Gromov-Witten theory and mirror symmetry

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Chapter 1

Gromov-Witten invariants

Roughly speaking, Gromov-Witten invariants is the number of algebraic curves (resp. pseduo-holomorphic curves) of a algebraic variety X (resp. complex manifold) passing through specific subvarieties (resp. submanifolds) under specific degree (homology class of the curves) and genus.

Definition 1.0.1. Let $\gamma_1, \ldots, \gamma_n \in H^*(X; \mathbb{Q})$ and let $\beta \in H^2(X; \mathbb{Q})$. The Gromov-Witten invariant of genus g degree β curves is

$$\langle \gamma_1 \dots \gamma_n \rangle_{g,n,\beta} = \int_{[\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir}} ev_1^*(\gamma_1) \cup \dots \cup ev_n^*(\gamma_n).$$

Here, a point in $[\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir}$ is $[f:C\to X,1,\ldots,n]$:

a map from the genus g curve C to the variety X modulo the automorphism of C.

The evaluation map $ev_i : [\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir} \to X$ is given by

$$ev_i([f:C\to X,1,\ldots,n])=f(i).$$

Let $\overline{\mathcal{M}}_{g,n}$ denote the moduli space (Deligne-Mumford stack) of genus g curves with n marked points, and let $\overline{\mathcal{C}}_{g,n}$ be the universal family of $\overline{\mathcal{M}}_{g,n}$.

1.1 Kontsevich's approach

Atiyah-Bott localization formula tells us: if there is a torus action $\mathbb{T} = (\mathbb{C}^*)^n$ on X, then the fixed points of torus action could tells us some properties of X.

By the classifying space theory, $B\mathbb{T} = (\mathbb{C}P^{\infty})^{\times n}$, so $H^*(B\mathbb{T}) = \mathbb{Q}[\lambda_1, \dots, \lambda_n]$. Let $\mathcal{R}_{\mathbb{T}} = \mathbb{Q}(\lambda_1, \dots, \lambda_n)$. Let $X_{\mathbb{T}} = E\mathbb{T} \times_{\mathbb{T}} X$, the equivariant cohomology of X is defined by

$$H_{\mathbb{T}}^*(X) := H^*(X_{\mathbb{T}}) = H^*(E\mathbb{T} \times_{\mathbb{T}} X),$$

so naturally, $H_{\mathbb{T}}^*(X)$ is a $H_{\mathbb{T}}^*(pt) = H^*(B\mathbb{T})$ -module. The localization of $H_{\mathbb{T}}^*(X)$ means $H_{\mathbb{T}}^*(X) \otimes \mathcal{R}_{\mathbb{T}}$.

Theorem (Atiyah-Bott). Let $X^{\mathbb{T}}$ be fixed locus of \mathbb{T} , let Z_j be a connection component of $X^{\mathbb{T}}$, and let N_j be the normal bundle of Z_j in X. Let $i_j: Z_j \to X$ and let $i_{j!}: H^*_{\mathbb{T}}(Z_j) \to H^*_{\mathbb{T}}(X)$ be the pushforward defined by the Gysin map. Let $\alpha \in H^*_{\mathbb{T}}(X) \otimes \mathcal{R}_{\mathbb{T}}$, we have

$$\alpha = \sum_{j} \frac{i_{j!} i_{j}^{*} \alpha}{Euler_{T}(N_{j})},$$

In particular,

$$\int_{X_{\mathbb{T}}} \alpha = \sum_{j} \int_{(Z_{j})_{\mathbb{T}}} \frac{i_{j}^{*} \alpha}{Euler_{T}(N_{j})}.$$

Kontsevich's approach is to apply Atiyah-Bott localization formula in $[\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir}$ so that we can simplify the computation. We can lift the \mathbb{T} action on X to $[\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir}$ in the following way: let $t \in \mathbb{T}$, $[f:C \to X,1,\ldots,n] \in [\overline{\mathcal{M}}_{g,n}(X,\beta)]^{vir}$, $x \in X$

$$(t \cdot f)(x) = f(t \cdot x).$$

In special, we will consider $\overline{\mathcal{M}}_{g,n}(\mathbb{P}^r,d)$ in this section. As claimed before, we need to find $[f:C\to X,1,\ldots,n]\in\overline{\mathcal{M}}_{g,n}(\mathbb{P}^r,d)^{\mathbb{T}}$. The fixed points of \mathbb{P}^r is

$${q_i = [0:0:\dots:1:0:\dots:0]}_{0 \le i \le r}.$$

The coordinate curve l_{ij} connecting q_i, q_j has one dimensional degree of freedom \mathbb{C}^* (as an invariant component). The curve $C \in \overline{C}_{g,n}$ is stable (i.e. $\operatorname{Aut}(C) < \infty$) if and only if 2g - 2 + n > 0. If a components C' of C is mapped to l_{ij} , then C' has two points mapped to q_i, q_j respectively (equivalent to with two marked points in C'), so $2g - 2 + 2 \leq 0$ implies g = 0, i.e. $C' \cong \mathbb{P}^1$ (see Fig 1.1). Meanwhile, $f|_{C'}$ must be uniformly ramified, so $f|_{C'}(z) = z^e, \forall z \in C' \cong \mathbb{P}^1$, for some $e \in \mathbb{N}^*$.

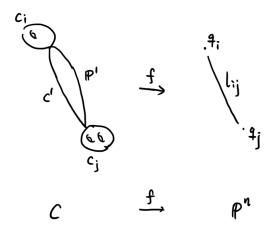


Figure 1.1: $f(C_i) = q_i$, $f(C') = l_{ij}$, $f(C_i) = q_i$

It is convenient to use a decorated graph $\vec{\Gamma} = (\Gamma, \vec{f}, \vec{d}, \vec{g}, \vec{s})$ (graph, maps, degrees, genus, marked points) to represent $[f: C \to X, 1, \dots, n] \in \overline{\mathcal{M}}_{g,n}(\mathbb{P}^r, d)^{\mathbb{T}}$. Let $\operatorname{val}(v)$, the valence of v, be the number of edges connecting vertex v, and let $n(v) = |s_v| + val(v)$. The stable map $[f: C \to X, 1, \dots, n]$ with fixed graph $\vec{\Gamma} = (\Gamma, \vec{f}, \vec{d}, \vec{g}, \vec{s})$ defines a substack

$$\overline{\mathcal{M}}_{\vec{\Gamma}} \subset \overline{\mathcal{M}}_{g,n}(\mathbb{P}^r, d).$$

On the other sides, consider

$$\varphi_{\vec{\Gamma}}: \prod_{\dim C_v=1} \overline{M}_{g_v, n(v)} \to \overline{\mathcal{M}}_{\vec{\Gamma}}.$$

If v, v' are connected by an edge e, then let $C_v, C_{v'}$ connected by a $C_e \cong \mathbb{P}^1$ associated with a degree d_e map to \mathbb{P}^r . Let \overline{M}_{Γ} be the product of above C_v, C_e . There is a group \mathbb{A}_{Γ} acting on \overline{M}_{Γ} . The group \mathbb{A}_{Γ} is defined by:

$$1 \to \prod_{edges} \mathbb{Z}/(d_e) \to \mathbb{A}_{\Gamma} \to Aut(\Gamma) \to 1.$$

and

$$\overline{\mathcal{M}}_{\vec{\Gamma}} = \overline{M}_{\Gamma}/\mathbb{A}_{\Gamma}.$$

Therefore, we know the \mathbb{T} -fixed locus of $\overline{\mathcal{M}}_{g,n}(\mathbb{P}^r,d)$ is $\overline{\mathcal{M}}_{\vec{\Gamma}}$. Let N_{Γ} be the normal bundle of $\overline{\mathcal{M}}_{\vec{\Gamma}}$ in $\overline{\mathcal{M}}_{g,n}(\mathbb{P}^r,d)$. Then there is an explicit formula for

the equivariant Euler class. Before doing that, we define some necessary notations. A flag F is a pair (v, e) such that e is an edge containing the vertex v. We put i(F) = v, j(F) the vertex of e different from v. Let

$$\omega_F = \frac{\lambda_{i(F)} - \lambda_{j(F)}}{d_e} \in H^2_{\mathbb{T}}(pt) = \mathbb{Q}[\lambda_0, \dots, \lambda_r],$$

which is the weight of T-action on $T_{q_{in}}C_e$.

Theorem 1.1.1 ($Euler_{\mathbb{T}}(N_{\Gamma})$). $Euler_{\mathbb{T}}(N_{\Gamma}) = e_{\Gamma}^F e_{\Gamma}^v e_{\Gamma}^e$, where

$$e_{\Gamma}^{F} = \prod_{n(i(F))\geq 3} (\omega_{F} - \psi_{F}) / \prod_{j\neq i(F)} (\lambda_{i(F)} - \lambda_{j}),$$

$$e_{\Gamma}^{v} = \prod_{v} \prod_{j\neq i_{v}} (\lambda_{i_{v}} - \lambda_{j}) \prod_{val(v)=2, s_{v}=\emptyset} (\omega_{F_{1}(v)} + \omega_{F_{2}(v)}) / \prod_{val(v)=1, s_{v}=\emptyset} \omega_{F(v)}$$

$$e_{\Gamma}^{e} = \prod_{e} \frac{(-1)^{d_{e}} (d_{e}!)^{2} (\lambda_{i} - \lambda_{j})^{2d_{e}}}{d_{e}^{2d_{e}}} \prod_{a+b=d_{e}, k\neq i, j} (\frac{a\lambda_{i} + b\lambda_{j}}{d_{e}} - \lambda_{k})$$

The proof is partially discussed in section 1.2.

1.2 Tangent-obstruction sequence

Consider
$$[f: C \to X, 1, ..., n] \in \overline{\mathcal{M}}_{\vec{\Gamma}}, \vec{\Gamma} = (\Gamma, \vec{f}, \vec{d}, \vec{g}, \vec{s})$$
. We put $V^1(\Gamma) := \{v \in V(\Gamma) : g_v = 0, val(v) = 1, |s_v| = 0\}$
 $V^2(\Gamma) := \{v \in V(\Gamma) : g_v = 0, val(v) = 2, |s_v| = 0\}$
 $V^{1,1}(\Gamma) := \{v \in V(\Gamma) : g_v = 0, val(v) = 1, |s_v| = 1\}$
 $V^s(\Gamma) := \{v \in V(\Gamma) : 2g_v - 2 + val(v) + |s_v| > 0\}$
 $y(v, e) := C_e \cap C_v$

The tangent-obstruction sequence is

$$0 \to Aut(C, 1, \dots, n)$$

$$\to Def(f) \to Def(C, 1, \dots, n, f) \to Def(C, 1, \dots, n)$$

$$\to Ob(f) \to Ob(C, 1, \dots, n, f) \to 0,$$

$$0 \to Hom(\Omega_C(p_1 + \dots + p_n), \mathcal{O}_C)$$

$$\to H^0(C, f^*T_X) \to T^1 \to Ext^1(\Omega_C(p_1 + \dots + p_n), \mathcal{O}_C)$$

$$\to H^1(C, f^*T_X) \to T^2 \to 0.$$

For simplicity:

$$0 \to B_1 \to B_2 \to T^1 \to B_4 \to B_5 \to T^2 \to 0.$$

The $N^{vir} = T^{1,m} - T^{2,m}$ (m means moving part).

$$Euler_{\mathbb{T}}(N^{vir}) = \frac{Euler_{\mathbb{T}}(B_2^m)Euler_{\mathbb{T}}(B_4^m)}{Euler_{\mathbb{T}}(B_1^m)Euler_{\mathbb{T}}(B_5^m)}.$$

(1) $Euler_{\mathbb{T}}(B_2^m)/Euler_{\mathbb{T}}(B_5^m)$. The normalization sequence of C is:

$$0 \to \mathcal{O}_C \to \bigoplus_{v \in V^s(\Gamma)} \mathcal{O}_{C_v} \oplus \bigoplus_{e \in E(\Gamma)} \mathcal{O}_{C_e}$$
$$\to \bigoplus_{v \in V^2(\Gamma)} \mathcal{O}_{y_v} \oplus \bigoplus_{(e,v) \in F^s(\Gamma)} \mathcal{O}_{y(e,v)} \to 0.$$

Take $\otimes f^*T_X$:

$$0 \to H^{0}(C, f^{*}T_{X}) \to \bigoplus_{v \in V^{s}(\Gamma)} H^{0}(C_{v}, f^{*}T_{X}) \oplus \bigoplus_{e \in E(\Gamma)} H^{0}(C_{e}, f^{*}T_{X})$$

$$\to \bigoplus_{v \in V^{2}(\Gamma)} T_{f(y_{v})}X \oplus \bigoplus_{(e,v) \in F^{s}(\Gamma)} T_{f(y(e,v))}X$$

$$\to H^{1}(C, f^{*}T_{X}) \to \bigoplus_{v \in V^{s}(\Gamma)} H^{1}(C_{v}, f^{*}T_{X}) \oplus \bigoplus_{e \in E(\Gamma)} H^{1}(C_{e}, f^{*}T_{X}) \to 0.$$

$$H^1(C_v, f^*T_X) = H^1(C_v, \mathcal{O}_{C_v}) \otimes T_{f(C_v)} X \cong H^0(C_v, \omega_{C_v})^{\vee} \otimes T_{f(C_v)} X$$

Here $H^0(C_v, \omega_{C_v})$ is Hodge bundle \mathbb{E} . By splitting principle, assume $\mathbb{E} = L_1 \oplus \cdots \oplus L_g$, then

 $H^0(C_v, f^*T_X) = T_{f(C_v)}X$

$$e(\mathbb{E}^{\vee} \otimes \mathbb{C}_{1}) = \prod_{i=1}^{g} c_{1}(L_{i}^{\vee} \otimes \mathbb{C}_{1}) = \prod_{i=1}^{g} c_{1}(L_{i}^{\vee}) + c_{1}(\mathbb{C}_{1})$$

$$= \prod_{i=1}^{g} (-c_{1}(L_{i}) + u) = \sum_{k=1}^{g} (-1)^{k} c_{k}(\mathbb{E}) u^{g-k} = \sum_{k=1}^{g} (-1)^{k} \lambda_{k} u^{g-k} =: \Lambda_{g}^{\vee}(u)$$

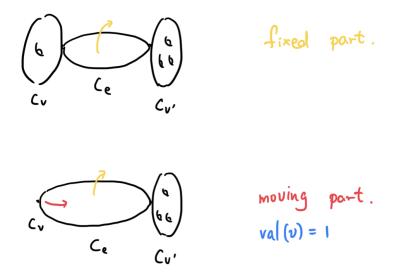


Figure 1.2: automorphism of (C, 1, ..., n)

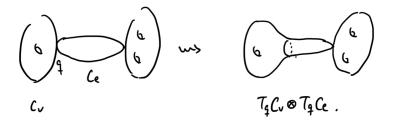


Figure 1.3: deformation of $(C, 1, \ldots, n)$

- (2) $Euler_{\mathbb{T}}(B_4^m)/Euler_{\mathbb{T}}(B_1^m)$.
- (2.1) $B_1 = Aut(C, 1, ..., n) = Hom(\Omega_C(p_1 + ... + p_n), \mathcal{O}_C)$: We should classify what is moving and what is fixed. Basically, we have

$$B_1^m = \bigoplus_{v \in V^1(\Gamma), (e,v) \in F(\Gamma)} T_{y(e,v)} C_e.$$

(2.2) $B_4 = Def(C, 1, ..., n) = Ext^1(\Omega_C(p_1 + ... + p_n), \mathcal{O}_C)$: \mathbb{P}^1 has just 1 complex structure, so we consider $g(C) \geq 1$. If we don't change node q, C will stay in the same class in $\overline{\mathcal{M}}_{g,n}$. Hence we must resolve the node, and

geometrically, resolution depends on $T_qC_v\otimes T_qC_e$. So basically we have

$$B_4^m = \bigoplus_{v \in V^2(\Gamma), E_v = (e, e')} T_{y_v} C_e \otimes T_{y_v} C_{e'} \oplus \bigoplus_{(e, v) \in F^s(\Gamma)} T_{y(e, v)} C_v \otimes T_{y(e, v)} C_e$$

Returning to the special case $X = \mathbb{P}^r$, we can get the theorem 1.1.1.

1.3 Aspinwall Morrison formula; Faber Pandaripande formula

In this section, we will use Kontsevich's approach to compute the multiple cover contribution of rigidly embedding curves \mathbb{P}^1 in a Calabi-Yau threefold X.

The geometry picture is this. The normal bundle N of $\mathbb{P}^1 \subset X$ is rank 2 and splits on \mathbb{P}^1 . Because X is Calabi-Yau and $c_1(\mathbb{P}^1) = 2$, the normal bundle is of degree 2. Embedded \mathbb{P}^1 's in a Calabi-Yau threefold (not necessary lines) with normal bundle $\mathcal{O}(-1) \oplus \mathcal{O}(-1)$ are called rigid. The degree 2 Gromov-Witten invariant of a generic quintic has two contributions:

- (1) rigid conics curves in X;
- (2) lines with double cover, so this part is related to $\overline{\mathcal{M}}_0(\mathbb{P}^1,2)$.

We want to compute the contribution of part (2). This problem finally leads to:

$$N_{g,d} = \int_{\overline{\mathcal{M}}_{g,0}(\mathbb{P}^1,d)} e(R^1 \pi_* f^* N),$$

where

$$\overline{C}_{g,0}(\mathbb{P}^1, d) \xrightarrow{f} \mathbb{P}^1$$

$$\downarrow^{\pi} \quad \text{and } N = \mathcal{O}(-1) \oplus \mathcal{O}(-1).$$

$$\overline{\mathcal{M}}_{g,0}(\mathbb{P}^1, d)$$

The decorated graphs $\vec{\Gamma} = (\Gamma, \vec{f}, \vec{d}, \vec{g}, \vec{s})$ in $\overline{\mathcal{M}}_{g,0}(\mathbb{P}^1, d)^{\mathbb{T}}$ are of the type in Figure 1.4. We can choose different lifts on $\mathcal{O}(-1) \oplus \mathcal{O}(-1)$ so that only $\vec{\Gamma} = (\Gamma, \vec{f}, \vec{d}, \vec{g}, \vec{s})$ with 1 edge contributing $N_{g,d}$.

(1) g = 0 (Aspinwall Morrison formula): $N_{0,d} = 1/d^3$;

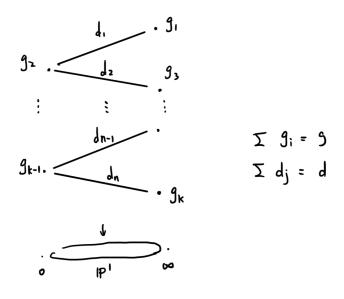


Figure 1.4: $\overline{\mathcal{M}}_{q,0}(\mathbb{P}^1,d)^{\mathbb{T}}$

(2) $g \ge 1$ (Faber-Pandharipande):

$$N_{g,d} = \sum_{g_1 + g_2 = g} \frac{1}{d} \int_{\overline{\mathcal{M}}_{g_1,1}} \lambda_{g_1} \psi^{2g_1 - 2} d^{2g_1 - 1}$$

$$\times \int_{\overline{\mathcal{M}}_{g_2,1}} \lambda_{g_2} \psi^{2g_2 - 2} d^{2g_2 - 1} = \sum_{g_1 + g_2 = g} b_{g_1} b_{g_2} d^{2g - 3}$$

$$b_0 = 0; b_g = \int_{\overline{\mathcal{M}}_{g,1}} \lambda_{g_2} \psi^{2g - 2} \quad (g > 0)$$

$$\sum_{g_2 = 0}^{\infty} b_g t^{2g} = \frac{t/2}{\sin t/2}.$$

Then use the Laurent series of $\cot t$, we have

$$N_{1,d}=\frac{1}{12d},$$

$$N_{g,d}=d^{2g-3}\frac{|B_{2g}|}{2g\cdot(2g-2)!}=|\chi(\overline{\mathcal{M}}_g)|\frac{d^{2g-3}}{(2g-3)!},\quad g\geq 2,$$
 where B_g is the Bernoulli number in $\frac{x}{e^x-1}$.

Chapter 2

Quantum Cohomology

2.1 quantum product

The quantum cohomology is a variation of classical cohomology. Let $T_0 = 1, T_1, \ldots, T_p, T_{p+1}, \ldots, T_m \in H^*(X)$ be a basis of $H^*(X)$ as a \mathbb{Q} -vector space $(T_1, \ldots, T_p \in H^2(X))$. Let $\beta \in H^2(X)$, $\gamma = \sum_{i=0}^m t_i T_i$, we define quantum potential as

$$F_0^X(t_0, \dots, t_m) = \sum_{n,\beta} \frac{1}{n!} \langle \gamma^n \rangle_{0,n,\beta}^X Q^{\beta}$$

$$= \frac{1}{6} \int_X (\sum_{i=0}^m t_i T_i)^3 + \sum_{\beta=0, n \ge 4} \langle T_0^{n_0} \dots T_m^{n_m} \rangle_{0,n,0} \prod_{i=1}^m \frac{t_i^{n_i}}{n_i!}$$

$$+ \sum_{\beta>0,n} Q^{\beta} \langle T_0^{n_0} \dots T_m^{n_m} \rangle_{0,n,\beta} \prod_{i=p+1}^m \frac{t_i^{n_i}}{n_i!}.$$

By string equations and divisor equations,

$$F_0^X(t_0, \dots, t_m) = \frac{1}{6} \int_X (\sum_{i=0}^m t_i T_i)^3 + \sum_{\beta = 0, n \ge 4} \langle T_1^{n_1} \dots T_m^{n_m} \rangle_{0, n, 0} \prod_{i=1}^m \frac{t_i^{n_i}}{n_i!} + \sum_{\beta > 0, n} Q^{\beta} q_1^{\int_{\beta} T_1} \dots q_p^{\int_{\beta} T_p} \langle T_{p+1}^{n_{p+1}} \dots T_m^{n_m} \rangle_{0, n, \beta} \prod_{i=p+1}^m \frac{t_i^{n_i}}{n_i!},$$

where $q_i = e^{t_i}$.

$$F_{ijk} := \frac{\partial^3 F_0^X}{\partial t_i \partial t_j \partial t_k} = \sum_{n,\beta} \frac{1}{n!} \langle T_i T_j T_k \gamma^n \rangle_{0,n+3,\beta}^X Q^{\beta}$$

$$= \int_{X} T_{i} T_{j} T_{k} + \sum_{\beta=0, n \geq 1} \langle T_{i} T_{j} T_{k} T_{1}^{n_{1}} \dots T_{m}^{n_{m}} \rangle_{0, n+3, 0} \prod_{i=1}^{m} \frac{t_{i}^{n_{i}}}{n_{i}!}$$

$$+ \sum_{\beta>0, n} Q^{\beta} q_{1}^{\int_{\beta} T_{1}} \dots q_{p}^{\int_{\beta} T_{p}} \langle T_{i} T_{j} T_{k} T_{p+1}^{n_{p+1}} \dots T_{m}^{n_{m}} \rangle_{0, n+3, \beta} \prod_{i=p+1}^{m} \frac{t_{i}^{n_{i}}}{n_{i}!}, \quad q_{i} = e^{t_{i}}.$$

Let $g_{ij} = (T_i, T_j)$ means the Poincare pair of T_i, T_j . The big quantum product is defined as

$$(T_i *_t T_j, T_k) := F_{ijk},$$

in other words,

$$T_i *_t T_j = \sum_{e,f} F_{ije} g^{ef} T_f.$$

It is known that the quantum product is a generalization of intersection theory: given T_i, T_j, T_k , they contribute to the quantum product if there exists \mathbb{P}^1 touching their Poincare dual classes at the same time. Extend the t_i in quantum multiplication linearly, then the $\mathbb{Q}[[t_0, \ldots, t_m]]$ -module $H^*(X) \otimes_{\mathbb{Q}} \mathbb{Q}[[t_0, \ldots, t_m]]$ is the big quantum cohomology QH(X).

The associativity of quantum product is formulated as WDVV equation:

$$F_{ija}g^{ab}F_{bkl} = F_{ila}g^{ab}F_{bjk}.$$

It is proved by a forgetful map $\pi : \overline{\mathcal{M}}_{0,4}(X,\beta) \to \overline{\mathcal{M}}_{0,4}$. One should notice that $\overline{\mathcal{M}}_{0,4} \cong \mathbb{P}^1$, so the boundary divisor $D(12|34) \sim D(13|24)$ and

$$\int_{[\overline{\mathcal{M}}_{0,4}(X,\beta)]^{vir}\cap\pi^*D(12|34)} ev_1^*(T_i)ev_2^*(T_j)ev_3^*(T_k)ev_4^*(T_l)\prod_{i=5}^{n+4} ev_i^*(\gamma)$$

$$= \int_{[\overline{\mathcal{M}}_{0,4}(X,\beta)]^{vir} \cap \pi^*D(13|24)} ev_1^*(T_i)ev_2^*(T_j)ev_3^*(T_k)ev_4^*(T_l) \prod_{i=5}^{n+4} ev_i^*(\gamma).$$

A useful trick is to separate $[\overline{\mathcal{M}}_{0,4}(X,\beta)]^{vir} \cap \pi^*D(12|34)$ by

$$\coprod_{n_1+n_2=n,\beta_1+\beta_2=\beta} [\overline{\mathcal{M}}_{0,n_1+3}(X,\beta_1) \times \overline{\mathcal{M}}_{0,n_2+3}(X,\beta_2)]^{vir} \cap (ev \times ev)^*[\Delta],$$

$$PD[\Delta] = g^{ab}T_a \otimes T_b$$

then we get

$$\sum_{n_1+n_2=n} \sum_{\beta_1+\beta_2=\beta} \langle T_i T_j T_a \gamma^{n_1} \rangle_{0,n_1+3,\beta_1} g^{ab} \langle T_b T_k T_l \gamma^{n_2} \rangle_{0,n_2+3,\beta_2}$$

$$= \sum_{n_1+n_2=n} \sum_{\beta_1+\beta_2=\beta} \langle T_i T_k T_a \gamma^{n_1} \rangle_{0,n_1+3,\beta_1} g^{ab} \langle T_b T_j T_l \gamma^{n_2} \rangle_{0,n_2+3,\beta_2}.$$

This is the essential part in the proof of associativity of quantum product.

Remark 2.1.1. It deserves to notice that the quantum product is defined by rational curves, so its usage mainly concentrates in genus 0 GW-invariants. The difficulty to define quantum product via higher genus curves is that there is no so good associativity as the genus 0 case. It must be a good work if we can find a way to give a quantum product via higher genus curves with associativity like now.

The small quantum product is defined by

$$T_i *_s T_j = T_i *_t T_j |_{t_{p+1} = \dots = t_m = 0}, 0 \le i, j \le m.$$

Precisely, let

$$\overline{F}_{ijk} = F_{ijk}|_{t_{p+1}=\dots=t_m=0} = \int_X T_i T_j T_k + \sum_{\beta>0} Q^{\beta} q_1^{\int_{\beta} T_1} \dots q_p^{\int_{\beta} T_p} \langle T_i T_j T_k \rangle_{0,3,\beta},$$

then

$$T_i *_s T_j = \overline{F}_{ije} g^{ef} T_f, \quad 1 \le e, f \le m.$$

Extend q_i linearly, the $\mathbb{Q}[[q_1,\ldots,q_p]]$ -module $H^*(X)\otimes_{\mathbb{Q}}\mathbb{Q}[[q_1,\ldots,q_p]]$ is defined as the small quantum cohomology $QH^s(X)$.

Example 2.1.2. $QH^{s}(\mathbb{P}^{m}) = \mathbb{Q}[H,q]/(H^{m+1}-q)$, where $H \in H^{2}(\mathbb{P}^{m},\mathbb{Q})$, $q = e^{t_{1}}$.

2.2 quantum differential equation

We can view the vector space H(X) as a Riemannian manifold M with standard flat metric g_{ij} given by Poincare pairing. The quantum product $*_t$ could be use to define a connection (called Dubrovin connection, or Givental connection) ∇^z , which is different from the Levi-Civita connection induced by its Riemannian metric.

Definition 2.2.1. (Dubrovin connection) Let $X, Y \in \Gamma(M, TM)$, ∇ be the Levi-Civita connection w.r.t g. The Dubrovin connection ∇^z is defined by

$$\nabla_X^z Y := \nabla_X Y - \frac{1}{z} X *_t Y.$$

The WDVV equation shows ∇^z is flat. i.e. $Rm^z = 0$.

Definition 2.2.2 (quantum differential equation). Let $\sigma \in \Gamma(M, TM)$, the equation $\nabla^z \sigma = 0$ is the quantum differential equation. The fundamental solution of quantum differential equation is $(m+1) \times (m+1)$ matrix s(z,t) $(t = (t_0, \ldots, t_m)) = (a_{ij})$, such that each column defines a solution

$$\sigma_j(t) = \sum_{i=0}^m a_{ij}(t) \frac{\partial}{\partial t_i}.$$

Now we want to find the solution of quantum differential equation. Let $(S(z)T_a, T_b) = g_{ab} + \langle \langle \frac{T_a}{z-\varphi_1}, T_b \rangle \rangle_{0,2}$, where

$$\langle\langle \frac{T_a}{z-\psi_1}, T_b \rangle\rangle_{0,2} = \sum_{n \geq 0, \beta, (n,\beta) \neq (0,0)} \frac{1}{n!} \langle \frac{T_a}{z-\psi_1} T_b \gamma^n \rangle_{0,n+2,\beta}.$$

Proposition 2.2.3. The $S_a = (S(z)T_a, T_b)g^{bc}\partial_c$ is a flat section with respect to ∇^z .

Chapter 3

Mirror Symmetry

I plan to follow Givental's approach to give a proof of genus 0 mirror symmetry of hypersurfaces in \mathbb{P}^n .

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