To: Professor Lu

From: Christopher Jin, Frederic Liu, Gordon Lou, Owen Chen, Yanjun He

Subject: ME1042 Lab 01 Forced and Free Vibrations

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On $2^{\rm nd}$ and $16^{\rm th}$ September, Gordon Lou, Owen Chen, Frederic Liu, Yanjun He, and I conducted our first experiment in the course Mechanical Measurements 2, from which we have explored the fundamentals of mechanical vibrations, including natural frequency and damping on a rigid beam.

In the first part of experiment, we researched the influence of added mass on the natural frequency of the cantilever beam and also examine the impact of frequency ratio on the amplitude and phase lag of vibration. In order to achieve these two goals, we recorded the natural frequency with different masses added to the cantilever beam and also the amplitude and phase lag of vibration with different frequencies of force applied to cantilever beam into Table 4, whose results are shown in Figure 1, 2, and 3, respectively.

Firstly, we conduct the error analysis in this experiment. Through our measurements as shown in Figure 1, It is obvious that the measured value is less than the theoretical value. After calculation, there is about 5% error between the measured value and the theoretical value and the measured value is always less than the theoretical value. The theoretical frequency is given by Equation 6. From this, we can know that the first possibility of error is caused by that the actual I_A is larger than the theoretical I_A . If the mass of the connecting part is ignored when measuring the mass of each part of the system, or the reading is smaller than the actual value, I_A will be smaller. Another possibility is that the actual k is smaller than the measured k. For example, when measuring the spring k value, the reading is too large, which will also lead to theoretical frequency greater than measured frequency (Sorokin & Ershova, 2003). In addition, we regard beam as a rigid body in Equation 6. However, in fact, the beam also has a certain elasticity. If the elasticity of the beam is considered, the k of the system will be smaller than that of the spring.

Secondly, we compare the natural frequency and damping ratio for two conditions - undamped and fully shut. For the damping ratio, the damp of the fully shut case is mainly come from the viscous damper. For the undamped case, it comes from the air fraction. In theory, because the damping liquid has a damping constant much larger than the air, the fully shut case should have a damping ratio larger than that of the undamped case. From Table 1, the recorded data conforms our expectation from the theory. The damping ratio of the fully shut case is 0.029, which is 14.5 times larger than that of the undamped case. For the natural frequency, in theory, from Equation 1, the damped natural frequency becomes lower as the damping ratio increases. However, in our results, the two natural frequencies are nearly the same. We hold the opinion that contradict is mainly because the damping ratio is too small in the order of magnitudes, which has low influence on the damped natural frequency (So et al., 2000). In the experiment we use the same beam for both the undamped and the full shut case, which means the system would have the same natural frequency. Assuming the natural frequency of the undamped beam in ideal conditions is ω_n , and substitute the damping ratio into Equation 1, we can obtain the damped natural frequency for the undamped case is: $\omega_d = 0.999998\omega_n$, while for the fully shut case is $\omega_d = 0.999579\omega_n$. The two natural frequencies vary after four decimal places, so we can say that the effect of the damper has low effect on the damped natural frequency.

Thirdly, we find the magnification factor for a speed ratio of 1.0 for two sets of results and compare them. For forced vibration, the magnification factor is shown in Equation 2. When

r approaches to 0, the magnification factor β approaches to 1. Therefore, when speed ratio is small, $\beta \approx 1$. In other words, $X \approx X_{st}$. And from Figure 2, we can know that when r approaches to 0.75, the amplitude is converging to some vale and the value is already stable. Also, damper system ($\zeta > 0$) will converge first than the undamped system ($\zeta \approx 0$). Therefore, we can use $X_{r=0.75,fully\ shut}$ to replace X_{st} . According to the previous deduction, the magnification factors for the condition of undamped and fully shut in this experiment can be calculate using Equation 4, which are equal to 135.6 and 13.4, respectively. The difference between these two magnification factors are huge, which is because when r approaches to 1, the magnification factor is shown in Equation 5. Therefore, at r=1, the magnification factor β only depends on the damping ratio. In undamped case, the damping ratio $\zeta_{undamped} \approx 0$, which will make the magnification factor extremely high. But in fully shut case, the damping ratio $\zeta_{fully\ shut}$ is much bigger than 0, which will decrease the magnification factor of this system dramatically. That's why $\beta_{undamped} \gg \beta_{fully\ shut}$ at r=1.0. And we also use Equation 5 to calculate the theoretical magnification factors, which are shown in Table 1.

Finally, we give our own comments on the amplitude and phase lag values vs speed ratio plot. In undamped and forced vibration case, the amplitude as shown in Equation 8 should approach infinity when r is close to 1. The result is consistent with theoretical analysis. According to the Figure 2, the amplitude of undamped case approaches to 0 when the frequency ratio is further away from 1; The amplitude becomes very large and increases rapidly when the frequency ratio approaches 1. Since the recorded data points are discrete and the test equipment has more or less damping, there will not be infinite situation in real experiment. In damped and forced vibration case, the amplitude is shown in Equation 9. Since the denominator has minimum value when the value of r is the same as Equation 10, the maximum amplitude occurs at r slightly greater than 1. The result is consistent with theoretical analysis. The amplitude of damped case also approaches to 0 when the frequency ratio is further away from 1 and it has maximum value when the frequency ratio is about 1.01. The maximum value is $2.019 \mathrm{~cm}$ based on the experimental data. The phase lag as shown in Equation 11 should increase from 0° to 180° and it will increase rapidly as r approaches 1 (Calim, 2009). Also, the lower the damping ratio, the higher the slope around r = 1. The figure drawn by the two groups of data as shown in Figure 3 is both consistent with the theoretical analysis. The slope of undamped curve is obviously larger than that of damped curve, and the ordinate scale ranges from 180° to 360° , which is equivalent to the scale range from 0° to 180° .

After this experiment, we have understood the fundamentals of natural frequency and damping on a rigid beam by observing oscillation affects, which promotes the theories we have learned in class. In the future, if we have the opportunity to research civil engineering and geotechnical monitoring, today's experiment will provide us with tremendous help regarding to its usefulness and significance.

In the second part of experiment, we researched the influence of added mass on the natural frequency of the cantilever beam of simply supported beam and also examine the impact of frequency ratio on the amplitude and phase lag of vibration.

Firstly, we find the effective mass using Rayleigh's improved solution as shown in Equation 12 and then calculate the theoretical oscillation frequency using Equation 13 whose results are shown in Table 3. Compared the calculation to the measured experimental value, we can find that the theoretical natural frequencies are always slightly smaller than the measured one, which may be because the blocks we added have a certain amount of wear, resulting in the actual mass of the block being less than 400 grams. Therefore, the frequency we measured is greater than the frequency calculated.

Secondly, we compare the natural frequency and damping ratio for two conditions - undamped and fully shut in the simply supported case. For the damping ratio of the simply supported beam experiment, comparing the result of the fully shut with the undamped case, the result shows the same tendency as the beam and spring experiment, which is because the

damping constant of the damping liquid is larger than that of the air. However, when comparing two sets of the damping ratios between the two experiments, our group find that the simply supported case shows a larger damping ratio than that of the beam and spring case whether or not the system is damped. As discussed in the previous part, the effective mass of the simply supported beam is shown in Equation 12, which is lower than the total mass of the system. According of the definition of the in Equation 16, damping ratio is negatively correlated with the effective mass. Thus, with the same damping constant, the decreasing of the effective mass is the main cause of the increasing of the damping ratio. For the natural frequency, from Equation 1, the damped natural frequency becomes lower as the damping ratio increases. This time, as the damping ratio becomes larger, the decrease in the recorded value of the natural frequency becomes more prominent, the value of the fully shut case is $0.12 \, \mathrm{rad/s}$ smaller than that of the undamped case.

Thirdly, we find the magnification factor for a speed ratio of 1.0 for two sets of results in the simply supported case and compare them. We use Equation 5 to calculate the theoretical magnification factors, which are shown in Table 2. At r=1, the magnification factor β only depends on the damping ratio. In undamped case, the damping ratio $\zeta_{undamped} \approx 0$, which will make the magnification factor extremely high. But in fully shut case, the damping ratio $\zeta_{fully\ shut} > 0$, which will decrease the magnification factor of this system dramatically. That's why $\beta_{undamped} \gg \beta_{fully\ shut}$ at r=1.0.

After this experiment, we have learnt how to analyze the beam model in the case of free vibration and simple support, calculate the theoretical natural frequency of the beam and compare it with the measured frequency, and find several reasons for the error and studied forced vibration and damping, which promotes the theories we have learned in class. In the future, if we have the opportunity to research from mechanical vibration to dynamic system (Liu & Li, 2020), today's experiment will provide us with tremendous help regarding to this usefulness and significance.

Table 1: Data Table for Damping Condition Calculations.

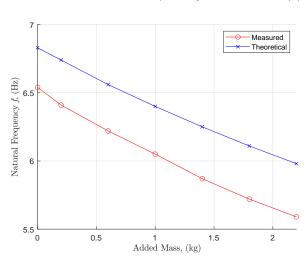
Case	Natural Frequency (Hz)	Damping Ratio	Magnification Ratio
Undamped	6.43	0.002	250
Fully Shut	6.43	0.029	17.2

Table 2: Data Table for Damping Condition Calculations in Simply Supported Case.

Case	Natural Frequency (Hz)	Damping Ratio	Magnification Ratio
Undamped	15.00	0.0035	143
Fully Shut	14.88	0.041	12.2

Equation for the relationship between the damped natural frequency and the damping ratio:

$$\omega_d = \omega_n \sqrt{1 - \zeta^2} \tag{1}$$



10
9
8
7
7
6
0.7
0.8
0.9
1
1.1
1.2
1.3
1.4
Frequency Ratio, $\frac{\omega}{\omega_n}$

Figure 2: The Relationship between Amplitude and Frequency Ratio during excitation damping.

Figure 1: Frequency vs Added Mass for Both Theoretical and Measured Values.

Equation for magnification factor:

$$\beta = \frac{1}{\sqrt{(1 - r^2)^2 + (2\zeta r)^2}} \tag{2}$$

Equation for speed ratio:

$$r = \frac{\omega}{\omega_n} \tag{3}$$

Equation for magnification factor in this experiment:

$$\beta = \frac{X_{r=1.0}}{X_{st}} \approx \frac{X_{r=1.0}}{X_{r=0.75,fully\ shut}}$$
 (4)

Equation for magnification factor at natural frequency:

$$\beta = \lim_{r \to 1} \frac{1}{\sqrt{(1 - r^2)^2 + (2\zeta r)^2}} = \frac{1}{2\zeta}$$
 (5)

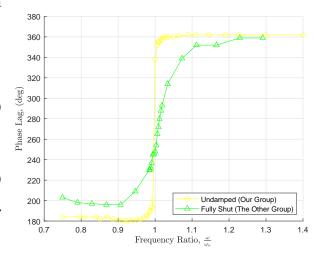


Figure 3: The Relationship between Phase Lag and Frequency Ratio during excitation damping.

Equation for theoretical frequency for can- Equation for effective mass (Rayleigh): tilever beam:

$$\omega = \sqrt{\frac{kl_{spring}^2}{I_A}} \tag{6}$$

in which

$$I_{A} = \left[\frac{1}{3}m_{beam}l_{beam}^{2}\right] + \left[\left(\frac{m_{spring}}{3} + m_{fixing}\right)l_{spring}^{2}\right] + \left[m_{exciter}l_{exciter}^{2}\right]$$
(7)

Equation for amplitude for vibration of beam:

$$X = \left| \frac{\delta_{st}}{1 - r^2} \right| \tag{8}$$

Equation for amplitude for vibration of beam in damped and forced vibration case:

$$X = \frac{1}{\sqrt{(1-r^2)^2 + (2\zeta r)^2}}$$

$$= \frac{1}{\sqrt{(4\zeta^2 + 1)r^4 - 2r^2 + 1}}$$
(9)

in which X reaches minimum when

$$r = \frac{1}{\sqrt{4\zeta^2 + 1}}\tag{10}$$

Equation for phase lag for vibration of beam:

$$\tan \phi = \frac{2\zeta r}{1 - r^2} \tag{11}$$

Added Mass (kg)	ed Mass (kg) Total Exciter Mass (kg) Effective Mass (kg)		Natural Frequency, f (Hz)		$1/f^2$
ridaca ividss (kg)	Total Exerci Wass (Rg)	Effective ivides (kg)	Measured	Theoretical	1/ J
0	4.2	5.0	15.53	15.50	0.00415
0.2	4.4	5.2	15.35	15.20	0.00424
0.6	4.8	5.6	14.78	14.64	0.00458
1.0	5.2	6.0	14.24	14.15	0.00493
1.4	5.6	6.4	13.77	13.70	0.00527
1.8	6.0	6.8	13.38	13.29	0.00559
2.2	6.4	7.2	13.01	12.91	0.00591

$$m_{effective} = m_{exciter} + \frac{17}{35} m_{beam}$$
 (12)

Equation for improved natural frequency:

$$f = \frac{1}{2\pi} \sqrt{\frac{6EI_{beam}}{m_{effective}l_3^3}} \tag{13}$$

Equation for forced vibration:

$$I_A \ddot{\theta} + kx_S l_{spring} + c l_{damper} \dot{x}_d = l_{exciter} Q \sin \Omega t$$
(14)

$$\ddot{\theta} + 2\gamma\dot{\theta} + \omega^2\theta = \frac{l_{exciter}Q}{I_A}Q\sin\Omega t \qquad (15)$$

Equation for damping ratio calculation:

$$\zeta = \frac{c}{2\sqrt{m_{eff}k_{eff}}}\tag{16}$$

Equation for stiffness of simply supported beam:

$$k_{beam} = \frac{48EI_{beam}}{l_{beam}^3} \tag{17}$$

(18)

Equation for period and frequency of simply supported beam:

 $T = \frac{1}{f} = 2\pi \sqrt{\frac{m_{exciter} l_{beam}^3}{48EL_{beam}}}$

Table 4: Recorded Data for Amplitude and Phase Angle with respect to Speed Ratio

Speed	Speed	X_{pp}	Phase
Ratio	(Hz)		Angle
0.750	4.82	0.2804	184
0.795	5.11	0.2804	184
0.809	5.20	0.2804	184
0.840	5.40	0.2804	184
0.851	5.47	0.2804	182
0.870	5.60	0.3488	184
0.885	5.69	0.3774	182
0.899	5.78	0.4415	181
0.910	5.85	0.4878	181
0.921	5.92	0.5696	180
0.930	5.98	0.6622	181
0.940	6.04	0.7682	181
0.950	6.11	0.9366	181
0.960	6.17	1.166	182
0.970	6.24	1.576	183
0.976	6.28	1.987	185
0.980	6.30	2.395	185
0.985	6.34	3.422	189
0.988	6.36	4.393	190
0.991	6.37	5.828	193
0.992	6.38	6.865	194
0.994	6.39	10.51	200
1.000	6.43	18.55	338
1.006	6.47	7.594	352
1.008	6.48	6.049	355
1.009	6.49	6.049	355
1.011	6.50	6.669	356
1.015	6.53	3.665	357
1.020	6.56	2.803	358
1.023	6.58	2.417	360
1.030	6.62	1.943	360
1.037	6.67	1.559	360
1.050	6.75	1.170	360
1.065	6.85	0.9360	361
1.089	7.00	0.6998	362
1.110	7.14	0.5718	362
1.146	7.37	0.4614	362
1.191	7.67	0.3731	362
1.227	7.89	0.3245	362
1.299	8.35	0.2605	362
1.400	9.00	0.2274	362

References

- Çalım, F. F. (2009). Free and forced vibrations of non-uniform composite beams. *Composite Structures*, 88(3), 413–423.
- Liu, C.-S. & Li, B. (2020). Forced and free vibrations of composite beams solved by an energetic boundary functions collocation method. *Mathematics and Computers in Simulation*, 177, 152–168.
- So, R., Zhou, Y., & Liu, M. (2000). Free vibrations of an elastic cylinder in a cross flow and their effects on the near wake. *Experiments in Fluids*, 29(2), 130–144.
- Sorokin, S. & Ershova, O. (2003). Forced and free vibrations of rectangular sandwich plates with parametric stiffness modulation. *Journal of sound and vibration*, 259(1), 119–143.

Appendices

1 Code for plot in the lab

Input Matlab source for plotting three graphs in this lab:

```
1 clc; clf; clear all;
2 % Clear all the things that are remained in the before command.
3 %% Problem 3.1.1
  hold on;
   % Begin to plot the figure in one graph.
   M = [0 \ 0.2:0.4:2.2]; % Added mass.
   fm = [6.54 \ 6.41 \ 6.22 \ 6.05 \ 5.87 \ 5.72 \ 5.59]; % Measured natural frequency.
   ft = [6.83 \ 6.74 \ 6.56 \ 6.40 \ 6.25 \ 6.11 \ 5.98]; % Theoretical natural frequency.
   plot(M, fm, 'ro-'); % Plot added mass with measured natural frequency.
   plot(M, ft, 'bx-'); % Plot added mass with theoretical natural frequency.
10
11
   grid on;
                % Open the grid on the graph.
12 xlim([0 2.2]); % Set the limit of x axis to [0 2.2].
13 xlabel('Added Mass, (kg)','interpreter','latex');
14 % Set the x label to the Added Mass, (kg).
15 ylabel('Natural Frequency\it \ f\rm, (Hz)','interpreter','latex');
16 % Set the y label to the Natural Frequency\it \ f\rm, (Hz).
17 legend('Measured','Theoretical');
18 % Set the Legend to 'Measured', 'Theoretical'.
19 hold off;
20 % Stop continuing plot the figure.
21 %% Problem 3.1.4.1
22 figure; %Open another figure.
23 hold on;
24 % Begin to plot the figure in one graph.
25 \text{ sr1} = [
26
        0.75
27
        0.795
28
        0.809
29
        0.84
30
        0.851
31
        0.87
32
        0.885
33
        0.899
34
        0.91
35
        0.921
36
        0.93
37
        0.94
38
        0.95
39
        0.96
40
        0.97
41
        0.976
42
        0.98
43
        0.985
44
        0.988
45
        0.991
46
        0.992
47
        0.994
48
49
        1.006
50
        1.008
51
        1.009
52
       1.011
53
       1.015
54
        1.02
55
        1.023
56
       1.03
57
        1.037
```

```
58
         1.05
 59
         1.065
 60
         1.089
 61
         1.11
 62
         1.146
 63
         1.191
 64
         1.227
 65
         1.299
 66
         1.4
 67
         ]; % Speed ratio for undamped.
     amp1 = [
 68
 69
         0.2804
70
         0.2804
 71
         0.2804
 72
         0.2804
 73
         0.2804
 74
         0.3488
75
         0.3774
 76
         0.4415
 77
         0.4878
 78
         0.5696
 79
         0.6622
 80
         0.7682
 81
         0.9366
 82
         1.166
 83
         1.576
 84
         1.987
 85
         2.395
86
         3.422
87
         4.393
88
         5.828
 89
         6.865
 90
         10.51
 91
         18.55
 92
         7.594
 93
         6.049
 94
         5.408
 95
         4.669
 96
         3.665
 97
         2.803
 98
         2.417
99
         1.943
100
         1.559
101
         1.17
102
         0.936
103
         0.6998
104
         0.5718
105
         0.4614
106
         0.3731
107
         0.3245
108
         0.2605
109
         0.2274
110 ]/2; % Amplitude for undamped.
111 plot(sr1, amp1, 'mo-'); % Plot the speed ratio with amplitude for undamped.
112 \text{ sr2} = [
113
         0.75
114
         0.79
115
         0.829
116
         0.868
         0.907
117
118
         0.946
119
         0.984
120
         0.986
```

```
121
       0.988
122
        0.991
123
        0.994
124
        0.997
125
126
        1.003
127
        1.006
128
        1.009
129
        1.012
130
        1.016
131
        1.019
132
        1.034
133
       1.073
134
       1.112
135
        1.166
136
        1.229
137
        1.291
        ]; % Speed ratio for fully shut.
138
139
   amp2 = [
140
       0.1368
141
        0.17
142
        0.212
143
        0.2826
144
        0.4018
145
        0.6534
146
        1.316
147
        1.413
148
        1.446
149
        1.546
150
        1.634
151
       1.755
152
       1.833
153
       1.92
154
       1.964
155
       2.019
156
       2.019
157
       1.975
158
       1.943
159
       1.557
160
       0.8784
161
        0.6314
        0.4414
162
163
        0.331
164
        0.2694
165
        ]/2;
               % Amplitude ratio for fully shut.
166 plot(sr2, amp2, 'c*-'); % Plot the speed ratio with amplitude for fully shut.
167 grid on; % Open the grid on the graph.
168 \text{ xlim}([0.7 \ 1.45]); % Set the limit of x axis to [0.7 \ 1.45].
170 % Set the x label to the Frequency Ratio, \frac{\infty}{\infty} {\omega_n} 
171 ylabel('Amplitude, (cm)','interpreter','latex');
172 % Set the y label to the Amplitude, (cm).
173 legend('Undamped (Our Group)', 'Fully Shut (The Other Group)');
174 % Set the Legend to 'Undamped (Our Group)', 'Fully Shut (The Other Group)'.
175 hold off;
176 % Stop continuing plot the figure.
177 %% Problem 3.1.4.1
178 figure; %Open another figure.
179 hold on;
180 % Begin to plot the figure in one graph.
181 pha1 = [
182
        184
183
        184
```

```
184
         184
185
         184
186
         182
187
         184
188
         182
189
         181
190
         181
191
         180
192
         181
193
         181
194
         181
195
         182
196
         183
197
         185
198
         185
199
         189
200
         190
201
         193
202
         194
203
         200
204
         248
205
         262
206
         265
207
         265
208
         266
209
         267
210
         268
211
         270
212
         270
213
         270
         270
214
215
         271
216
         272
217
         272
218
         272
219
         272
220
         272
221
         272
222
         272
223
         ]; % Phase lag for undamped.
224 plot(srl, phal, 'yo-'); % Plot the speed ratio with phase lag for undamped.
225 pha2 = [
226
        203
227
         198
228
         197
229
         196
230
         196
231
         209
232
         230
233
         230
234
         232
235
         237
         245
236
237
         246
238
         248
239
         254
240
         265
241
         272
242
         280
243
         288
244
         293
245
         314
246
         339
```

```
247
        352
248
        352
249
        359
250
        359
251
        ]; % Phase lag for fully shut.
252 plot(sr2, pha2, 'g^-'); % Plot the speed ratio with phase lag for fully shut.
253 grid on; % Open the grid on the graph.
254 % xlim([0.7 1.45]);
255 xlabel('Frequency Ratio, <math>\frac{\ode{n}}{\ode{n}}{\ode{n}}^{\circ},'interpreter','latex');
256 % Set the x label to the Frequency Ratio, \frac{\infty}{\infty} {\omega_n}.
257 ylabel('Phase Lag, ($\deg$)','interpreter','latex');
258 % Set the y label to the Phase Lag, (\del{lag}).
259 legend('Undamped (Our Group)', 'Fully Shut (The Other Group)', 'Location', 'SouthEast'
        → );
260 % Set the Legend to 'Undamped (Our Group)', 'Fully Shut (The Other Group)'.
261 hold off;
262 % Stop continuing plot the figure.
```