## Cahn-Hilliard-Navier-Stokes model

Governing equations

$$\begin{split} \hat{\varrho} \left( \mathbf{u}_t + \mathbf{u} \cdot \boldsymbol{\nabla} \mathbf{u} \right) - \hat{\nu} \Delta \mathbf{u} + \boldsymbol{\nabla} p &= \hat{\varrho} \mathbf{g} - \varphi \boldsymbol{\nabla} \mu_{\mathrm{ch}}, \\ \boldsymbol{\nabla} \cdot \mathbf{u} &= 0, \\ \varphi_t + \mathbf{u} \cdot \boldsymbol{\nabla} \varphi &= \boldsymbol{\nabla} \cdot \left( M \boldsymbol{\nabla} \mu_{\mathrm{ch}} \right), \\ \mu_{\mathrm{ch}} &= -\varepsilon \sigma \Delta \varphi + \frac{\sigma}{\varepsilon} \boldsymbol{\Psi}'(\varphi). \end{split}$$

with zero boundary conditions for velocity  $\mathbf{u}\big|_{\partial\Omega} = 0$ , zero boundary flux of phase  $\nabla \varphi \cdot \mathbf{n}\big|_{\partial\Omega} = 0$  and also zero boundary flux of the chemical potential  $M(\varphi)\nabla \mu_{\mathrm{ch}} \cdot \mathbf{n}\big|_{\partial\Omega} = 0$ . Initial conditions were predescribes as no velocity  $\mathbf{u}(t=0) = \mathbf{0}$  and circular bubble in the bottom half of the rectangle containing a different phase  $\varphi(t=0) = \varphi_0$ .

The degenerate mobility is considered to take the simple form  $M(\varphi) = M_0 (1 - \varphi)^2 \varphi^2$ , and the potential is of the same form, i.e.  $\Psi(\varphi) = (1 - \varphi)^2 \varphi^2$ .

We considered the density, viscosity and other quantities depending on the phase to take the form

$$\hat{\varrho}(\varphi) = \begin{cases} \hat{\varrho}_1 &, & \varphi < 0 \\ \hat{\varrho}_1 + (\hat{\varrho}_2 - \hat{\varrho}_1)\varphi &, & 0 \le \varphi \le 1, \\ \hat{\varrho}_2 &, & \varphi > 1. \end{cases}$$

Boundary conditions were chosen such that most of the terms in the weak formulation zero-out and we are left with this simplistic version

$$\int_{\Omega} \hat{\varrho} \mathbf{u}_{t} \cdot \mathbf{v} + \int_{\Omega} \hat{\varrho} \left( \mathbf{u} \cdot \nabla \mathbf{u} \right) \cdot \mathbf{v} + \hat{\nu} \nabla \mathbf{u} : \nabla \mathbf{v} - \int_{\Omega} \hat{\varrho} \mathbf{g} \cdot \mathbf{v} + \varphi \mu_{\mathrm{ch}} \nabla \cdot \mathbf{v} + \int_{\Omega} q \nabla \cdot \mathbf{u} - p \nabla \cdot \mathbf{v} \\
+ \int_{\Omega} \varphi_{t} \psi + \int_{\Omega} \left( \mathbf{u} \cdot \nabla \varphi \right) \psi + \int_{\Omega} M \nabla \mu_{\mathrm{ch}} \cdot \nabla \psi + \int_{\Omega} \mu_{\mathrm{ch}} \nu_{\mathrm{ch}} - \varepsilon \sigma \nabla \varphi \cdot \nabla \nu_{\mathrm{ch}} - \frac{\sigma}{\varepsilon} \Psi'(\varphi) \nu_{\mathrm{ch}} = 0$$

For stability reasons, this systems needs to be computed on a mesh with high detail and some backward method, like the backward Euler. The resulting functional renders

$$\begin{split} \int_{\Omega} \frac{\varphi_{i+1} - \varphi_{i}}{\mathrm{d}t} \psi + \int_{\Omega} \left(\theta \hat{\varrho}(\varphi_{i+1}) + (1 - \theta) \hat{\varrho}(\varphi_{i})\right) \frac{\mathbf{u}_{i+1} - \mathbf{u}_{i}}{\mathrm{d}t} \cdot \mathbf{v} \\ &+ \theta \left(F_{\mathrm{phase}} + F_{\mathrm{Navier-Stokes}}\right) \left(\mathbf{u}_{i+1}, \ p_{i+1}, \ \varphi_{i+1}, \ \mu_{\mathrm{ch}\, i+1}\right) \\ &+ (1 - \theta) \left(F_{\mathrm{phase}} + F_{\mathrm{Navier-Stokes}}\right) \left(\mathbf{u}_{i}, \ p_{i}, \ \varphi_{i}, \ \mu_{\mathrm{ch}\, i}\right) \\ &+ \int_{\Omega} q \mathbf{\nabla} \cdot \mathbf{u} + \int_{\Omega} \mu_{\mathrm{ch}} \nu_{\mathrm{ch}} - \varepsilon \sigma \mathbf{\nabla} \varphi \cdot \mathbf{\nabla} \nu_{\mathrm{ch}} - \frac{\sigma}{\varepsilon} \Psi'(\varphi) \nu_{\mathrm{ch}} = 0 \end{split}$$