbf{R}_n)$$
$$= \frac{1}{\sqrt{N}} \sum_{\mathbf{R}_l \in B}^{N} e^{i\mathbf{k}\cdot(\mathbf{R}_l+\mathbf{R}_m)} \Phi_j(\mathbf{r}-\mathbf{R}_l)$$
$$= \Phi_{jk}(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{R}_m}$$

3. Tight-binding fonction d'onde

$$\psi(\mathbf{r};\mathbf{k}) = \sum_{j=1}^{n} C_{j}(\mathbf{k}) \Phi_{j\mathbf{k}}(\mathbf{r})$$

4. <u>Tight-binding éléments de matrice</u>

$$H_{jj'}(\mathbf{k}) = \langle \Phi_{j\mathbf{k}} | H | \Phi_{j'\mathbf{k}} \rangle$$
 transfer integral matrix $S_{jj'}(\mathbf{k}) = \langle \Phi_{j\mathbf{k}} | \Phi_{j'\mathbf{k}} \rangle$ overlap integral matrix

Principe de variation pour l'équation de Schrodinger.

Theorem: La fonctionelle
$$E[\psi] \equiv \frac{\int d\mathbf{r} \frac{h^2}{2m} (\nabla \psi(\mathbf{r}))^2 + V(\mathbf{r}) |\psi(\mathbf{r})|^2}{\int d\mathbf{r} |\psi(\mathbf{r})|^2}$$

est minimisé pour <u>tous</u> function d'onde qui satisfie l'équation de Schrodinger.

D'après Ashcroft et Mermin, "Solid State Physics", HRW, 1976

Preuve:

Définir
$$F[\psi, \phi] = \int d\mathbf{r} \left(\frac{h^2}{2m} (\nabla \psi(\mathbf{r})) \cdot (\nabla \phi^*(\mathbf{r})) + V(\mathbf{r}) \psi(\mathbf{r}) \phi^*(\mathbf{r}) \right)$$

 $(\psi, \phi) = \int d\mathbf{r} [\psi(\mathbf{r}) \phi^*(\mathbf{r})]$

Soit $H \psi = E \psi$ il s'ensuite que

$$F[\psi+\delta\psi,\psi+\delta\psi] = E[(\psi,\psi)+(\psi,\delta\psi)+(\delta\psi,\psi)]+O(\delta\psi)^{2}$$
$$(\psi+\delta\psi,\psi+\delta\psi) = (\psi,\psi)+(\psi,\delta\psi)+(\delta\psi,\psi)+O(\delta\psi)^{2}$$

$$E[\psi + \delta \psi] = \frac{F[\psi + \delta \psi, \psi + \delta \psi]}{(\psi + \delta \psi, \psi + \delta \psi)} = E + O(\delta \psi)^{2}$$

5. <u>Tight-binding energies</u>

Parce-que le calcul "tight-binding" est faite avec un ensemble restreint de fonctions de base, il s'agit d'un ansatz et non un calcul complet. Alors, on emploi le principe de variation:

$$E_{0}(\mathbf{k}) \leq E_{TB}(\mathbf{k}) = \frac{\langle \Psi_{\mathbf{k}} | H | \Psi_{\mathbf{k}} \rangle}{\langle \Psi_{\mathbf{k}} | \Psi_{\mathbf{k}} \rangle} = \frac{\sum_{j,j'} \overline{C}_{j} H_{jj'}(\mathbf{k}) C_{j'}}{\sum_{j,j'} \overline{C}_{j} S_{jj'}(\mathbf{k}) C_{j'}}$$

Minimizer:

$$\frac{\partial E_{TB}(\mathbf{k})}{\partial \bar{C}_{j}} = 0 = \frac{\sum_{j} H_{jj'}(\mathbf{k}) C_{j'}}{\sum_{j,j'} \bar{C}_{j} S_{jj'}(\mathbf{k}) C_{j'}} - \frac{\sum_{j,j'} \bar{C}_{j} H_{jj'}(\mathbf{k}) C_{j'}}{\sum_{j,j'} \bar{C}_{j} S_{jj'}(\mathbf{k}) C_{j'}} \frac{\sum_{j} S_{jj'}(\mathbf{k}) C_{j'}}{\sum_{j,j'} \bar{C}_{j} S_{jj'}(\mathbf{k}) C_{j'}}$$

$$\Rightarrow \sum_{j} H_{jj'}(\mathbf{k}) C_{j'} = \lambda(\mathbf{k}) \sum_{j} S_{jj'}(\mathbf{k}) C_{j'}$$

$$C_{j} \neq 0 \Rightarrow \det(H(\mathbf{k}) - \lambda(\mathbf{k}) S(\mathbf{k})) = 0$$

Example: trans-polyacétylène

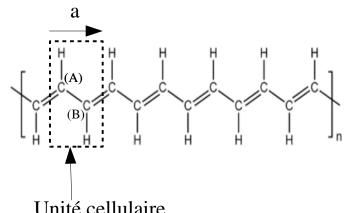
Vecteur de reseau: $\mathbf{a} = a \hat{\mathbf{x}}$

Vecteur de reseau reciproque: $b = \frac{2\pi}{a} \hat{x}$

Example: *trans*-polyacétylène

Vecteur de reseau: $a = a \hat{x}$

Vecteur de reseau reciproque: $b = \frac{2\pi}{3} \hat{x}$



Unité cellulaire

sp² hybridization \rightarrow σ -liens dans le plan, π -lien (p-p) dehors le plan

On cherche les niveau pour les π -electrons:

$$\begin{split} \Phi_{A}(\boldsymbol{r};\boldsymbol{k}) &= \frac{1}{\sqrt{N}} \sum_{\boldsymbol{R}_{n} \in B}^{N} e^{i\boldsymbol{k} \cdot \boldsymbol{R}_{n}} \varphi_{2 p_{z}}(\boldsymbol{r} - \boldsymbol{R}_{n}) = \frac{1}{\sqrt{N}} \sum_{n=0}^{N} e^{ikna} \varphi_{2 p_{z}}(\boldsymbol{r} - na \,\hat{\boldsymbol{x}}) \\ \Phi_{B}(\boldsymbol{r};\boldsymbol{k}) &= \frac{1}{\sqrt{N}} \sum_{\boldsymbol{R}_{n} \in B}^{N} e^{i\boldsymbol{k} \cdot (\boldsymbol{R}_{n} + \boldsymbol{r}_{AB})} \varphi_{2 p_{z}}(\boldsymbol{r} - \boldsymbol{R}_{n} + \boldsymbol{r}_{AB}) = \frac{1}{\sqrt{N}} \sum_{n=0}^{N} e^{ikna + ikR_{x}} \varphi_{2 p_{z}}(\boldsymbol{r} - \boldsymbol{r}_{AB} - na \,\hat{\boldsymbol{x}}) \end{split}$$

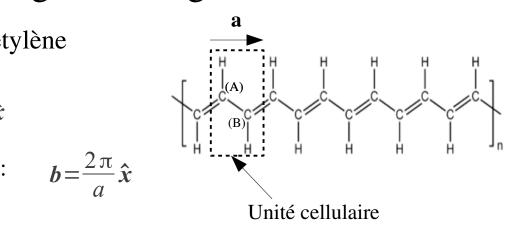
$$\mathbf{k} = \frac{2\pi}{a} \frac{p_i}{N} \hat{\mathbf{x}}, \quad 0 \le p_i < N$$
 $0 \le k \le 2\pi/a$

Example: trans-polyacétylène

Vecteur de reseau: $\mathbf{a} = a \hat{\mathbf{x}}$

Vecteur de reseau reciproque:

$$\boldsymbol{b} = \frac{2\pi}{a} \hat{\boldsymbol{x}}$$



$$\Phi_{A}(\mathbf{r};\mathbf{k}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}_{n} \in B}^{N} e^{i\mathbf{k}\cdot\mathbf{R}_{n}} \phi_{2p_{z}}(\mathbf{r} - \mathbf{R}_{n}) = \frac{1}{\sqrt{N}} \sum_{n=0}^{N} e^{ikna} \phi_{2p_{z}}(\mathbf{r} - na\,\mathbf{\hat{x}})$$

$$\Phi_{B}(\mathbf{r};\mathbf{k}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}_{n} \in B}^{N} e^{i\mathbf{k}\cdot(\mathbf{R}_{n}+\mathbf{r}_{AB})} \Phi_{2p_{z}}(\mathbf{r}-\mathbf{R}_{n}+\mathbf{r}_{AB}) = \frac{1}{\sqrt{N}} \sum_{n=0}^{N} e^{ikna+ikR_{x}} \Phi_{2p_{z}}(\mathbf{r}-\mathbf{r}_{AB}-na\,\hat{\mathbf{x}})$$

$$H_{AA}(\mathbf{k}) = \langle \Phi_{A\mathbf{k}} | H | \Phi_{A\mathbf{k}} \rangle = \frac{1}{N} \sum_{n,m=0}^{N} e^{ik(n-m)a} \langle \Phi_{2p_z}(\mathbf{r} - ma\mathbf{\hat{x}}) | H | \Phi_{2p_z}(\mathbf{r} - na\mathbf{\hat{x}}) \rangle$$

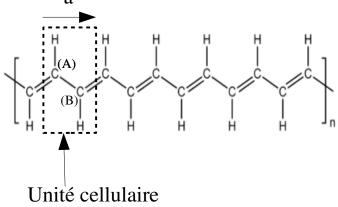
$$= \langle \Phi_{p_z}(\mathbf{r}) | H | \Phi_{p_z}(\mathbf{r}) \rangle + \underbrace{\frac{1}{N} \sum_{n \neq m}^{N} e^{ik(n-m)a} \langle \Phi_{2p_z}(\mathbf{r} - ma\mathbf{\hat{x}}) | H | \Phi_{2p_z}(\mathbf{r} - na\mathbf{\hat{x}}) \rangle}_{\text{faible}}$$

 $\approx \epsilon_{2n}$

Example: trans-polyacétylène

Vecteur de reseau: $a = a \hat{x}$

Vecteur de reseau reciproque: $b = \frac{2\pi}{a} \hat{x}$



$$H_{AB}(k\,\hat{\boldsymbol{x}}) = \langle \Phi_{Ak} | H | \Phi_{Bk} \rangle = \frac{1}{N} \sum_{n,m=0}^{N} e^{ik(n-m)+ik\,r_{AB,x}} \langle \Phi_{2\,p_z}(\boldsymbol{r}-ma\,\hat{\boldsymbol{x}}) | H | \Phi_{2\,p_z}(\boldsymbol{r}-\boldsymbol{r}_{AB}-na\,\hat{\boldsymbol{x}}) \rangle$$

$$= e^{ik\,R_x} \langle \Phi_{2\,p_z}(\boldsymbol{r}) | H | \Phi_{2\,p_z}(\boldsymbol{r}-\boldsymbol{r}_{AB}) \rangle + e^{ik(r_{AB,x}-a)} \langle \Phi_{2\,p_z}(\boldsymbol{r}) | H | \Phi_{2\,p_z}(\boldsymbol{r}-\boldsymbol{r}_{AB}+\boldsymbol{a}) \rangle$$

$$+ \frac{1}{N} \sum_{m=0;\,n\neq0,-1}^{N} e^{ik\,r_{AB,x}+ik(n-m)} \langle \Phi_{2\,p_z}(\boldsymbol{r}-ma\,\hat{\boldsymbol{x}}) | H | \Phi_{2\,p_z}(\boldsymbol{r}-\boldsymbol{r}_{AB}-na\,\hat{\boldsymbol{x}}) \rangle$$
faible

if
$$r_{ABx} = a/2$$

 $\approx 2t \cos(ka/2), \quad t = \langle \phi_{2p_z}(\mathbf{r}) | H | \phi_{2p_z}(\mathbf{r} - a/2) \rangle$

Example: trans-polyacétylène

Vecteur de reseau:

$$a = a \hat{x}$$

Vecteur de reseau reciproque:

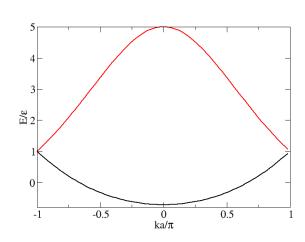
$$b = \frac{2\pi}{a} \hat{x}$$
Unité cellulaire

$$H_{ij} = \begin{pmatrix} \epsilon_{2p} & 2t\cos(ka/2) \\ 2t\cos(ka/2) & \epsilon_{2p} \end{pmatrix} \qquad S_{ij} = \begin{pmatrix} 1 & 2s\cos(ka/2) \\ 2s\cos(ka/2) & 1 \end{pmatrix}$$

$$0 = det \begin{pmatrix} \epsilon_{2p} - E & 2(t - sE)\cos(ka/2) \\ 2(t - sE)\cos(ka/2) & \epsilon_{2p} - E \end{pmatrix}$$

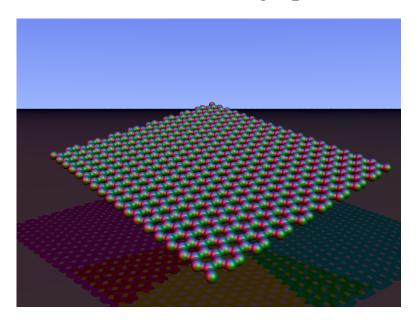
$$= (\epsilon_{2p} - E)^2 - 4(t - sE)^2 \cos^2(ka/2)$$

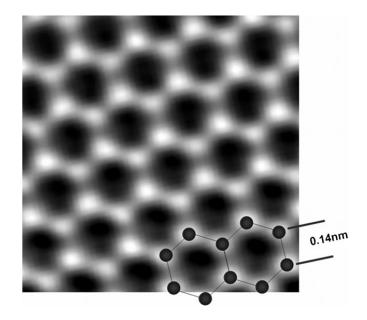
$$E_{\pm}(k) = \frac{\epsilon_{2p} \pm 2t \cos(ka/2)}{1 \pm 2s \cos(ka/2)}, -\frac{\pi}{a} < k < \frac{\pi}{a}$$



Graphene 1

graphène = un seul feuillet de graphite

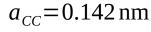


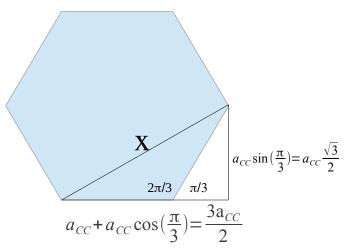


Structure électronique d'un atome de carbone = $1s^2 2s^2 2p^2$ coeur = $1s^2$ 4 électrons de valence = $2s^2 2p^2$

Chaque atome de carbone offre 3 orbitales atomiques sp^2 et une orbitale $2p_z$ Les orbitales atomiques sp^2 forment les liens σ Les orbitales atomiques $2p_z$ forment les liens π

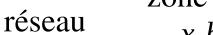
Graphene 2

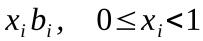


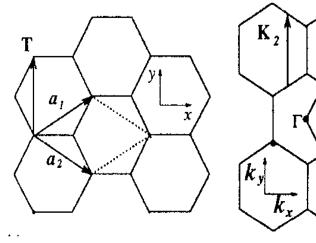


$$x = \sqrt{\frac{9 a_{CC}^2}{4} + \frac{3 a_{CC}^2}{4}} = \sqrt{3} a_{CC}$$

zone de Brillouin







$$a_{1,2} = \left(\frac{\sqrt{3}}{2}a, \pm \frac{a}{2}\right)$$
 $b_{1,2} = \left(\frac{2\pi}{\sqrt{3}a}, \pm \frac{2\pi}{a}\right)$

$$a = ||\boldsymbol{a}_1|| = ||\boldsymbol{a}_2|| = 0.246 \,\mathrm{nm}$$

R. Saito, G. Dresselhaus & M. S. Dresselhaus,

Physical Properties of Carbon Nanotubes (Imperial College Press, London, 1998)

Graphene Tight-binding: π-bands

$$H_{AA} = H_{BB} = \epsilon_{2p}$$

$$H_{AB} = t(e^{i \mathbf{k} \cdot \mathbf{R}_1} + e^{i \mathbf{k} \cdot \mathbf{R}_2} + e^{i \mathbf{k} \cdot \mathbf{R}_3})$$

$$\equiv tf(\mathbf{k})$$

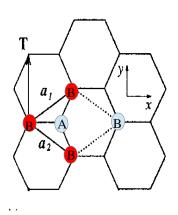
$$f(\mathbf{k}) = e^{-ik_x a/\sqrt{3}} + 2e^{ik_x a/2\sqrt{3}}\cos(\frac{k_y a}{2})$$

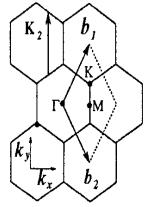
$$S_{AB} = sf(\mathbf{k})$$

$$E_{\pi\pi} = \frac{\epsilon_{2p} \pm tw(\mathbf{k})}{1 \pm sw(\mathbf{k})}$$

réseau

zone de Brillouin





$$\boldsymbol{a}_{1,2} = \left(\frac{\sqrt{3}}{2}a, \pm \frac{a}{2}\right)$$

$$\boldsymbol{a}_{1,2} = \left(\frac{\sqrt{3}}{2}a, \pm \frac{a}{2}\right)$$
 $\boldsymbol{b}_{1,2} = \left(\frac{2\pi}{\sqrt{3}a}, \pm \frac{2\pi}{a}\right)$

$$w(\mathbf{k}) = \sqrt{\left[f(\mathbf{k})\right]^2} = \sqrt{1 + 4\cos\left(\frac{\sqrt{3}k_x a}{2}\right)\cos\left(\frac{k_y a}{2}\right) + 4\cos^2\left(\frac{k_y a}{2}\right)}$$

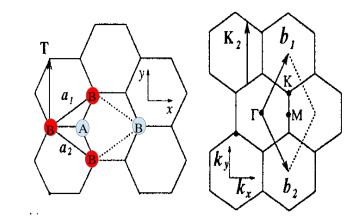
 $s = 0 \Leftrightarrow$ "Slater-Koster approximation"

R. Saito, G. Dresselhaus & M. S. Dresselhaus, Physical Properties of Carbon Nanotubes (Imperial College Press, London, 1998)

Graphene Tight-binding: σ-bands

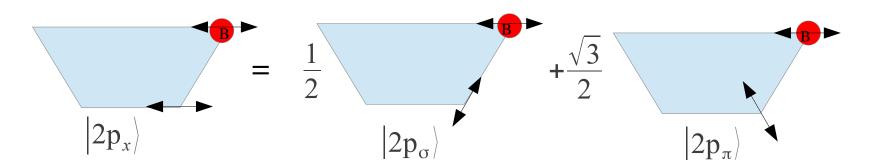
Basis: $(2s^{A}, 2p_{x}^{A}, 2p_{y}^{A}, 2s^{B}, 2p_{x}^{B}, 2p_{y}^{B})$

$$H_{AA} = \begin{pmatrix} \epsilon_{2s} & 0 & 0 \\ 0 & \epsilon_{2p} & 0 \\ 0 & 0 & \epsilon_{2p} \end{pmatrix}, \quad S_{AA} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



Alors,
$$|2p_x\rangle = \cos\left(\frac{\pi}{3}\right)|2p_\sigma\rangle + \cos\left(\frac{2\pi}{3}\right)|2p_\pi\rangle = \frac{1}{2}|2p_\sigma\rangle + \frac{\sqrt{3}}{2}|2p_\pi\rangle$$

Car l'état de l = 1 se transforme comme un vecteur sous rotations.



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Graphene Tight-binding: σ-bands

Basis:
$$(2s^{A}, 2p_{x}^{A}, 2p_{y}^{A}, 2s^{B}, 2p_{x}^{B}, 2p_{y}^{B})$$

$$|2p_{x}\rangle = \frac{1}{2}|2p_{\sigma}\rangle + \frac{\sqrt{3}}{2}|2p_{\pi}\rangle$$

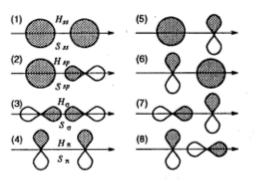
$$= \left(\frac{1}{2}\right)^{2} + \left(\frac{\sqrt{3}}{2}\right)^{2}$$

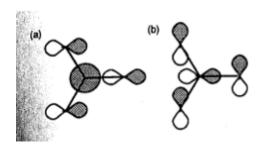
$$+ \left(\frac{1}{2}\right)\left(\frac{\sqrt{3}}{2}\right) + \left(\frac{1}{2}\right)\left(\frac{\sqrt{3}}{2}\right) + \left(\frac{1}{2}\right)\left(\frac{\sqrt{3}}{2}\right)$$

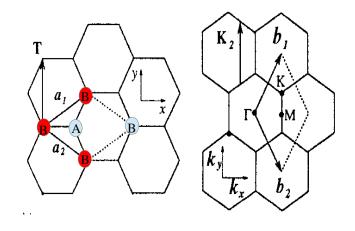
$$= \left(\frac{1}{4}H_{\sigma\sigma} + \frac{3}{4}H_{\pi\pi}\right)e^{i(k_{x}R_{x}+k_{y}R_{y})} = \left(\frac{1}{4}H_{\sigma\sigma} + \frac{3}{4}H_{\pi\pi}\right)e^{i\left(\frac{k_{x}a}{2\sqrt{3}} + \frac{k_{y}a}{2}\right)}, \quad H_{\sigma\pi} = 0$$

Graphene Tight-binding: σ-bands

Basis: $(2s^{A}, 2p_{x}^{A}, 2p_{y}^{A}, 2s^{B}, 2p_{x}^{B}, 2p_{y}^{B})$





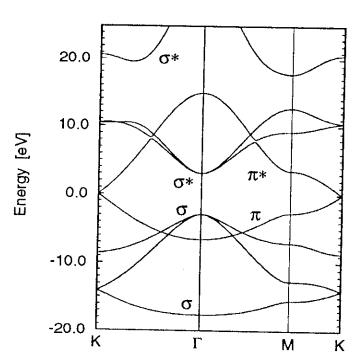


R. Saito, G. Dresselhaus & M. S. Dresselhaus, *Physical Properties of Carbon Nanotubes* (Imperial College Press, London, 1998)

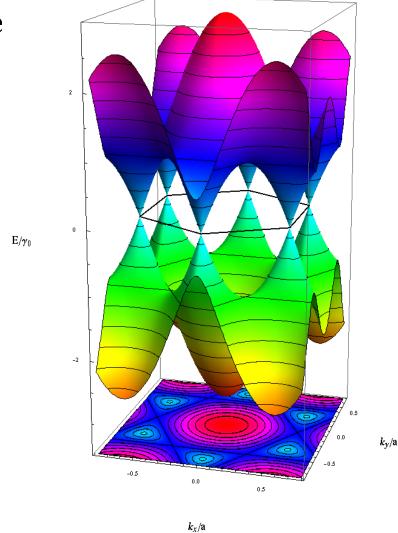
$$\langle 2 p_x^B | H | 2 p_x^A \rangle = \frac{1}{4} (H_{\sigma\sigma} + 3 H_{\pi\pi}) e^{i \left(\frac{k_x a}{2\sqrt{3}} + \frac{k_y a}{2}\right) i}$$

$$\langle 2 p_x^B | H | 2 p_y^A \rangle = \frac{i\sqrt{3}}{2} (H_{\sigma\sigma} + H_{\pi\pi}) e^{-ik_x a/2\sqrt{3}} \sin \frac{k_y a}{2}$$

$$\langle 2s^{B}|H|2p_{x}^{A}\rangle = H_{sp}\left(-e^{ik_{x}a/\sqrt{3}} + e^{ik_{x}a/(2\sqrt{3})}\right)$$



Graphene



component	H (eV)	S
ss	-6.7969	0.212
sp	-5.580	0.102
σσ	-5.037	0.146
ππ	-3.033	0.129
ε _s -ε _p	-8.868	

R. Saito, G. Dresselhaus & M. S. Dresselhaus, Physical Properties of Carbon Nanotubes (Imperial College Press, London, 1998)

Carbon Structures

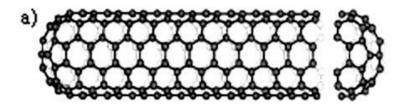
- Introduction
 - Histoire
 - Carbone
 - Hybridization
- Tight-binding calculations
 - Principe
 - Example: trans
 - Graphene π -liens
 - Graphene σ-liens
- Structure de carbone nanotubes
- Structure Electronique
- Synthesis
- Transport
- Multi-walled nanotubes

Nanotube structure

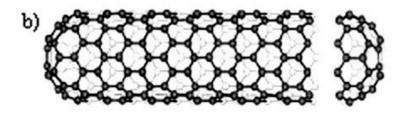
Rayon: 0.7-10 nm (pour la plupart < 2nm)

Longeur: 1-100 µm

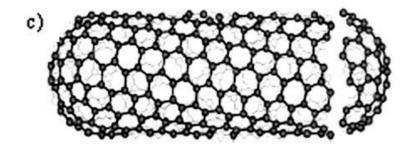
Chapeau ("cap"): fullerene avec 6 pentagons +plusieurs hexagon



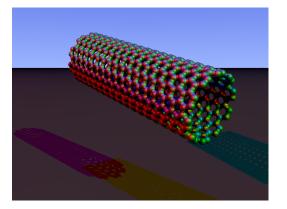
"armchair"

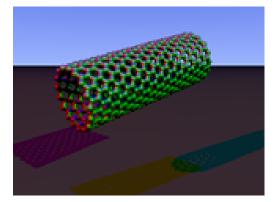


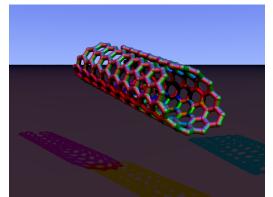
"zigzag"



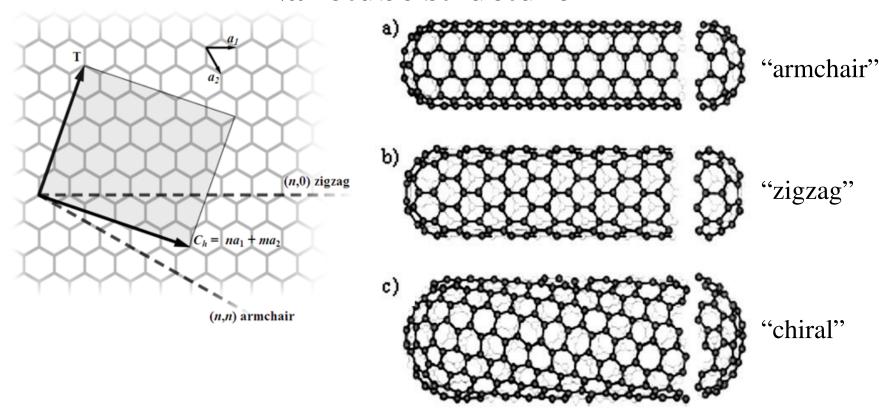
"chiral"







Nanotube structure

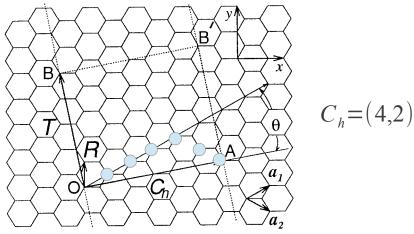


Division fondamentale:

- "achiral": possédant une symétrie miroir (seulement la "zigzag" et la "armchair"
- "chiral": pas de symétrie mirroir

Nanotube structure: vecteur chiral

Les points O,A et B,B' sont identifiés



Ligne OA est la vecteur chiral

Vecteur chiral:
$$C_h = n a_1 + m a_2 \Leftrightarrow (n, m), \quad 0 \le |m| \le n$$

N.B.:
$$a_1 \cdot a_1 = a_2 \cdot a_2 = a^2$$
, $a_1 \cdot a_2 = \frac{a^2}{2}$ $a = 0.144 \text{ nm} \times \sqrt{3} = 0.249 \text{ nm}$

Diamètre:
$$d_t = L/\pi = ||C_h||/\pi = \sqrt{n^2 + m^2 + nm} \times 0.079 \text{ nm}$$

Angle chiral:
$$\cos \theta = \frac{\boldsymbol{C}_h \cdot \boldsymbol{a}_1}{|\boldsymbol{C}_h||\boldsymbol{a}_1|} = \frac{2n + m}{2\sqrt{n^2 + m^2 + nm}}$$

R. Saito, G. Dresselhaus & M. S. Dresselhaus, *Physical Properties of Carbon Nanotubes* (Imperial College Press, London, 1998)

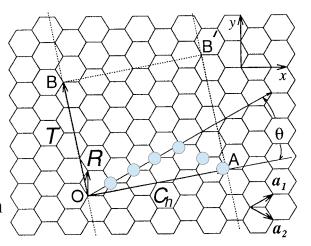
Nanotube structure : vecteur de translation

Les points O,A et B,B' sont identifiés

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$$C_h = (4,2)$$

Ligne OB est la vecteur de translation



La vecteur de translation définit le unité cellulaire le long de la nanotube.

Vecteur de translation: $T = t_1 a_1 + t_2 a_2 \Leftrightarrow (n, m)$

$$\boldsymbol{C}_h \cdot \boldsymbol{T} = 0 \Leftrightarrow t_1(2n+m) + t_2(n+2m) = 0 \Leftrightarrow t_1 = \frac{2m+n}{\gcd(2m+n, m+2n)}, \quad t_2 = -\frac{m+2n}{\gcd(2m+n, m+2n)}$$

2N

Dans l'example:
$$\gcd(8,10) = 2 \Leftrightarrow t_1 = \frac{8}{2} = 4, \ t_2 = -\frac{10}{2} = -5$$

Le nombre d'hexagônes par unite cellulaire est:
$$N = \frac{2L^2}{a^2 \gcd(2m+n, m+2n)}$$

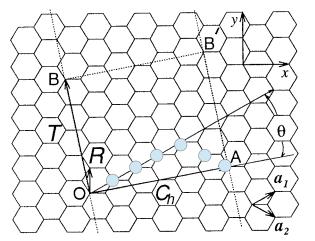
Le nombre d'atoms par unite cellulaire est:

Nanotube structure : vecteur de symétrie

Les points O,A et B,B' sont identifiés

R. Saito, G. Dresselhaus & M. S. Dresselhaus, *Physical Properties of Carbon Nanotubes* (Imperial College Press, London, 1998)

$$C_h = (4,2)$$



Les positions de tous les atomes dans la maille élémentaire peut être représenté en tant que n**R** avec les bords periodique.

Ligne OR est la vecteur de symétrie.

R est la site d'atom avec la plus petite non-zero projection sur C.

$$R = p a_1 + q a_2 \Leftrightarrow (p, q)$$
 $R \cdot \frac{C_h}{|C_h|} = \frac{p(m + \frac{n}{2}) + q(\frac{m}{2} + n)}{\sqrt{m^2 + n^2 + mn}} = \frac{p(2m + n) + q(m + n2)}{2\sqrt{m^2 + n^2 + mn}}$

La valeur minimale possible pour le numérateur est

$$p(2m+n)+q(m+2n)=\gcd(2m+n,2n+m)$$

Aussi
$$0 < \frac{R \cdot T}{|T|} < T \Leftrightarrow 0 < (p(t_1 + t_2/2) + q(t_2 + t_1/2)) < t_1^2 + t_2^2 + t_1 t_2$$

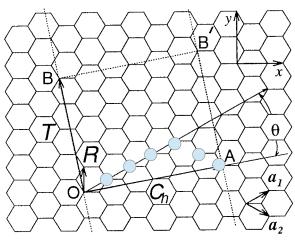
 $\Leftrightarrow 0 < mp - nq \le N$

Nanotube structure : vecteur de symétrie

Les points O,A et B,B' sont identifiés

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$$C_h = (4,2)$$



Ligne OR est la vecteur de symétrie.

R est la site d'atom avec la plus petite non-zero projection sur C.

$$\mathbf{R} = p \, \mathbf{a}_1 + q \, \mathbf{a}_2 \Leftrightarrow (p, q) \qquad \qquad \frac{\mathbf{R} \cdot \mathbf{C}_h}{|\mathbf{C}_h|} = \frac{|\mathbf{R} \times \mathbf{T}|}{|\mathbf{T}|} = \frac{(t_1 q - t_2 p)(\mathbf{a}_1 \times \mathbf{a}_2)}{|\mathbf{T}|} \Rightarrow t_1 q - t_2 p = 1$$

Aussi
$$0 < \frac{R \cdot T}{|T|^2} = \frac{|C_h \times R|}{|C||T|} = \frac{mp - nq}{N} < 1$$

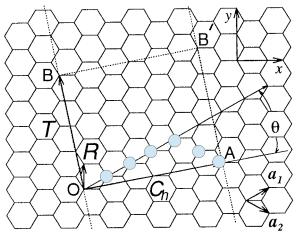
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Nanotube structure : vecteur de symétrie

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$$C_h = (4,2)$$



Ligne OR est la vecteur de symétrie.

Les positions de tous les atomes dans la maille élémentaire peut être représenté en tant que n**R** avec les bords periodique car

$$NR \cdot \frac{C_h}{|C_h|} = N \frac{|R \times T|}{|T|} = N \frac{|a_1 \times a_2|}{|T|} = |C_h|$$

$$R = p a_1 + q a_2 \Leftrightarrow (p, q)$$

$$t_1 q - t_2 p = 1$$

$$0 < mp - nq \le N$$

Structure

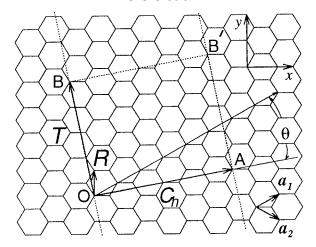
	zigzag	armchair	chiral
С	(n,0)	(n,n)	(n,m)
Т	(1,-2)	(1,-1)	$\left(\frac{2m+n}{d_R}, -\frac{2n+m}{d_R}\right)$
R	(1,-1)	(1,0)	W K
L/a	n	$\sqrt{3} n$	$\sqrt{m^2 + n^2 + nm}$
Т	$\sqrt{3}$	1	$\sqrt{3} L/d_R$
N	2n	2n	$2L^2/(a^2d_R)$

$$d_{R} = \gcd(2m+n, 2n+m)$$

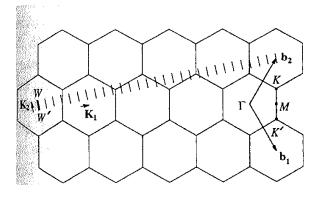
$$= \begin{pmatrix} d & \text{if } n-m \text{ is not a multiple of } 3d \\ 3d & \text{if } n-m \text{ is a multiple of } 3d \end{pmatrix}, \quad d = \gcd(n, m)$$

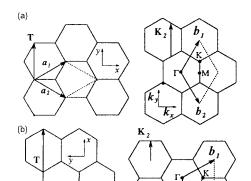
Nanotube structure

réseau



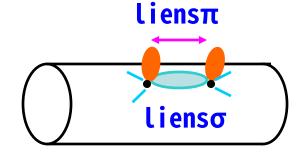
zone de Brillouin





« armchair » (n,n)

 \ll zigzag \gg (n,0)



R. Saito, G. Dresselhaus & M. S. Dresselhaus,

Physical Properties of Carbon Nanotubes (Imperial College Press, London, 1998)