

WEDGE NOTES

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Abstract. Three main contributions: stable DG formulations for general vertex-mapped wedges, and energy stable DG schemes for acoustic-elastic coupling in the presence of arbitrary heterogeneous media.

1. Introduction.

2. Vertex-mapped wedges. Properties:

1. $\frac{\partial r}{\partial xyz} J, \frac{\partial t}{\partial xyz} J \in P^1(\Delta) \otimes P^1([-1, 1])$.

2. $\frac{\partial s}{\partial xyz} J \in P^0(\Delta) \otimes P^2([-1, 1])$.

3. On triangular faces, $\mathbf{n}J^f \in P^0(\Delta)$.

4. On quadrilateral faces, $\mathbf{n}J^f \in P^1(\Delta)$.

Since $\frac{\partial}{\partial r}, \frac{\partial}{\partial t} : P^N(\Delta) \otimes P^N([-1, 1]) \rightarrow P^{N-1}(\Delta) \otimes P^N([-1, 1])$, this implies nodal collocation can be used to apply geometric factors.

Curvilinear wedges treated using quadrature-based skew-symmetric formulation; only need to interpolate over triangles.

3. Nodal elements. Important note: GQ nodes in the extruded direction.

Metric identities are satisfied for low order mappings of wedges; i.e.

$$\widehat{\nabla} \cdot \begin{pmatrix} r_x J & r_y J & r_z J \\ s_x J & s_y J & s_z J \\ t_x J & t_y J & t_z J \end{pmatrix} = \widehat{\nabla} \cdot (J\mathbf{G}) = 0.$$

where \mathbf{G} is the matrix of geometric factors. Even with this, GLL is not energy-stable.

3.1. Lift matrices under WADG. Under WADG, the computation of the RHS uses only reference differentiation and lift matrices.

For triangular faces, the lift matrix consists of $(N + 1)$ diagonal sub-matrices. These diagonal entries are all identical.

For quadrilateral faces, the lift matrix consists of block diagonal columns.

4. Acoustic-elastic coupling. Elastic wave equation (velocity-stress formulation)

$$\begin{aligned} \mathbf{C}^{-1} \frac{\partial \boldsymbol{\sigma}}{\partial t} + \nabla \cdot \mathbf{v} &= 0 \\ \rho \frac{\partial \mathbf{v}}{\partial t} + \text{sym}(\nabla \boldsymbol{\sigma}) &= 0. \end{aligned}$$

Assuming $\mu = 0$, reduction to the acoustic wave equation with $p = \text{tr}(\mathbf{S})$ and wavespeed $c = \sqrt{\frac{\lambda}{\rho}}$.

$$\begin{aligned} \frac{1}{c^2} \frac{\partial p}{\partial t} &= \nabla \cdot \mathbf{u} \\ \rho \frac{\partial \mathbf{u}}{\partial t} &= \nabla p. \end{aligned}$$

For acoustic-elastic interfaces, continuity conditions are given as

$$\begin{aligned} \mathbf{S}\mathbf{n} &= p\mathbf{n} \\ \mathbf{v} \cdot \mathbf{n} &= \mathbf{u} \cdot \mathbf{n}. \end{aligned}$$

We define the numerical flux \mathbf{f}^* as the sum of a central flux term (involving the average of left and right hand sides at an interface) and a penalization term.

On the acoustic side of the interface, this flux is

$$\langle (\mathbf{A}_n^T \boldsymbol{\sigma} - p\mathbf{n}) - \tau_u ((\mathbf{v} - \mathbf{u}) \cdot \mathbf{n}) \mathbf{n}, \boldsymbol{\tau} \rangle_{L^2(\partial D^k)} + \langle (\mathbf{v} - \mathbf{u}) \cdot \mathbf{n} - \tau_p \mathbf{n} \cdot (\mathbf{A}_n^T \boldsymbol{\sigma} - p\mathbf{n}), v \rangle_{L^2(\partial D^k)}$$

On the elastic side of the interface, the flux is

$$\langle (p\mathbf{n} - \mathbf{A}_n^T \boldsymbol{\sigma}) - \tau_v ((\mathbf{u} - \mathbf{v}) \cdot \mathbf{n}) \mathbf{n}, \mathbf{w} \rangle_{L^2(\partial D^k)} + \langle \mathbf{A}_n (\mathbf{u} - \mathbf{v}) - \tau_\sigma \mathbf{A}_n (p\mathbf{n} - \mathbf{A}_n^T \boldsymbol{\sigma}), \mathbf{q} \rangle_{L^2(\partial D^k)}$$

After integrating by parts the v and \mathbf{w} equations, taking $(v, \boldsymbol{\tau}) = (p, \mathbf{u})$ and $(\mathbf{w}, \mathbf{q}) = (\mathbf{v}, \boldsymbol{\sigma})$ yields flux terms which cancel.

Acoustic side central flux terms

$$\langle \mathbf{u} \cdot \mathbf{n} + \mathbf{v} \cdot \mathbf{n}, p \rangle_{L^2(\partial D^k)} + \langle \mathbf{A}_n^T \boldsymbol{\sigma} - p\mathbf{n}, \mathbf{u} \rangle_{L^2(\partial D^k)}$$

Elastic central flux terms

$$\langle \mathbf{A}_n^T \boldsymbol{\sigma} + p\mathbf{n}, \mathbf{v} \rangle_{L^2(\partial D^k)} + \langle \mathbf{A}_n (\mathbf{u} - \mathbf{v}), \boldsymbol{\sigma} \rangle_{L^2(\partial D^k)}$$

The terms

$$\langle \mathbf{u} \cdot \mathbf{n}, p \rangle_{L^2(\partial D^k)}, \quad \langle \mathbf{A}_n^T \boldsymbol{\sigma}, \mathbf{v} \rangle_{L^2(\partial D^k)}$$

are cancelled locally over each element. The remaining terms cancel after summing contributions from both sides of the acoustic-elastic interface.

Dissipative penalization terms can then be added to provide a stabilization. For the acoustic domain,

$$\langle \mathbf{n}^T \mathbf{S} \mathbf{n} - p, p \rangle_{L^2(\partial D^k)}, \quad \langle \mathbf{v} - \mathbf{u}, \mathbf{u} \rangle_{L^2(\partial D^k)}$$

Note: penalization does not require exact integration. It shouldn't even require explicit quadrature...

5. Numerical experiments. Test wedges, tets, wedge-tet hybrid meshes with other waves for due diligence.

5.1. Stoneley wave. Test acoustic-elastic coupling.

6. Future work. Need to load $O(N^3)$ geofacs ($O(N^2)$ geofacs per wedge). Can reduce triangular costs using Bernstein-Bezier elements

- Volume kernel: reduce from $O(N^5)$ to $O(N^4)$ computational complexity ($O(N^2)$ per triangle, cheaper application of geofacs using Bernstein polynomial multiplication, $O(N^2)$ in extruded direction but this cost is near-negligible due to the use of fast shared memory).
- Surface kernel: still $O(N^4)$. 1D interpolation operators are the same cost whether BB or nodal. Lift matrix cost can be reduced, but doesn't change asymptotic cost.
- Update kernel: for vertex-mapped wedges, can reduce from $O(N^5)$ to $O(N^4)$: $J \in P^1(\triangle)$ over each element, so can be applied using polynomial multiplication and projection down.