

Entropy stable Gauss collocation discontinuous Galerkin methods

Jesse Chan

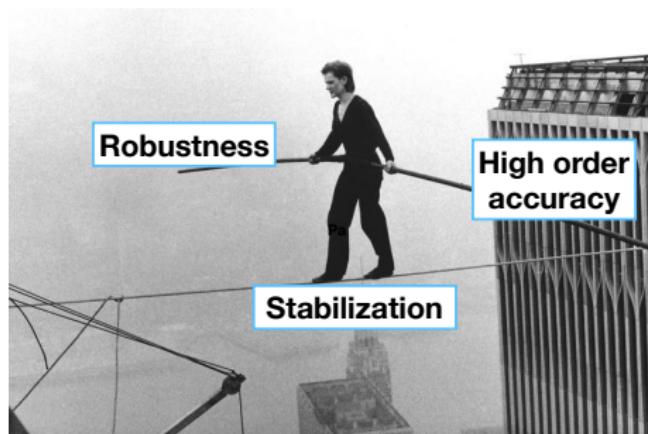
with DCDR Fernandez, Mark Carpenter (NASA Langley)

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SIAM CSE, Spokane, Washington
March 1, 2019

Discretely entropy stable schemes

Goal: address the stability and robustness of high order methods for nonlinear conservation laws.



- Aim for stability independently of artificial viscosity, limiters, discretization errors.
- *Mechanical* approach to stability based on algebraic properties of a discretization.
- Semi-discrete stability property for high order DG methods.

Finite volume methods: Tadmor, Chandrashekhar, Ray, Svard, Fjordholm, Mishra, LeFloch, Rohde, ...

High order tensor product elements: Fisher, Carpenter, Gassner, Winters, Kopriva, Hindenlang, Persson, Pazner, ...

High order general elements: Chen and Shu, Crean, Hicken, DCDR Fernandez, Zingg, ...

Entropy stability for nonlinear problems

- Generalizes energy stability to **nonlinear** systems of conservation laws (Burgers', shallow water, compressible Euler, MHD).

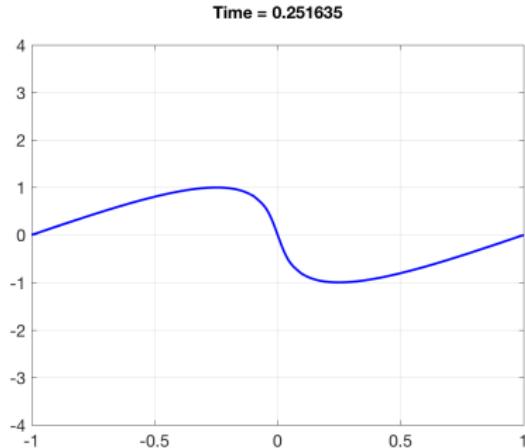
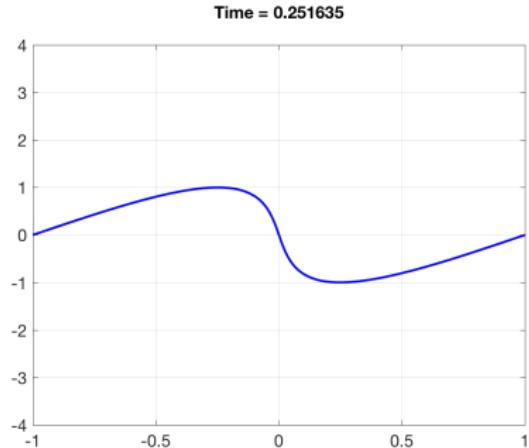
$$\frac{\partial \mathbf{u}}{\partial t} + \frac{\partial \mathbf{f}(\mathbf{u})}{\partial x} = 0.$$

- Continuous entropy inequality: given a convex **entropy** function $S(\mathbf{u})$ and “entropy potential” $\psi(\mathbf{u})$,

$$\int_{\Omega} \mathbf{v}^T \left(\frac{\partial \mathbf{u}}{\partial t} + \frac{\partial \mathbf{f}(\mathbf{u})}{\partial x} \right) = 0, \quad \boxed{\mathbf{v} = \frac{\partial S}{\partial \mathbf{u}}} \\ \Rightarrow \int_{\Omega} \frac{\partial S(\mathbf{u})}{\partial t} + \left(\mathbf{v}^T \mathbf{f}(\mathbf{u}) - \psi(\mathbf{u}) \right) \Big|_{-1}^1 \leq 0.$$

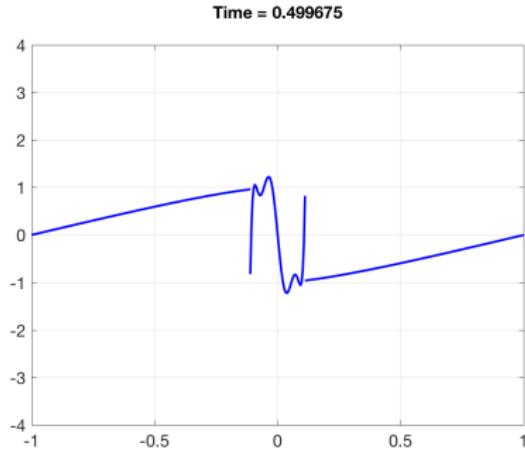
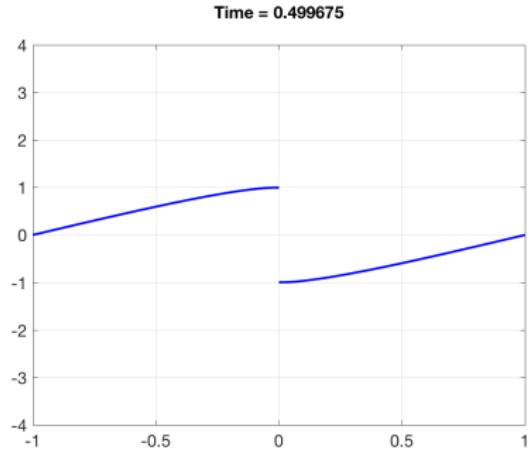
- Proof of entropy inequality relies on **chain rule**, integration by parts.

High order methods typically unstable for nonlinear PDEs



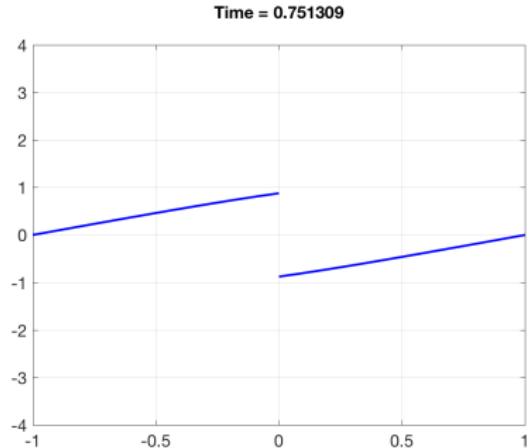
- High order methods tend to blow up for under-resolved solutions (shocks, turbulence), sensitive to discretization.
- Instability tied to loss of the **chain rule**.

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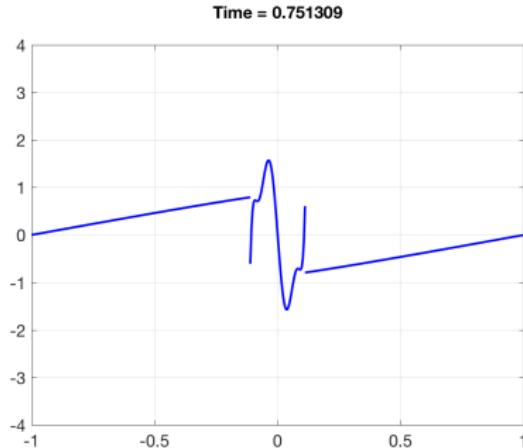


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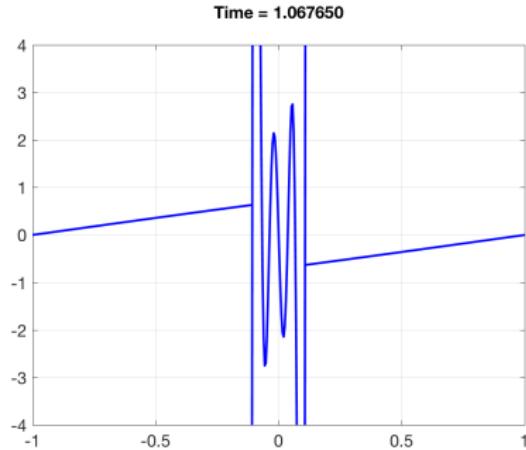
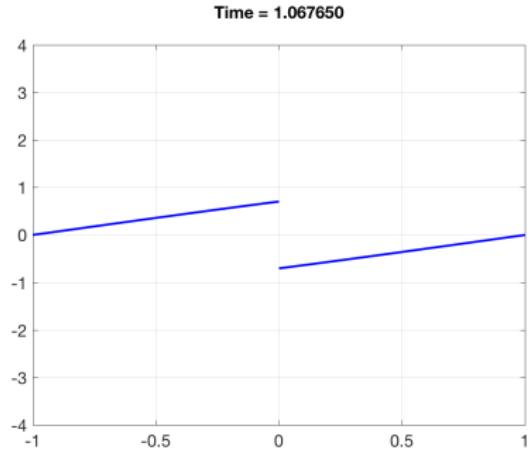
(a) Exact solution



(b) 8th order DG

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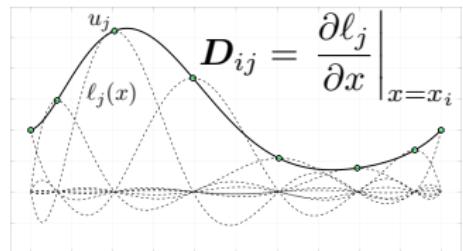
Talk outline

- 1 Summation-by-parts and high order DG
- 2 “Decoupled” block SBP operators
- 3 Numerical experiments
- 4 A mortar approach to non-conforming interfaces

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Nodal DG, summation-by-parts (SBP), flux differencing



$$\mathbf{B} = \begin{bmatrix} -1 & & & \\ & 0 & & \\ & & \ddots & \\ & & & 1 \end{bmatrix}$$

- Mimic **integration by parts** algebraically

$$\mathbf{Q} = \mathbf{B} - \mathbf{Q}^T, \quad \mathbf{Q} = \mathbf{M}\mathbf{D}, \quad \mathbf{M} \text{ diagonal mass matrix.}$$

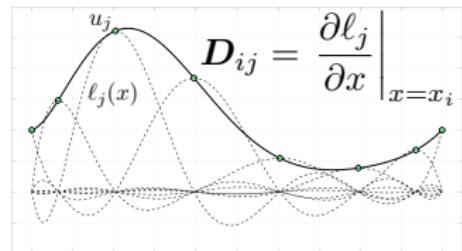
- Typical collocation scheme:

$$\mathbf{M} \frac{du}{dt} + \mathbf{Qf}(u) = 0 \implies \mathbf{M}_{ii} \frac{du_i}{dt} + \sum_j \mathbf{Q}_{ij} \mathbf{f}(u_j) = 0.$$

- Flux differencing: recover standard form if $\mathbf{f}_S(u_i, u_j) = \frac{1}{2}(u_i + u_j)$.

$$\mathbf{M}_{ii} \frac{du_i}{dt} + \sum_j \mathbf{Q}_{ij} 2\mathbf{f}_S(u_i, u_j) = 0 \implies \boxed{\mathbf{M} \frac{du}{dt} + 2(\mathbf{Q} \circ \mathbf{F}_S) \mathbf{1} = 0}.$$

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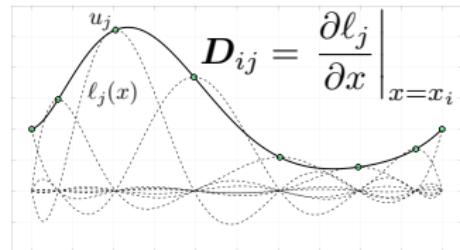
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Entropy stable schemes: a brief derivation

- Skew-symmetric formulation on each element:

$$\boldsymbol{M} \frac{d\boldsymbol{u}}{dt} + 2(\boldsymbol{Q} \circ \boldsymbol{F}_S) \boldsymbol{1} = 0$$

- Trick: use Tadmor's entropy conservative numerical flux for $\boldsymbol{f}_S, \boldsymbol{f}^*$

$$\boldsymbol{f}_S(\boldsymbol{u}, \boldsymbol{u}) = \boldsymbol{f}(\boldsymbol{u}), \quad (\text{consistency})$$

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- Proof of entropy conservation: multiply by \boldsymbol{v}^T

$$\boldsymbol{v}^T \boldsymbol{M} \frac{d\boldsymbol{u}}{dt} + \boldsymbol{v}^T \left((\boldsymbol{Q} - \boldsymbol{Q}^T) \circ \boldsymbol{F}_S \right) \boldsymbol{1} + \boldsymbol{v}^T \boldsymbol{B} \boldsymbol{f}^* = 0.$$

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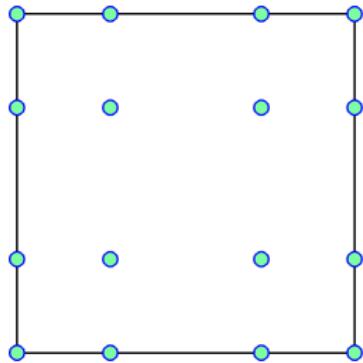
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Choices of collocation nodes



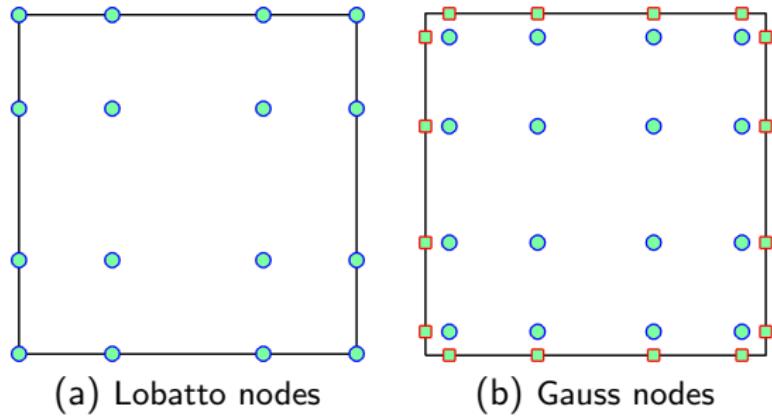
(a) Lobatto nodes

- Existing high order entropy stable schemes based on Lobatto nodes.
 - Gauss nodes: advantages in increasing quadrature accuracy?
 - Challenge in extending to Gauss nodes: inter-element coupling terms.

Fisher, Carpenter, et al. (2013, 2014), Gassner, Winters, Hindenlang, Kopriva (2016, 2018), ...

Parsani et al. (2016), *Entropy Stable Staggered Grid Discontinuous Spectral Collocation Methods*

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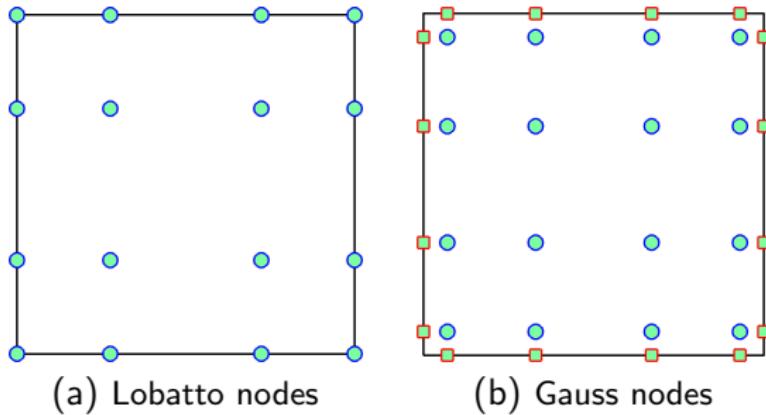


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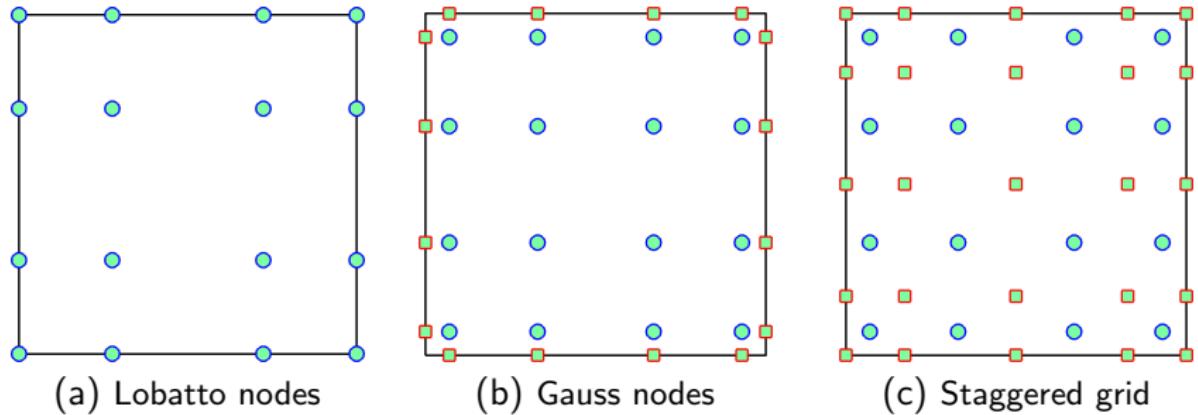


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Generalized SBP operators

- Differentiation matrix \mathbf{D} : evaluates derivatives at Gauss points.
- Define boundary interpolation matrix \mathbf{E}

$$\mathbf{E} = \begin{bmatrix} \ell_1(-1), & \dots & \ell_{N+1}(-1) \\ \ell_1(1), & \dots & \ell_{N+1}(1) \end{bmatrix}$$

- Generalized SBP property: let $\mathbf{Q} = \mathbf{MD}$,

$$\mathbf{Q} + \mathbf{Q}^T = \mathbf{E}^T \mathbf{B} \mathbf{E}, \quad \mathbf{B} = \begin{bmatrix} -1 & \\ & 1 \end{bmatrix}$$

$$\implies \int_{-1}^1 \frac{\partial I_N u}{\partial x} I_N v + \int_{-1}^1 I_N u \frac{\partial I_N v}{\partial x} = (I_N u)(I_N v)|_{-1}^1.$$

where I_N is the degree N polynomial interpolation operator

Entropy stability on multiple elements

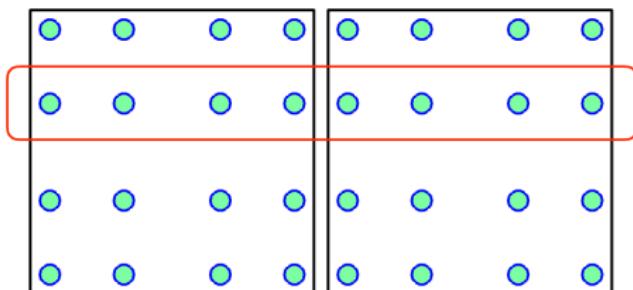


Figure: Coupling between Gauss nodes on neighboring elements.

- Re-deriving the entropy stable formulation with GSBP operators:

$$\boldsymbol{M} \frac{d\boldsymbol{u}}{dt} + 2(\boldsymbol{Q} \circ \boldsymbol{F}_S) \boldsymbol{1} = 0.$$

- The presence of the interpolation matrix \boldsymbol{E} increases inter-element coupling, complicates BC imposition.

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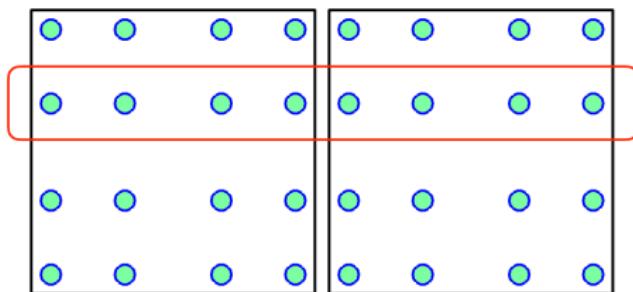


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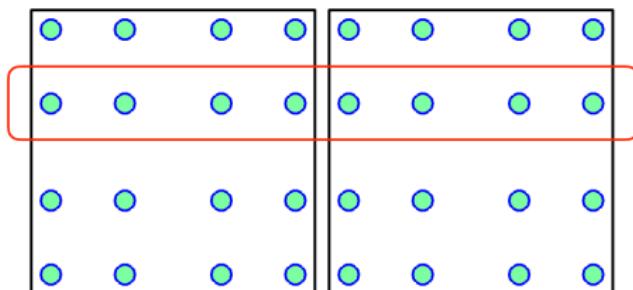


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- The presence of the interpolation matrix \boldsymbol{E} increases inter-element coupling, complicates BC imposition.

Talk outline

- 1 Summation-by-parts and high order DG
- 2 "Decoupled" block SBP operators
- 3 Numerical experiments
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A “decoupled” SBP operator

- Goal: SBP property without the boundary interpolation operator \mathbf{E}

$$\mathbf{Q}_N = \begin{bmatrix} \mathbf{Q} - \frac{1}{2}\mathbf{E}^T\mathbf{B}\mathbf{E} & \frac{1}{2}\mathbf{E}^T\mathbf{B} \\ -\frac{1}{2}\mathbf{B}\mathbf{E} & \frac{1}{2}\mathbf{B} \end{bmatrix},$$

- If $\mathbf{Q} + \mathbf{Q}^T = \mathbf{E}^T\mathbf{B}\mathbf{E}$, then \mathbf{Q}_N satisfies the SBP property

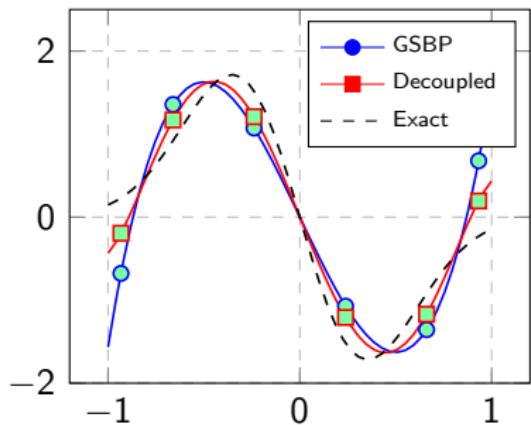
$$\boxed{\mathbf{Q}_N + \mathbf{Q}_N^T = \begin{bmatrix} \mathbf{0} \\ \mathbf{B} \end{bmatrix}} \sim \boxed{\int_{-1}^1 \frac{\partial I_N u}{\partial x} v + u \frac{\partial I_N v}{\partial x} = uv|_{-1}^1}.$$

- \mathbf{Q}_N approximates $f \frac{\partial g}{\partial x}$ by \mathbf{u} using data at $\mathbf{x} = [\mathbf{x}_{\text{vol}}, \mathbf{x}_{\text{face}}]$

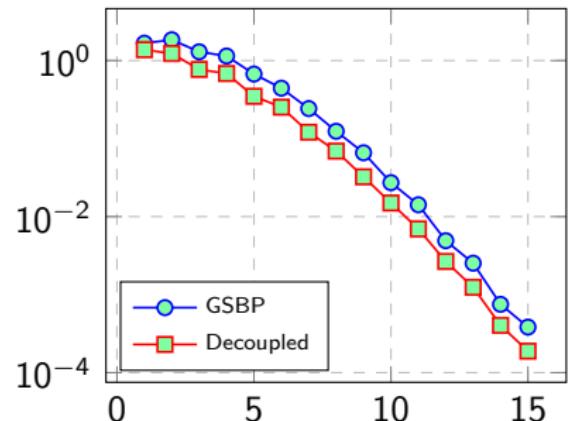
$$\mathbf{M}\mathbf{u} = \begin{bmatrix} \mathbf{I} \\ \mathbf{E} \end{bmatrix}^T \text{diag}(\mathbf{f}) \mathbf{Q}_N \mathbf{g}, \quad \mathbf{f}_i, \mathbf{g}_i = f(\mathbf{x}_i), g(\mathbf{x}_i).$$

- Reduces to traditional SBP operator under appropriate quadrature.

Decoupled SBP operators add boundary corrections



(a) Derivative approximations

(b) L^2 error w.r.t. degree N

- Equivalent to a variational problem for a polynomial $u(\mathbf{x}) \approx f \frac{\partial g}{\partial \mathbf{x}}$.

$$\int_{-1}^1 u(\mathbf{x})v(\mathbf{x}) = \int_{-1}^1 f \frac{\partial I_N g}{\partial \mathbf{x}} v + (g - I_N g) \frac{(fv + I_N(fv))}{2} \Big|_{-1}^1.$$

Entropy stable schemes using decoupled SBP operators

- Replace SBP operator with decoupled SBP operator

$$\mathbf{M} \frac{d\mathbf{u}}{dt} + \left((\mathbf{Q} - \mathbf{Q}^T) \circ \mathbf{F}_S \right) \mathbf{1} + \mathbf{B} \mathbf{f}^* = 0, \quad (\text{standard SBP})$$

- \mathbf{F}_S is the matrix of flux evaluations between solution values at *both* volume and face nodes

$$(\mathbf{F}_S)_{ij} = \mathbf{f}_S(\tilde{\mathbf{u}}_i, \tilde{\mathbf{u}}_j), \quad \tilde{\mathbf{u}} = \begin{bmatrix} \mathbf{u} \\ \mathbf{u}_f \end{bmatrix}$$

- Face values of \mathbf{u}_f : **interpolate entropy variables** (entropy projection). Resulting scheme is entropy stable.

Chan (2018). *On discretely entropy conservative and entropy stable discontinuous Galerkin methods*.

Parsani et al. (2016), *Entropy Stable Staggered Grid Discontinuous Spectral Collocation Methods*

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- \mathbf{F}_S is the matrix of flux evaluations between solution values at *both* volume and face nodes

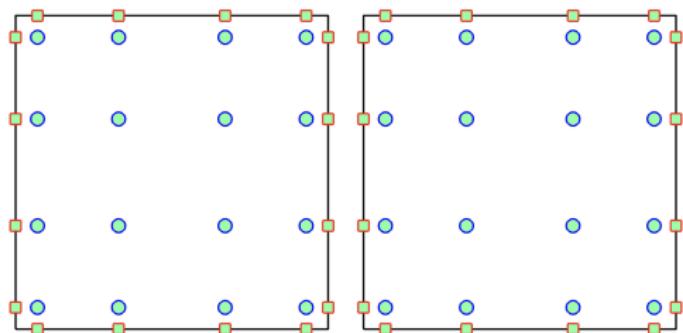
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Higher dimensions and curved meshes



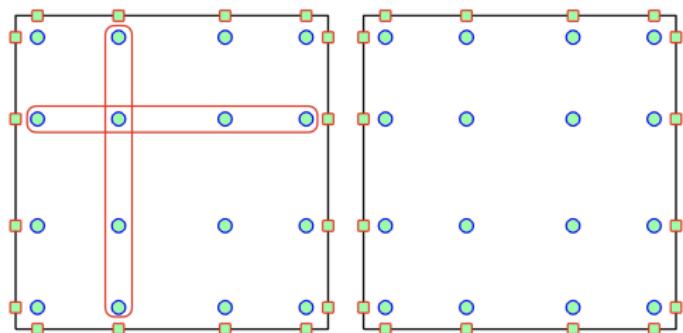
$$\left(\underbrace{\begin{bmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B}^T & \mathbf{C} \end{bmatrix}}_{\mathbf{Q}_N^i - (\mathbf{Q}_N^i)^T} \circ \underbrace{\begin{bmatrix} \mathbf{F}_S^{vv} & \mathbf{F}_S^{vf} \\ \mathbf{F}_S^{fv} & \mathbf{F}_S^{ff} \end{bmatrix}}_{\mathbf{F}_S} \right) \mathbf{1}$$

- Multi-dimensional operators defined using Kronecker products

$$\mathbf{Q}^1 = \mathbf{Q}_{1D} \otimes \mathbf{M}_{1D}, \quad \mathbf{E} = \begin{bmatrix} \mathbf{E}_{1D} \otimes \mathbf{I}_2 \\ \mathbf{I}_2 \otimes \mathbf{E}_{1D} \end{bmatrix}, \quad \mathbf{B}^1 = \begin{bmatrix} \mathbf{B}_{1D} \otimes \mathbf{I}_2 \\ \mathbf{0} \end{bmatrix}.$$

- Curved meshes: discretize geometric terms in split form, ensure GCL.
- Compute geometric terms at Lobatto points, interpolate to Gauss.

Higher dimensions and curved meshes



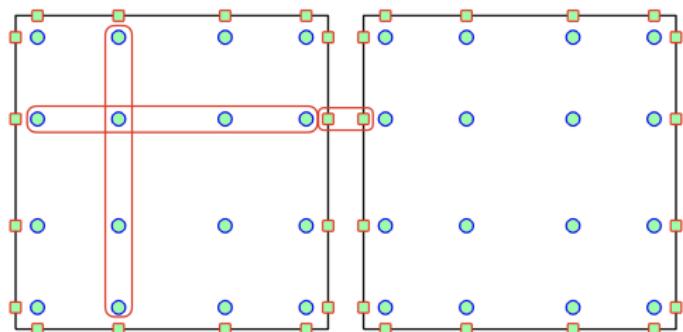
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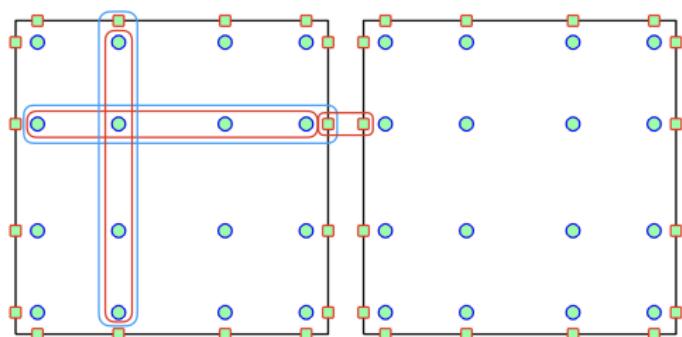
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Higher dimensions and curved meshes



$$\left(\begin{bmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B}^T & \mathbf{C} \end{bmatrix} \circ \begin{bmatrix} \mathbf{F}_S^{vv} & \mathbf{F}_S^{vf} \\ \mathbf{F}_S^{fv} & \mathbf{F}_S^{ff} \end{bmatrix} \right) \mathbf{1}$$

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Accurate quadrature can reduce error on curved meshes

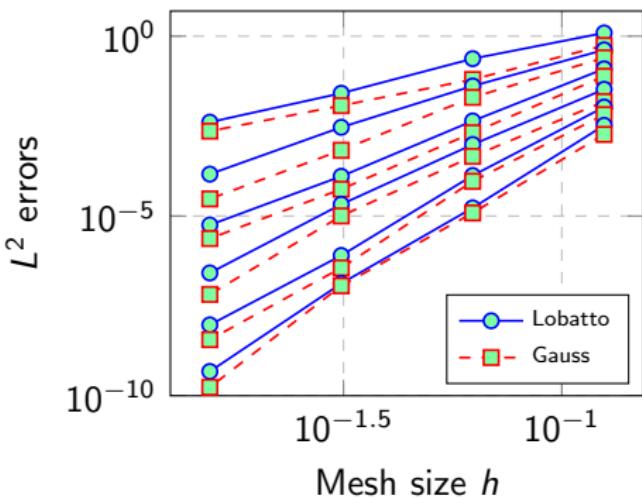
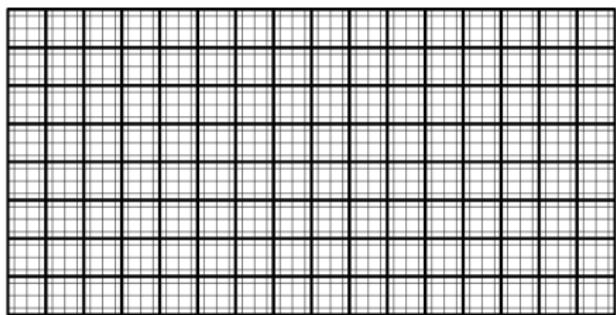


Figure: L^2 errors for the 2D isentropic vortex at time $T = 5$ for degree $N = 2, \dots, 7$ Lobatto and Gauss collocation schemes (similar behavior in 3D).

Accurate quadrature can reduce error on curved meshes

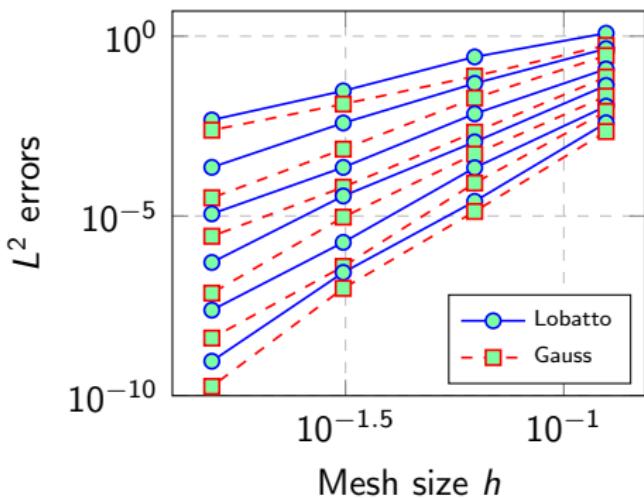
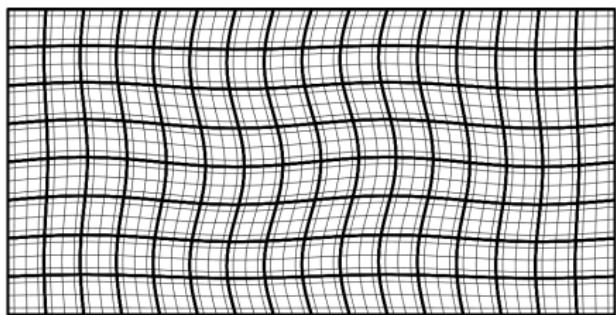


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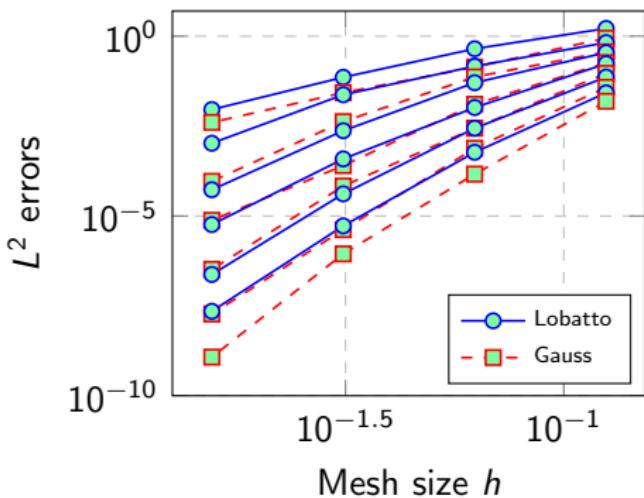
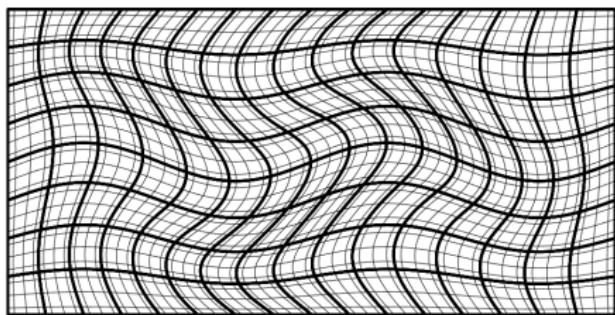


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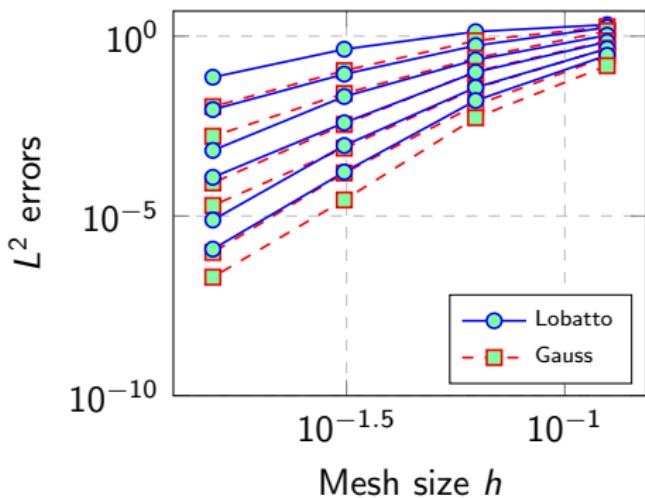
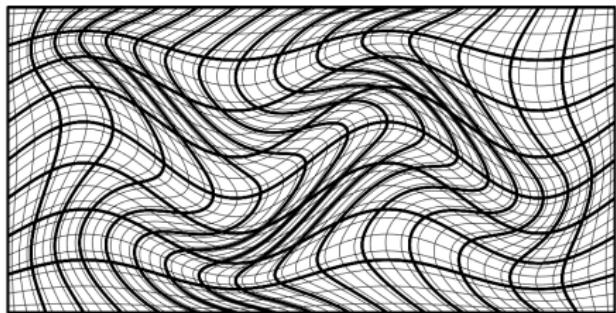


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3D inviscid Taylor-Green vortex

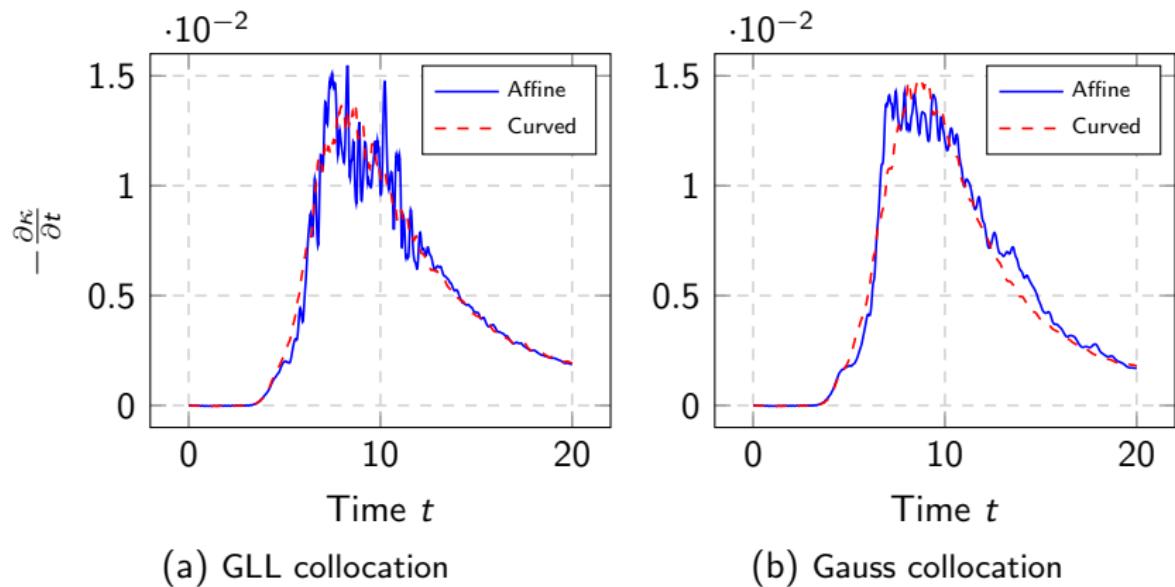


Figure: Kinetic energy dissipation rate for entropy stable GLL and Gauss collocation schemes with $N = 7$ and $h = \pi/8$.

Shock vortex interaction

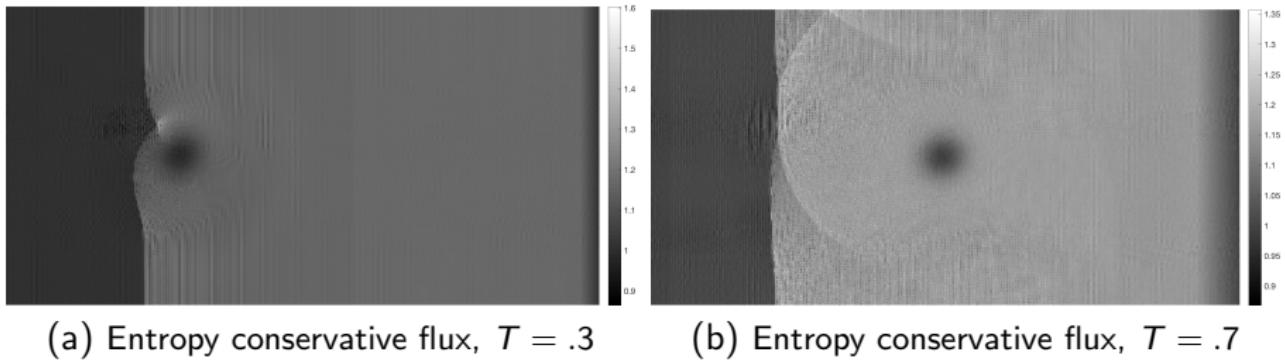


Figure: Shock vortex interaction problem using high order entropy stable Gauss collocation schemes with $N = 4, h = 1/100$.

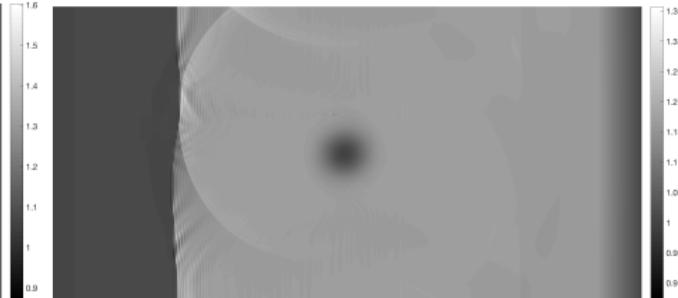
Jiang, Shu (1998). *Efficient Implementation of Weighted ENO Schemes*.

Winters, Derigs, Gassner, and Walch (2017). *A uniquely defined entropy stable matrix dissipation operator for high Mach number ideal MHD and compressible Euler simulations*.

Shock vortex interaction



(a) Lax-Friedrichs flux, $T = .3$



(b) Lax-Friedrichs flux, $T = .7$

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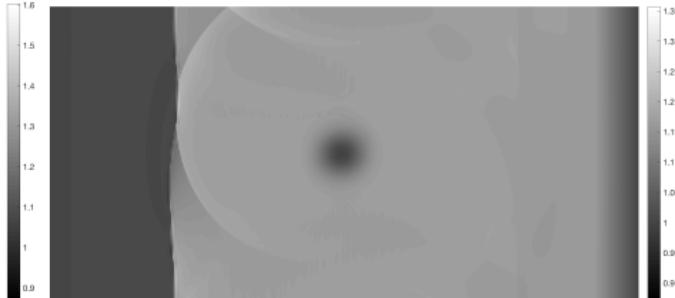
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Shock vortex interaction



(a) Matrix dissipation flux, $T = .3$



(b) Matrix dissipation flux, $T = .7$

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Shock vortex interaction

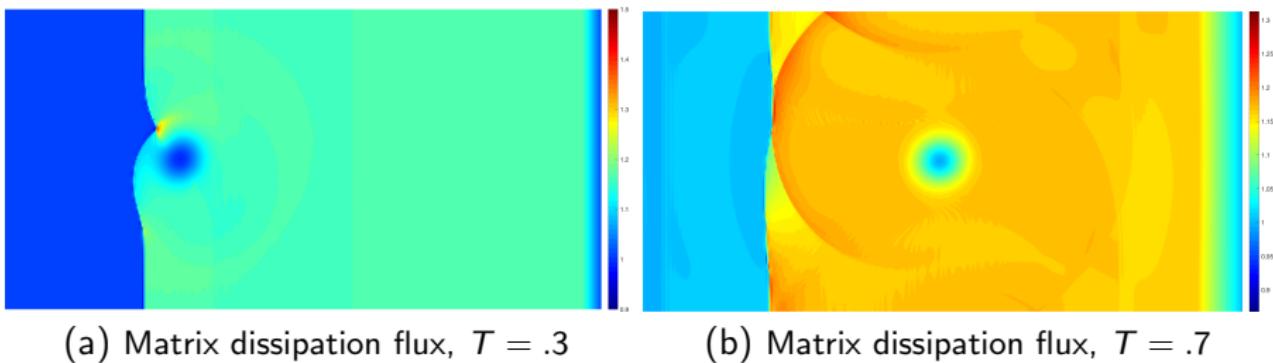
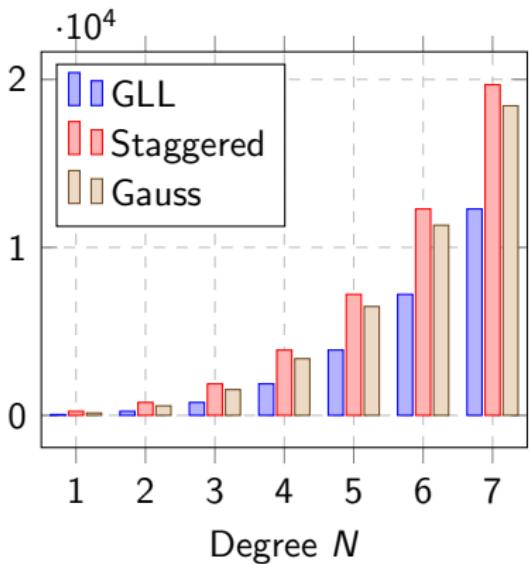


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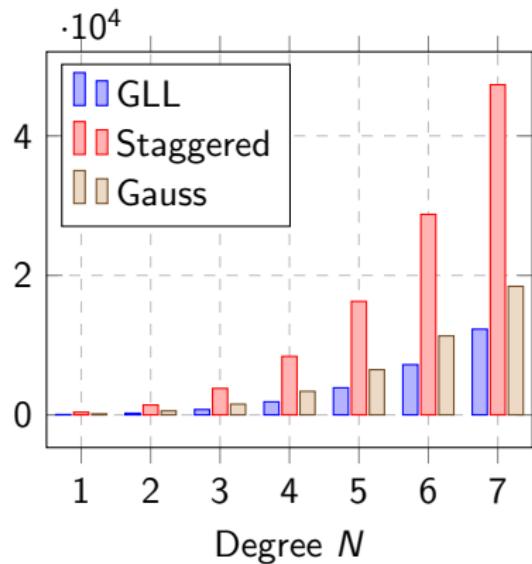
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Cost estimates: flux evaluations and matrix operations



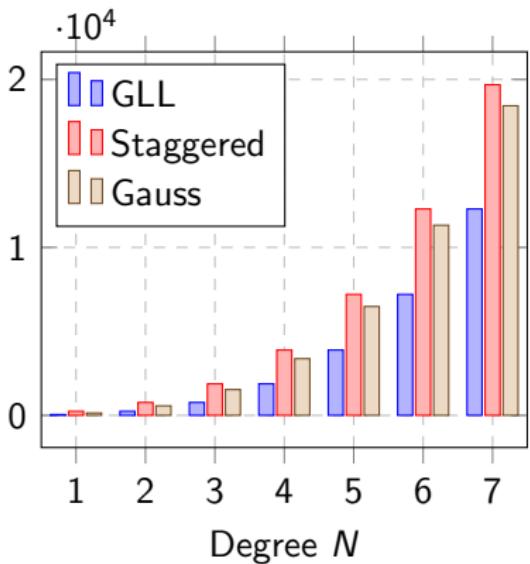
(a) Two-point flux evaluations



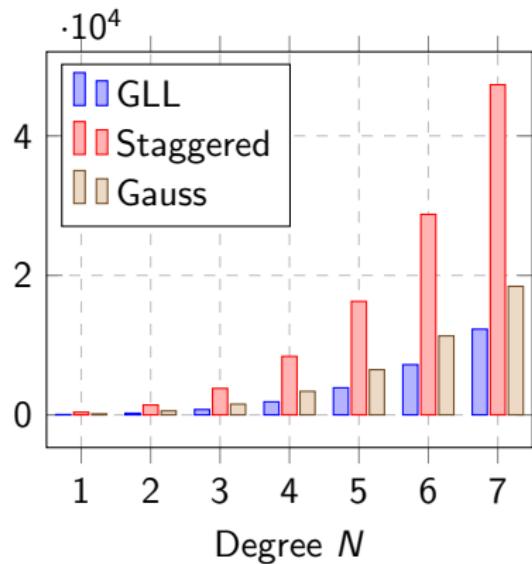
(b) Est. matrix op. FLOPS

Note: does not significantly reduce total flux evaluations.

Cost estimates: flux evaluations and matrix operations



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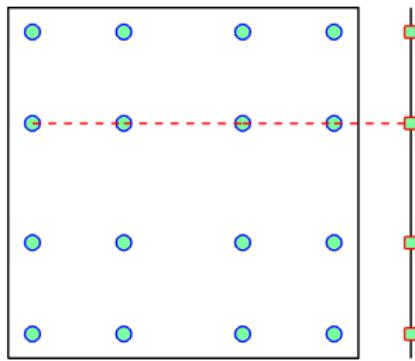
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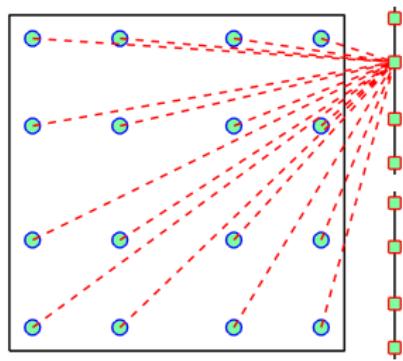
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Non-conforming interfaces (with DCDR Fernandez)



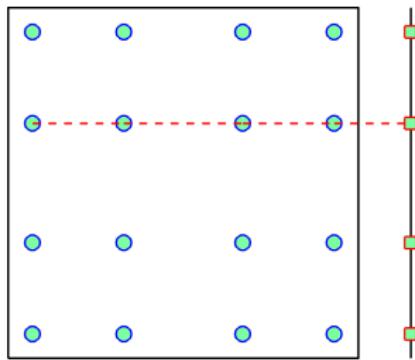
(a) Conforming surface quadrature nodes



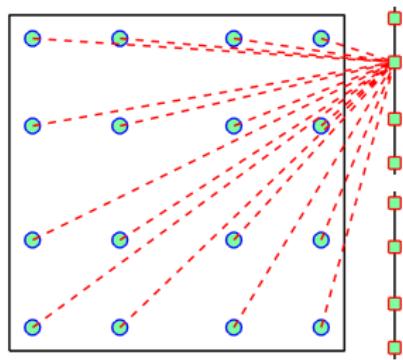
(b) Non-conforming surface nodes

- Volume/surface nodes interact through $f_S(\mathbf{u}_i, \mathbf{u}_j)$ and interpolation.
- Fix: weakly couple conforming+non-conforming faces using a mortar.

Non-conforming interfaces (with DCDR Fernandez)



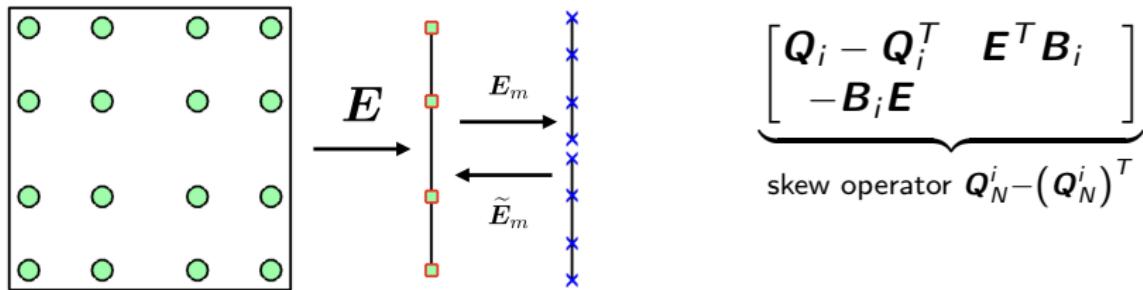
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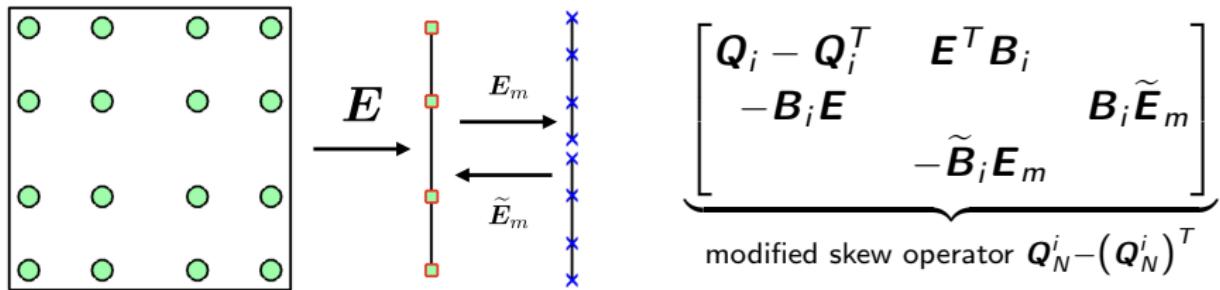
A decoupled SBP approach to mortars



- Define appropriate interpolation operators $\mathbf{E}_m, \tilde{\mathbf{E}}_m$ between conforming and non-conforming (mortar) nodes.
- Modify the skew-symmetric formulation as follows:

$$\mathbf{M} \frac{d\mathbf{u}}{dt} + \sum_{i=1}^d \begin{bmatrix} \mathbf{I} \\ \mathbf{E} \end{bmatrix}^T \left[\begin{array}{cc} \mathbf{Q}_i - \mathbf{Q}_i^T & \mathbf{E}^T \mathbf{B}_i \\ -\mathbf{B}_i \mathbf{E} & \end{array} \right] + \mathbf{E}^T \mathbf{B}_i \mathbf{f}_i^* = 0.$$

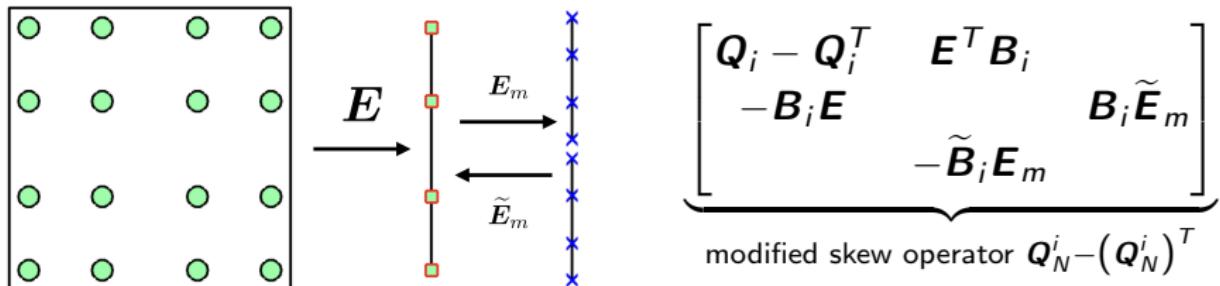
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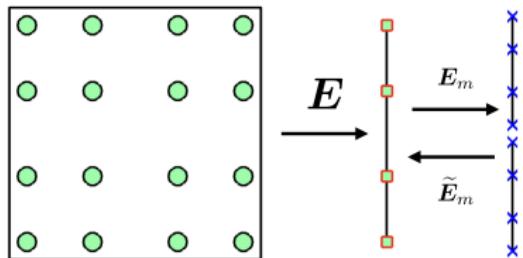
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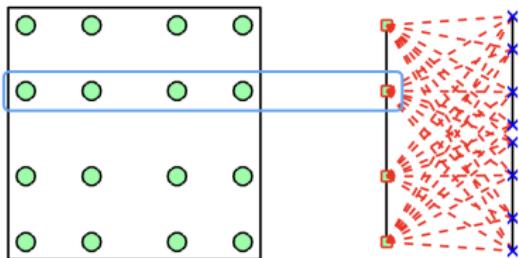
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Mortar implementation



(a) Mortar operators



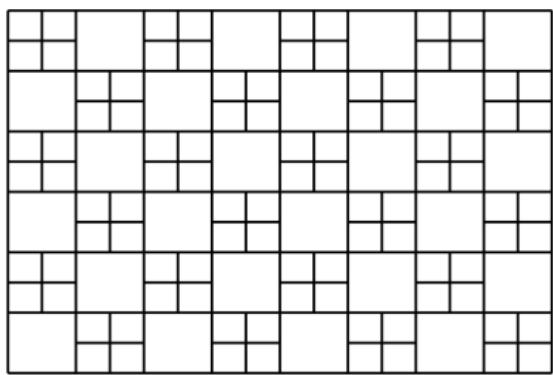
(b) Mortar coupling

$$\boldsymbol{M} \frac{d\boldsymbol{u}}{dt} + \sum_{i=1}^d \begin{bmatrix} \boldsymbol{I} \\ \boldsymbol{E} \end{bmatrix}^T \left(\begin{bmatrix} \boldsymbol{Q}_i - \boldsymbol{Q}_i^T & \boldsymbol{E}^T \boldsymbol{B}_i \\ -\boldsymbol{B}_i \boldsymbol{E} \end{bmatrix} \circ \boldsymbol{F}_S \right) \boldsymbol{1} + \boldsymbol{E}^T \boldsymbol{B}_i \tilde{\boldsymbol{f}}_i^* = 0$$

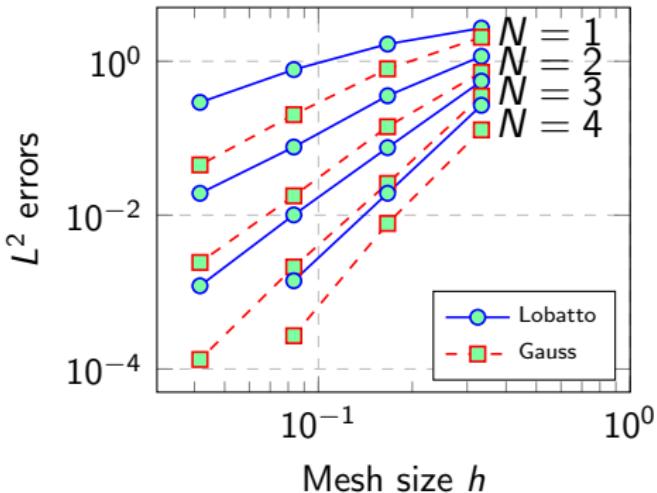
$$\tilde{\boldsymbol{f}}_i^* = \tilde{\boldsymbol{E}}_m \boldsymbol{f}_i^* + \left(\tilde{\boldsymbol{E}}_m \circ \boldsymbol{F}_S^{sm} \right) \boldsymbol{1} - \tilde{\boldsymbol{E}}_m (\boldsymbol{E}_m \circ \boldsymbol{F}_S^{ms}) \boldsymbol{1}$$

Can reformulate as an entropy stable correction to the numerical flux.

Numerical results: non-conforming meshes



(a) Coarse non-conforming mesh



(b) Sub-optimal rates if under-integrated

The skew-symmetric formulation guarantees entropy stability for both Lobatto and Gauss quadratures, but Gauss is more accurate.

Summary and future work

- Gauss collocation DG can improve accuracy on curved meshes.
- Can treat non-conforming interfaces using a mortar-based coupling.
- This work is supported by DMS-1719818 and DMS-1712639.

Thank you! Questions?



Chan (2019). *Skew-symmetric entropy stable modal discontinuous Galerkin formulations*.

Chan, Del Rey Fernandez, Carpenter (2018). *Efficient entropy stable Gauss collocation methods*.

Chan, Wilcox (2018). *On discretely entropy stable weight-adjusted DG methods: curvilinear meshes*.

Chan (2018). *On discretely entropy conservative and entropy stable discontinuous Galerkin methods*.

3D isentropic vortex

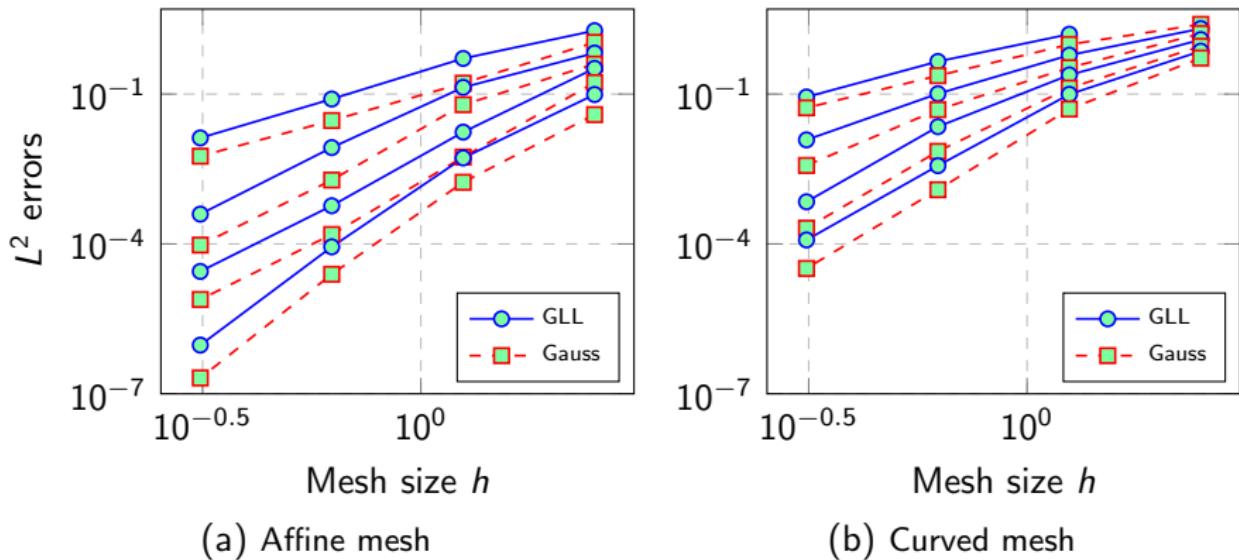
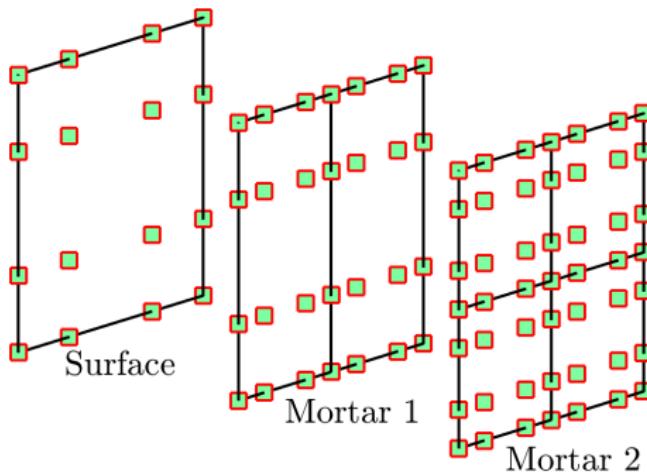


Figure: L^2 errors for the 3D isentropic vortex for $N = 2, \dots, 5$.

Entropy stable mortars in 3D



The mortar-based approach can reduce the number of flux evaluations for 3D non-conforming meshes.