

Fig. 1. – Spatial distribution of the activator at  $t = 100$  time units in the presence of noise: (a) complex spiral dynamics for  $D = 0.80$  and (b) spiral meandering for  $D = 2.08$ . Noise strength is given by  $\sigma = 0.28$  in the two cases.

disappeared through the boundaries of the system. Therefore, the deterministic model either displays stable spiral waves or decays to the quiescent state after a short transient, depending on the initial conditions considered.

The situation is radically different in the presence of parametric noise. Two snapshots of the spatial distribution of the activator variable  $u_i(t)$  taken at  $t = 100$  time units are given in fig. 1, for two different values of the spatial coupling  $D > D_{\text{th}}$ . In both cases the system evolves from random initial conditions. Close to the propagation threshold (fig. 1a) the system exhibits a long-lived complex state which closely resembles the spiral-chaos regime found in deterministic systems with more complicated local dynamics [10]. This state has been observed to survive for at least  $t = 3000$  time units. Further above threshold (fig. 1b), the complexity of the dynamical state is reduced, consisting only of a small number of spirals, either single or coupled in pairs (as in the figure), whose core undergoes a slow meandering (drifting of spirals due to noise has already been observed and characterised in the complex Ginzburg-Landau

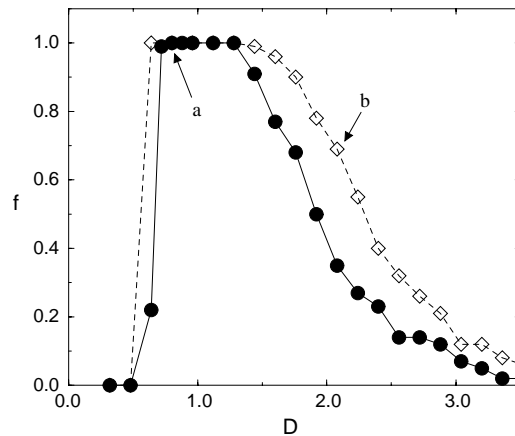


Fig. 2. – Fraction  $f$  of realizations (out of 100) that do not decay to the rest state after a reference time  $t_{\text{ref}} = 100$ , for increasing values of the diffusion coefficient  $D$ . Two noise strengths are considered:  $\sigma = 0.28$  (empty diamonds) and  $\sigma = 0.40$  (solid circles). Labels  $a$  and  $b$  correspond to the non-decaying situations of fig. 1.