

ME 1071: Applied Fluids

Lecture 11 Introduction to Compressible Flow

Spring 2021

Weekly Study Plan





Weeks	Dates	Lectures
1	Mar. 9	Course Introduction, Fluids Review
2	Mar. 16	Chapter 8: Internal Incompressible Viscous Flow
3	Mar. 23	Chapter 8: Internal Incompressible Viscous Flow
4	Mar. 30	Chapter 8/Exam I Review
5	Apr. 6	Exam I
6	Apr. 13	Chapter 9: External Incompressible Viscous Flow
7	Apr. 20	Chapter 9: External Incompressible Viscous Flow
8	Apr. 25	Chapter 11: Flow in Open Channels
9	Apr. 27	Chapter 11: Flow in Open Channels
10	May. 11	Exam II Review
11	May. 18	Exam II
12	May. 25	Chapter 12: Introduction to Compressible Flow
13	Jun. 1	Chapter 12: Introduction to Compressible Flow
14	Jun. 8	Chapter 12: Introduction to Compressible Flow
15	Jun. 15	Chapter 5: CFD Related Topics
16	Jun. 22	Final Exam Review (Final Exam at Jun. 26)

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Outlines





- > Isentropic Flow with Area Variation
- **▶** Calculation of Normal Shock-Wave Properties
- Measurement of Velocity in Compressible Flows

Isentropic Flow with Area Variation

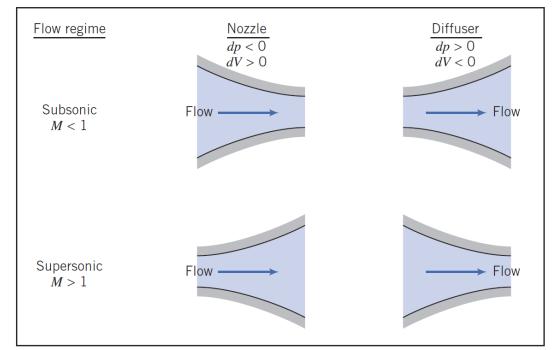




Flow variation induced by area change

Area-velocity relation

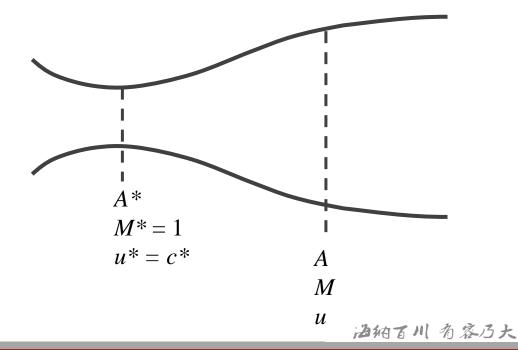
$$\frac{dA}{A} = (M^2 - 1)\frac{du}{u}$$



Nozzle and diffuser shapes as a function of initial Mach number.

Area-Mach number relation

$$\left(\frac{A}{A^*}\right)^2 = \frac{1}{M^2} \left[\frac{2}{k+1} \left(1 + \frac{k-1}{2} M^2 \right) \right]^{(k+1)/(k-1)}$$



Compressible Flow with Area Variation





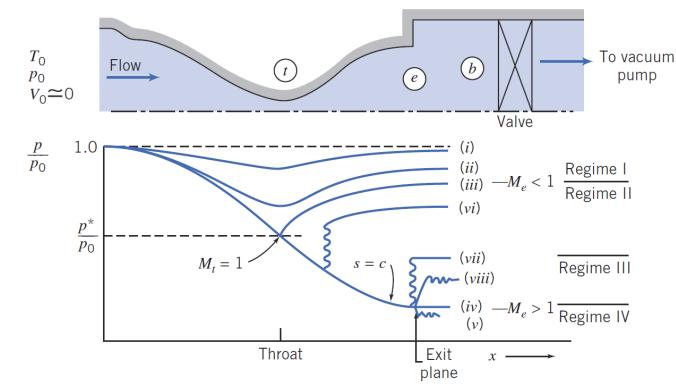
- Compressible Flow in a Convergent-Divergent Nozzle
 - Once there is shock wave, the flow is not isentropic.

$$\frac{T_0}{T} = 1 + \frac{k-1}{2}M^2$$

$$\frac{p_0}{p} = \left(1 + \frac{k-1}{2}M^2\right)^{k/(k-1)}$$

$$\frac{
ho_0}{
ho} = \left(1 + \frac{k-1}{2}M^2\right)^{1/(k-1)}$$

How to calculate the properties of shock waves?



Pressure distributions for flow in a converging-diverging nozzle for different back pressures.

Outlines





- **▶** Isentropic Flow with Area Variation
- > Calculation of Normal Shock-Wave Properties
- Measurement of Velocity in Compressible Flows



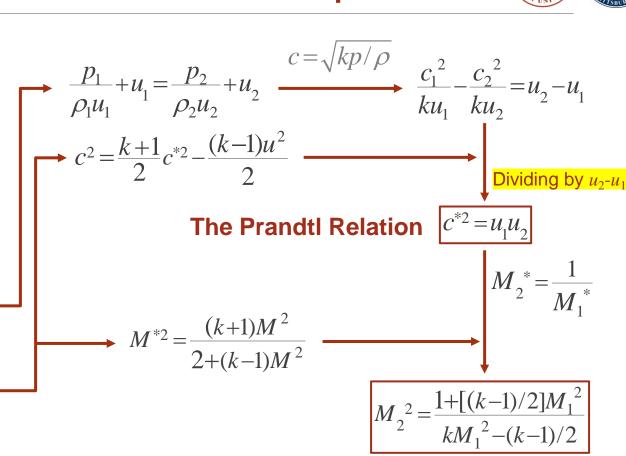


Assumptions

- Steady flow, $\partial /\partial t = 0$.
- Adiabatic, $\dot{q} = 0$.
- No viscous effects.
- No body forces, $\mathbf{f} = 0$.
- + One-dimensional flow, $M_1 >= 1$

$egin{array}{c} p_1 \ ho_1 \ T_1 \ M_1 \end{array}$	b	2	$p_2 \\ \rho_2 \\ T_2 \\ M_2$
u_1 $p_{0,1}$ $h_{0,1}$ $T_{0,1}$			u_2 $p_{0,2}$ $h_{0,2}$ $T_{0,2}$
= 1	L		s_2

	*1	
Continuity	$\rho_1 u_1 = \rho_2 u_2$	
Momentum	$p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2$	
Energy	$h_1 + \frac{u_1^2}{2} = h_2 + \frac{u_2^2}{2}$	
Enthalpy	$h_2 = c_p T_2$	
Equation of state	$p_2 = \rho_2 R T_2$	



The relation between the Mach numbers ahead and behind the normal shock wave.





The Mach numbers across a normal shock wave

$$M_{2}^{2} = \frac{1 + [(k-1)/2]M_{1}^{2}}{kM_{1}^{2} - (k-1)/2} \begin{cases} \text{Invalid} & \text{if } M_{1} < 1 \\ M_{2} = 1 & \text{infinitely weak shock wave, Mach wave} \\ M_{2} < 1 & \text{if } M_{1} > 1 \\ M_{2} \rightarrow \sqrt{\frac{k-1}{2k}} = 0.378 & \text{if } M_{1} \rightarrow \infty \end{cases}$$

The ratios of thermodynamics properties across a normal shock wave

Continuity
$$\rho_1 u_1 = \rho_2 u_2$$
 $\rho_2 = \frac{u_1}{u_2}$ $\rho_2 = \frac{u_1^2}{u_1 u_2} = \frac{u_1^2}{v_1^2} = \frac{u_1^2}{v_2^2} = M_1^{*2}$ Momentum $p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2$ $p_2 - p_1 = \rho_1 u_1^2 - \rho_2 u_2^2 = \rho_1 u_1^2 (1 - \frac{u_2}{u_1})$

$$\frac{p_2 - p_1}{p_1} = \frac{\rho_1}{p_1} u_1^2 (1 - \frac{u_2}{u_1}) = \frac{k \rho_1}{k p_1} u_1^2 (1 - \frac{u_2}{u_1}) = \frac{k u_1^2}{a_1^2} (1 - \frac{u_2}{u_1}) \longrightarrow \frac{p_2}{p_1} = 1 + \frac{2k}{k+1} (M_1^2 - 1)$$

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The ratios of thermodynamics properties across a normal shock wave

$$\begin{cases} h_2 = c_1 T_2 \end{cases}$$

$$\frac{h_2}{h_1} = \frac{T_2}{T_1} = \frac{p_2}{p_1} \frac{\rho_1}{\rho_2}$$

Equation of state
$$p = \rho RT$$
 $h_2 = r_2 =$

The five unknowns p_2 , u_2 , ρ_2 , h_2 , and T_2 are explicitly solved.

• The ratios of thermodynamics properties in limiting case $(M_1 \rightarrow \infty)$

$$\lim_{M_1 \to \infty} M_2 = \sqrt{\frac{1 + [(k-1)/2]M_1^2}{kM_1^2 - (k-1)/2}} = \sqrt{\frac{k-1}{2k}} = 0.378$$

$$\lim_{M_1 \to \infty} \frac{\rho_2}{\rho_1} = \frac{u_1}{u_2} = \frac{(k+1)M_1^2}{2 + (k-1)M_1^2} = \frac{k+1}{k-1} = 6$$

$$\lim_{M_1 \to \infty} \frac{p_2}{p_1} = 1 + \frac{2k}{k+1} (M_1^2 - 1) = \infty$$

$$\lim_{M_1 \to \infty} \frac{h_2}{h_1} = \frac{T_2}{T_1} = \left[1 + \frac{2k}{k+1} (M_1^2 - 1) \right] \frac{2 + (k-1)M_1^2}{(k+1)M_1^2} = \infty$$

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The entropy change across the normal shock wave

$$s_{2} - s_{1} = c_{p} \ln \frac{T_{2}}{T_{1}} - R \ln \frac{p_{2}}{p_{1}} \xrightarrow{M_{1} >= 1} s_{2} - s_{1} = c_{p} \ln \left[1 + \frac{2k}{k+1} (M_{1}^{2} - 1) \right] \frac{2 + (k-1)M_{1}^{2}}{(k+1)M_{1}^{2}} \right] - R \ln \left[1 + \frac{2k}{k+1} (M_{1}^{2} - 1) \right] \ge 0$$

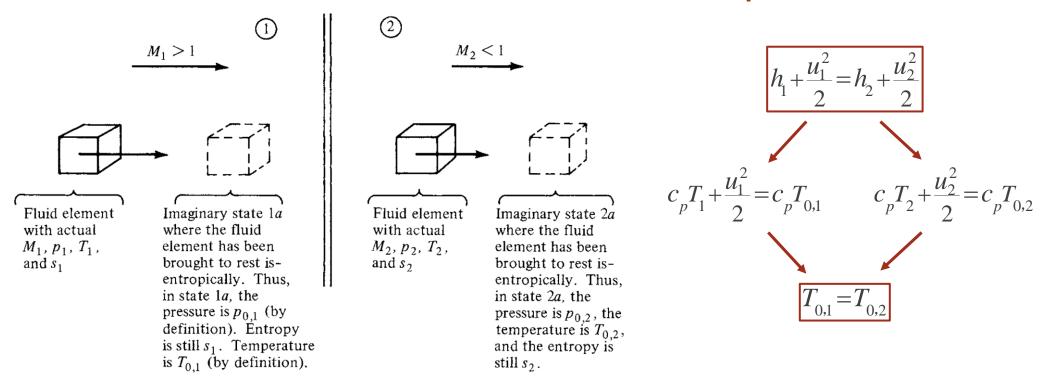
The second law of thermodynamics requires $M_1 \ge 1$.

- Within the shock wave large gradients in velocity and temperature occur.
- Friction and thermal conduction are significant.
- Therefore, the entropy increases across the shock wave.





Total conditions across the normal shock wave: total temperature



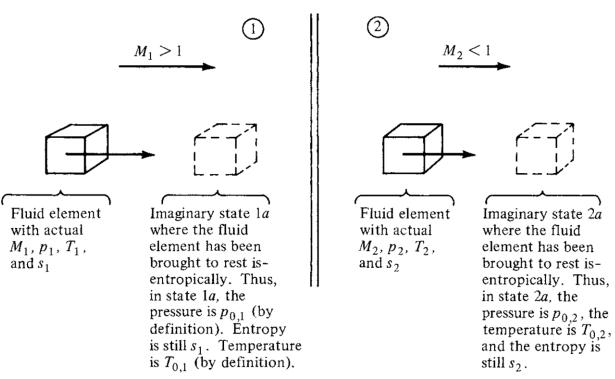
- Total temperature is constant across a stationary normal shock wave.
- Consistent with the conclusion derived for the adiabatic calorically perfect gas.

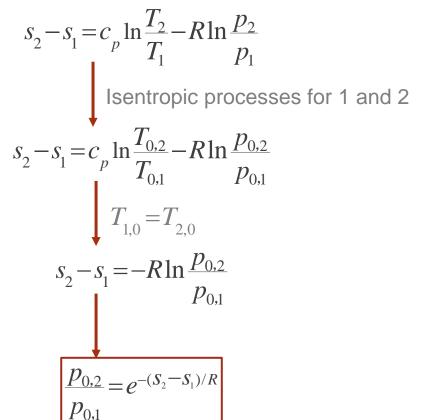
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Total conditions across the normal shock wave: total pressure



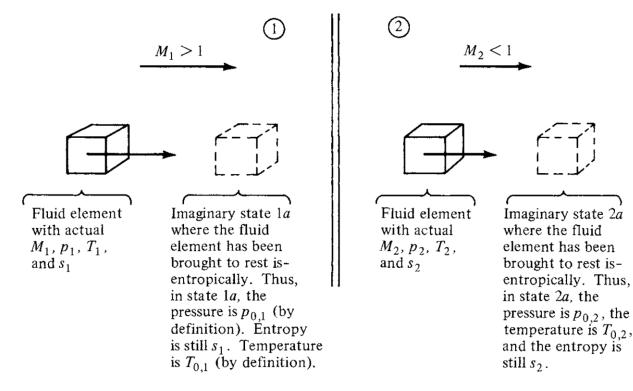


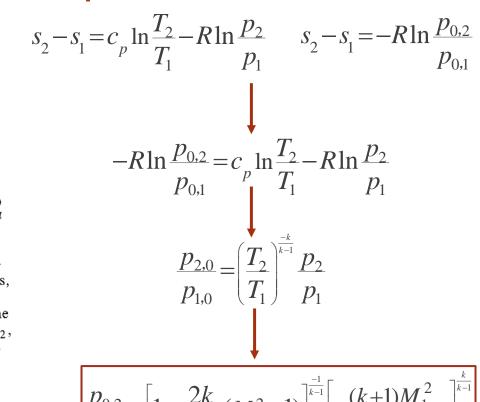
• Total pressure decreases across a normal shock wave.





Total conditions across the normal shock wave: total pressure





Total pressure decreases across a normal shock wave.

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Assumptions

- Steady flow, $\partial /\partial t = 0$.
- Adiabatic, $\dot{q} = 0$.
- No viscous effects.
- No body forces, $\mathbf{f} = 0$.
- + One-dimensional flow, $M_1 >= 1$

$p_1 \\ \rho_1 \\ T_1 \\ M_1$	b	2	$\begin{array}{c} p_2 \\ \rho_2 \\ T_2 \\ M_2 \end{array}$
u_1 $p_{0,1}$ $h_{0,1}$ $T_{0,1}$			u_2 $p_{0,2}$ $h_{0,2}$ $T_{0,2}$
<u>1</u>	<u>_</u>		2

M 2 _	$\frac{1+[(k-1)/2]M_1^2}{1+[(k-1)/2]M_1^2}$
1 v1 ₂ –	$kM_1^2 - (k-1)/2$

$$\frac{\rho_2}{\rho_1} = \frac{u_1}{u_2} = \frac{(k+1)M_1^2}{2 + (k-1)M_1^2}$$

$$p_{0,2} = e^{-(S_2 - S_1)/R}$$

 $T_{0,1} = T_{0,2}$

Continuity
$$\rho_1 u_1 = \rho_2 u_2$$
Momentum $p_1 + \rho_1 u_1^2 = p_2 + \rho_2 u_2^2$ Energy $h_1 + \frac{u_1^2}{2} = h_2 + \frac{u_2^2}{2}$ Enthalpy $h_2 = c_p T_2$ Equation of state $p_2 = \rho_2 R T_2$

$$\frac{p_2}{p_1} = 1 + \frac{2k}{k+1} (M_1^2 - 1)$$

$$\frac{h_2}{h_1} = \frac{T_2}{T_1} = \left[1 + \frac{2k}{k+1}(M_1^2 - 1)\right] \frac{2 + (k-1)M_1^2}{(k+1)M_1^2}$$

$$s_2 - s_1 = c_p \ln \left\{ \left[1 + \frac{2k}{k+1} (M_1^2 - 1) \right] \frac{2 + (k-1)M_1^2}{(k+1)M_1^2} \right\} - R \ln \left[1 + \frac{2k}{k+1} (M_1^2 - 1) \right] \ge 0$$





The variations of properties across a normal shock wave

$$M_2^2 = \frac{1 + [(k-1)/2]M_1^2}{kM_1^2 - (k-1)/2}$$

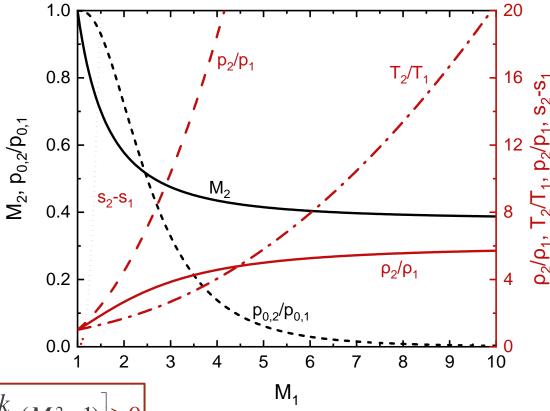
$$T_{0,1} = T_{0,2}$$

$$\frac{\rho_2}{\rho_1} = \frac{u_1}{u_2} = \frac{(k+1)M_1^2}{2 + (k-1)M_1^2}$$

$$\frac{p_{0,2}}{p_{0,1}} = e^{-(S_2 - S_1)/R}$$

$$\frac{p_2}{p_1} = 1 + \frac{2k}{k+1} (M_1^2 - 1)$$

$$\frac{h_2}{h_1} = \frac{T_2}{T_1} = \left[1 + \frac{2k}{k+1}(M_1^2 - 1)\right] \frac{2 + (k-1)M_1^2}{(k+1)M_1^2}$$



$$s_2 - s_1 = c_p \ln \left\{ \left[1 + \frac{2k}{k+1} (M_1^2 - 1) \right] \frac{2 + (k-1)M_1^2}{(k+1)M_1^2} \right\} - R \ln \left[1 + \frac{2k}{k+1} (M_1^2 - 1) \right] \ge 0$$

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Compressible Flow with Area Variation





Compressible Flow in a Convergent-Divergent Nozzle

Once there is shock wave, the flow is not isentropic.

$$\frac{T_0}{T} = 1 + \frac{k-1}{2}M^2$$

$$\frac{p_0}{p} = \left(1 + \frac{k-1}{2}M^2\right)^{k/(k-1)}$$

$$\frac{
ho_0}{
ho} = \left(1 + \frac{k-1}{2}M^2\right)^{1/(k-1)}$$

To vacuum Flow (b)pump $\left(e\right)$ Valve $\begin{array}{c} (ii) \\ (iii) \end{array} - M_e < 1$ Regime III

The flows before and after the normal shock wave are still respectively isentropic, except that all the properties will jump at the shock wave.

Pressure distributions for flow in a converging-diverging nozzle for different back pressures.

Exit

plane

Throat

Outlines

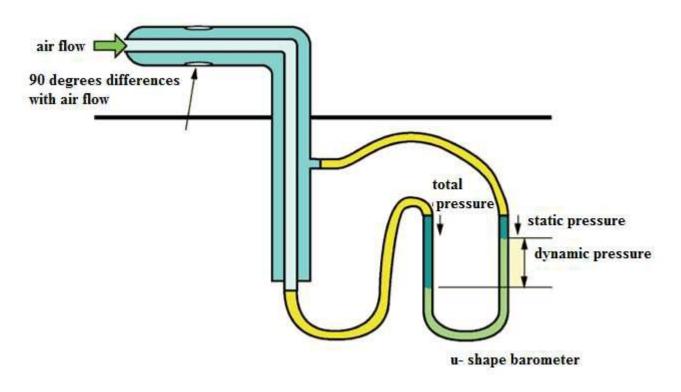




- > Introduction
- **▶** The Basic Normal Shock Equations
- Speed of Sound
- Special Forms of the Energy Equation
- ➤ When Is a Flow Compressible?
- Calculation of Normal Shock-Wave Properties
- > Measurement of Velocity in Compressible Flows



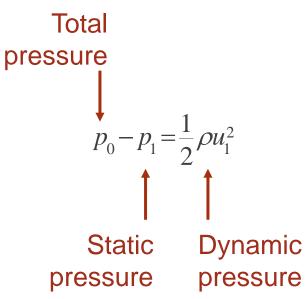
Velocity measurement using a pitot tube



For incompressible flow

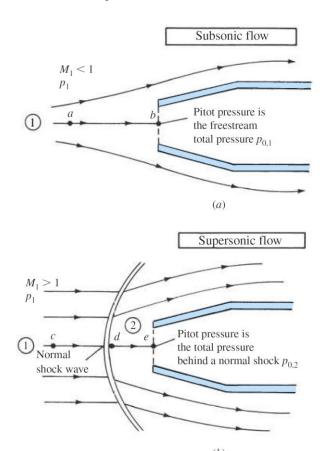
Bernoulli's Equation

$$p_1 + \frac{1}{2}\rho_1 u_1^2 = p_2 + \frac{1}{2}\rho_2 u_2^2$$





Velocity measurement using a pitot tube



- For subsonic compressible flow
- $a \rightarrow b$ is isentropic process.

$$\frac{p_{0,1}}{p_1} = \left(1 + \frac{k-1}{2}M_1^2\right)^{k/(k-1)}$$

$$\downarrow$$

$$M_1^2 = \frac{2}{k-1} \left[\left(\frac{p_{0,1}}{p_1}\right)^{(k-1)/k} - 1 \right]$$

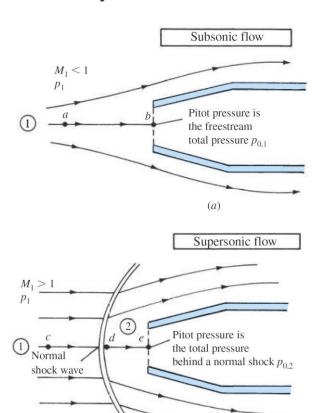
$$\downarrow$$

$$u_1^2 = \frac{2c_1^2}{k-1} \left[\left(\frac{p_{0,1}}{p_1}\right)^{(k-1)/k} - 1 \right]$$

Comparing to incompressible flow, additional knowledge of c_1 is needed for subsonic compressible flow.



Velocity measurement using a pitot tube



- For supersonic compressible flow (method 1)
- $c \rightarrow$ the point before the shock wave is isentropic process.
- $d \rightarrow e$ is also isentropic process.

However, the point before the shock wave \rightarrow d is nonisentropic.

$$\frac{p_{0,2}}{p_1} = \frac{p_{0,2}}{p_2} \frac{p_2}{p_1}$$

$$\frac{p_{0,2}}{p_2} = \left(1 + \frac{k-1}{2}M_2^2\right)^{k/(k-1)} M_2^2 = \frac{1 + \left[(k-1)/2\right]M_1^2}{kM_1^2 - (k-1)/2} \frac{p_2}{p_1} = 1 + \frac{2k}{k+1}(M_1^2 - 1)$$

$$\downarrow \qquad \qquad \downarrow$$

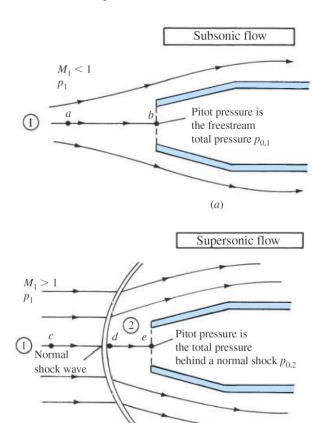
$$\frac{p_{0,2}}{p_1} = \left(\frac{(k+1)^2 M_1^2}{4kM_1^2 - 2(k-1)}\right)^{k/(k-1)} \left(1 + \frac{2k}{k+1}(M_1^2 - 1)\right)$$

The Rayleigh Pitot tube fomula

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Velocity measurement using a pitot tube



For supersonic compressible flow (method 2)

 $d \rightarrow e$ is isentropic process.

The point before the shock wave \rightarrow e can be calculated using the results of the properties across the shock wave.

$$\frac{p_{0,2}}{p_1} = \frac{p_{0,2}}{p_{0,1}} \frac{p_{0,1}}{p_1}$$

$$\frac{p_{0,2}}{p_{0,1}} = \left[1 + \frac{2k}{k+1}(M_1^2 - 1)\right]^{\frac{-1}{k-1}} \left[\frac{(k+1)M_1^2}{2 + (k-1)M_1^2}\right]^{\frac{k}{k-1}} \qquad \frac{p_{0,1}}{p_1} = \left[1 + \frac{k-1}{2}M_1^2\right]^{k/(k-1)}$$

$$\frac{p_{0,2}}{p_1} = \left[\frac{(k+1)^2 M_1^2}{4kM_1^2 - 2(k-1)}\right]^{k/(k-1)} \left[1 + \frac{2k}{k+1}(M_1^2 - 1)\right]$$

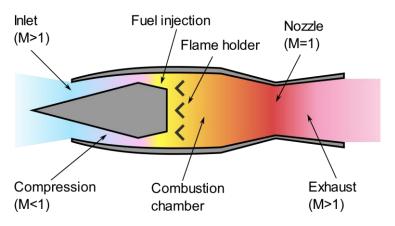
The Rayleigh Pitot tube formula

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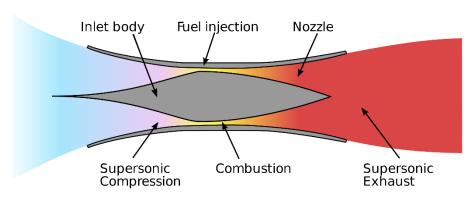




Examples 1 & 2



Ramjet (冲压发动机, ram撞击)



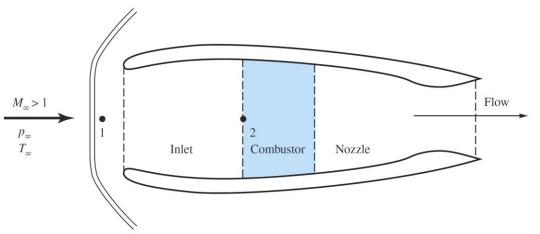
Scramjet (超燃冲压发动机)

- Ramjet is a form of airbreathing jet engine that uses the engine's forward motion to compress incoming air without a compressor.
- Scramjet (supersonic combustion ramjet) is a variant of a ramjet airbreathing jet engine in which combustion takes place in supersonic airflow.





Examples 1 & 2



The ramjet is flying at Mach 2 at a standard altitude of 10 km, where the air pressure and temperature are 2.65 imes 10⁴ N/m² and 223.3 K, respectively. Calculate the air temperature and pressure at point 2 when the Mach number at that point is 0.2.

399 K, 1.42 atm

Repeat Example 1, except for a freestream Mach number M_{∞} = 10. Assume that the ramjet has been redesigned so that the Mach number at point 2 remains equal to 0.2.

4690 K, 32.7 atm

$$\frac{T_{0,1}}{T_{\infty}} = 1 + \frac{k-1}{2} M_{\infty}^{2}$$

$$\frac{p_{0,1}}{p_{\infty}} = \left(1 + \frac{k-1}{2}M_{\infty}^{2}\right)^{k/(k-1)}$$

$$\frac{T_{0,1}}{T_{\infty}} = 1 + \frac{k-1}{2} M_{\infty}^{2} \qquad \frac{p_{0,1}}{p_{\infty}} = \left[1 + \frac{k-1}{2} M_{\infty}^{2}\right]^{k/(k-1)} \qquad \frac{p_{0,1}}{p_{0,\infty}} = \left[1 + \frac{2k}{k+1} (M_{\infty}^{2} - 1)\right]^{\frac{-1}{k-1}} \left[\frac{(k+1)M_{\infty}^{2}}{2 + (k-1)M_{\infty}^{2}}\right]^{\frac{k}{k-1}}$$

$$T_{1,0} = T_{2,0}$$

$$\frac{T_{0,2}}{T_2} = 1 + \frac{k-1}{2}M_2^2$$

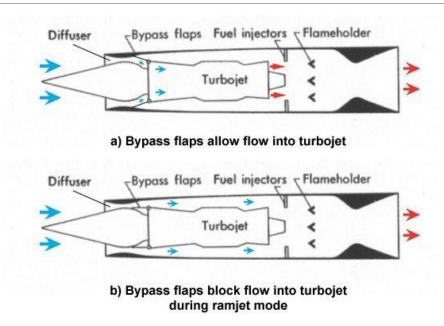
$$\frac{p_{0,2}}{p_2} = \left(1 + \frac{k-1}{2}M_2^2\right)^{k/(k-1)}$$





• Example 3





Consider the Lockheed SR-71 Blackbird flying at a standard altitude of 25 km. The pressure measured by a Pitot tube on this airplane is 3.88×10^4 N/m². Calculate the velocity of the airplane. At an altitude of 25 km, p = 2.5273×10^3 N/m² and T = 216.66 K.

$$\frac{p_{0,1}}{p_{0,\infty}} = \left[1 + \frac{2k}{k+1} (M_{\infty}^2 - 1)\right]^{\frac{-1}{k-1}} \left[\frac{(k+1)M_{\infty}^2}{2 + (k-1)M_{\infty}^2}\right]^{\frac{k}{k-1}} \qquad a = \sqrt{kRT} \qquad V = Ma = 1003 \text{ m/s}$$





Problem 12.54

Air flows from a large tank (p = 650 kPa, T = 550°C) through a converging nozzle, with a throat area of 600 mm², and dischargers to the atmosphere. Determine the mass flow rate of the flow for isentropic flow through the nozzle.

$$\frac{p_t}{p_0} = \frac{101}{650} = 0.155 < 0.528 \rightarrow choked \ flow, M_t = 1.0$$

$$\frac{T_0}{T_t} = \left[1 + \frac{k-1}{2}M_t^2\right]$$

$$\frac{p_0}{p_t} = \left[1 + \frac{k-1}{2}M_t^2\right]^{\frac{k}{k-1}}$$

$$\dot{m} = \rho_t V_t A_t$$





Problem 12.58

Air flows from a large tank ($p = 650 \ kPa$, $T = 550 ^{\circ}$ C) through a converging nozzle, with a throat area of 600 mm², and dischargers to the atmosphere. Determine the mass flow rate of the flow for isentropic flow through the nozzle.

$$\sum F_x = 0 = p_{1,g}A_1 - p_{2,g}A_2 - R_x = -V_1(\dot{m}) + V_2(\dot{m})$$

$$R_x = p_{1,g}A_1 - p_{2,g}A_2 + \dot{m}(V_1 - V_2)$$

$$\frac{p_t}{p_0} = \frac{101}{171} = 0.6257 > 0.528 \rightarrow not \ choked \ flow$$

$$\frac{p_0}{p_t} = \left[1 + \frac{k-1}{2}M_t^2\right]^{\frac{k}{k-1}} \qquad \frac{p_0}{p_1} = \left[1 + \frac{k-1}{2}M_1^2\right]^{\frac{k}{k-1}} \qquad \frac{p_0}{p_2} = \left[1 + \frac{k-1}{2}M_2^2\right]^{\frac{k}{k-1}}$$

$$\frac{T_0}{T_t} = \left[1 + \frac{k-1}{2}M_t^2\right] \qquad \frac{T_0}{T_1} = \left[1 + \frac{k-1}{2}M_1^2\right] \qquad \frac{T_0}{T_2} = \left[1 + \frac{k-1}{2}M_2^2\right]$$





Problem 12.59

Air enters a converging-diverging nozzle at 2MPa and 313 K. At the exit of the nozzle, the pressure is 200 kPa. Assume adiabatic, frictionless flow through the nozzle. The throat area is 20 cm². What is the area at the nozzle exit? What is the mass flow rate of the air?

$$\frac{p_0}{p_e} = \left[1 + \frac{k-1}{2} M_e^2\right]^{\frac{k}{k-1}}$$

$$\frac{A}{A^*} = \frac{1}{M_e} \left[\frac{1 + \frac{k-1}{2} M_e^2}{\frac{k+1}{2}} \right]^{\frac{k+1}{2(k-1)}}$$

$$\dot{m} = A_e p_0 \sqrt{\frac{k}{RT_0}} \left(\frac{2}{k+1}\right)^{\frac{k+1}{2(k-1)}}$$





Problem 12.64

Air discharges through a convergent-divergent nozzle which is attached to a large reservoir. At a point in the nozzle in a normal shock wave is detected across which the absolute pressure jumps from 69 to 207 kPa. Calculate the pressures in the throat of the nozzle and in the reservoir.

$$\frac{p_2}{p_1} = 1 + \frac{2k}{k+1} (M_1^2 - 1)$$

$$\frac{p_0}{p_1} = \left[1 + \frac{k-1}{2}M_1^2\right]^{\frac{k}{k-1}}$$

$$\frac{p_t}{p_0} = 0.528$$





This assignment is due by 6pm on June 10th.

- Upload your solution to BB.
- You can either type your solution out using a word editor like Microsoft Word or clearly hand write the information and then copy/scan your solution to a digital file.