



# ME1020

## Mechanical vibrations

### Lecture 10

Multi DOF system vibration 2 (Free undamped vibration)



# Objectives

- Analyze the free vibration of a 2-DOF undamped system and generalized the results to  $n$ -DOF undamped systems
- Describe the characteristics of the principal mode shapes and beat phenomenon in free vibration

# Revision of matrix algebra

Consider a matrix  $A$  defined by

$$A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

❖ The inverse of  $A$  is (provided  $\det\{A\} \neq 0$ ):

$$A^{-1} = \frac{1}{\det\{A\}} \begin{bmatrix} d & -b \\ -c & a \end{bmatrix}$$

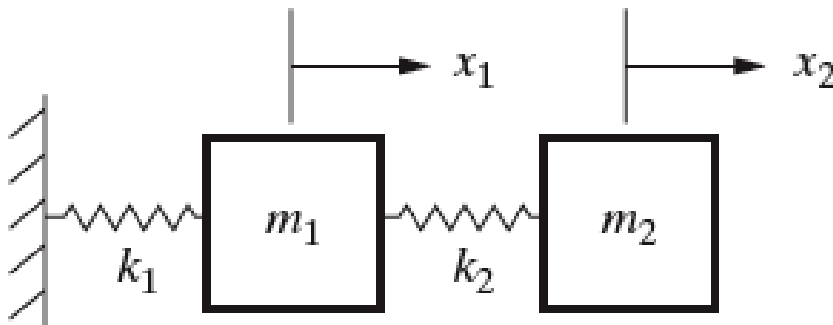
❖ Where the determinant of matrix  $A$  is defined as

$$\det(A) = \det\left\{\begin{bmatrix} a & b \\ c & d \end{bmatrix}\right\} = ad - bc$$

❖ A square matrix is singular if its determinant is zero

# Example 1

Derive the equations of motion for the following 2 DOF system with mass  $m_1 = 9$  kg,  $m_2 = 1$  kg, spring  $k_1 = 24$  N/m, and  $k_2 = 3$  N/m



Total kinetic energy:

$$T = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2$$

$$T = 4.5 \dot{x}_1^2 + \frac{1}{2} \dot{x}_2^2$$

▪ Note  $q_1 = x_1$ ,  $q_2 = x_2$

$$\frac{\partial T}{\partial \dot{q}_1} = \frac{\partial T}{\partial \dot{x}_1} = 9\dot{x}_1$$

$$\frac{\partial T}{\partial q_1} = \frac{\partial T}{\partial x_1} = 0$$

$$\frac{\partial T}{\partial \dot{q}_2} = \frac{\partial T}{\partial \dot{x}_2} = \dot{x}_2$$

$$\frac{\partial T}{\partial q_2} = \frac{\partial T}{\partial x_2} = 0$$

# Example 1

Potential energy:

$$U = \frac{1}{2}k_1x_1^2 + \frac{1}{2}k_2(x_2 - x_1)^2$$

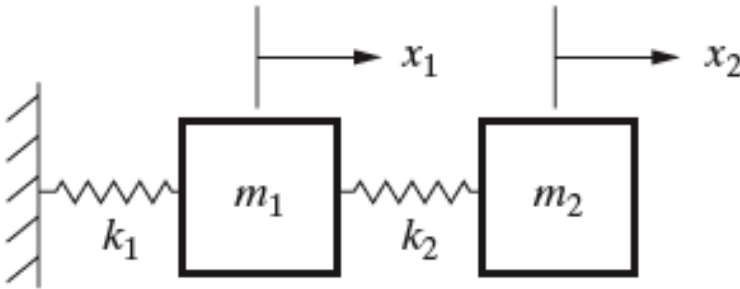
$$U = 12x_1^2 + \frac{3}{2}(x_2 - x_1)^2$$

$$\frac{\partial U}{\partial q_1} = \frac{\partial U}{\partial x_1} = 24x_1 - 3(x_2 - x_1)$$

$$\frac{\partial U}{\partial q_2} = \frac{\partial U}{\partial x_2} = 3(x_2 - x_1)$$

No dissipation function

No generalized forces



# Example 1

For  $k = 1, 2$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_k} \right) - \frac{\partial T}{\partial q_k} + \frac{\partial D}{\partial \dot{q}_k} + \frac{\partial U}{\partial q_k} = Q_k$$

❖ For  $k = 1$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{x}_1} \right) - \frac{\partial T}{\partial x_1} + \frac{\partial D}{\partial \dot{x}_1} + \frac{\partial U}{\partial x_1} = Q_1$$

$$\frac{\partial}{\partial t} (9\dot{x}_1) + 24x_1 - 3(x_2 - x_1) = 0$$

❖ For  $k = 2$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{x}_2} \right) - \frac{\partial T}{\partial x_2} + \frac{\partial D}{\partial \dot{x}_2} + \frac{\partial U}{\partial x_2} = Q_2$$

$$\frac{\partial}{\partial t} (\dot{x}_2) + 3(x_2 - x_1) = 0$$

# Example 1

$$\frac{\partial}{\partial t}(9\dot{x}_1) + 24x_1 - 3(x_2 - x_1) = 0$$

$$\frac{\partial}{\partial t}(\dot{x}_2) + 3(x_2 - x_1) = 0$$

- ❖ The 2 equations are simplified to:

$$9\ddot{x}_1 + 27x_1 - 3x_2 = 0$$

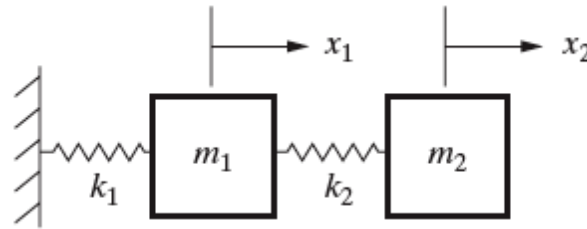
$$\ddot{x}_2 - 3x_1 + 3x_2 = 0$$

- ❖ Express in matrix form:

$$\begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Leftrightarrow [m][\ddot{x}] + [k][x] = [0]$$

- ❖ Note: system is undamped

# Free vibration (2-DOF) undamped



Consider the free vibration of the system given in example 1:

$$\begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Leftrightarrow [m][\ddot{x}] + [k][x] = [0]$$

Without damping and forces, the system should vibrate continuously (at the natural frequencies) under initial condition disturbance. Assume a solution of the form

$$[x(t)] = \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix} \cos(\omega t + \phi) = [u]e^{j\omega t}$$

- ❖  $[u]$  is a non zero vector of constants to be determined along with  $\omega$
- ❖  $j = \sqrt{-1}$  and  $e^{j\omega t} = \cos \omega t + j \sin \omega t$  represents harmonic motion



# Free vibration (2-DOF) undamped

$$[x(t)] = [u]e^{j\omega t}$$
$$[\ddot{x}(t)] = -\omega^2[u]e^{j\omega t}$$

Substitute into

$$[m][\ddot{x}(t)] + [k][x(t)] = [0]$$
$$-\omega^2[m][u]e^{j\omega t} + [k][u]e^{j\omega t} = [0]$$
$$\{-\omega^2[m] + [k]\}[u]e^{j\omega t} = [0]$$
$$\{-\omega^2[m] + [k]\}[u] = [0]$$

- ❖ Note:  $[u] \neq 0$  and this means that inverse of  $\{-\omega^2[m] + [k]\}$  cannot exist, which implies that  $\det\{-\omega^2[m] + [k]\} = 0$
- ❖ For the 2-DOF system, we can obtain two values for  $\pm\omega$  by solving

$$\det\{-\omega^2[m] + [k]\} = 0$$

# Free vibration (2-DOF) undamped

To illustrate the process, we will utilize the values for  $[m]$  and  $[k]$ :

$$[m] = \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{and} \quad [k] = \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix}$$

❖  $\det\{-\omega^2[m] + [k]\} = 0$  will give the characteristic equation

$$\det\left\{-\omega^2 \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix}\right\} = \det\left\{\begin{bmatrix} 27 - 9\omega^2 & -3 \\ -3 & 3 - \omega^2 \end{bmatrix}\right\} = 0$$

$$(27 - 9\omega^2)(3 - \omega^2) - 9 = 0$$

$$9\omega^4 - 54\omega^2 + 72 = 0$$

$$\omega^4 - 6\omega^2 + 8 = 0 \quad \text{or} \quad (\omega^2 - 2)(\omega^2 - 4) = 0$$

❖ We get two values for  $\pm\omega$  and these are  $\omega_1 = \pm\sqrt{2}$  and  $\omega_2 = \pm 2$  rad/s after solving the characteristic equation

# Free vibration (2-DOF) undamped

❖ The two values for  $\pm\omega$  are  $\omega_1 = \pm\sqrt{2}$  and  $\omega_2 = \pm 2$  rad/s

❖ We can use these values to solve for  $[u]$  using

$$\{-\omega^2[m] + [k]\}[u] = [0]$$

❖ For  $\omega_1 = \pm\sqrt{2}$  rad/s:

$$\left\{-2 \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix}\right\} [u]_1 = [0]$$

$$\begin{bmatrix} 9 & -3 \\ -3 & 1 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$9X_1 - 3X_2 = 0 \quad \text{and} \quad -3X_1 + X_2 = 0$$

❖ We cannot solve for  $X_1$  and  $X_2$  but can get their ratio  $X_1/X_2 = 1/3$

❖ For  $X_2 = 1$ ,  $X_1 = 1/3$  and

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1/3 \\ 1 \end{bmatrix}_1$$

# Free vibration (2-DOF) undamped

❖ The two values for  $\pm\omega$  are  $\omega_1 = \pm\sqrt{2}$  and  $\omega_2 = \pm 2$  rad/s

❖ We can use these values to solve for  $[u]$  using

$$\{-\omega^2[m] + [k]\}[u] = [0]$$

❖ For  $\omega_2 = \pm 2$  rad/s

$$\left\{-4 \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix}\right\} [u]_2 = [0]$$

$$\begin{bmatrix} -9 & -3 \\ -3 & -1 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$-9X_1 - 3X_2 = 0 \quad \text{and} \quad -3X_1 - X_2 = 0$$

❖ We cannot solve for  $X_1$  and  $X_2$  but can get their ratio  $X_1/X_2 = -1/3$

❖ For  $X_2 = 1$ ,  $X_1 = -1/3$  and

$$[u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} -1/3 \\ 1 \end{bmatrix}_2$$

# Free vibration (2-DOF) undamped

- ❖ We have  $\omega_1 = \pm\sqrt{2}$  rad/s and  $\omega_2 = \pm 2$  rad/s
- ❖ We also have

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1/3 \\ 1 \end{bmatrix}_1 \text{ and } [u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} -1/3 \\ 1 \end{bmatrix}_2$$

- ❖ The assumed solution is  $[x(t)] = [u]e^{j\omega t}$  which is a combination of

$$[x(t)] = (ae^{j\omega_1 t} + be^{-j\omega_1 t})[u]_1 + (ce^{j\omega_2 t} + de^{-j\omega_2 t})[u]_2$$

- ❖ Where  $a$ ,  $b$ ,  $c$ , and  $d$  are to be determined based on initial conditions
- ❖ Applying Euler's equation, an alternate simpler form of solution is

$$[x(t)] = A \sin(\omega_1 t + \phi_1) [u]_1 + B \sin(\omega_2 t + \phi_2) [u]_2$$

- ❖ Where  $A$ ,  $B$ ,  $\phi_1$ , and  $\phi_2$  are to be determined based on initial conditions
- ❖ In this form it is clear that the natural frequencies are  $\omega_1$  and  $\omega_2$
- ❖  $[u]_1$  and  $[u]_2$  are called the mode shapes

# Free vibration (2-DOF) undamped

To illustrate the process, we will continue to obtain the free response for the system with initial conditions  $x_1(0) = 1 \text{ mm}$ ,  $x_2(0) = \dot{x}_1(0) = \dot{x}_2(0) = 0$

$$[x(t)] = A \sin(\omega_1 t + \phi_1) [u]_1 + B \sin(\omega_2 t + \phi_2) [u]_2$$

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = A \sin(\sqrt{2}t + \phi_1) \begin{bmatrix} 1/3 \\ 1 \end{bmatrix}_1 + B \sin(2t + \phi_2) \begin{bmatrix} -1/3 \\ 1 \end{bmatrix}_2$$

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} \frac{1}{3}A \sin(\sqrt{2}t + \phi_1) - \frac{1}{3}B \sin(2t + \phi_2) \\ A \sin(\sqrt{2}t + \phi_1) + B \sin(2t + \phi_2) \end{bmatrix}$$

Differentiate to get the velocities

$$\begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{3}A \cos(\sqrt{2}t + \phi_1) - \frac{2}{3}B \cos(2t + \phi_2) \\ \sqrt{2}A \cos(\sqrt{2}t + \phi_1) + 2B \cos(2t + \phi_2) \end{bmatrix}$$

# Free vibration (2-DOF) undamped

A time  $t = 0$ :  $x_1(0) = 1$  mm,  $x_2(0) = x_2(0) = \dot{x}_1(0) = \dot{x}_2(0) = 0$

$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} \frac{1}{3}A \sin(\phi_1) - \frac{1}{3}B \sin(\phi_2) \\ A \sin(\phi_1) + B \sin(\phi_2) \end{bmatrix}$$

$$\begin{bmatrix} 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \frac{\sqrt{2}}{3}A \cos(\phi_1) - \frac{2}{3}B \cos(\phi_2) \\ \sqrt{2}A \cos(\phi_1) + 2B \cos(\phi_2) \end{bmatrix}$$

There are four equations:

$$\frac{1}{3}A \sin(\phi_1) - \frac{1}{3}B \sin(\phi_2) = 1$$

$$A \sin(\phi_1) + B \sin(\phi_2) = 0$$

$$\frac{\sqrt{2}}{3}A \cos(\phi_1) - \frac{2}{3}B \cos(\phi_2) = 0$$

$$\sqrt{2}A \cos(\phi_1) + 2B \cos(\phi_2) = 0$$

$$\left. \begin{array}{l} \frac{\sqrt{2}}{3}A \cos(\phi_1) - \frac{2}{3}B \cos(\phi_2) = 0 \\ \sqrt{2}A \cos(\phi_1) + 2B \cos(\phi_2) = 0 \end{array} \right\} \begin{array}{l} \frac{2\sqrt{2}}{3}A \cos(\phi_1) = 0 \\ \phi_1 = \pi/2 \\ \phi_2 = \pi/2 \end{array}$$

# Free vibration (2-DOF) undamped

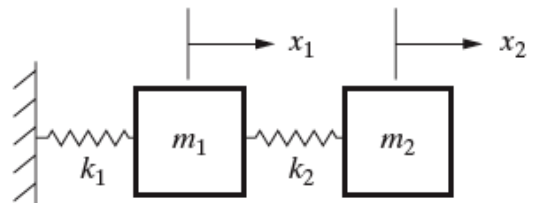
Found  $\phi_1 = \pi/2$  and  $\phi_2 = \pi/2$

$$\frac{1}{3}A \sin(\phi_1) - \frac{1}{3}B \sin(\phi_2) = \frac{1}{3}A - \frac{1}{3}B = 1$$

$$A \sin(\phi_1) + B \sin(\phi_2) = A + B = 0$$

Solve to get:  $2A = 3$  or  $A = 1.5$  mm and  $B = -1.5$  mm

The system is:

$$\begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 27 & -3 \\ -3 & 3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$


The solution is:

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} 0.5 \sin(\sqrt{2}t + \pi/2) + 0.5 \sin(2t + \pi/2) \\ 1.5 \sin(\sqrt{2}t + \pi/2) - 1.5 \sin(2t + \pi/2) \end{bmatrix}$$

❖ The response contains both natural frequencies of the system



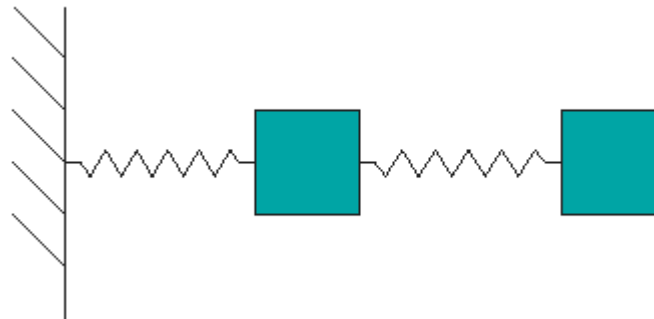
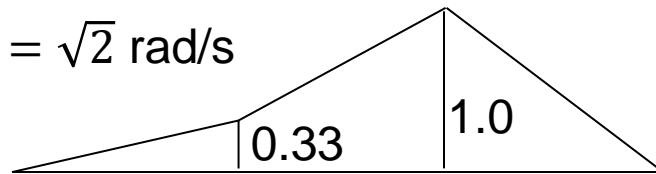
# Principal modes

The principal modes for the 2-DOF system discussed are the mode shapes

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1/3 \\ 1 \end{bmatrix}_1 \quad \text{and} \quad [u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} -1/3 \\ 1 \end{bmatrix}_2$$

Consider the first mode where  $[x]_1 = A_1 \sin(\sqrt{2}t) [u]_1$

$$\omega_1 = \sqrt{2} \text{ rad/s}$$

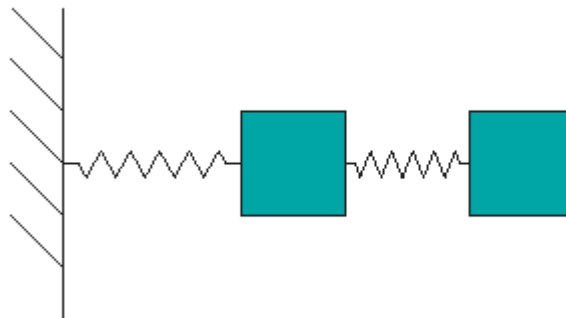
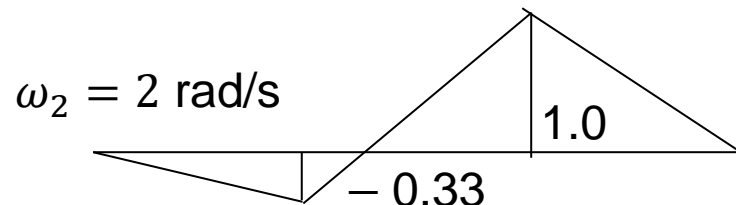


# Principal modes

The principal modes for the 2-DOF system discussed are the mode shapes

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1/3 \\ 1 \end{bmatrix}_1 \quad \text{and} \quad [u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} -1/3 \\ 1 \end{bmatrix}_2$$

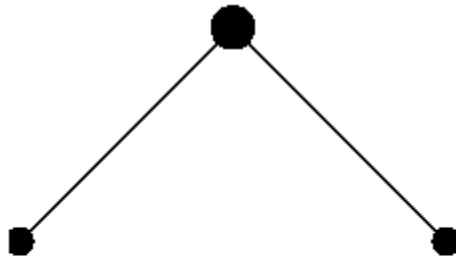
Consider the second mode where  $[x]_2 = A_2 \sin(2t) [u]_2$



# Mode shapes

A 1 DOF system has one mode of vibration and 1 natural frequency

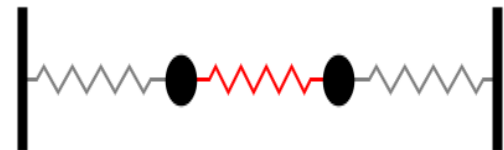
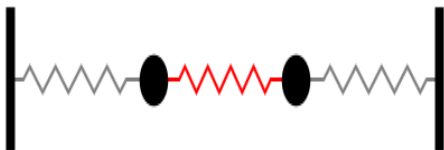
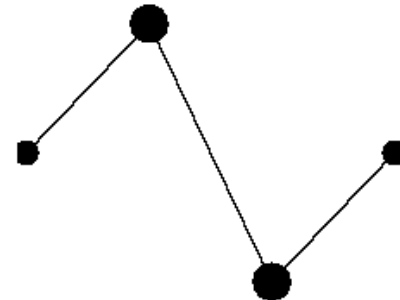
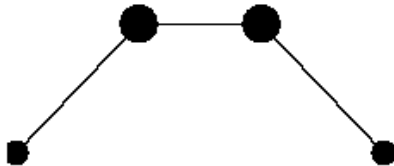
$$\omega_n = \sqrt{\frac{k}{m}}$$



# Mode shapes

A 2 DOF system has two modes of vibration and 2 natural frequencies

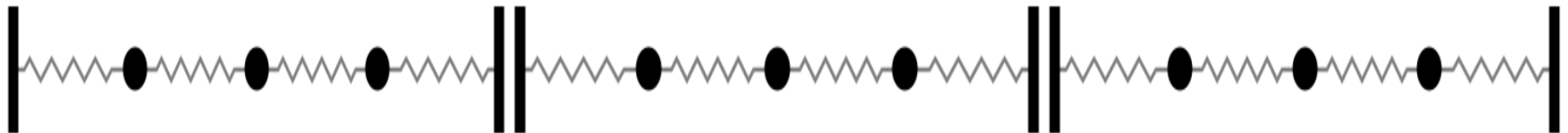
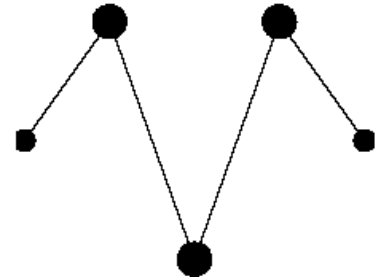
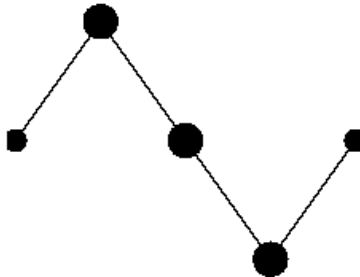
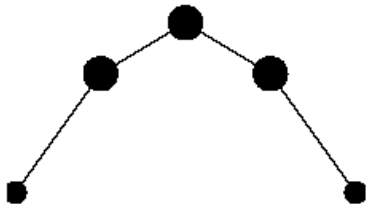
$$\omega_n \neq \sqrt{\frac{k}{m}}$$



# Mode shapes

A 3 DOF system has three modes of vibration and 3 natural frequencies

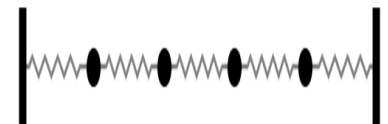
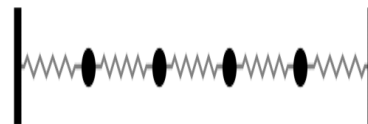
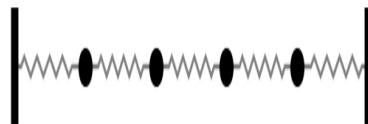
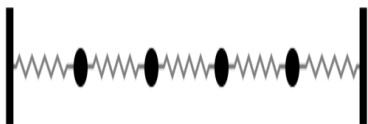
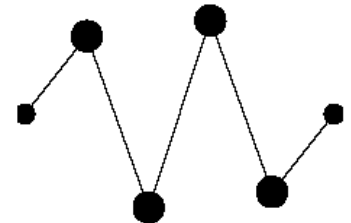
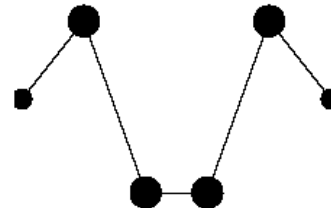
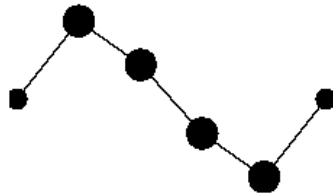
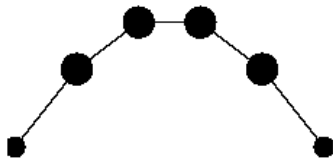
$$\omega_n \neq \sqrt{\frac{k}{m}}$$



# Mode shapes

A 4 DOF system has four modes of vibration and 4 natural frequencies

$$\omega_n \neq \sqrt{\frac{k}{m}}$$



# Mode shapes

A Discrete system consists of a finite number of masses or rigid bodies. The DOF of a discrete system can be viewed as:

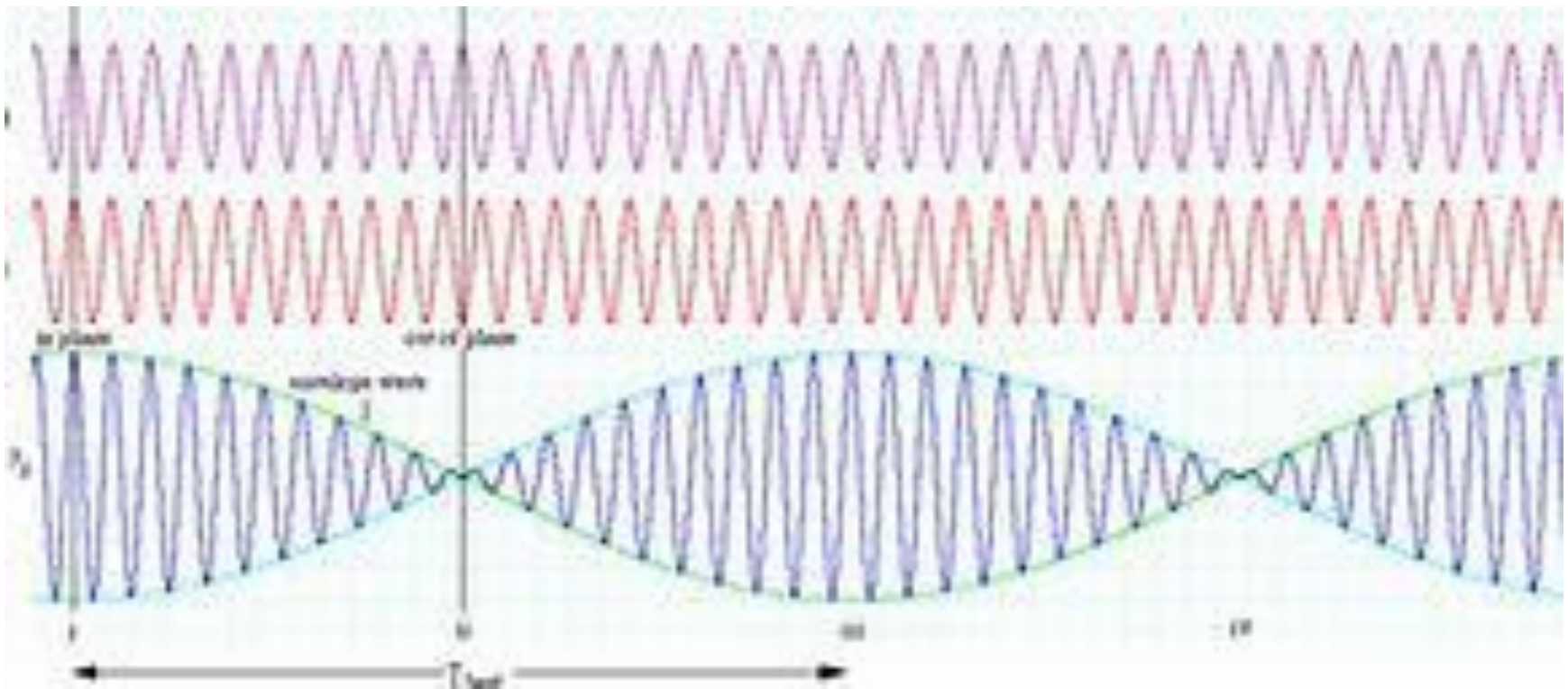
$$\text{DOF} = \sum (\text{number of possible motion of mass } i)$$

- ❖ A Continuous system has distributed mass properties.
- ❖ So far we have modeled the system using discrete masses. The mode shape concept can be extended to a continuous system



# Beats

For a 2-DOF free undamped vibration system, the beats phenomenon occurs when the 2 natural frequencies are near each other





# Beats

Spring:

$n=50.5$   $D_m=60\text{mm}$   $d=3\text{mm}$

$m_s=0.545\text{kg}$

$D=75.65\text{N/m}$   $D^*=86.91\text{Nmm}$

Mass bob:

$m=3.375\text{kg}$   $I\approx 0.00383\text{kgm}^2$

Eigen angular frequency:

$\omega_{01}\approx\omega_{02}=4.62\text{ rad/sec}$



# Free vibration (n-DOF) undamped

The free vibration analysis for the 2-DOF system can be extended to a  $n$ -DOF:

- ❖ Determine the  $n$  number of natural frequencies  $\omega_i$  for  $i = 1, \dots, n$  using

$$\det\{-\omega^2[m] + [k]\} = 0$$

- ❖ For  $i = 1, \dots, n$  use natural frequency  $\omega_i$  to determine the mode shape  $[u]_i$

$$\{-\omega_i^2[m] + [k]\}[u]_i = [0]$$

- ❖ The general solution is of the form

$$[x(t)] = \sum_{i=1}^n [u]_i A_i \sin(\omega_i t + \phi_i)$$

- ❖ Initial conditions can be specified for  $[x(0)] = \begin{bmatrix} x_1(0) \\ \vdots \\ x_n(0) \end{bmatrix}$  and  $[\dot{x}(0)] = \begin{bmatrix} \dot{x}_1(0) \\ \vdots \\ \dot{x}_n(0) \end{bmatrix}$

- ❖ Solve for the integration constants  $A_i$  and  $\phi_i$  using

$$[x(0)] = \sum_{i=1}^n [u]_i A_i \sin(\phi_i) \text{ and } [\dot{x}(0)] = \sum_{i=1}^n [u]_i \omega_i A_i \cos(\phi_i)$$

# Repeated frequencies

- ❖ If the characteristic equation has  $m$  repeated roots, then there are  $m$  natural frequencies with modes that coincide. This means that only  $n - m$  of the linear algebraic equations from which the mode shape is calculated are independent. Thus the  $m$  elements of the mode shape can be arbitrary chosen and the most general mode shape involves  $m$  arbitrary constants. These are used to specify the  $m$  linearly independent mode shape

$$[u]_i, [u]_{i+1}, \dots, [u]_{i+m}$$

and the general solution can be applied

- ❖ Note: the method used to compute the natural frequencies and mode shapes presented is not the most efficient way to solve vibration problems

# Example 2

Determine the natural frequencies and mode shapes for the following system:

$$\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 2k & 0 \\ 0 & 2k \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

❖ Determine the 2 natural frequencies  $\omega_i$  for  $i = 1, 2$  using

$$\det\{-\omega^2[m] + [k]\} = 0$$

$$\det\left\{-\omega^2 \begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} + \begin{bmatrix} 2k & 0 \\ 0 & 2k \end{bmatrix}\right\} = \det\left\{\begin{bmatrix} 2k - m\omega^2 & 0 \\ 0 & 2k - m\omega^2 \end{bmatrix}\right\} = 0$$
$$(2k - m\omega^2)(2k - m\omega^2) = 0$$

We get two values for  $\pm\omega$  and these are  $\omega_1 = \omega_2 = \pm\sqrt{2k/m}$  rad/s

The natural frequencies are repeated

# Example 2

- ❖ We found the 2 natural frequencies are  $\omega_1 = \omega_2 = \pm\sqrt{2k/m}$  rad/s
- ❖ For  $\omega = \sqrt{2k/m}$  rad/s:

$$\{-\omega^2[m] + [k]\}[u] = [0]$$

$$\begin{bmatrix} 2k - m\omega^2 & 0 \\ 0 & 2k - m\omega^2 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = [0]$$
$$0(X_1) + 0(X_2) = 0$$

- ❖ This implies that  $X_1$  and  $X_2$  can take on any values
- ❖ Two convenient pairs of mode shapes (perpendicular to each other) are

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}_1 \quad \text{and} \quad [u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}_2$$

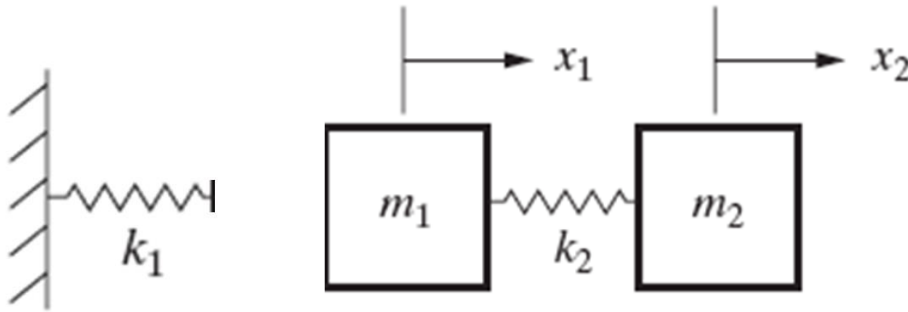
These can be used to solve for the free vibration when given the initial conditions

# Unrestrained systems

- ❖ Note that the natural frequencies are obtained after solving the characteristics equation
- ❖ The characteristic equation only have real roots when mass and stiffness matrices are symmetry
- ❖ The system is unstable if the characteristics equation has complex roots
- ❖ If the stiffness matrix is singular (i.e.  $\det\{[k]\} = 0$ ), then one of the natural frequency is zero
- ❖ The system is unrestrained when one of the natural frequency is zero

# Example 3

Analyze the free vibration of example 1 if the spring  $k_1$  is now disconnected from mass  $m_1$  as shown (note: mass = 9 kg,  $m_2 = 1$  kg, spring  $k_2 = 3$  N/m)



Total kinetic energy:

$$T = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2$$

$$T = 4.5 \dot{x}_1^2 + \frac{1}{2} \dot{x}_2^2$$

▪ Note  $q_1 = x_1$ ,  $q_2 = x_2$

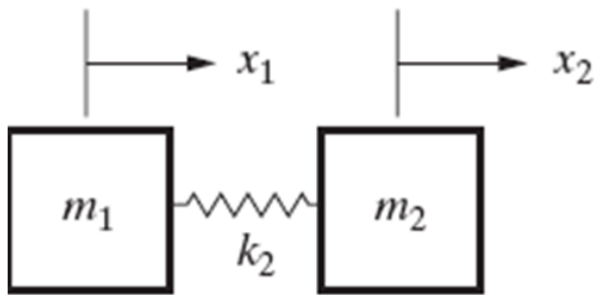
$$\frac{\partial T}{\partial \dot{q}_1} = \frac{\partial T}{\partial \dot{x}_1} = 9\dot{x}_1$$

$$\frac{\partial T}{\partial q_1} = \frac{\partial T}{\partial x_1} = 0$$

$$\frac{\partial T}{\partial \dot{q}_2} = \frac{\partial T}{\partial \dot{x}_2} = \dot{x}_2$$

$$\frac{\partial T}{\partial q_2} = \frac{\partial T}{\partial x_2} = 0$$

# Example 3



Potential energy:

$$U = \frac{1}{2} k_2 (x_2 - x_1)^2$$

$$U = \frac{3}{2} (x_2 - x_1)^2$$

$$\frac{\partial U}{\partial q_1} = \frac{\partial U}{\partial x_1} = -3(x_2 - x_1)$$

$$\frac{\partial U}{\partial q_2} = \frac{\partial U}{\partial x_2} = 3(x_2 - x_1)$$

No dissipation function

No generalized forces



# Example 3

For  $k = 1, 2$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_k} \right) - \frac{\partial T}{\partial q_k} + \frac{\partial D}{\partial \dot{q}_k} + \frac{\partial U}{\partial q_k} = Q_k$$

❖ For  $k = 1$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{x}_1} \right) - \frac{\partial T}{\partial x_1} + \frac{\partial D}{\partial \dot{x}_1} + \frac{\partial U}{\partial x_1} = Q_1$$

$$\frac{\partial}{\partial t} (9\dot{x}_1) - 3(x_2 - x_1) = 9\ddot{x}_1 - 3(x_2 - x_1) = 0$$

❖ For  $k = 2$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{x}_2} \right) - \frac{\partial T}{\partial x_2} + \frac{\partial D}{\partial \dot{x}_2} + \frac{\partial U}{\partial x_2} = Q_2$$

$$\frac{\partial}{\partial t} (\dot{x}_2) + 3(x_2 - x_1) = \ddot{x}_2 + 3(x_2 - x_1) = 0$$

# Example 3

$$\begin{aligned}9\ddot{x}_1 - 3(x_2 - x_1) &= 0 \\ \ddot{x}_2 + 3(x_2 - x_1) &= 0\end{aligned}$$

❖ Express in matrix form:

$$\begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 3 & -3 \\ -3 & 3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Leftrightarrow [m][\ddot{x}] + [k][x] = [0]$$

❖ Determine the 2 natural frequencies  $\omega_i$  for  $i = 1, 2$  using

$$\det\{-\omega^2[m] + [k]\} = 0$$

$$\det\left\{-\omega^2 \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 3 & -3 \\ -3 & 3 \end{bmatrix}\right\} = \det\left\{\begin{bmatrix} 3 - 9\omega^2 & -3 \\ -3 & 3 - \omega^2 \end{bmatrix}\right\} = 0$$
$$(3 - 9\omega^2)(3 - \omega^2) - 9 = 0$$

$$9\omega^4 - 30\omega^2 = 0 \quad \text{or} \quad \omega^2(9\omega^2 - 30) = 0$$

We get two values for  $\pm\omega$  and these are  $\omega_1 = 0$  and  $\omega_2 = \pm 1.826$  rad/s

Note the system is restrained

# Example 3

- ❖ The two values for  $\pm\omega$  are  $\omega_1 = 0$  and  $\omega_2 = \pm 1.826$  rad/s
- ❖ We can use these values to solve for  $[u]$  using

$$\{-\omega^2[m] + [k]\}[u] = [0]$$

- ❖ For  $\omega_1 = 0$  rad/s:

$$\left\{ \begin{bmatrix} 3 & -3 \\ -3 & 3 \end{bmatrix} \right\} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = [0]$$

$$[u]_1 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_1 = \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1$$

- ❖ For  $\omega_2 = \pm 1.826$  rad/s

$$\left\{ -3.334 \begin{bmatrix} 9 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 3 & -3 \\ -3 & 3 \end{bmatrix} \right\} [u]_2 = [0]$$

$$\begin{bmatrix} -27 & -3 \\ -3 & -0.334 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad \text{or} \quad [u]_2 = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix}_2 = \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2$$

# Example 3

To illustrate the process, we will continue to obtain the free response for the system with initial conditions  $x_1(0) = 1$  mm,  $x_2(0) = \dot{x}_1(0) = \dot{x}_2(0) = 0$

- ❖ The solution of the free vibration of the undamped and unrestrained system is a combination of the modes

$$\begin{aligned} [x(t)] &= A \sin(\omega_1 t + \phi_1) [u]_1 + B \sin(\omega_2 t + \phi_2) [u]_2 \\ \begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} &= A \sin(\phi_1) \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1 + B \sin(1.826t + \phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2 \\ \begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \end{bmatrix} &= 1.826B \cos(1.826t + \phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2 \end{aligned}$$

Applying the initial conditions at  $t = 0$

$$\begin{aligned} \begin{bmatrix} 1 \\ 0 \end{bmatrix} &= A \sin(\phi_1) \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1 + B \sin(\phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2 \\ \begin{bmatrix} 0 \\ 0 \end{bmatrix} &= 1.826B \cos(\phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2 \end{aligned}$$

# Example 3

$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} = A \sin(\phi_1) \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1 + B \sin(\phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2$$

$$\begin{bmatrix} 0 \\ 0 \end{bmatrix} = 1.826B \cos(\phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2$$

These consist of four equations

$$1 = A \sin(\phi_1) - 0.111B \sin(\phi_2)$$

$$0 = A \sin(\phi_1) + B \sin(\phi_2)$$

$$0 = -0.2027B \cos(\phi_2)$$

$$0 = 1.826B \cos(\phi_2)$$

$$\left. \begin{array}{l} \phi_2 = \pi/2 \\ \text{or } 3\pi/2 \end{array} \right\}$$

Using  $\phi_2 = \pi/2$  or  $3\pi/2$

$$1 = A \sin(\phi_1) \mp 0.111B$$

$$0 = A \sin(\phi_1) \pm B$$

$$\left. \begin{array}{l} 1 = \pm 1.111B \\ B = \pm 0.9 \end{array} \right\}$$

Using  $B = \pm 0.9$ , we get  $A \sin(\phi_1) = \mp 0.9$

# Example 3

- ❖ The solution of the free vibration of the undamped and unrestrained system is

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = A \sin(\phi_1) \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1 + B \sin(1.826t + \phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2$$

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \mp 0.9 \begin{bmatrix} 1 \\ 1 \end{bmatrix}_1 \pm 0.9 \sin(1.826t + \phi_2) \begin{bmatrix} -0.111 \\ 1 \end{bmatrix}_2$$

$$\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix} = \mp 0.9 \pm \sin(1.826t + \phi_2) \begin{bmatrix} -0.1 \\ 0.9 \end{bmatrix}$$

Where plus (minus ) sign is to be used when  $\phi_2$  is taken as  $\pi/2$  ( $3\pi/2$ )