Universidad de Oviedo

Grado en Física

Métodos Numéricos y sus Aplicaciones a la Física

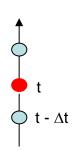
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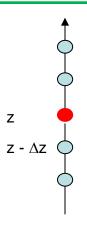
Diferencias Finitas en el dominio del tiempo

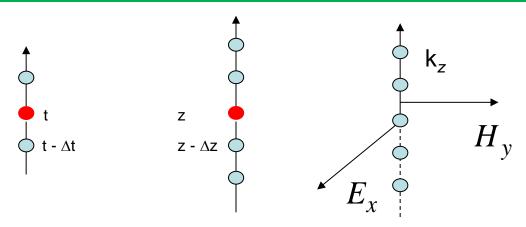
FDTD

$$\mu_0 \frac{\partial \vec{H}}{\partial t} = - \left(\nabla \times \vec{E} \right)$$

$$\varepsilon \frac{\partial \vec{E}}{\partial t} = (\nabla \times \vec{H}) - \vec{J}_{ex}$$

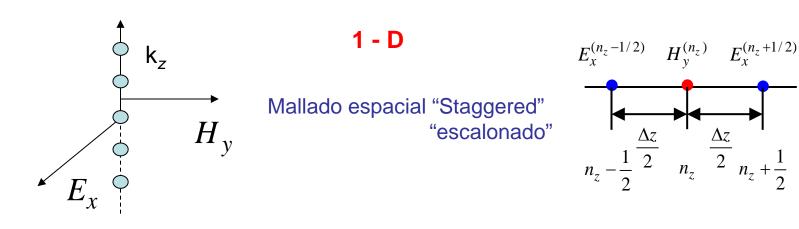






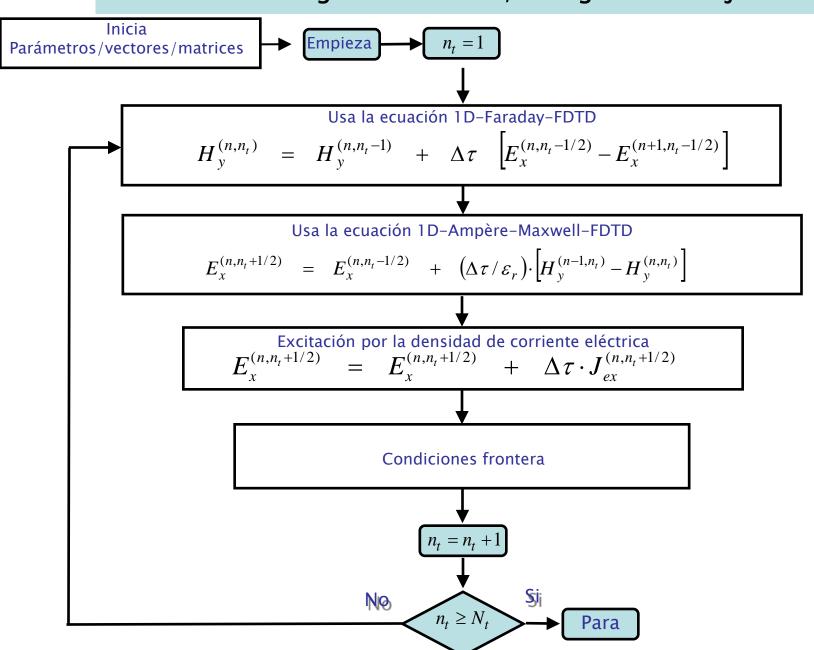
$$\frac{\partial H_{y}}{\partial t} = - \frac{1}{\mu_{o}} \frac{\partial E_{x}}{\partial z}$$

$$\varepsilon \frac{\partial E_x}{\partial t} = \frac{\partial H_y}{\partial z} - J_{ex}$$

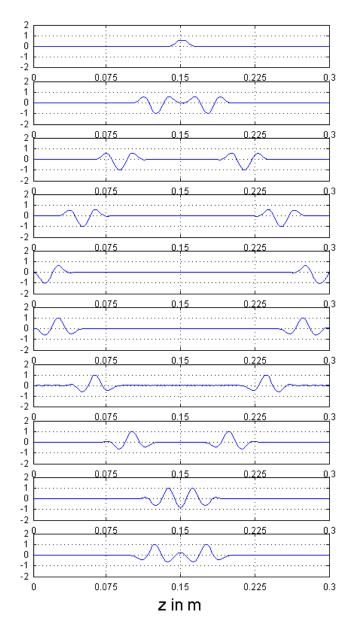


Algorítmo temporal tipo "Leap-frog" **Plano** "salto de rana" temporal $(n_z + 1)$ (n_z) (n_t) $\left(n_z + \frac{1}{2}\right)$

1-D Algoritmo FDTD / Diagrama de flujo

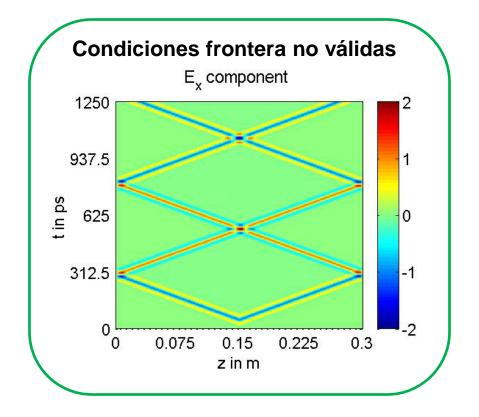


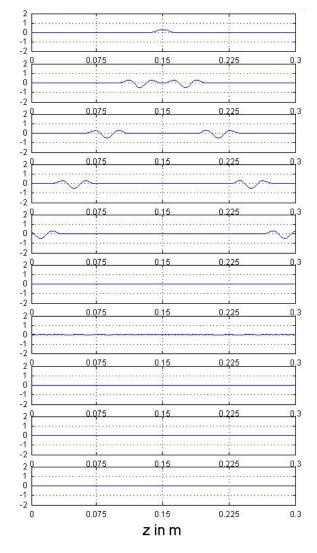
$E_x(z,t_0)$ a distintos t_0



Condiciones frontera

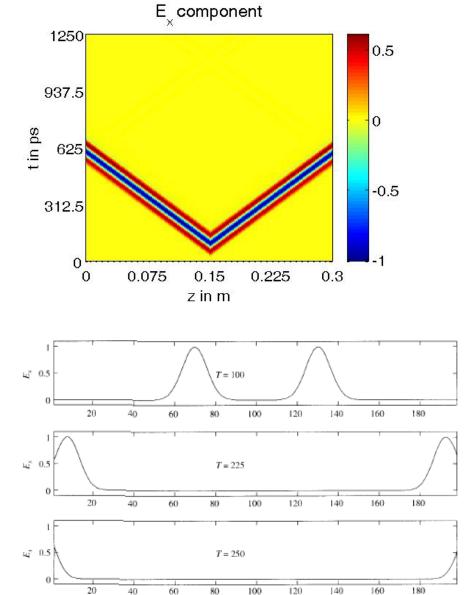
$$\left. \begin{array}{l}
 E_x^{(1,n_t)} = 0 \\
 E_x^{(N_z,n_t)} = 0
 \end{array} \right\} \quad 1 \le n_t \le N_t$$





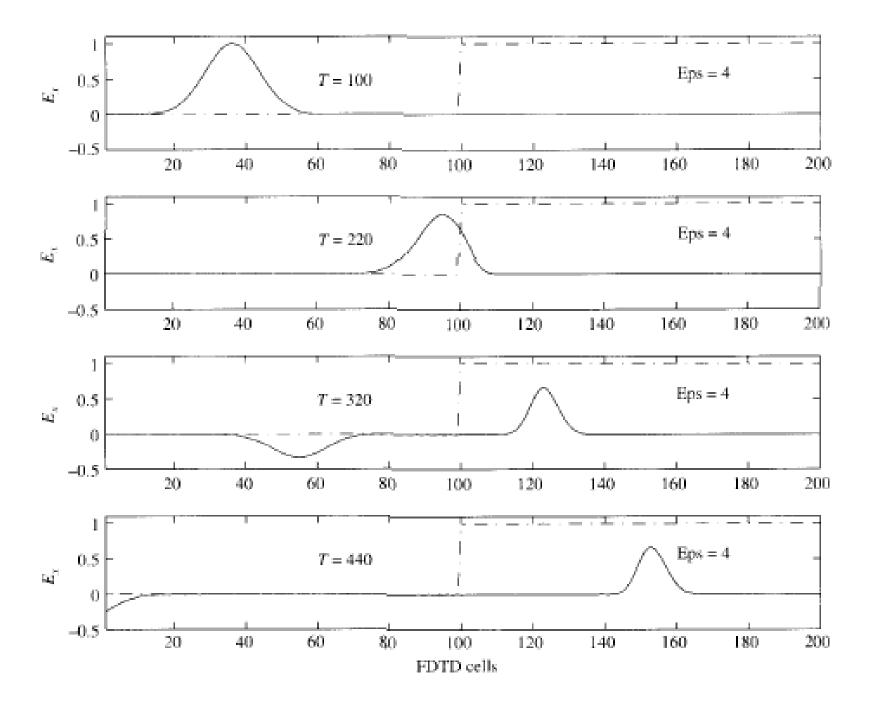
Si $c\Delta t = \Delta z/2$ una onda plana necesita dos pasos de tiempo, $2n_t$, para viajar a través de una celda de espaciado Δz

$$\left. \begin{array}{l}
 E_x^{(1,n_t)} = E_x^{(2,n_t-2)} \\
 E_x^{(N_z,n_t)} = E_x^{(N_z-1,n_t-2)}
 \end{array} \right\} \quad 1 \le n_t \le N_t$$



¡ Guarda $E_x(2)$, $E_x(N-1)$ para dos tiempos y almacénalo en $E_x(1)$, $E_x(N)$!

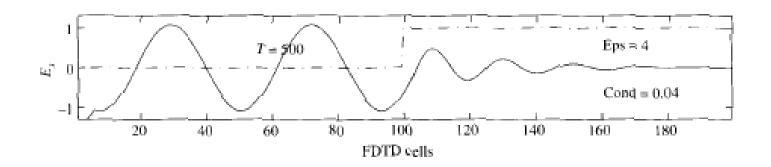
FDTD cells

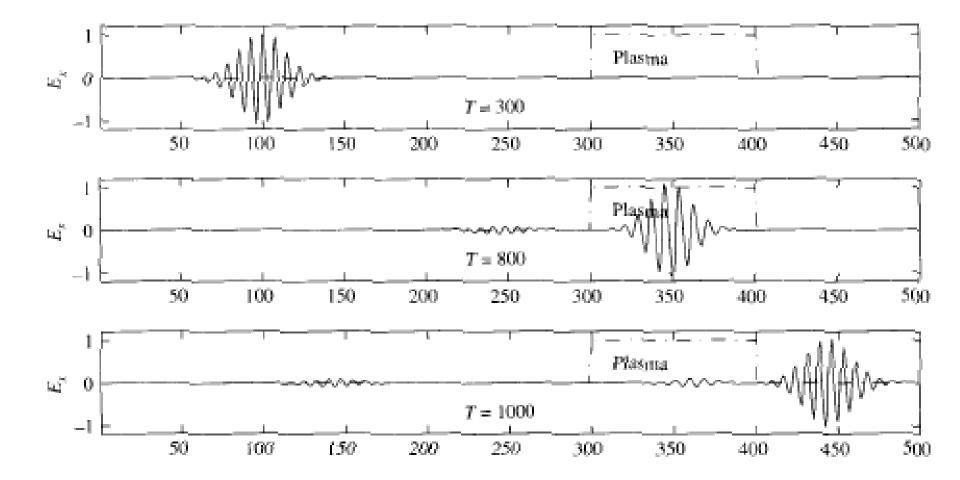


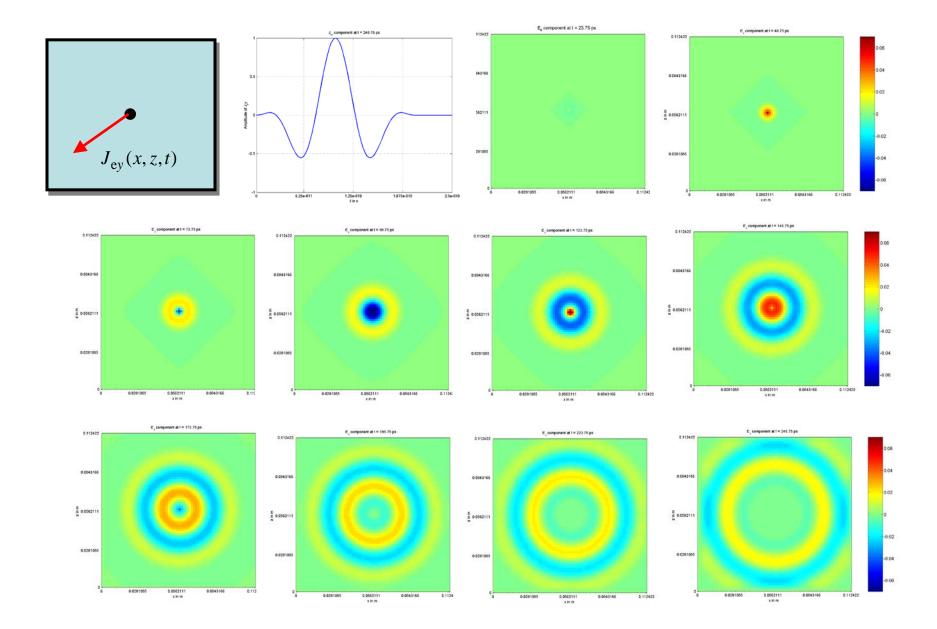
Tipos de fuentes blandas

Anchura de la celda

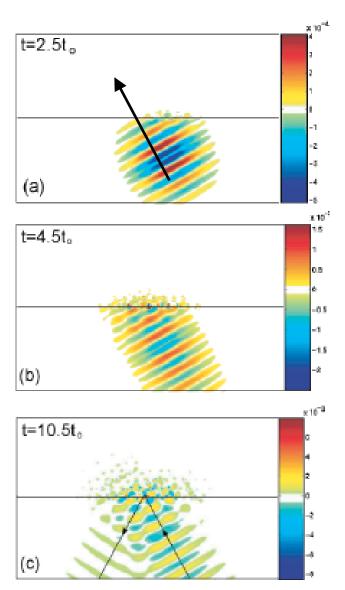
Para una buena resolución espacial $\Delta x = \lambda /10$

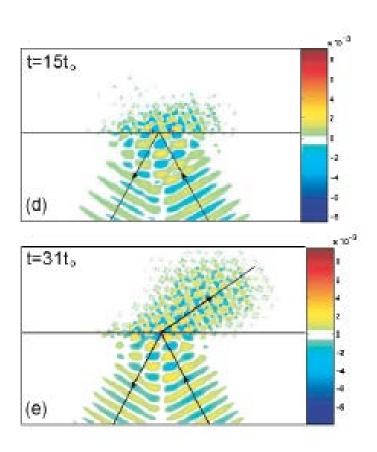


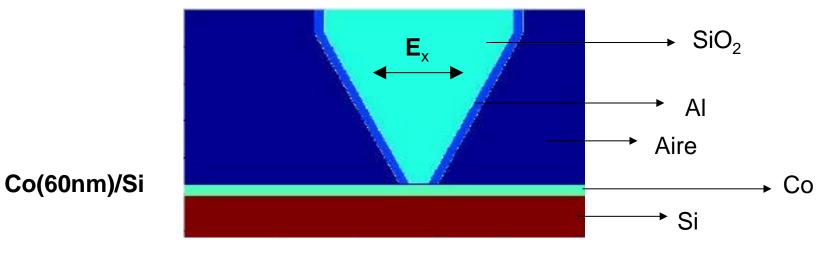




Refraction in Media with a Negative Refractive Index

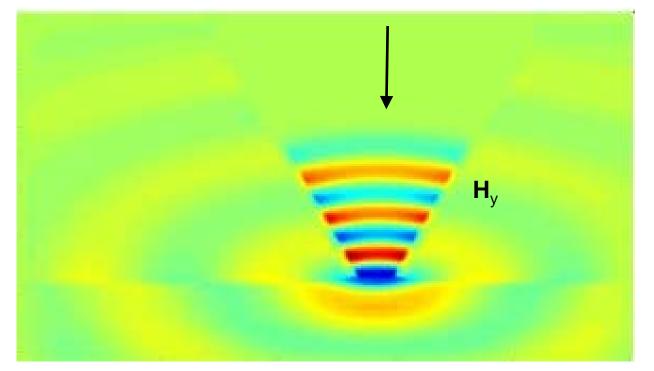






 $\lambda = 633 \text{ nm}$ abertura = 96 nm $d_{punta-muestra} = 20 \text{ nm}$

1100x1100 celdas² Área celda= 2x2 nm²



Evolución temporal de una función de onda

Método explícito (FDTD)

$$i\hbar\frac{\partial\psi(r,t)}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2\psi(r,t) + V(r,t)\cdot\psi(r,t)$$

$$\frac{\partial \psi(r,t)}{\partial t} = i \frac{\hbar}{2m} \nabla^2 \psi(r,t) - \frac{i}{\hbar} V(r,t) \cdot \psi(r,t),$$

$$\psi(r,t) = \psi_{\text{real}}(r,t) + i\psi_{\text{imag}}(r,t). \qquad \frac{\partial \psi_{\text{real}}(r,t)}{\partial t} = -\frac{\hbar}{2m} \cdot \nabla^2 \psi_{\text{imag}}(r,t) + \frac{1}{\hbar} V \psi_{\text{imag}}(r,t) \qquad [1]$$

$$\frac{\partial \psi_{\text{imag}}(r,t)}{\partial t} = \frac{\hbar}{2m} \cdot \nabla^2 \psi_{\text{real}}(r,t) - \frac{1}{\hbar} V \psi_{\text{real}}(r,t). \qquad [2]$$

$$\frac{\psi_{\text{real}}^{n}(k) - \psi_{\text{real}}^{n-1}(k)}{\Delta t} = -\frac{\hbar}{2m} \frac{\psi_{\text{imag}}^{n-1/2}(k+1) - 2\psi_{\text{imag}}^{n-1/2}(k) + \psi_{\text{imag}}^{n-1/2}(k-1)}{(\Delta z)^{2}} + \frac{1}{\hbar}V(k) \cdot \psi_{\text{imag}}^{n-1/2}(k), \quad [1]$$

Método semi-implícito (Crank-Nicolson)

Am. J. Phys. 50, 902 (1982);

$$i\hbar \frac{\partial}{\partial t} \Psi(x,t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x,t) + V(x)\Psi(x,t).$$

$$a = i \hbar \quad b = -a^{2}/2m$$

$$\frac{b}{(\Delta x)^{2}} \Psi(x_{i+1},t_{j}) + \left(\frac{2a}{\Delta t} - \frac{2b}{(\Delta x)^{2}} - V(x_{i})\right) \Psi(x_{i},t_{j}) + \frac{b}{(\Delta x)^{2}} \Psi(x_{i-1},t_{j}) + b \left[\Psi(x_{i+1},t_{j-1}) - 2\Psi(x_{i},t_{j-1}) + \Psi(x_{i+1},t_{j-1})\right]$$

$$= \left(\frac{2a}{\Delta t} + V(x_{i})\right) \Psi(x_{i},t_{j-1})$$