

# Electronics of Radio

Notes on David Rutledge's book

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# Basic concepts

- Potential difference ( $V, \phi$ ):  $\phi = \int_a^r E \cdot ds$ , energy per charge,  $1V = 1 \text{ J/s}$
- Kirkoff 1:  $\sum_{loop} V_i = 0$  (Conservation of energy)
- Kirkoff node:  $\sum_{node} I_i = 0$  (Conservation of charge)
- $V(t) = V_p \cos(\omega t), \omega = 2\pi f, I(t) = I_p \cos(\omega t), \omega = 2\pi f$
- Instantaneous power:  $P(t) = V(t)I(t) = V_p I_p \cos^2(\omega t)$
- Average power:  $P_a = \int_0^{1/f} V(t)I(t)dt = V(t)I(t) = \int_0^{2\pi/\omega} V_p I_p \cos^2(\omega t)dt = \frac{V_p I_p}{2}$
- Band names:

Name	Frequency
VLF	3-30kHz
LW	20-300kHz
MW	300kHz-3MHz
HF	3MHz-30MHz
VHF	30-300MHz

Name	Frequency
UHF	300MHz-1GHz
uW	1-30GHz
milliW	30-300GHz
submilliF	>300GHz

# Signals

- Gain (G) expressed in decibels:  $G = 10 \log_{10}(P_{out}/P_{in})$
- Mixer:
  - $V(t) = \cos(\omega_1 t) \cos(\omega_2 t) = \frac{1}{2}[\cos(\omega_+ t) + \cos(\omega_- t)], \omega_+ = \omega_1 + \omega_2, \omega_- = \omega_1 - \omega_2$
- Modulation

Name	Equation
AM	$V(t) = a(t)\cos(\omega_c t)$
FM	$V(t) = V_c\cos((\omega_c + a(t))t)$
FSK	$V(t) = V_c\cos(\omega_1 t), \text{ if } 1$ $V(t) = V_c\cos(\omega_0 t), \text{ if } 0$
PSK	$V(t) = +V_p\cos(\omega t), \text{ if } 1$ $V(t) = -V_p\cos(\omega t), \text{ if } 0$

# Resistors, capacitors, inductors

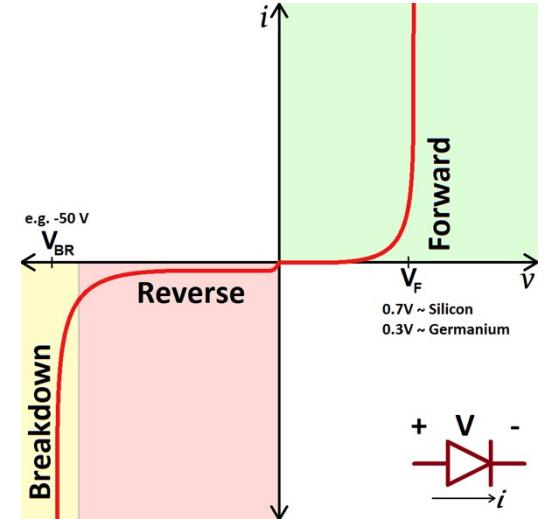
- Resistors
  - Analytic model:  $IR = V$
  - Energy dissipated:  $E = \int_{t_i}^{t_f} IV dt = \int_{t_i}^{t_f} I^2 R dt$
- Capacitors
  - Analytic model:  $CV = q, C \frac{dV}{dt} = i$
  - Capacitor Energy stored:  $E = \int_{t_i}^{t_f} CV \frac{dV}{dt} dt = \frac{1}{2} CV^2$
- Inductors
  - Analytic model:  $V = L \frac{di}{dt}$
  - Inductor Energy stored:  $E = \int_{t_i}^{t_f} IV dt = \int_{t_i}^{t_f} LI \frac{dI}{dt} dt = \frac{1}{2} LI^2$



Credit: Make Electronics

# Diodes, transformers

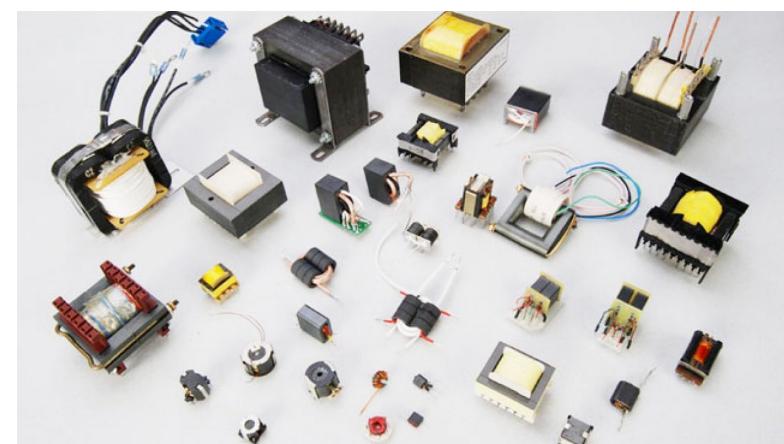
- Diodes
  - Devices that allow current to flow only in one direction
  - Silicon diodes, for example have, essentially infinite resistance if  $V_{ac} < 0$ , that is if the cathode is at a higher potential than the anode and very low resistance if  $V_{ac} > .7V$ .
  - The cathode is usually labelled with a band



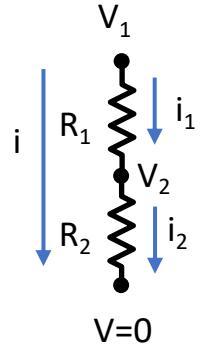
Credit: Make Electronics

- Transformers

- AC only:  $\frac{N_2}{N_1} = \frac{V_2}{V_1}$



# Simple circuit analysis with Kirchhoff

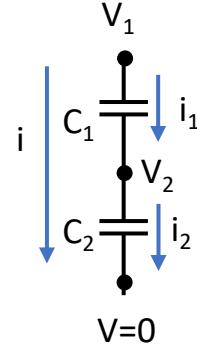


- $R_{eq}$  is the equivalent resistance, replacing the top left circuit with a single resistance.
- By Kirchhoff's node rule,  $i_1 = i_2 = i$ , so
- $\frac{V_1 - V_2}{R_1} = \frac{V_2}{R_2} = \frac{V_1}{R_{eq}}$  thus  $\frac{R_1}{R_{eq}} V_1 = V_1 - V_2$  and  $\frac{R_2}{R_{eq}} V_1 = V_2$ . Adding, we get  $\frac{R_1}{R_{eq}} V_1 + \frac{R_2}{R_{eq}} V_1 = V_1$ . Dividing by  $V_1$  and solving, we get  $R_1 + R_2 = R_{eq}$



- Again let  $R_{eq}$  is the equivalent resistance, replacing the bottom left circuit with a single resistance.
- By Kirchhoff's node rule,  $i_1 + i_2 = i$ , so
- $\frac{V_1}{R_1} + \frac{V_1}{R_2} = \frac{V_1}{R_{eq}}$ .
- Solving, we get.  $\frac{R_1 R_2}{R_1 + R_2} = R_{eq}$

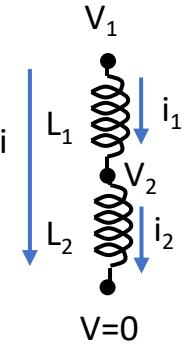
- $C_{eq}$  is the equivalent capacitance, replacing the top right circuit with a single capacitor.
- By Kirchhoff's node rule,  $i_1 = i_2 = i$ , so
- $C_1 \frac{d(V_1 - V_2)}{dt} = C_2 \frac{d(V_2)}{dt} = C_{eq} \frac{dV_1}{dt}$
- $\frac{C_{eq}}{C_1} \frac{d(V_1)}{dt} = \frac{d(V_1 - V_2)}{dt}$  and  $\frac{C_{eq}}{C_2} \frac{d(V_1)}{dt} = \frac{d(V_2)}{dt}$
- Adding and cancelling the  $\frac{d(V_1)}{dt}$ , we get
- $\frac{C_{eq}}{C_1} + \frac{C_{eq}}{C_2} = 1$  and solving, we get.  $\frac{C_1 C_2}{C_1 + C_2} = C_{eq}$



- $C_{eq}$  is the equivalent capacitance, replacing the bottom right circuit with a single capacitor.
- By Kirchhoff's node rule,  $i_1 + i_2 = i$
- $C_{eq} \frac{dV_1}{dt} = C_1 \frac{dV_1}{dt} + C_2 \frac{dV_1}{dt}$ , so
- $C_{eq} = C_1 + C_2$

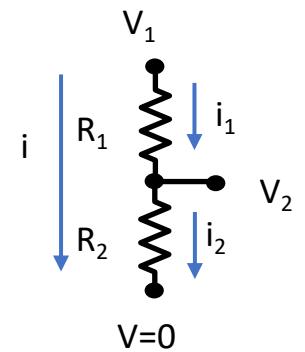


# Simple circuit analysis with Kirchhoff



- Let  $L_{eq}$  be the equivalent inductance, replacing the top left circuit with a single inductor.
- By Kirchhoff's node rule,  $i_1 = i_2 = i$ , so
- $L_{eq} \frac{di}{dt} = V_1$ ,  $L_1 \frac{di_1}{dt} = V_1 - V_2$ ,  $L_1 \frac{di_2}{dt} = V_2$
- $V_1 = L_{eq} \frac{di}{dt} = L_1 \frac{di}{dt} + L_2 \frac{di}{dt}$  and
- $L_{eq} = L_1 + L_2$

- Let  $L_{eq}$  be the equivalent inductance, replacing the bottom left circuit with a  $\frac{di}{dt} = \frac{V_1}{L_{eq}}$ ,
- $\frac{di_1}{dt} = \frac{V_1}{L_1}$ ,  $\frac{di_2}{dt} = \frac{V_1}{L_2}$ ,
- single inductor.
- By Kirchhoff's node rule,  $i_1 + i_2 = i$ , so
- $\frac{V_1}{L_{eq}} = \frac{V_1}{L_1} + \frac{V_1}{L_2}$  and
- $L_{eq} = \frac{L_1 L_2}{L_1 + L_2}$



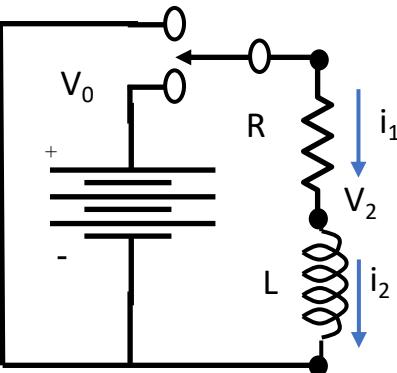
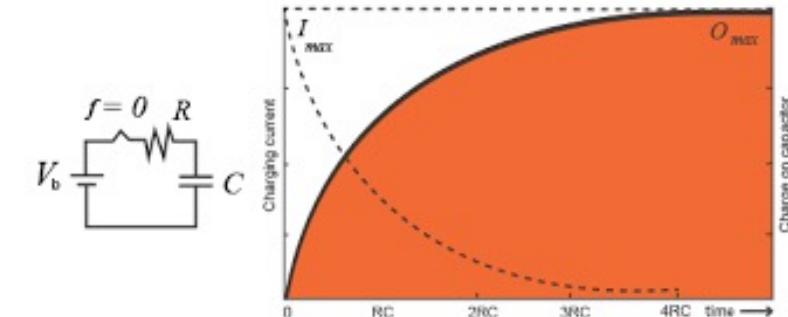
- The circuit on the right, is useful and is called a *voltage divider*.
- $i = i_1 = i_2$  so  $\frac{V_1 - V_2}{R_1} = \frac{V_2}{R_2}$ ,  $V_1 - V_2 = \frac{R_1}{R_2} V_2$
- Thus,  $V_1 = (1 + \frac{R_1}{R_2})V_2$  and so
- $V_2 = \frac{R_2}{R_1 + R_2} V_1$

# RC/RL circuit analysis with Kirchhoff



- RC behavior: charging

- $V_0 - V_2 = i_1 R = V_R, i_1 = \frac{V_R}{R}$
- $i_2 = C \frac{dV_2}{dt}, V_C = V_2$
- $i_1 = i_2, V_C = V_0 - V_R$
- $\frac{V_R}{R} = C \frac{dV_C}{dt}, RC \frac{dV_C}{dt} = V_0 - V_C$ , or  $RC \frac{dV_C}{dt} + V_C = V_0$
- Solution is  $V_C = V_0 - V_0 e^{-\frac{t}{RC}}$



- RL behavior: charging

- $V_0 - V_2 = i_1 R = V_R$
- $V_L = V_2 = L \frac{di_2}{dt}$
- $i_1 = i_2, V_R = V_0 - V_L$ , so  $L \frac{d}{dt} \frac{V_0 - V_L}{R} = V_L$
- $\frac{L}{R} \frac{dV_L}{dt} + V_L = 0$
- Solution is  $V_L = V_0 e^{-\frac{Rt}{L}}$



# Phasors

- $V(t) = RI(t)$
- $V(t) = L\dot{I}(t)$
- $I(t) = C\dot{V}(t)$
- Suppose  $V(t) = A\cos(\omega t + \theta)$  and  $I(t) = B\cos(\omega t + \phi)$ . If  $\phi > \theta$ , we say the current leads the voltage.
- $V(t) = \operatorname{Re}(e^{j(\omega t + \theta)})$ , and  $I(t) = \operatorname{Re}(e^{j(\omega t + \phi)})$
- Now define  $V = Ae^{j\theta}$  and  $I = Be^{j\phi}$ , so  $|V| = A$ ,  $|I| = B$ ,  $\angle V = \theta$ , and  $\angle I = \phi$ .  $V$  and  $I$  are called phasors and do not include time. Note that  $V(t) = \operatorname{Re}(Ve^{j\omega t})$  and  $I(t) = \operatorname{Re}(Ie^{j\omega t})$ .
- Note that  $I = CVj\omega$ , for a capacitor and  $V = LIj\omega$ , for an inductor

# Circuit analysis and impedance

- Impedance unifies the “simple” ohms law with capacitance and inductance.
- $Z = R$ , for resistors,  $Z = j\omega L$ , for inductors and  $Z = \frac{1}{j\omega C}$ , for capacitors.
- In general,  $Z = R + jX$  and all the ohm like laws hold for resistors, capacitors and inductors .
  - $Z_{eq} = Z_1 + Z_2$  for two components with impedance  $Z_1, Z_2$  connected in series
  - $Z_{eq} = \frac{Z_1 Z_2}{Z_1 + Z_2}$  for two components with impedance  $Z_1, Z_2$  connected in parallel
- For example, for a resistor and capacitor in series has impedance  $Z = R + \frac{1}{j\omega C}$

# Phasors, impedance and power

- For the circuit on the right,  $Z = R + \frac{1}{j\omega C}$  is the impedance for the resistor and capacitor in series.
- The phasor  $I = \frac{V_0}{Z}$  and the phasor  $V = \frac{I}{j\omega C} = \frac{V_0}{1+j\omega RC}$
- Further,  $|I| = \frac{V_0}{|Z|}$ ,  $\angle I = \angle \frac{V_0}{|Z|}$  and  $|V| = \frac{|I|}{|j\omega C|} = \left| \frac{V_0}{1+j\omega RC} \right|$
- For phasors  $V, I$ , define the complex power as  $P = \frac{V\bar{I}}{2} = Z \frac{I\bar{I}}{2} = R \frac{|I|^2}{2} + jX \frac{|I|^2}{2}$ ; the first term is the real power, the second is called the *reactive power*.
- The average power is  $P_a = \text{Re}(P) = \text{Re}\left(\frac{V\bar{I}}{2}\right)$ . We define the reactive power as  $P_r = \text{Im}(P)$ .
- $P_r = \omega(E_L - E_C)$ , where  $E_L$  and  $E_C$  are respectively, the energy stored in the inductor and capacitor respectively.



# Q and phasors

- Consider the series resonance on the right.  $Z_{LCR} = R + j\left(\omega L - \frac{1}{\omega C}\right)$
- The phasor,  $I = \frac{V_0}{Z_{LCR}}$ , and the phasor  $V_R = \frac{V_0}{Z_{LCR}}Z_R$ , where  $Z_R = R$ .
- So  $V_R = \frac{RC\omega V_0}{RC\omega + j(LC\omega^2 - 1)}$ .
- $|V_R|$  is maximum when  $\omega^2 LC = 1$ . Put  $\omega_0 = \frac{1}{\sqrt{LC}}$ . When  $\omega = \omega_0$ ,  $|V_R| = V_R = V_0$ .
- $|V_R| = \frac{V_0}{\sqrt{2}}$ , when  $X = R$ . Note that the power through R when  $X = R$  is half the power through R when  $X = 0$  or  $\omega = \omega_0$ .
- Let the frequencies where  $R = \pm X$  be denoted  $\omega_u$  and  $\omega_l$ , where  $\omega_u > \omega_l$ .
- We define  $Q = \frac{X}{R} = \frac{\omega_0 L}{R} = \frac{1}{\omega_0 C R}$ .
- Solving for  $\omega_u$  and  $\omega_l$ , we get  $\frac{L\omega_u}{\omega_0} - \frac{\omega_0}{C\omega_u} = R$  and  $\frac{L\omega_l}{\omega_0} - \frac{\omega_0}{C\omega_l} = -R$ , or, in terms of  $Q$ ,
- $\frac{\omega_u - \omega_0}{\omega_0} = \frac{1}{Q}$  and  $\frac{\omega_l - \omega_0}{\omega_0} = -\frac{1}{Q}$ . In fact,  $\omega_0 = \sqrt{\omega_u \omega_l}$ , and so  $\frac{\omega_u - \omega_l}{\omega_0} = \frac{1}{Q}$ .
- Thus  $Q = \frac{\omega_0}{\omega_u - \omega_l} = \frac{\omega_0}{\Delta\omega}$
- From the definition of  $P_a$ , earlier,  $Q = \omega_0 \frac{E}{P_a}$ , where  $E$  is the total energy stored in  $L$  and  $C$ , which is in turn the peak  $E_L$  and peak  $E_C$  at resonance.



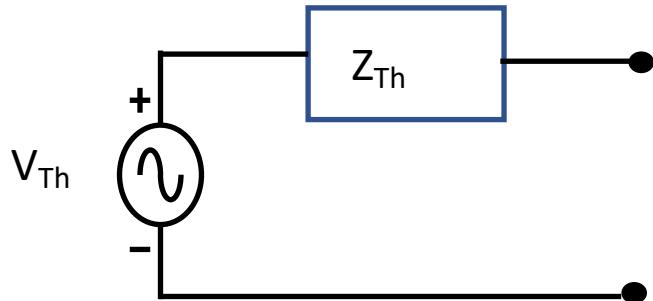
# Phasors, impedance and power

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- Further,  $|I| = \frac{V_0}{|Z|}$ ,  $\angle I = \angle \frac{V_0}{|Z|}$  and  $|V| = \frac{|I|}{|j\omega C|} = \left| \frac{V_0}{1+j\omega RC} \right|$
- For phasors  $V, I$ , define the complex power as  $P = \frac{V\bar{I}}{2} = Z \frac{I\bar{I}}{2} = R \frac{|I|^2}{2} + jX \frac{|I|^2}{2}$ ; the first term is the real power, the second is called the *reactive power*.
- The average power is  $P_a = \text{Re}(P) = \text{Re}\left(\frac{V\bar{I}}{2}\right)$ . We define the reactive power as  $P_r = \text{Im}(P)$ .
- $P_r = \omega(E_L - E_C)$ , where  $E_L$  and  $E_C$  are respectively, the energy stored in the inductor and capacitor respectively.



# Thevenin and Norton

- Thevenin: Any combination of linear sources and passive elements terminating in two terminals is equivalent to a pure voltage source in series with an impedance



- Norton: Any combination of linear sources and passive elements terminating in two terminals is equivalent to a pure current source in parallel with a conductance



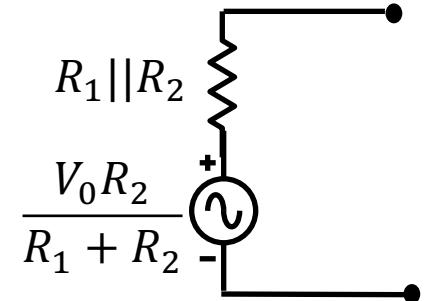
- Similar theorems for two terminal input and output devices (with transfer function)

# Thevenin and Norton

- We can use lookback resistance to calculate the Thevenin equivalent resistance and ideal source.
- To find the lookback resistance, short the source and apply the usual laws.
  - Here  $R_s = R_1 \parallel R_2$
- To find the new ideal source, notice  $R_1$  and  $R_2$  form a voltage divider.
  - The new source voltage is  $\frac{V_0 R_2}{R_1 + R_2}$



Is equivalent to



# Exercise 1: Resistors

1. Consider (A). Find the formula for power in the load. Find the  $R_l$  that maximizes the power to the load.

- $V_l = \frac{R_l}{R_s + R_l} V_0, I_l = \frac{V_0}{R_s + R_l}$ .

- $P_l = V_l I_l = \frac{R_l}{(R_s + R_l)^2} V_0^2$ , which is maximum when  $R_l = R_s$

2. Find the Thevenin and Norton parameters fore (B).

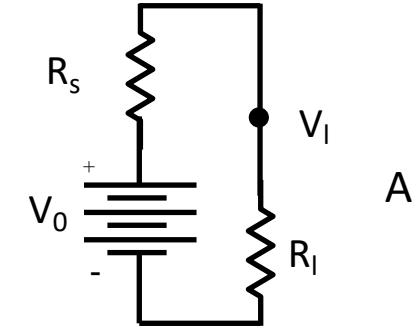
- $V_{Th} = \frac{R_3}{R_1 + R_3} V_0$

- $R_{Th} = R_2 + R_1 || R_3$

3. Find the Thevenin and Norton parameters fore (C).

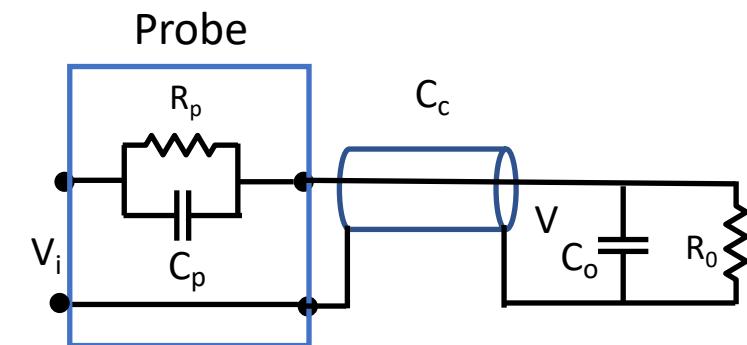
- $V_{Th} = \frac{R_3}{R_2 + R_3} V_0$

- $R_{Th} = R_2 || R_3$



# Exercise 3: Capacitors

1. In the circuit on the right,  $V_0$  is a 2 volt pp ideal square wave source of frequency 20Hz,  $R_s = 50\Omega$ ,  $R = 300k\Omega$  and  $C = 10 nF$ . Period is 50 msec
2. What is the voltage,  $V$ , at the output? The scope has an input resistance of  $1M\Omega$ .
  - About a volt at peak
3. Let  $t_2$ , the time to discharge to 0V. Calculate  $\tau$  and  $t_2$ .
  - $\tau = 3 \times 10^5 \times 10^{-8} \text{ sec} = 3 \text{ msec}$
  - $t_{12} \approx 1.5ms$
4. Capacitance on the scope prevents the delay from being 0. Measure the new  $t_2$  with these changes.
5. Given  $C_0$  and  $C_p$  and  $R_p$ .
  - $C_0 = 100pf/m$ ,  $C_o = 50pF$ ,  $C_p = 10pF$
6. Now calculate the new  $t_{12}$ .
  - $\tau = 6\mu\text{-sec}$



# Exercise 4: Diode detectors

- For AM,  $V(t) = V_c \cos(\omega_c t) + a(t) \cos(\omega_c t)$ , Define the modulation depth  $m = \frac{a_p}{V_c}$
- In circuit on the right,  $R = 10k\Omega$ ,  $C = 10 nF$
- Set function generator for  $f_c = 1MHz$ ,  $V_{c,pp} = 5V$ ,  $f_m = 1kHz$ ,  $m = .7$ 
  1. Calculate  $\tau$  for the RC circuit.  $\tau = 10^4 \times 10^{-8} \text{ sec} = .1ms$ .
    - $T_m$  is period of modulating signal.  $T_m = 10^{-3} \text{ sec} = 1ms$ . So  $\tau \ll T_m$
    - $T_c$  is period of modulating signal.  $T_c = 10^{-6} \text{ sec} = 1\mu s$ .  $\tau \gg T_c$
    - As you change  $f_m$  does the frequency of  $V_o$  track it? (It better)
  2. Compare the max voltage of the AM signal to the max of  $V_o$ .
    - $V_o, p \approx .8V$ ,  $V_{i,p} \approx 1.4V$
  3. What happens when we make  $m = 1.0$



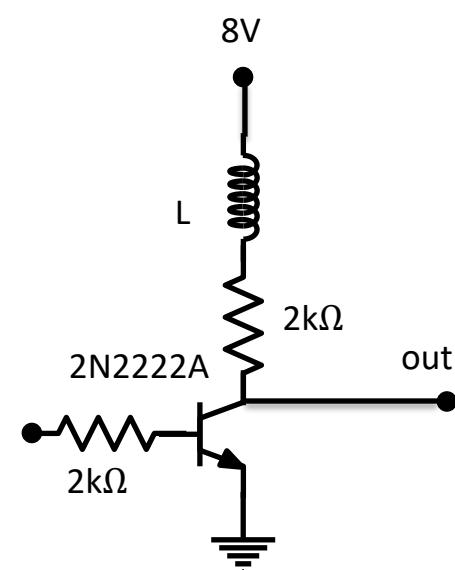
# AM Modulator for previous exercise

- I didn't have a signal generator that produced an AM signal, so I used the modulator on the right with the indicated inputs to produce the AM needed for the detector in the previous exercise.



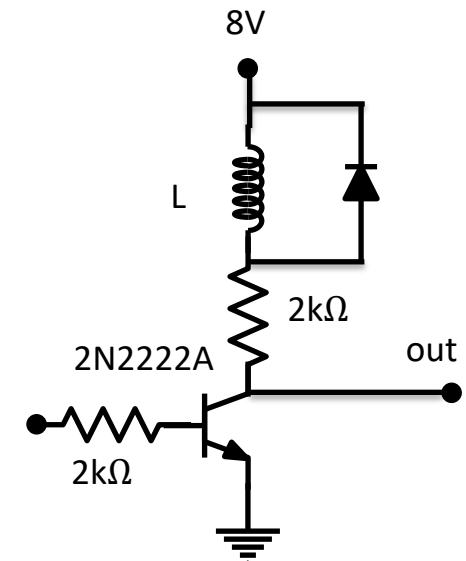
# Exercise 5: Inductors

- Set function generator for square wave with 5V  $V_{pp}$ , a Thevenin equivalent source resistance of  $R_{Th} = 50\Omega$  and frequency 1kHz. Connect a load,  $R_L = 100\Omega$  load, L=1mH
  - Observe square wave with rounded corners, measure the time,  $t_2$  to decay to 0
    - About  $20\mu\text{sec}$
  - In the top circuit, calculate inductor current and the expected delay,  $t_2$ 
    - $Z_{eq} = 150 + jL\omega$ ,  $\omega = 2\pi \times 10^3$ ,  $V_i = Re(V_{i,p}e^{j\omega t})$
    - As phasors,  $iZ_{eq} = V_i$ ,  $|i| \sqrt{150^2 + (\omega L)^2} = V_{i,p}$ ,  $|i| = \frac{V_{i,p}}{\sqrt{150^2 + (2\pi)^2}}$ ,  $\theta = \angle i = \arctan(-\frac{2\pi}{150})$ ,  $\theta \approx -2.4 \text{ rad} = -15^\circ$
    - $V_o = Re\left(\frac{100V_{i,p}}{\sqrt{150^2 + (2\pi)^2}} e^{j(\omega t + \theta)}\right)$ ,  $|V_o| = 1.6V$ ,
    - $\tau_{RL} = \frac{10^{-3}}{100} \text{ sec} \approx 10 \mu\text{sec}$
  - In the second circuit, use 2 scope channels: one at input, one at output.
    - $1\mu\text{sec rise time}$ . Ringing at 10MHz.  $\frac{1}{\sqrt{LC}} = 62.8 \times 10^6$ .  $C = \frac{10^3}{(62.8 \times 10^6)^2} \approx .25\text{pF}$
    - Note: I made the pull-up 100K.



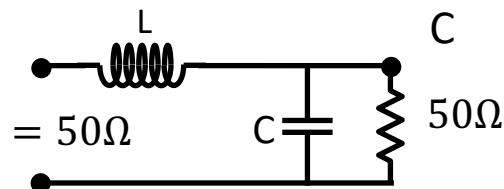
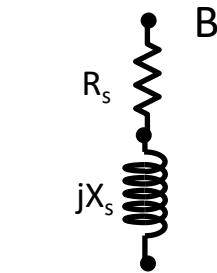
# Exercise 6: Diodes and snubbers

- Add indicated snubber diode.
  - 1. Swing up is nearly immediate with snubber
  - 2. Ringing at 10MHz.  $\frac{1}{\sqrt{LC}} = 62.8 \times 10^6$ .  $C = \frac{10^3}{(62.8 \times 10^6)^2} \approx .25 \mu F$
  - 3. What is its effect on ringing?
    - Ringing is uniform at 5 MHz
  - 4. Diode should be on when transistor is off.
- 
- Note: I made the pull-up 100K.



# Exercise 7: Parallel to Series conversion

- For series:  $Z_s = R_s + j\omega L$ ,  $Q_s = \frac{X_s}{R_s}$
  - For parallel:  $\frac{1}{Z_p} = \frac{1}{R_p} + \frac{1}{j\omega L}$ , so  $Z_p = \frac{j\omega L R_p}{R_p + j\omega L}$  and  $Q_p = \frac{R_p}{X_p}$
  - If  $Q_p = Q_s$ ,  $X_p X_s = R_p R_s$
- If circuits (A) and (B) have the same impedance, what is the relationship between  $R_p, X_p$  and  $R_s, X_s$ ?
    - $\frac{1}{Z_p} = \frac{1}{R_p} + \frac{1}{jX_s}$ ,  $Z_p = \frac{jR_p X_p}{R_p + jX_p} = \frac{X_p^2 R_p + jR_p^2 X_p}{R_p^2 + X_p^2}$ ,  $Z_s = R_s + jX_s$
    - $R_s = \frac{X_p^2 R_p}{R_p^2 + X_p^2}$ ,  $X_s = \frac{R_p^2 X_p}{R_p^2 + X_p^2}$ ,  $R_s = X_p \frac{X_p R_p}{R_p^2 + X_p^2}$ ,  $X_s = R_p \frac{X_p R_p}{R_p^2 + X_p^2}$ , set  $\rho = \frac{X_p R_p}{R_p^2 + X_p^2}$  for later reference
    - This shows the Q's must be equal as stated above.
  - Find a formula for  $X_s$ , for large  $Q = Q_p = Q_s$  and small  $Q = Q_p = Q_s$
  - Use circuit (C) to transfer a  $50\Omega$  load (circuit C) to a  $5\Omega$  load. What is  $X_C$  at 7 MHz?  $Z_i = 50\Omega$   
What are C and L in that model?
    - Use the parallel to series conversion to make a series equivalent circuit consisting of C and the  $50\Omega$  with  $R_s = 5\Omega$



# Exercise 8: Series resonance

- For the circuit on the right,  $C = 8 - 50\text{pf}$ ,  $L = 15\mu\text{H}$  forming a bandpass filter.  $R = 100\Omega$
- If  $C = 34\text{pf}$ , the resonant frequency is  $\omega = \frac{1}{\sqrt{35 \times 10^{-12} \times 15 \times 10^{-6}}} = \frac{10^9}{\sqrt{525}} \approx 44.2$ , so the resonant frequency is  $\frac{44.2}{2\pi} \approx 7.07\text{MHz}$

1. Tune the resonant frequency to  $7\text{MHz}$  and find  $f_u$ ,  $f_l$  and  $\Delta f$  and thus  $Q$ .

- $f_u = 7.67\text{MHz}$ ,  $f_l = 6.47\text{MHz}$ ,  $Q = \frac{f}{\Delta f} = \frac{7}{1.2} = 5.8$

2. Compute what these values should be

- $Z_{eq} = R + j(\omega L - \frac{1}{\omega C})$
- As phasors,  $i = |i|e^{j\theta}$ ,  $|i| = \frac{V_{i,0}}{\sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2}}$ ,  $\theta = -\arctan(\frac{\omega L - \frac{1}{\omega C}}{R})$
- $V_0 = iR$ , Power through R at  $\omega$  is  $P(\omega) = |i|^2|R = \frac{|V_{i,0}|^2 R}{R^2 + (\omega L - \frac{1}{\omega C})^2}$ . At resonance,  $P(\omega_r) = \frac{|V_{i,0}|^2}{R}$ . To find half power, set  $\frac{1}{2} = (\frac{|V_{i,0}|^2 R}{R^2 + (\omega L - \frac{1}{\omega C})^2}) / (\frac{|V_{i,0}|^2}{R})$ , or  $R = \omega L - \frac{1}{\omega C}$ .
- Solving gives  $f_u = 7.67\text{MHz}$ ,  $f_l = 6.53\text{MHz}$ ,  $Q = 6.1$
- General formulas:  $\omega_u = \frac{R}{2L} + \sqrt{(\frac{R}{2L})^2 + \frac{1}{LC}}$ ,  $\omega_l = -\frac{R}{2L} + \sqrt{(\frac{R}{2L})^2 + \frac{1}{LC}}$



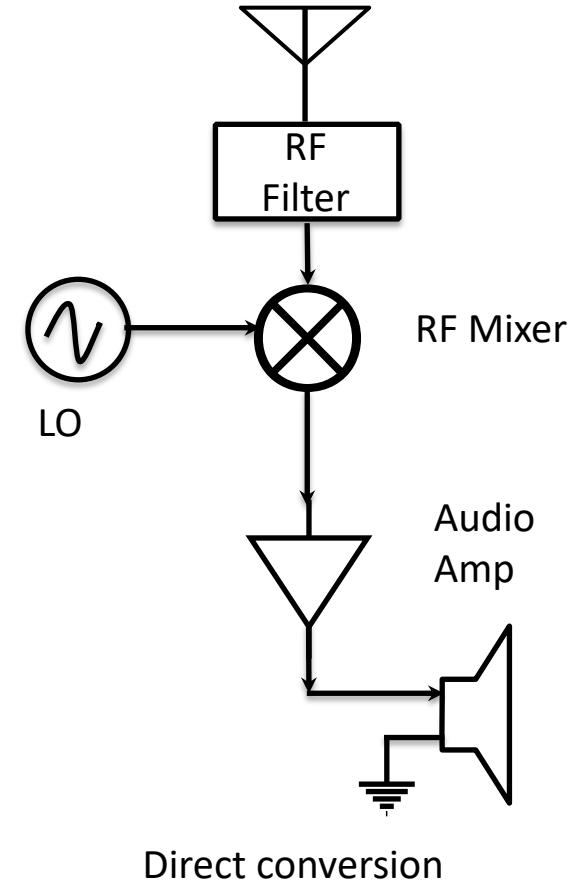
# Exercise 9: Parallel resonance

- $L = A_l N^2$ ,  $A_l = 4 \frac{nH}{turn^2}$  for T37-2 core so for 28 turns,  $L_6 = 3.1\mu H$
- 1. Again, find the resonant frequency, the frequencies corresponding to a 3db falloff, the bandwidth and the Q of this circuit. This circuit is in the transmit oscillator
  - For  $C_{38}$ ,  $L_6$ ,  $R = 100\Omega$ , network:  $Q = 100 \sqrt{\frac{10^{-10}}{3.1 \times 10^{-6}}} = 5.6$ .
  - $BW = \frac{f_r}{Q} = \frac{7MHz}{5.6} = 1.25MHz$ .  $f_u = f_r + \frac{BW}{2} = 7.625MHz$ ,  $f_l = f_r - \frac{BW}{2} = 6.375MHz$ .
  - General formulas:  $Q = R \sqrt{\frac{C}{L}}$ ,  $BW = \frac{f_r}{Q}$ ,  $f_u = f_r + \frac{BW}{2}$ ,  $f_l = f_r - \frac{BW}{2}$

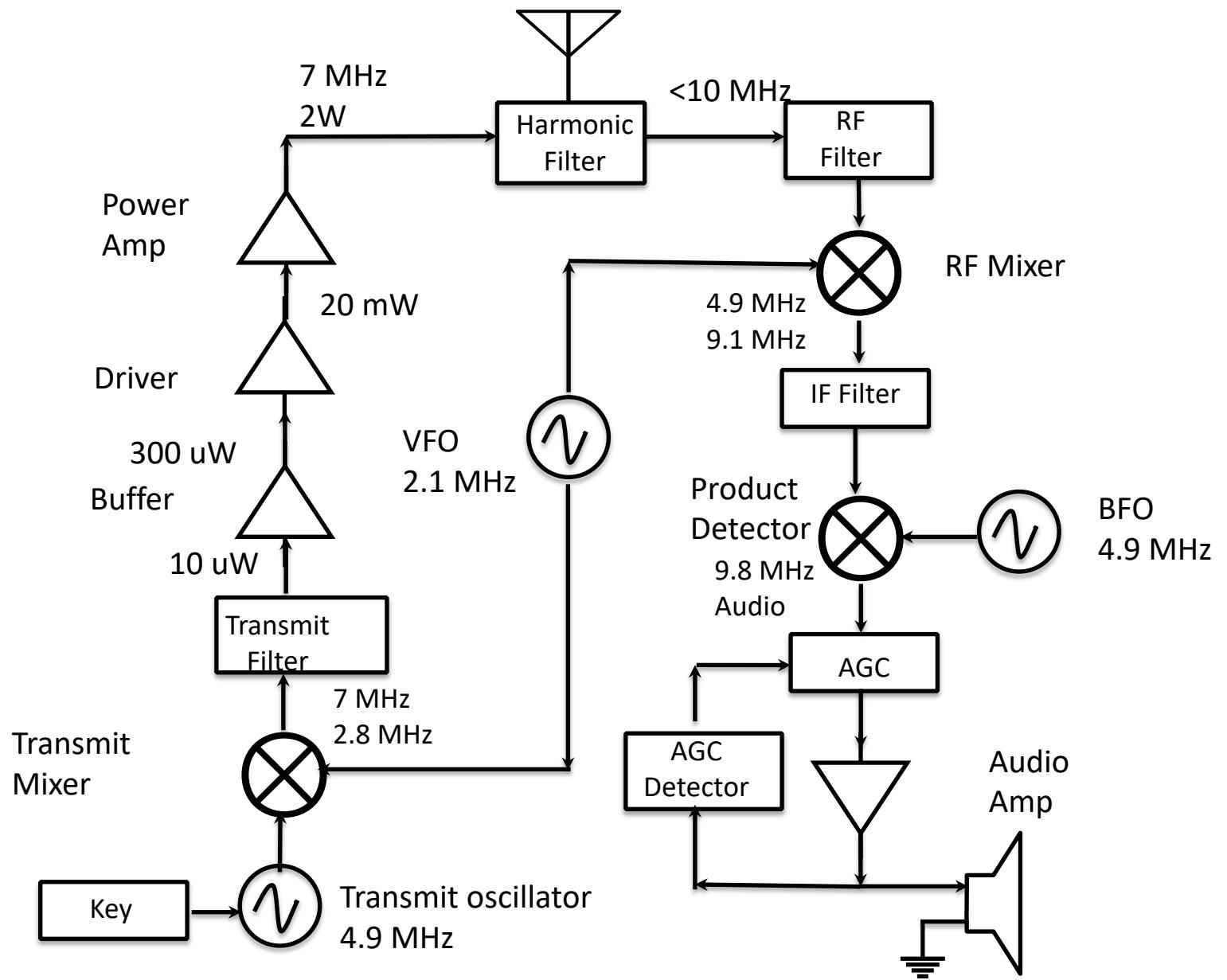


# Direct conversion and superhet receivers

- Image frequency
  - $\omega_{rf} = \omega_{LO} - \omega_a$
  - $\omega_i = \omega_{LO} + \omega_a$
- Superheterodyne designs
  - $\omega_{rf} = \omega_{IF} + \omega_{VFO}$
  - $\omega_{vi} = \omega_{IF} - \omega_{VFO}$
  - $\omega_{IF} = \omega_{BFO} - \omega_a$
  - $\omega_{bi} = \omega_{BFO} + \omega_a$
  - $\omega_{usb} = \omega_{VFO} + \omega_{BFO} + \omega_a$
  - $\omega_{lsb} = \omega_{VFO} + \omega_{BFO} - \omega_a$

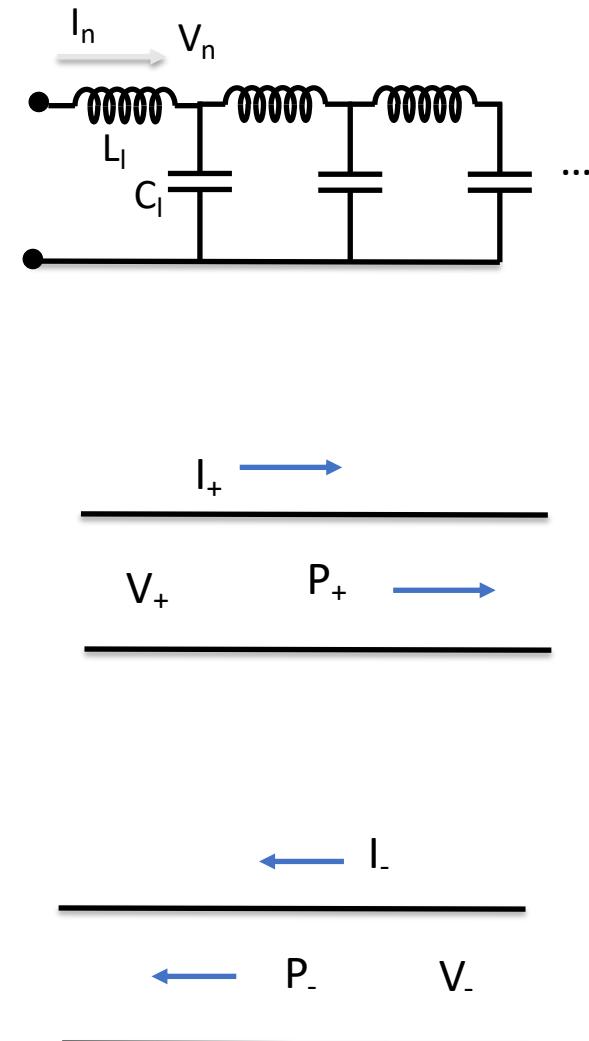


# Norcal 40A



# Transmission Lines

- $V_{n+1} - V_n = -L_l \frac{\partial I_{n+1}}{\partial t}, L = \frac{L_l}{l}$
- $I_{n+1} - I_n = -C_l \frac{\partial V_n}{\partial t}, C = \frac{C_l}{l}$
- $\frac{\partial^2 V}{\partial z^2} = LC \frac{\partial^2 V}{\partial t^2}$  and  $\frac{\partial^2 I}{\partial z^2} = LC \frac{\partial^2 I}{\partial t^2}$
- Solution is  $V(z - vt)$ ,  $v = \frac{1}{\sqrt{LC}}$ , which is the velocity for the forward wave
- $V' = vLI'$ ,  $\frac{V}{I} = \sqrt{\frac{L}{C}}$ ,  $Z_0 = \sqrt{\frac{L}{C}}$ ,  $Z_0$  is the forward or “characteristic” impedance
- Another solution is  $V(z + vt)$ ,  $v = \frac{1}{\sqrt{LC}}$  which is the velocity for the reverse wave
- $Z_0 = \frac{V_+}{I_+}$ ,  $-Z_0 = \frac{V_-}{I_-}$ ,  $V = V_+ + V_-$ ,  $-Z_0$  is the backwards looking impedance
- $P_+(t) = \frac{V_+^2}{Z_0}$ ,  $P_-(t) = -\frac{V_-^2}{Z_0}$



# Transmission Lines - continued

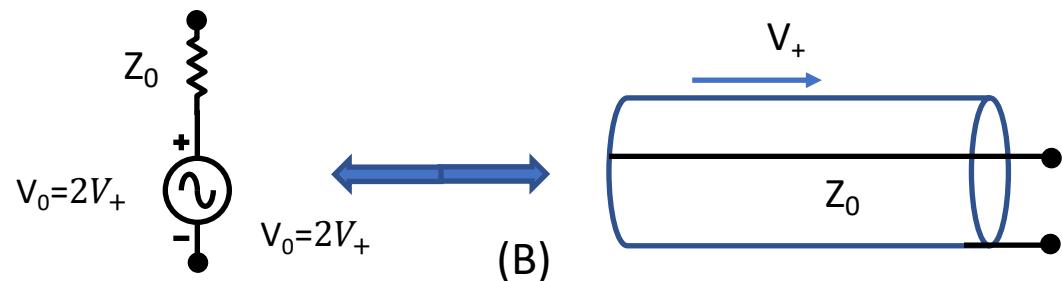
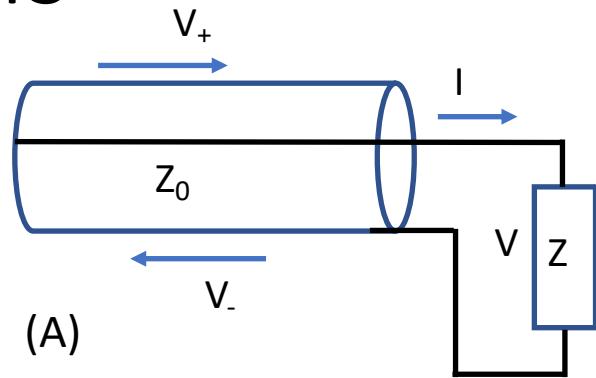
- $V(z - vt) = A\cos(\omega t - \beta z)$ ,  $v = \frac{\omega}{\beta}$ ; in phasor notation,  $V = Ae^{-j\beta z}$ .
- To compute power, let  $V_+ = Ae^{-j\beta z}$ ,  $V_- = Ae^{j\beta z}$
- Complex power is  $P = \frac{V\bar{I}}{2}$ ,  $P_+ = \frac{V_+\bar{I}_+}{2} = \frac{|V_+|^2}{2Z_0}$ ,  $P_- = \frac{V_-\bar{I}_-}{2} = -\frac{|V_-|^2}{2Z_0}$ , with  $\frac{V}{I} = Z_0$
- Suppose over a transmission line,  $Z$  is the distributed impedance/m,  $Y$  is the distributed admittance/m and suppose the forward wave is  $Ae^{j(\omega t - jk)}$ , and its phasor is  $V = Ae^{-jkz}$ . Let  $Z = \frac{V}{I}$ ,
- $\frac{dV}{dz} = -ZI$ ,  $\frac{dI}{dt} = -YV$ . ... (1)
- Put  $jk = \alpha + \beta j$  (to account for attenuation), then  $jk = \sqrt{ZY}$  and the forward phasor becomes  $e^{(-\alpha z - j\beta z)}$ .  $\alpha_{dB/m}$  is a transmission loss.  $\alpha_{dB/m} = 8.686\alpha_{nepers/m}$ . By differentiating (1), we get  $jkV = ZI$ ,  $-jkI = YV$ . Solutions are  $jk = \sqrt{ZY}$ ,  $Z_0 = \frac{V}{I} = \sqrt{\frac{Z}{Y}}$ , all complex
- Example: If  $Z = R + j\omega L$ ,  $Y = j\omega C + G$  for the transmission line, then  $jk = \sqrt{(j\omega L + R)(j\omega C + G)}$  and  $Z_0 = \sqrt{\frac{(j\omega L + R)}{(j\omega C + G)}}$  (positive real root)

# Transmission Lines - dispersion

- $\alpha$  and  $\nu$  can vary with frequency; this is dispersion.
- Heaviside: Adjust parameters so  $\frac{R}{L} = \frac{G}{C}$ , then  $\alpha$  and  $\nu$  are constants and we get:
  - $jk = j\omega\sqrt{LC}(1 + \frac{R}{j\omega L})$  and  $\nu = \frac{\omega}{\beta} = \frac{1}{\sqrt{LC}}$ ,  $\alpha = \sqrt{RG}$
  - We also get  $Z_0 = \sqrt{\frac{L}{C}}$  as with a lossless line.
  - If  $\omega L \gg R$ 
    - $G = 0$  and  $Z_0 = \sqrt{\frac{(j\omega L+R)}{j\omega C}} \approx \sqrt{\frac{L}{C}}$
    - If  $R \gg \omega L$ 
      - $jk = \sqrt{\frac{(j\omega L+R)}{j\omega C}} \approx \sqrt{j\omega RC}$ , and  $\alpha = \sqrt{\frac{\omega RC}{2}}$ ,  $\nu = \sqrt{\frac{2\omega}{RC}}$
- For first transatlantic cable,  $L = 460 \frac{nH}{m}$ ,  $C = 75 \frac{pF}{m}$ ,  $f = 12 \text{Hz}$ ,  $R = 7 \frac{m\Omega}{m}$ ,  $l = 3600 \text{ km}$ ,  $\alpha = \sqrt{\frac{\omega RC}{2}} = 4.4 \times 10^{-3} \text{nepers/m}$ ,  $\alpha l = 140 \text{dB}$ 
  - $\alpha l \approx 140 \text{dB}$  and highly dispersive

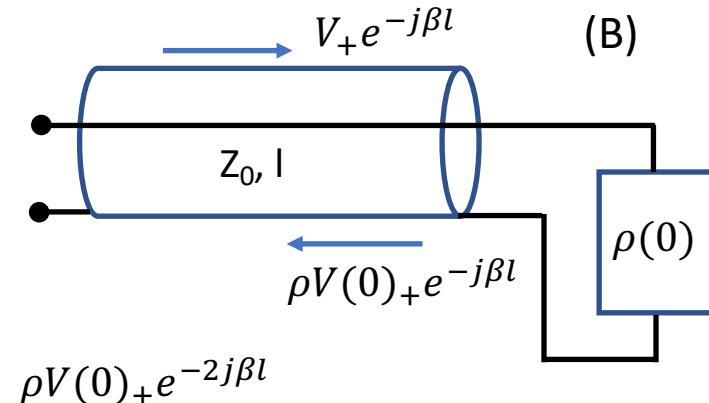
# Transmission Lines-reflections

- $\rho = \frac{V_-}{V_+}$ ,  $\rho V_+ = V_-$ ,  $V = V_+ + V_- = (1 + \rho)V_+$
- $\rho_i = \frac{i_-}{i_+} = -\rho$
- $\tau = \frac{V}{V_+} = 1 + \rho = \frac{2Z}{Z+Z_0}$ ,  $V = 2V_+$
- Consider the circuit in the upper right (A):  $V = V_+ + V_-$ ,  $I = I_+ + I_-$ ,  $Z = \frac{V}{I}$
- $Z = \frac{V}{I} = \frac{V_+ + V_-}{I_+ + I_-}$
- $\frac{Z}{Z_0} = \frac{1+\rho}{1-\rho}$ ,  $\rho = \frac{Z-Z_0}{Z+Z_0}$ ,  $\rho_{open-circuit} = 1$ .
  
- For (B):
- Lookback resistance is  $R_s = Z_0$ , short circuit for (B) is  $i_s = \frac{V_0}{Z_0}$
- Thevenin equivalent for open circuit is (B)
- $P_+ = \frac{V_+^2}{2Z_0} = \frac{V_0^2}{8R_s}$ , This is the total available power



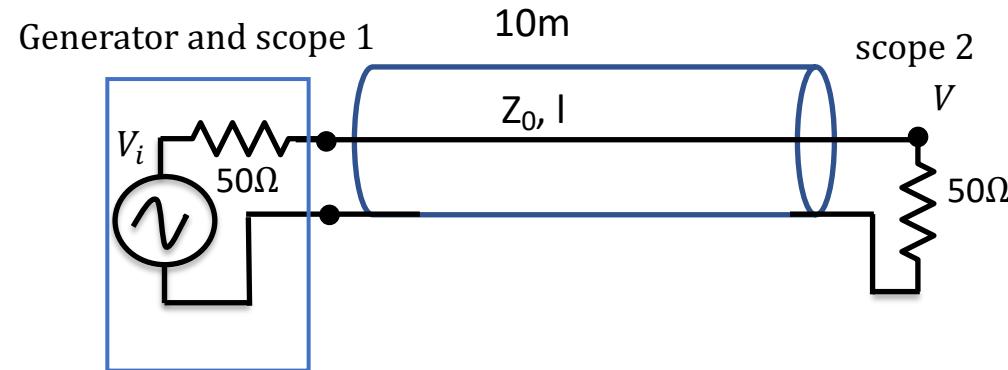
# Transmission Lines – resonance and Q

- For (A) on right,  $V_+ = \frac{V_0}{2}$ ,  $V = V_+ + V_- = V_0 e^{-j\beta l}$ ,  $V_- = \frac{V_0}{2} e^{-2j\beta l}$ 
  - $V_g = V_0 e^{-j\beta l} \cos(\beta l) = \frac{V_0}{2} (1 + e^{-2j\beta l})$ ,  $V_g(\frac{\lambda}{4}) = 0$
  - $I_g = \frac{V_+ - V_-}{Z_0} = jI_s e^{-j\beta l} \sin(\beta l)$ .
  - $X = \frac{V_g}{jI_g} = \frac{Z_0}{\tan(\beta l)}$
- $Q = \omega \frac{E}{P_a}$ ,  $E = \frac{lP_+}{v}$ ,  $P_a = P_+ - P_+ e^{-2\alpha l} \rho(0) \approx 2\alpha l P_+$ ,  $Q = \frac{\beta}{2\alpha}$
- In (B) to the right, the coefficient of reflection is  $\rho(0)$  and the generator absorbs the reverse wave.  $V = V_+ + V_- = V_0 e^{-j\beta l}$ .
  - $V_f = \rho(0)V_+e^{-j\beta l}$ ,  $V_r = \rho(0)V_+e^{-2j\beta l}$
  - $\rho(l) = \frac{V_-}{V_+} = \rho(0)e^{-2j\beta l}$  is the reflection coefficient at generator.
  - $\rho(\frac{\lambda}{2}) = \rho(0)$ ,  $\rho(\frac{\lambda}{4}) = -\rho(0)$
  - $\frac{Z(\lambda/4)}{Z_0} = \frac{Z_0}{Z(0)}$ ,  $Z = \frac{Z}{Z_0}$ ,  $y = \frac{1}{z}$ ,  $Z(\frac{\lambda}{4}) = -\frac{1}{z(0)}$
  - Normalized:  $Z(\lambda/4) = \frac{1}{z(0)}$
  - $Z_0 = \sqrt{Z(\lambda/4)Z(0)}$ ,  $Z_0 = \sqrt{R_S R_L}$

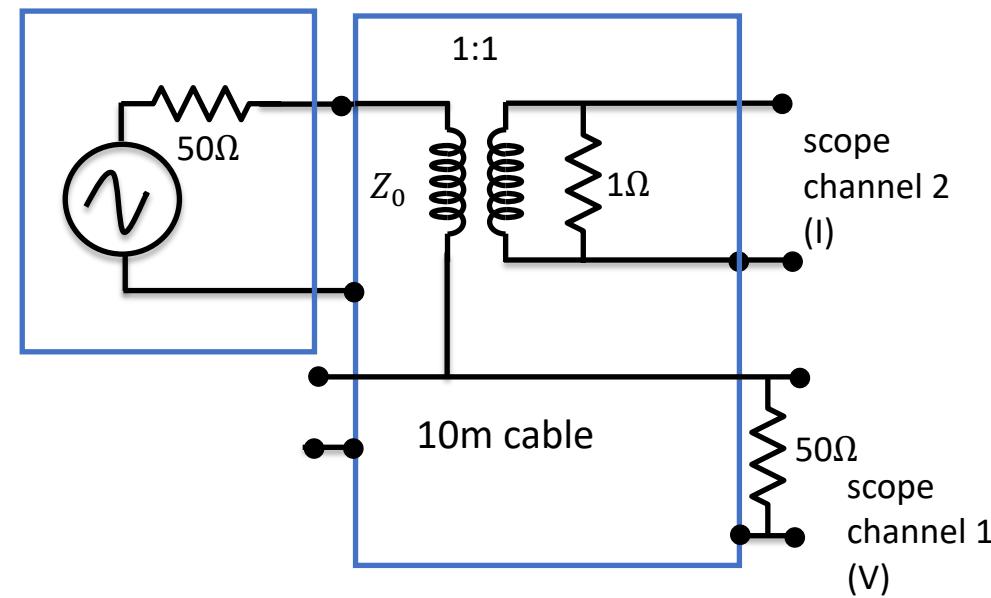


# Exercise 10: Coax

- We'll measure the velocity of waves in RG58U by connecting one channel of the scope to the input and one to the output.
- Measure the velocity,  $v$ , in 10m coax at 7MHz. Try different frequencies. Use  $50\text{ns}$ ,  $5V_{pp}$  using square waves at 20kHz.
  - Do the same with an antenna
  - Calculate  $Z_0$  with  $50\Omega$  termination for the circuit on the right.
  - Remove the  $50\Omega$  and measure the  $V$  and use it and  $Z_0$  to calculate  $L$ , and  $C$  for the coax
    - Measured speed is  $v = 2 \times 10^8 \text{ m/s}$ .  $Z_0 = 50\Omega$ . For high impedance,  $Z_0 = \sqrt{\frac{L}{C}}$  and  $v = \frac{1}{\sqrt{LC}}$ . So,  $Z_0^2 C = L$  and  $v^2 = \frac{1}{LC}$ , so  $Z_0^2 C^2 v^2 = 1$ .  $C = \frac{1}{Z_0 v} = 10^{-10} \text{ F}$ .  $2500 \times 10^{-10} \text{ F} = L = 250 \text{ nH}$ , which is what we use in the next problem.

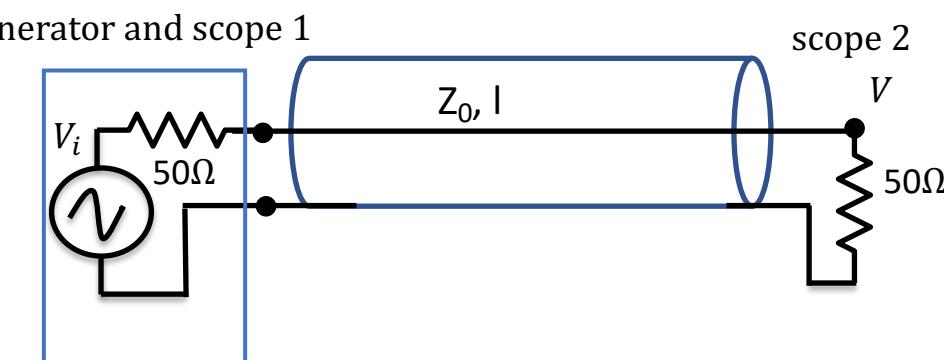


Function generator



# Exercise 11: Waves

- Suppose we want to send voice over 100km of coax,  $Z_L = 50\Omega$ ,  $l = 100\text{km}$ .
- Measure the SWR which is the ratio of the maximum to minimum output
    - $V = V_+ + V_-$ ,  $\rho = \frac{Z - Z_0}{Z + Z_0}$ ,  $Z = 50\Omega$ , we get  $Z_0$  from the previous exercise.
    - $|V_{max}| = |V_f| + |V_r| = (1 + \rho)|V_f|$ ,  $|V_{min}| = (1 - \rho)|V_f| = |V_f|$ .  $SWR = \frac{|V_{max}|}{|V_{min}|} = \frac{1 + \rho}{1 - \rho}$ ,
  - If  $L = 250 \frac{nH}{m}$ ,  $C = 100\text{pf}/m$  and the distributed resistance at voice is  $50 \text{ m}\Omega/m$ , calculate total dB loss at 500, 1000 and 2000Hz using the high frequency approximation.
    - $Z(f) = j\omega L + R = 50 \times 10^{-3} + j \cdot 2\pi f \cdot 250 \times 10^{-9}$
    - $Y(f) = j\omega C + \frac{1}{R} = \frac{1}{50 \times 10^{-3}} + j 2\pi f \cdot 10^{-10}$
    - $Z_0(f) = \sqrt{\frac{Z(f)}{Y(f)}}$
    - $Z_0(500) = 400\Omega$ ,  $Z_0(1000) = 282\Omega$ ,  $Z_0(500) = 200\Omega$ ,
    - High resistance approximation:  $\alpha(f) = \sqrt{\frac{\omega RC}{2}}$ ,
    - $\alpha(500) = 8.8 \times 10^{-5}$ ,  $\alpha(1000) = 12.6 \times 10^{-5}$
    - $\alpha(2000) = 17.6 \times 10^{-5} \times 10^5$
    - For 100km, loss is  $\alpha \times 10^5$



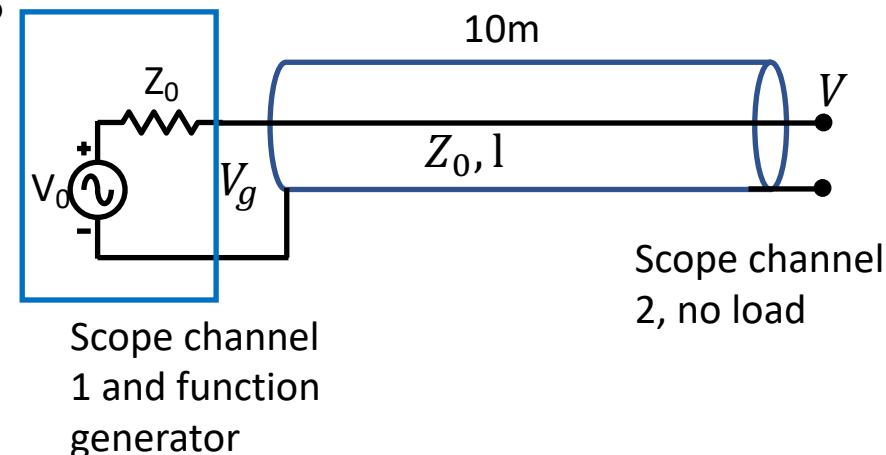
# Exercise 11: Waves

3. Add a 100mH inductor every 1km. Now what's the loss?

- $Z(f) = j\omega L + R = 50 \times 10^{-3} + j \cdot 2\pi f \cdot 10^{-4}$ ,  $Z_0(f) = \sqrt{\frac{Z(f)}{Y(f)}}$
- $Z_0(f) = \sqrt{\frac{Z(f)}{Y(f)}}$
- $Z_0(500) = 318\Omega$ ,  $Z_0(1000) = 317\Omega$ ,  $Z_0(2000) = 316\Omega$
- High reactance approximation:  $\alpha(f) = \frac{R}{2Z_0}$ ,  $Z_0 = \sqrt{\frac{L}{C}} = 1000\Omega$
- $\alpha(f) = \frac{R}{2Z_0(f)}$ ,  $\alpha(500) = \alpha(1000) = \alpha(2000) = \frac{5 \times 10^{-2}}{2000} = 5.5 \times 10^{-5}$  nepers/m
- For 100km, loss is  $\alpha \times 10^5 = 5.5$  or  $5.5 \times 8.868 \approx 49dB$

# Exercise 12: Resonance

- RG58U has a capacitance of about  $100 \frac{pF}{m}$ . Let  $\alpha$  be the attenuation constant and  $\beta$  be the phase
1. Derive an expression for  $|\frac{V_g}{V}|$  and assuming  $\alpha$  is small by finding the first resonance where  $V_g$  is minimum.
    - $V_g = V_0 e^{-j\beta l} \cos(\beta l)$ ,  $V = V_0 \exp(-j\beta l)$ .  $|\frac{V_g}{V}| = \cos(\beta l)$ . So, at  $l = \frac{\lambda}{2}$ ,  $|V_g| = |V|$
  2. Find  $\alpha$  and the wave velocity by finding the resonant frequency (without the load,  $1V_{pp}$ ) and noting the time delay with a scope on the input and output. Use  $|\frac{V_g}{V}|$  to calculate  $\alpha$ .
    - $|V|$  will be maximum at resonant frequency with unterminated line.
    - $|V_g(l)|$  is minimum when  $l = \frac{\lambda_r}{4}$  and  $\beta l = \frac{\pi}{4}$ . This gives  $\beta$ .
    - At  $l = \frac{\lambda}{2}$ ,  $|\frac{V_g}{V}| = e^{-\alpha(\lambda/2)}$
  3. Use this to calculate the velocity,  $v$ . How large is the frequency shift caused?
    - $v = \frac{\omega_r}{\beta}$ . [ $v$  should be about  $2 \times 10^8 \text{ m/s}$ ]
  4. Find, as usual,  $f_u$ ,  $f_l$ , and  $Q$ .
    - $Q = \frac{\beta}{2\alpha}$
    - $Q = \frac{f_r}{BW}$ , so  $BW = \frac{f_r}{Q}$ .  $f_u = f_r + \frac{BW}{2}$ , and  $f_l = f_r - \frac{BW}{2}$



# Filters

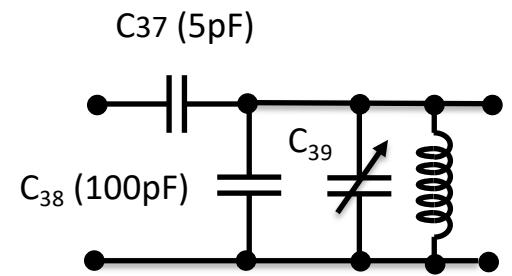


- Circuits on the left are called ladder filters.
- Low pass (Butterworth equivalent): Tabled values for inductors and capacitors based on frequency and dB drop-off.
- Can convert low pass into bandpass.
- For low pass to high pass
- Butterworth:  $L = \frac{P_i}{P} = 1 + (\frac{f}{f_c})^{2n}$ ,  $f_c$  is 3dB cutoff
- Chebyshev:  $L = 1 + \alpha C_n^2 (\frac{f}{f_c})^{2n}$ ,  $f_c$  is 3dB cutoff
- Normalized reactance's:  $a_i = \sin(\frac{(2i-1)\pi}{2n})$
- Ripple loss:  $1 + \alpha = 10^{\frac{Lr}{10}}$
- $\beta = \sinh(\frac{\tanh^{-1}(1/\sqrt{1+\alpha})}{n})$ ,  $c_i = \frac{a_i a_{i-1}}{c_{i-1}(\beta^2 + \sin^2(\frac{(i-1)\pi}{n}))}$
- Example: cutoff at 10MHz, 4<sup>th</sup> order, 50ohm output, 3dB cutoff,  $L(20MHz) = 6n = 24dB$ ,  $a_1 = 0.765$ ,  $X_1 = x_1 Z_0 = 38\Omega$ ,  $L_1 = \frac{X_1}{\omega_c} = 610nH$ ,  $b_2 = a_2 = 1.848$ ,  $B_2 = b_2/Z_0$ ,  $C_2 = B_2/\omega_c$
- Yuck!



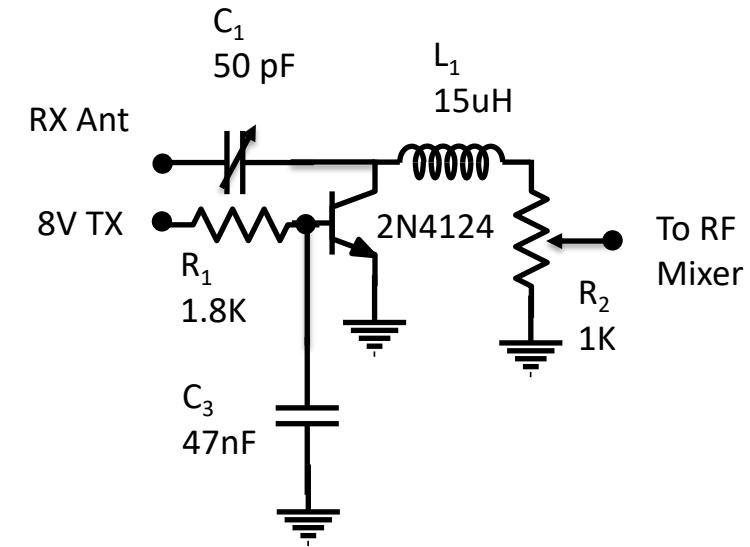
# Norcal transmit bandpass filter

- $C_{39} = 50\text{pF}$ ,
- $L_6$  is 36 turns #28 on T37-2 which has  $A_l = 4 \frac{\text{nH}}{\text{turn}^2}$ ,  $L_6 = A_l \cdot 36^2 = 3.1\mu\text{H}$
- $Z_2 = -\frac{j}{(C_{38}+C_{39})\omega_o}$ ,  $Z_3 = jL_6\omega_o$ ,  $Z_1 = \frac{j}{C_{37}\omega_o}$
- $Z_{2,3-eq} = \frac{jL_6\omega_0}{L_6(C_{38}+C_{39})\omega_0^2-1}$
- Resonance is when  $Z_{2,3-eq} \rightarrow \infty$ ,  $\omega_o^2 = \frac{1}{(C_{38}+C_{39})L_6} \approx \frac{10^{18}}{465}$ , when almost all the voltage drop is across  $Z_{2,3-eq}$   $\omega_o = \frac{10^9}{\sqrt{465}} \approx 50.8 \times 10^6$ ,  $f_0 = \frac{\omega_o}{2\pi} \approx 7.1 \text{ MHz}$
- Q of filter is:  $Q_s = \frac{X_s}{R_s}$ .  $R_s$  comes from the other components and must be measured
- Note that  $Z_{2,3-eq}$  is small for the other modulation product



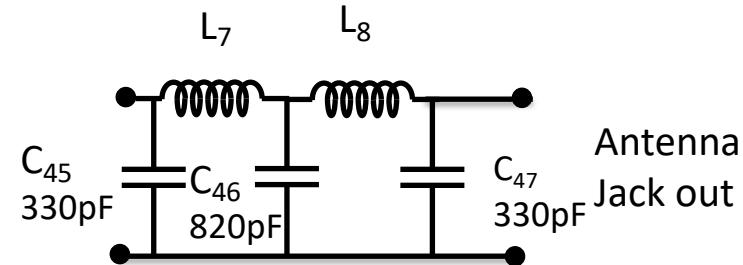
# Norcal RF Filter with switch

- We calculated the output of the  $C_1$ ,  $L_1$  bandpass filter in exercise 1.
- When the transistor conducts, the antenna is shorted to ground through  $C_1$  and  $C_3$ . What power is delivered to  $R_2$  when this happens?



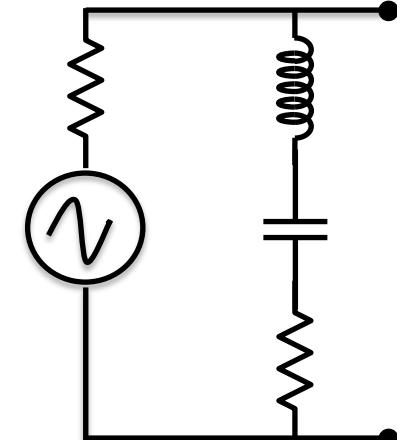
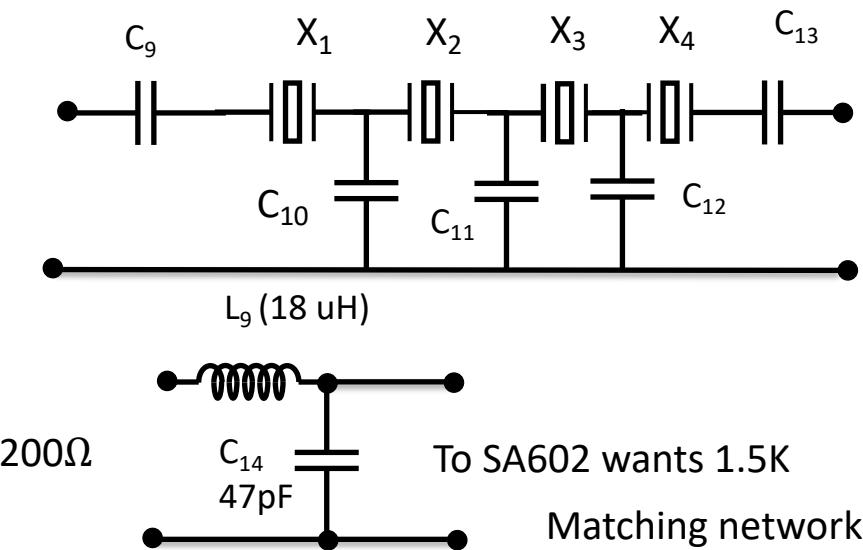
# Exercise 13: Norcal Harmonic Filter

- $L_7, L_8$  use T37-2 core, 18 turns, 1.3 $\mu$ H. Use  $50\Omega$  termination and set function generator at 10Vpp.
  1. Compute and compare loss at 7MHz and 14MHz.
  2. From  $A_l = 5nH/turn^2$ , calculate  $L_7$  and  $L_8$ .
  3. What is the spur strength at 7, 14 and 28MHz? Measure and calculate.



# Exercise 14: Norcal IF Cohn Filter

- $X_1$  through  $X_4$  are 4.91 MHz
- $C_{10}$ ,  $C_{11}$ ,  $C_{12}$  are 270 pF
- Set function generator to 50mV<sub>pp</sub> from function generator
- Calculate R and X for filter
  - 1. Measure the resonant frequency of one of the crystals
  - 2. Calculate the parameters of the crystal



Equivalent circuit for  
crystal and generator

# Transformers

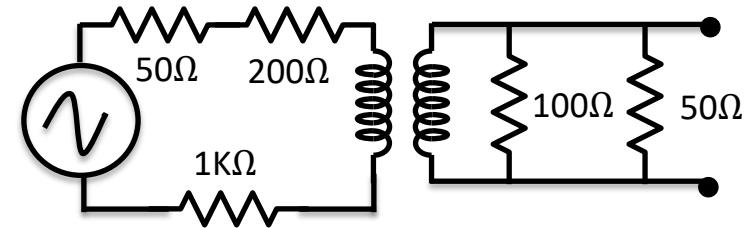
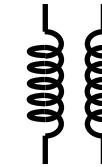
- For solenoid,  $\oint B \cdot ds = \mu_0 n I$  inside
- $LI = \Phi_B$ . Since there are  $n$  turns in the solenoid, over the solenoid,  $LI = \mu_0 n^2 I$ , so  $L = \mu_0 n^2$ .
- This is the source of  $L = A_l n^2$
- $V_s = \frac{N_s}{N_p} V_p$
- $Z_p = \left(\frac{N_p}{N_s}\right)^2 Z_s$

# Exercise 15: Norcal Driver Transformers

- $T_1$  is driver matcher uses FT 37-43

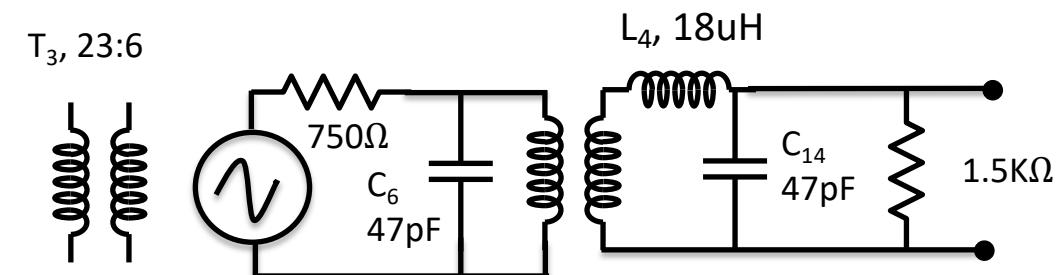
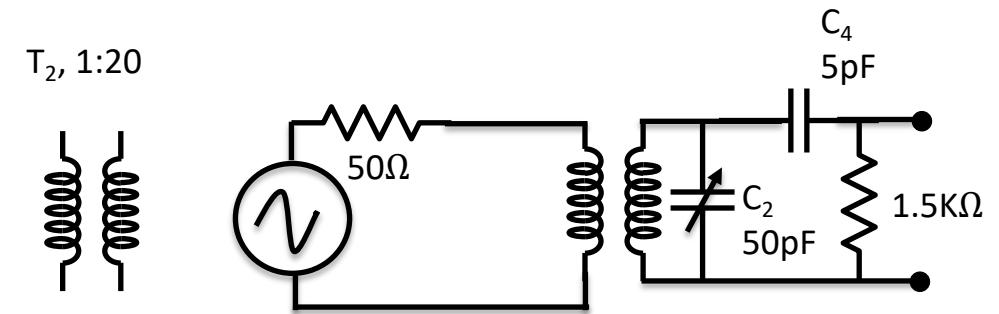
1. Measure the output V.
2. Calculate V
3. Measure the 3dB cutoff,  $f_c$ .
4. Use  $f_c$  to calculate  $A_t$

$T_1$ , 14:4



# Exercise 16: Norcal Tuned Transformers

- $T_2, T_3$  are IF matchers using FT 37-61
  - .5Vpp sine at 7MHz
1. Measure 3dB bandwidth
  2. Find  $P/P_+$



# Acoustics

- $P = P_0 + P_e, \rho = \rho_0 + \rho_e$
- 1. Gas moves and changes density: Displacement of undisturbed air is  $x$ . At time  $t$ , it's at  $x + \chi(x, t)$ , so  $\rho_0 \Delta x = \rho(x + \Delta x + \chi(x + \Delta x, t) - x - \chi(x, t))$ , or  $\rho_0 \Delta x = \rho(\frac{\partial \chi(x, t)}{\partial x} \Delta x + \Delta x)$ . So,  $\rho_e = -\rho_0 \frac{\partial \chi}{\partial x}$
- 2. Change in density causes change in pressure:  $P = f(\rho), P_0 + P_e = f(\rho_0 + \rho_e) = f(\rho_0) + \rho_e f'(\rho_0), f'(\rho_0) = \kappa = (\frac{dP}{d\rho})_0$ , or  $P_e = \kappa \rho_e$
- 3. Pressure differences cause motion:  $P(x, t) - P(x + \Delta x, t) = -\frac{\partial P_e}{\partial x} \Delta x$ , Newton's law gives  $\rho_0 \frac{\partial^2 \chi}{\partial t^2} = -\frac{\partial P_e}{\partial x} = -\kappa \frac{\partial \rho_e}{\partial x}$
- Substituting (1) into (3) gives  $\frac{\partial^2 \chi}{\partial t^2} = \kappa \frac{\partial^2 \chi}{\partial x^2}$ , put  $\kappa = \frac{1}{c_s^2}$
- Solution is  $\chi(x, t) = f(x - vt)$  [Different  $f$  than above].
- To find,  $\kappa = (\frac{dP}{d\rho})_0$ , note that the flow is adiabatic so  $PV^\gamma = C'$  and  $\rho$  varies inversely as  $V$ , so  $P = \rho^\gamma C$ , and finally, using  $PV = Nkt$ ,  $\kappa = (\frac{dP}{d\rho})_0 = \frac{\gamma k T}{n}$
- $L_p = 20 \log(\frac{P}{P_0}), P_0 = 20 \mu Pa$

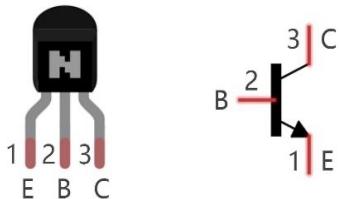
Sound	$L_p$	Power density
rustling leaves	10dB	1pW/m <sup>2</sup>
broadcast studio	20dB	1pW/m <sup>2</sup>
classroom	50dB	10nW/m <sup>2</sup>
heavy truck	90dB	1nW/m <sup>2</sup>
Shout at 1m	100dB	10mW/m <sup>2</sup>
jackhammer	110db	100mW/m <sup>2</sup>
jet takeoff at 50m	120dB	1W/m <sup>2</sup>

# Exercise 17: Tuned Speaker

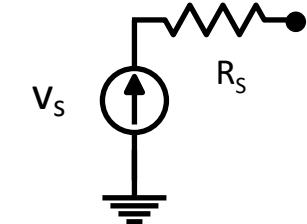
- Connect speaker to function generator 600Hz, 25mVrms.
1. Sound peaks at resonance. Find resonant frequency  $L_p$ .
  2. Measure  $f_l$ ,  $f_u$  by noting the 3dB loss. Calculate Q.
  3. Use voltmeter to find resonance with speaker (nominally 8ohm) to calculate impedance
  4. Calculate the resonant frequency from a transmission line equivalent circuit.

# Bipolar Transistors

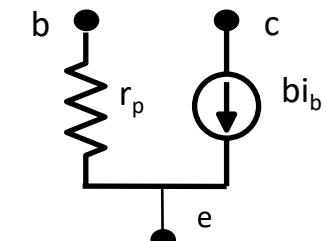
- NPN, PNP types
- Model
  - $V_f \approx .7V, V_s \approx .2V$
  - Conducts when  $V_{be} > V_f$
  - $i_c = \beta i_b$
  - $i_c = \alpha i_e$
  - $\beta = \frac{\alpha}{1-\alpha}$
  - $\beta \approx 100, \beta_r \approx 10$
- Switch
  - $G_s = \frac{i_b}{15mV}$
  - $R_S = 2\Omega$



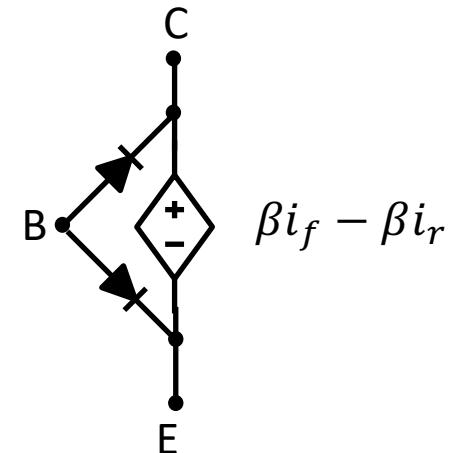
Region	$V_{be}$	$V_{bc}$	$V_{ce}$	$i_c$
active	$V_f$	$< V_f$	$> V_s$	$\beta i_b$
reverse	$< V_f$	$V_f$	$< -V_s$	$-(\beta_r + 1)i_b$
on(sat)	$V_f$	$V_f$	$> -V_s$	$> -(\beta_r + 1)i_b$ $< \beta i_b$
off	$< V_f$	$< V_f$	*	0



Bipolar source model



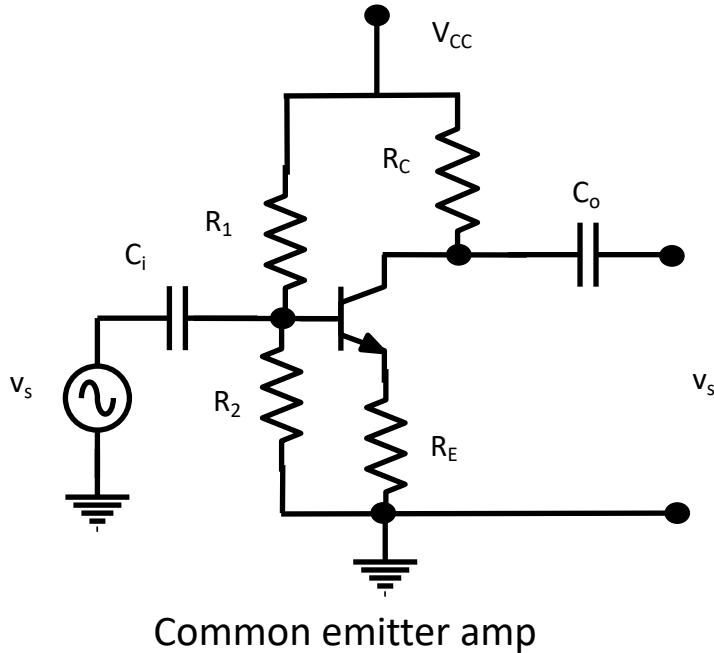
Bipolar equivalent circuit



Bipolar model

# BJT common emitter amplifier

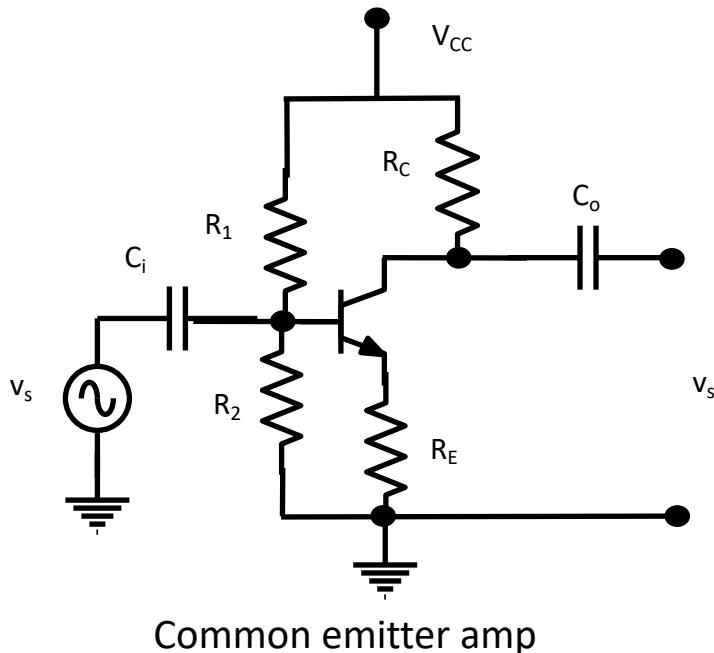
- Here's how to design a common emitter amplifier. We use a 2n3904 transistor with  $\beta=150$ . This circuit will work! Build it.



1. Pick the supply voltage  $V_{cc}=12V$ .
2. Choose a gain (amplification factor),  $A = 5$ .
3. Choose the “Q point” of the conducting transistor (4mA).
4.  $V_{cc} = (i_c \cdot R_C) + V_{ce} + i_e R_E \sim i_e \cdot (R_C + R_E) + V_{ce}$  with  $i_c=4\text{mA}$ . We get  $(R_C + R_E) = (V_{cc} - V_{ce})/(4\text{mA}) = 1.75 \text{ k}\Omega$ .
5. Since  $A = 5$  and  $A=R_C/R_E$ ,  $R_C= 5 R_E$  so  $R_E \sim 270 \Omega$  (this is a standard resistor value) and  $R_C= 1.5\text{k}\Omega$ .

Credit: Ward, Hands on Radio.

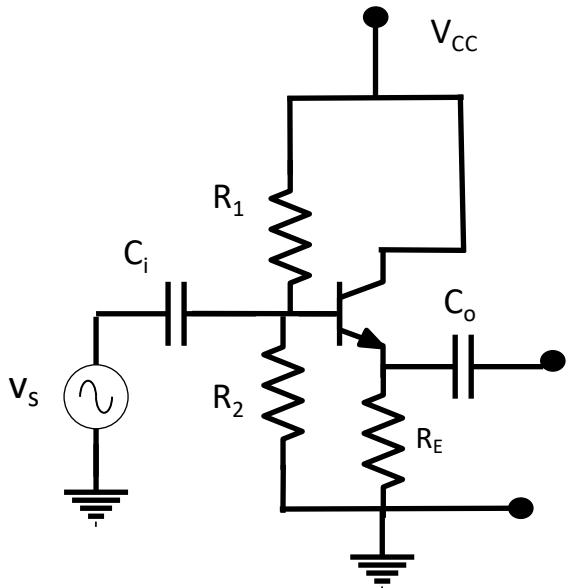
# BJT common emitter amplifier continued



6.  $i_b = 4\text{mA}/\beta = 27 \mu\text{A}$ .
  7. Since  $V_{be}$  must be greater than .7V throughout the input signal range, we want the voltage across  $R_2$  to satisfy  $V_{be} + i_c R_E = 1.8\text{V}$ .
  8. We insert a voltage divider consisting of  $R_1$  and  $R_2$ , so that  $R_1 = (12-1.8)/270 \mu\text{A} \sim 39 \text{k}\Omega$ .
  9.  $C_o$  and  $C_i$  are picked to offer small resistance to the frequency range we're interested in and  $C_o = C_i = 5 \mu\text{F}$ .
- I haven't explained why we want  $R_E$  but it provides thermal stability for the transistor over the range we care about. The fact that  $A=R_C/R_E$  can be calculated using Kirchhoff's laws.

Credit: Ward, Hands on Radio.

# BJT common collector amplifier

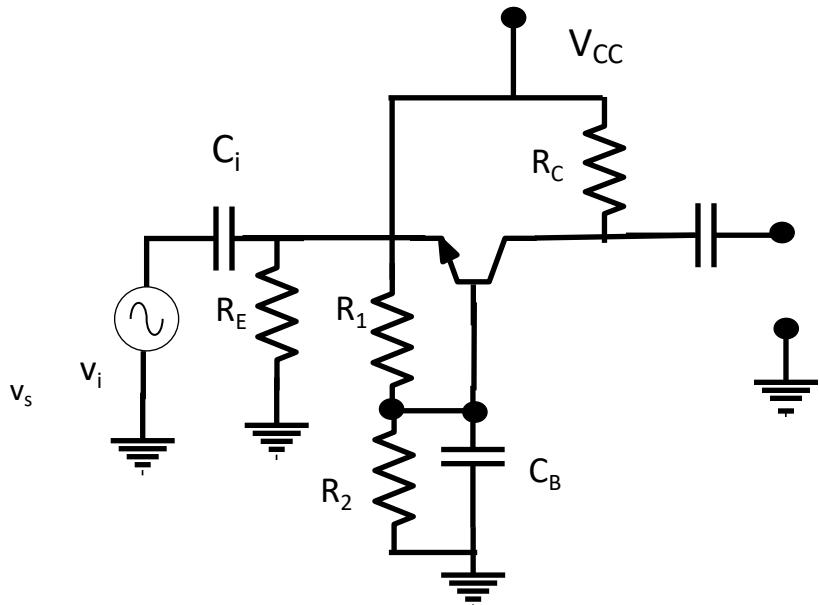


1.  $\beta = 150, A_V = 1, V_{cc} = 12V$
2. Q-pt:  $i_{ce} = 5mA, V_{ce,q} = 6V$  (rule of thumb),  $v_{be} = .7V$ .
3.  $i_{R_1 \rightarrow R_2} = 10i_b$  (ROT),  $V_{ce} = v_{be} + i_{ce,q}R_E, R_E = 1.2k\Omega, i_b = \frac{V_{ce,q}}{\beta} = 33\mu A$
4.  $V_{R_2} = V_{be} + i_C R_E = 6.7V, V_{R_1} = 5.3V$
5.  $R_2 = \frac{6.7}{330\mu A} = 20k\Omega, R_1 = \frac{5.3}{330\mu A} = 16k\Omega$
6.  $Z_{in} = R_1 || R_2 || (\beta + 1)R_E, R_{in} = 50\Omega, Z_{out} = 5\Omega$

Common collector amp (Emitter Follower)

Credit: Ward, Hands on Radio.

# BJT common base amplifier

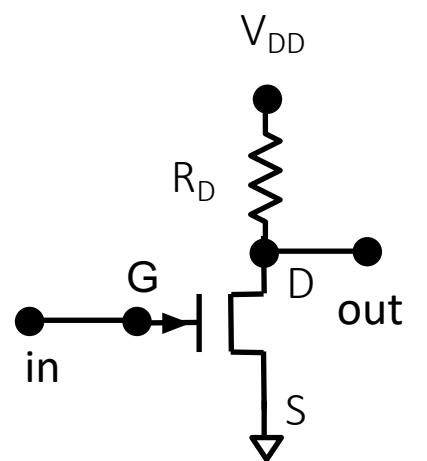


- $A_I = \frac{i_C}{i_E} = \frac{\beta}{\beta+1}, A_V = \frac{R_C || R_L}{r_e}, Z_{out} \approx R_C$
- 1.  $V_{CC} = 12V, V_{be} = .7V, R_E = 50\Omega, R_L = 1k\Omega, i_{ce,q} = 5mA, V_{ce,q} = 6V$
- 2.  $i_b = \frac{i_{ce,q}}{\beta} = 33\mu A, i_{R_1 \rightarrow R_2} = 10 i_b = 330\mu A$  (ROT)
- 3.  $V_{R_2} = V_{be} + i_C R_E = 6.7V, V_{R_1} = 5.3V$
- 4.  $R_1 = \frac{5.3}{330\mu A} = 16k\Omega, R_C = \frac{V_{cc}-i_{c,Q}R_E-V_{ce,Q}}{i_{c,Q}} = 1.35k\Omega$
- 5.  $A_V = \frac{R_C || R_L}{26/i_e} = 115$

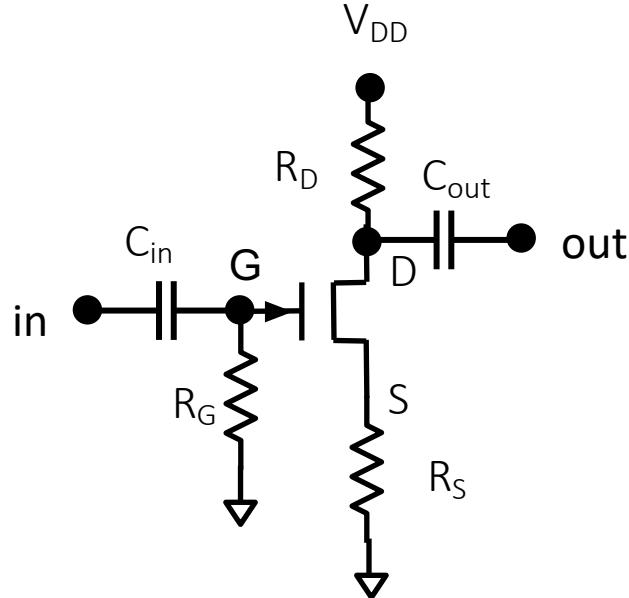
Common base amp

# JFETs

- JFET circuit model:  $I_{ds} = V_{ds} \left( \frac{2I_{dss}}{V_c^2} \right) \left( V_{gs} - V_c - \frac{V_{ds}}{2} \right)$
- $g_m = \frac{\Delta i_{ds}}{\Delta v_{gs}}$
- For circuit on right,  $g_m \Delta v_{gs} = \Delta i_{ds}$  and  $R_D \Delta i_{ds} = V_{out}$ , so  $-g_m R_D \Delta v_{gs} = V_{out}$
- Similar model for MOSFETs
- Op amp:  $V_{out} = A_{OL}(V_+ - V_-)$ , input resistance is very high



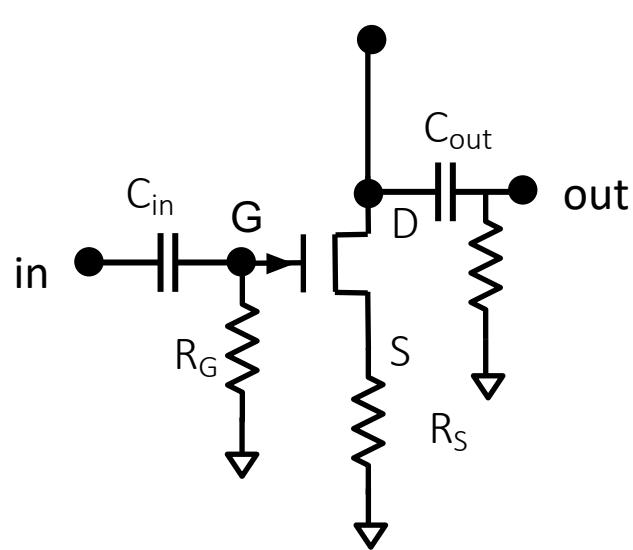
# JFET common source amplifier



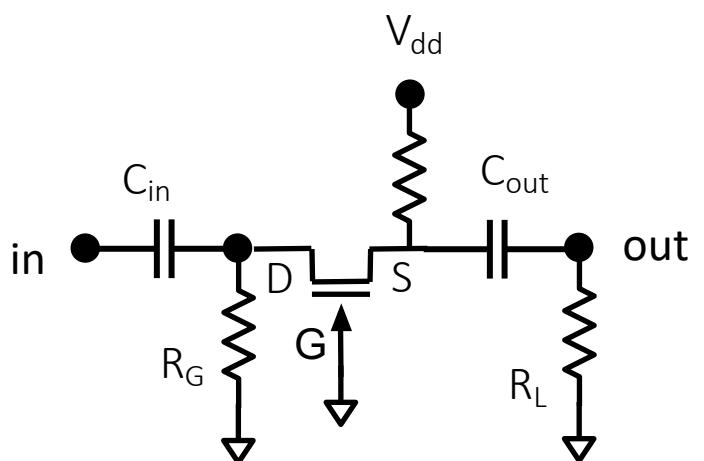
- $A_V = \frac{g_m R_D}{1 + g_m R_S} = -\frac{R_D}{R_L}, R_S = \frac{-V_P}{i_{dd}} \left(1 - \sqrt{\frac{i_{dd}}{i_{dss}}}\right)$ .  $g_m \approx 15mA/V$
- 1.  $V_{dd} = 12V, i_{dss} = 35mA, V_P = 3.0V, A_V = 10, i_{dd} = 10mA$
- 2. From equation above,  $R_S = 139\Omega, R_D = 10R_S = 1390\Omega$
- 3.  $A_V = -g_m(R_D || R_L)$

# JFET common drain amplifier

- Similar to BJT emitter follower

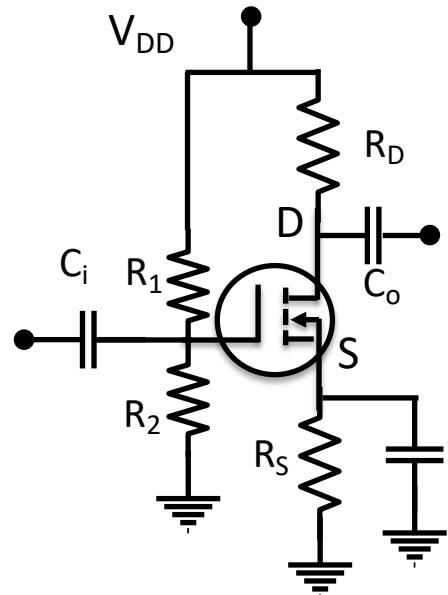


# JFET common gate amplifier



- $V_{DD} = 12V, i_{dss} = 60mA, V_P = -6, A_V = 10, R_L = 1k\Omega, R_S = 50\Omega$
  - $i_{d,q} = \frac{V_P}{2R_S^2 i_{dss}} \left( V_P + \sqrt{V_P^2 - 4R_S i_{dss} V_P} \right) - \frac{V_P}{R_S}$
1. Solve for  $R_D$ :  $10 = g_m \times R_D || R_L, R_D = 2k\Omega$
  2. Find  $i_{d,q} = \frac{V_P}{2R_S^2 i_{dss}} \left( V_P + \sqrt{V_P^2 - 4R_S i_{dss} V_P} \right) - \frac{V_P}{R_S} = 10mA$

# CMOS common emitter amplifier

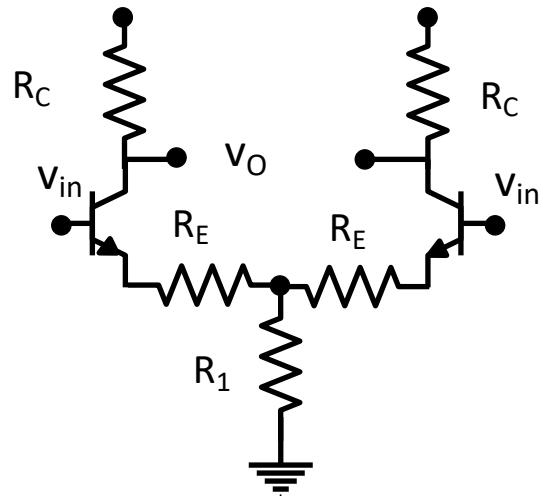


- Pick power
- $V_{DD} = i_D R_D + V_{DS} + i_D R_S$
- $V_{GS} = V_G - i_S R_S$
- $V_G = V_{DD} \frac{R_1}{R_1 + R_2}$
- $i_D = k(V_G - V_{TH})^2$
- Bias around  $\frac{V_{DD}}{3}$
- Pick gain,  $A = \frac{R_D}{R_S + \frac{1}{g_m}}$

# Differential Amplifier

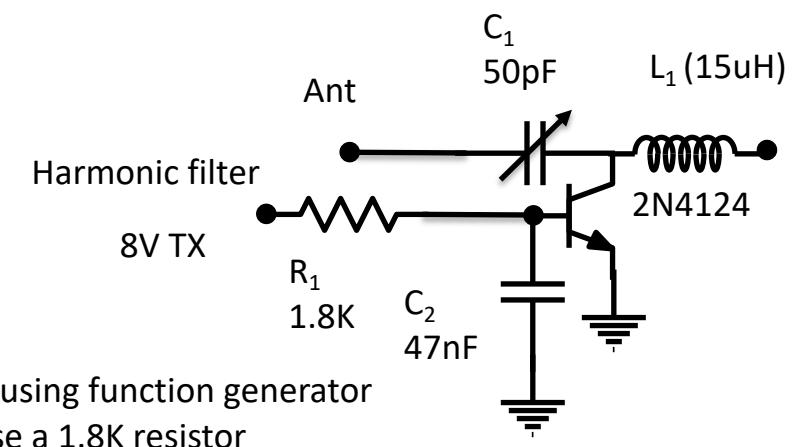
- Two port model
- $\begin{pmatrix} i_1 \\ i_2 \end{pmatrix} = \begin{bmatrix} y_{11} & y_{12} \\ y_{21} & y_{22} \end{bmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix}$
- Pick power  $\mp 12$
- Choose collector current ( $2mA$ ) by picking  $R_1$
- Pick gain,  $A = \frac{R_C}{2R_E}$

Differential amplifier



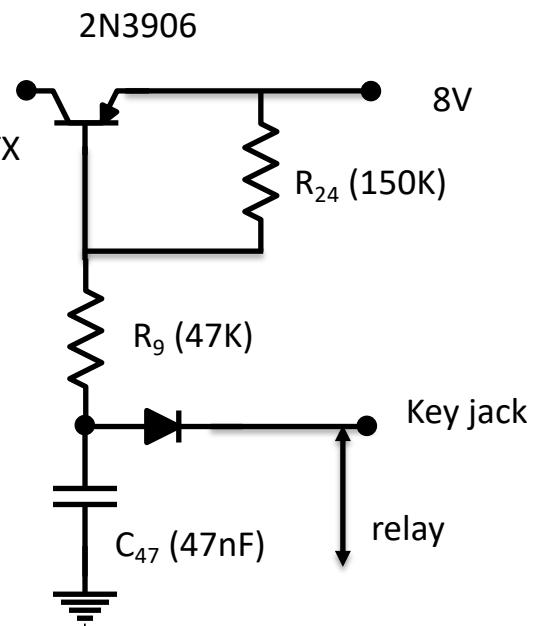
# Exercise 19: Norcal receiver switch

1. Consider the rising part of the base voltage waveform. Calculate slope.
2. Do the same for the falling part for voltage below .6V. Calculate  $t_2$ .
3. Measure switch attenuation
4. Measure the voltage with the switch on. Measure output voltage and calculate on-off rejection ratio  $R=20 \log(V_{\text{off}}/V_{\text{on}})$
5. Find the saturation resistance  $R_s$ .
6. Calculate the expected attenuation



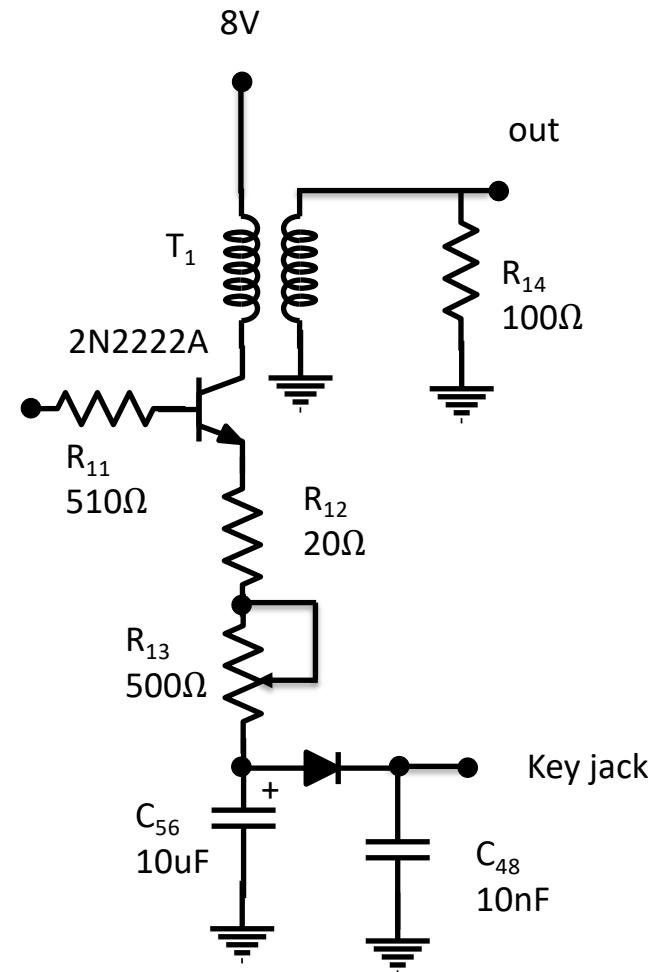
# Exercise 20: NorCal transmitter switch

1. Calculate voltage on  $C_{57}$ . Measure time for capacitor to charge half-way. Calculate what the time should be.
2. Calculate the approximate current  $i_c$  when Q4 is on. Assume base voltage on Q1 is 700 mV. Neglect saturation voltage on Q4. Calculate base current  $i_b$  required to produce this collector current assuming  $\beta = 100$ .
3. Calculate  $i_b$  at key down assuming a 700 mV drop in base-emitter of Q4 and at 600mV at D11
4. Sketch collector voltage at Q4 showing where transistor is saturated. What is the delay in going active?
5. Use the delay to measure  $\beta$ .



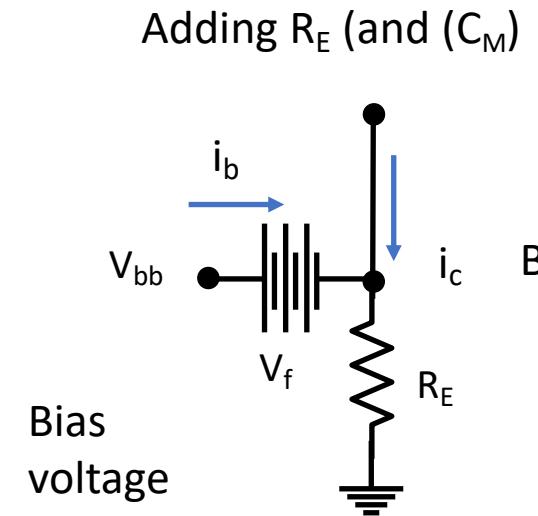
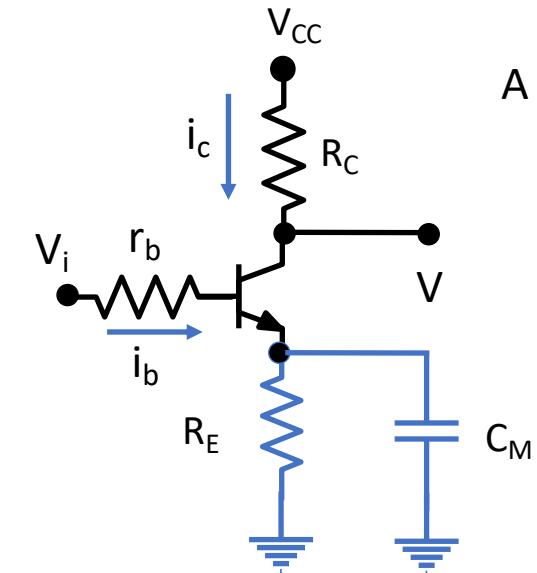
# Exercise 21: Norcal Driver

1. Measure the voltage gain  $G_v = \frac{v_o}{v_i}$  with R13 at minimum and maximum gain.
2. Calculate expected voltage gain at each setting.
3. 560ohm source resistance  $V_o = 2V$



# Emitter degeneration

- To the usual transistor circuit (A), on the right, we add  $R_E$ . (B) is an equivalent circuit.
- $V_{bb} \approx V_f + i_c R_E$ . Let  $V$  be the output AC and  $V_i$  be the input AC, then the gain is  $G = \frac{V}{V_i}$ .
- $V_i = i_b r_b + i_E R_E \approx i_C R_E$ ,  $Z_i = \frac{V_i}{i_b}$ ,
- $V = -i_c R_C$ . So  $G_v = -\frac{R_C}{R_E}$  (Doesn't depend on  $\beta$ ).
- $V_i \approx \beta i_b R_E$  and  $Z_i = \frac{V_i}{i_b}$ , so  $Z_i = \beta R_E$ .
- $C_M$  is called a Miller capacitor,  $i_m = j\omega(V_i - V) = j\omega C_M(1 + |G_v|)V_i$
- So with the Miller capacitor,  $Z_i = \beta R_E || (1 + |G_v|) C_M$
- $r_c \approx \frac{V_{early}}{i_c}$ ,  $r_c$  is the collector resistance
- $R_s' = R_s + r_b$ ,  $r_b$  is the base resistance
- $z_c = r_c || C_c$ ,  $C_c$  is specified in data sheet (8pF),  $z_c$  is the collector impedance
- $Z_o = \frac{V}{i_c}$ ,  $i = i_c - \beta i_b$ ,  $i_b = -\frac{i_c R_s}{R_s' + R_E}$ ,  $i = i_c(1 + \frac{\beta R_E}{R_s' + R_E})$
- $V = i z_c + i_c (R_s' || R_E)$
- $Z_o = \frac{V}{i_c} = z_c \left(1 + \frac{\beta R_E}{R_s' + R_E}\right) + R_s' || R_E$ .
- $|z_c| \gg R_E$ , so  $Z_o = z_c \left(1 + \frac{\beta R_E}{R_s' + R_E}\right)$

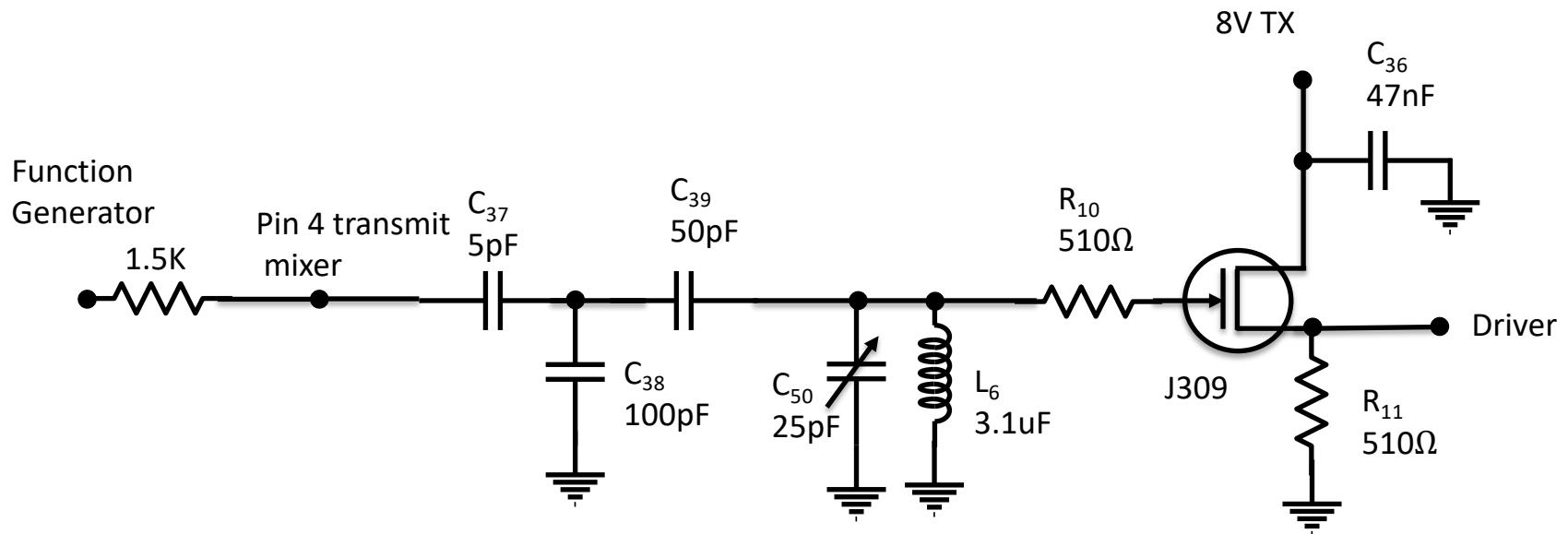


# Exercise 22: Emitter degeneration

- In Driver amplifier, add probe to  $R_{11}$ , this allows us to measure the AC voltage,  $V_i$ 
  1. Measure  $G_v = \frac{V}{V_i}$  with  $R_{13}$  turned fully counterclockwise
  2. Calculate the expected voltage gain for each setting
  3. Measure  $V_i$  at the maximal gain setting
  4. The open circuit voltage is  $V_0 = 2V$ , calculate  $V_i$  in terms of  $C_M$

# Exercise 23: Norcal Buffer amplifier

1. Measure the DC voltage at source of the JFET
2. Calculate the source and drain voltages you should expect
3. Measure the voltage gain
4. Find the transconductance you should expect
5. Calculate the available power  $P_+$  from the function generator through a 1.5Kohm load. Calculate gain in in dB

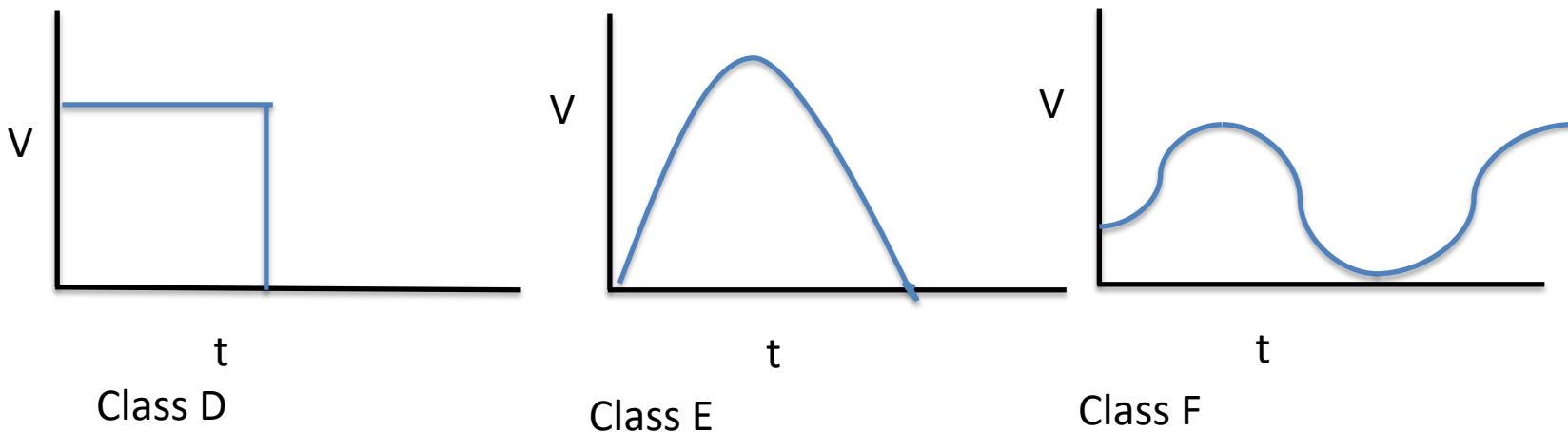
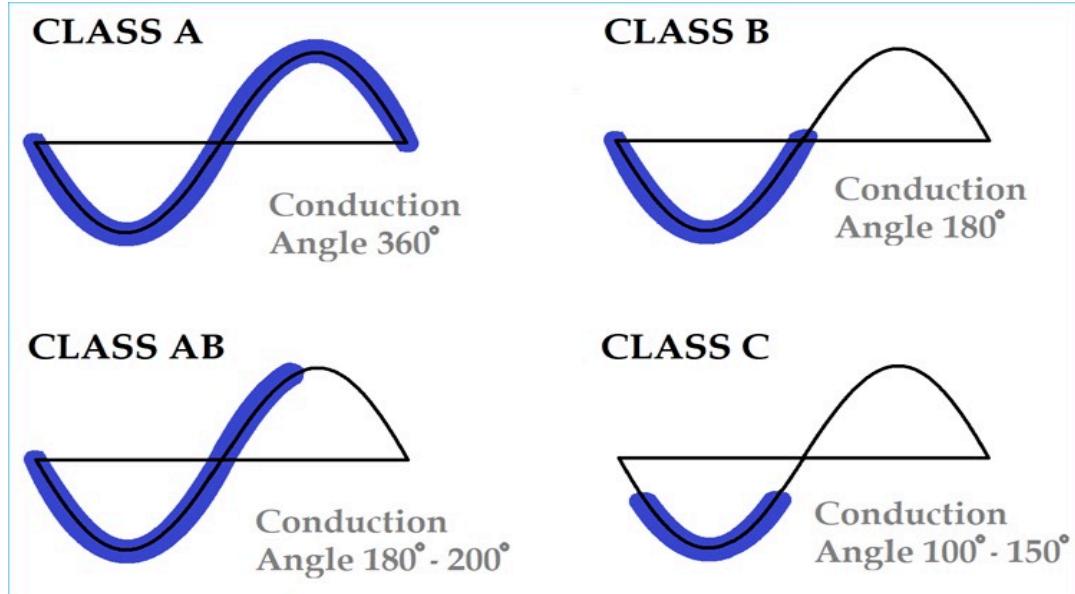


# Amplifier classes

Class	Efficiency	Characteristics
A	35%	Full bias
B	60%	Low bias
C	75%	Saturating
D	75%	Switch in pass-band
E	90%	Voltage switch
F	80%	Harmonic resonators

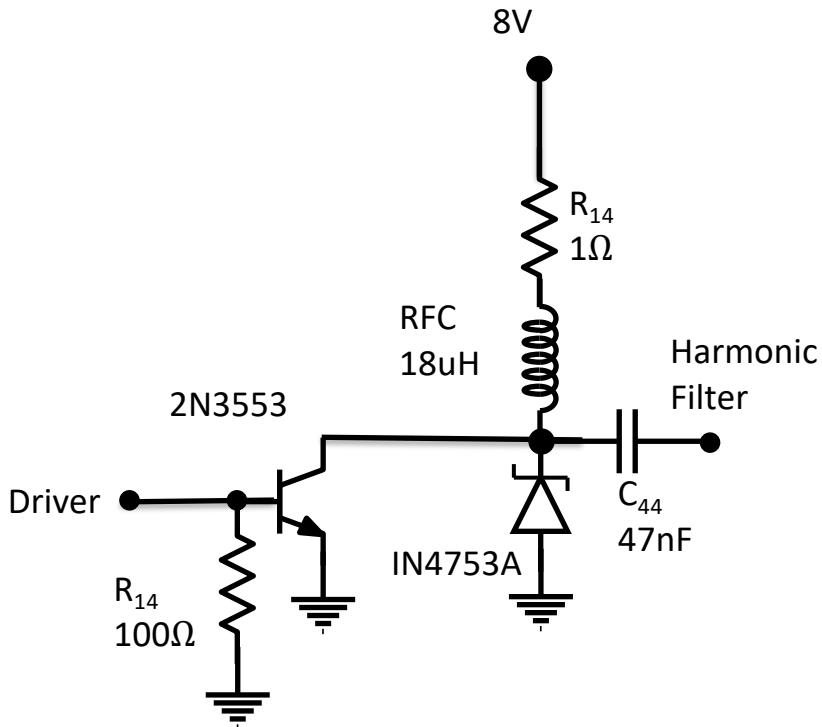
$$\eta = \frac{P}{P_0}, P_d = P_0 - P_i$$

$$P_d = P_a + P_{on}$$

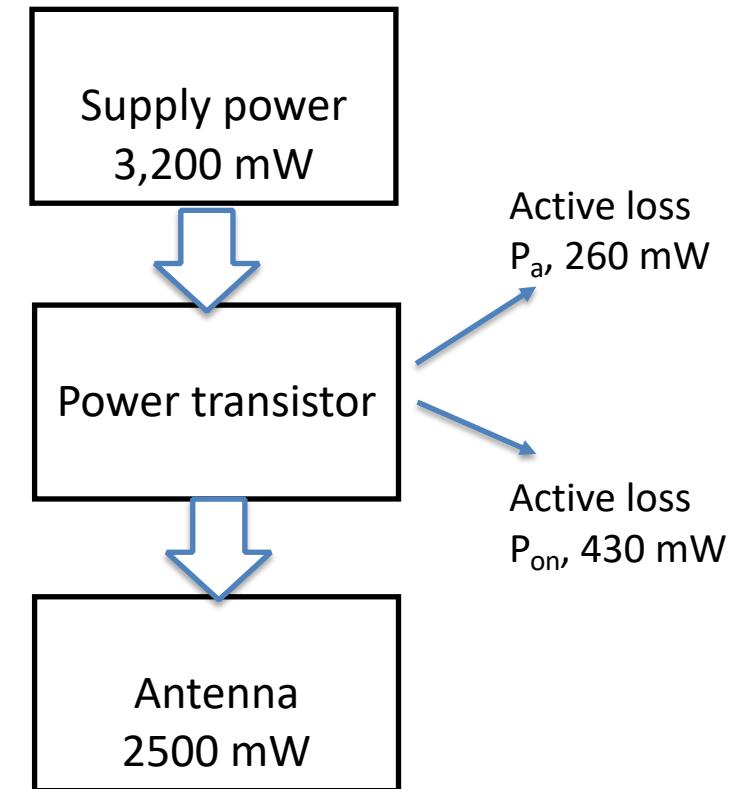


# Exercise 24: Norcal Power Amp

Norcal-40 Power amp is class C



- $R_t = \frac{T-T_0}{P_d}$
  - $T_0$  is ambient temperature,  $T$  is heat sink temperature
1. Calculate pp across 50ohm load required for output of 2W
  2. Find pp output voltages of 5, 10, 15, 20, 25 and 30V. Calculate power supply current subtracting 2mA for regulator
  3. Plot efficiency  $\eta = \frac{P}{P_0}$ . Plot dissipated power  $P_d = P_o - P_i$



# Thermal modelling

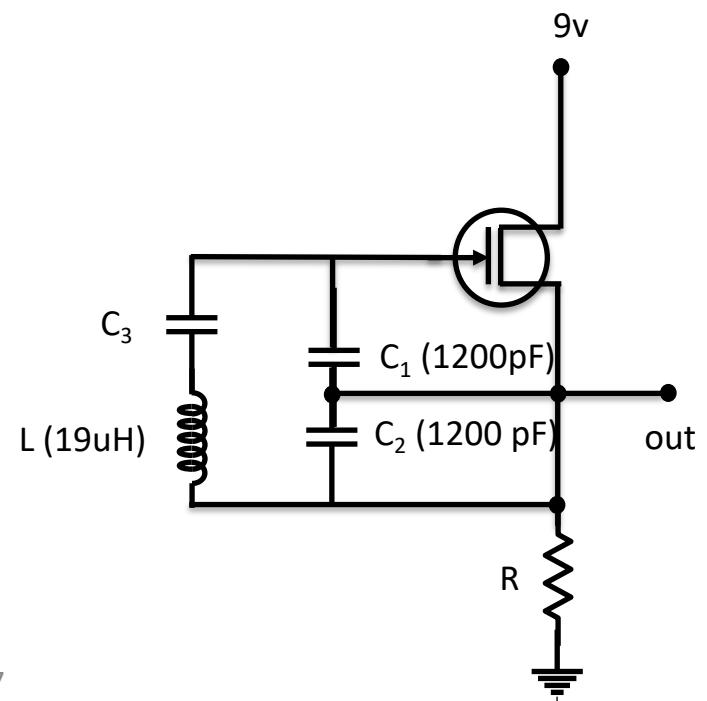
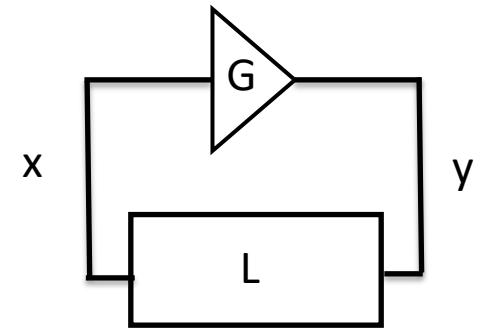
- $T$  is heat sink temperature,  $T_0$  is ambient temperature,  $P_d$  is power dissipated.
- $R_t = \frac{T-T_0}{P_d}$ ,  $R_t$  is the thermal resistance
- $C_t \dot{T} = P_d$ ,  $C_t$  is the thermal capacitance
- $R_j = \frac{T_j-T}{P_d}$ ,  $T_j$  is the junction temperature
- $f(t) + \tau f'(t) = f_\infty$ ,  $f(t) = f_0 e^{-\frac{t}{\tau}}$
- $P_d = \frac{T(t)-T_0}{R_t} + C_t T'(t)$ ,  $\tau = C_t R_t$ ,  $T_\infty = P_d R_t + T_0$
- $T(t) + \tau T'(t) = T_\infty$ ,  $\tau = C_t R_t$ .
- $T_\infty = P_d R_t + T_0$
- $T(t) = T_\infty - P_d R_t e^{-\frac{t}{\tau}}$
- $T_j(t) = T(t) + R_j P_d$

# Exercise 25: Thermal modelling

- For Motorola 2N3553,  $T_j = 25^\circ C/W$ 
  1. Measure ambient temperature
  2. Turn function generator until output is 30Vpp
  3. After 20 minutes, measure  $T_\infty$ . Use this to calculate  $R_t$  and  $T_j$
  4. Plot heat sink temperature vs time. Measure  $t_2$  and calculate  $C_t$

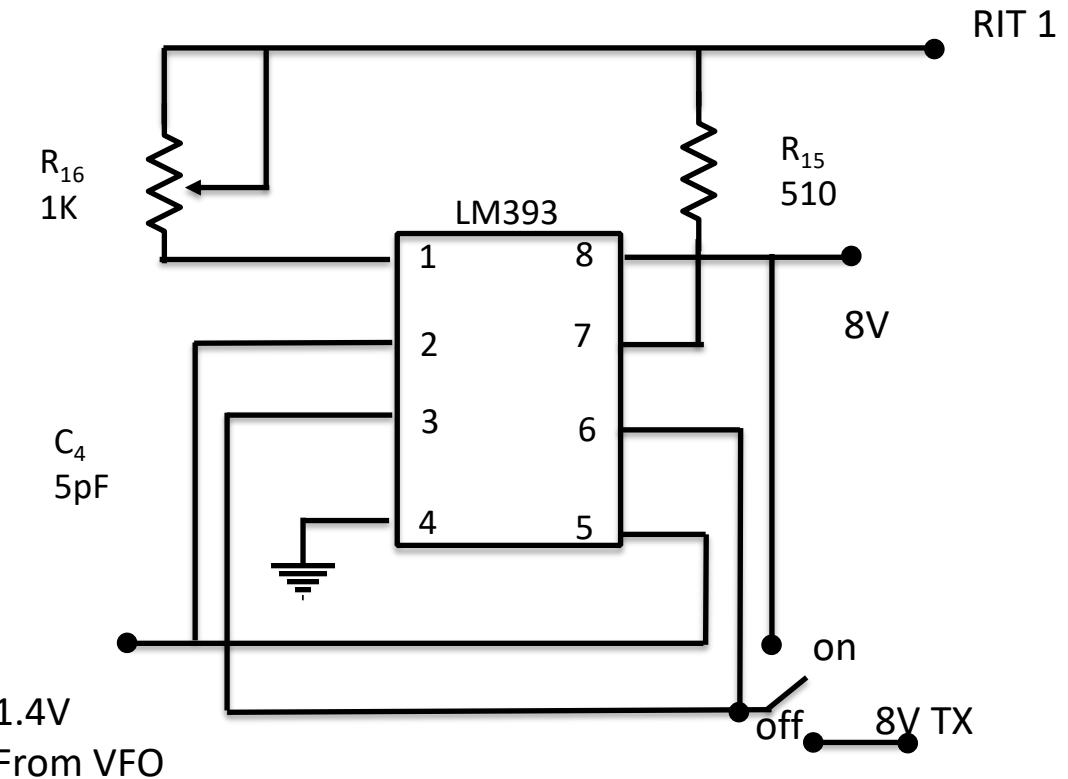
# Clapp oscillator

- Oscillation condition
  - $Gx = y$
  - $Ly=x$
  - $|G| = |L|$  and  $\angle G = \angle L$
- Clapp (circuit on right)
  - $i_d = g_m v_{gs}$
  - Resonance:  $-\frac{1}{j\omega_0 C_2} = j\omega_0 L + \frac{1}{j\omega_0 C_3} + \frac{1}{j\omega_0 C_1}$
  - $\omega_0 = \frac{1}{\sqrt{LC}}$ ,  $C = C_1 || C_2 || C_3$
  - At resonance,  $v_{gs} = R i_d \frac{C_1}{C_2}$ ,  $L = \frac{C_1}{R C_2}$
  - Oscillation continues if  $g_m > \frac{C_1}{R C_2}$
  - $v_{gs} = 2v_s$



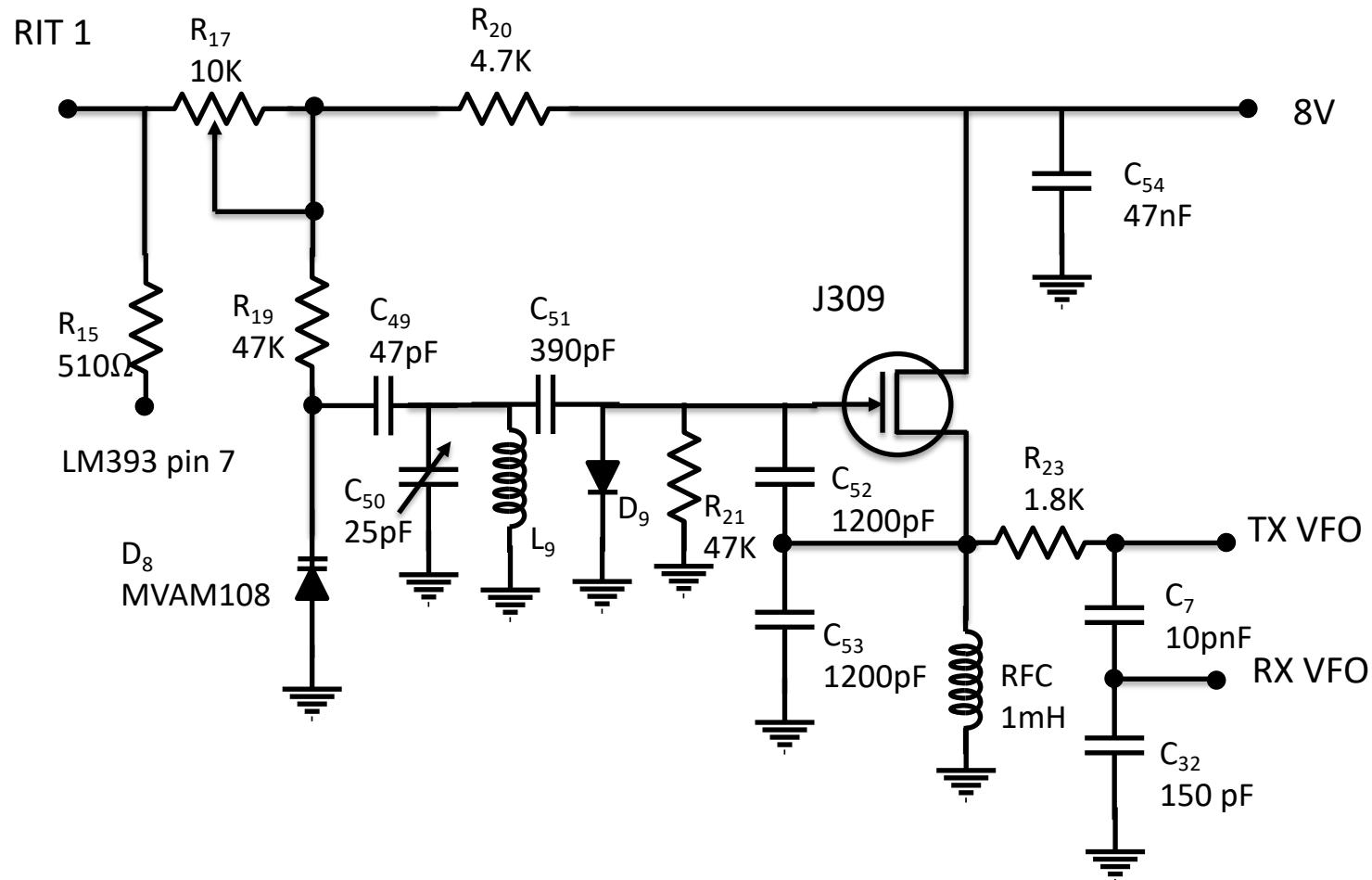
# Norcal Receiver Incremental Tuning (RIT)

- LM393 is a comparator



# Exercise 26: Norcal VFO

- L9: T68-7 62 turns
  - Check MVAM108 capacitor when  $R_{17}$  is high and low
  - Start resistor ( $R_{21}$ ) pulls gate to ground at start
  - When gain limiting diode (D9) conducts, it pulls gate negative
  - Oscillator keeps growing as long as  $g_m > 1/R$
1. Measure DC voltage across wiper in  $R_{17}$
  2. Calculate expected V for large signal oscillation
  3. How does this change if we consider the inductor and source-drain resistance
  4. How does the frequency change as  $R_{17}$  changes?
  5. Calculate the oscillation frequency and the loss ratio  $|V/V_1|$

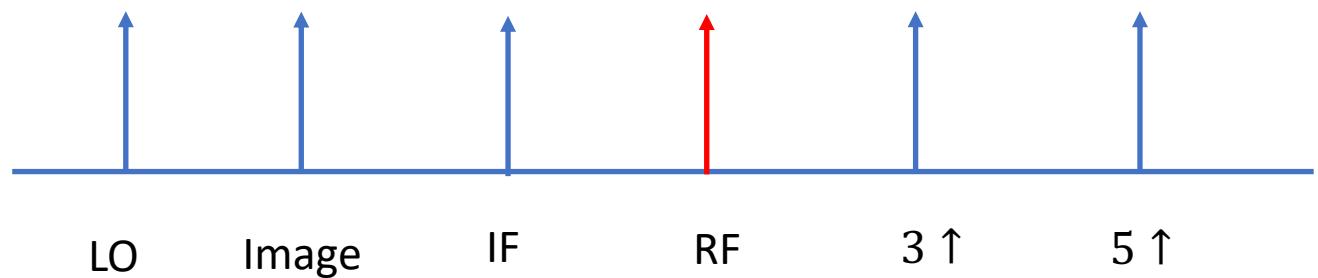


# Exercise 27: Gain limiting

1. Measure the voltage,  $V$ , on R23
2. In deriving the oscillation condition, we neglected the inductor resistance and drain source resistance,  $r_d$ . How does this affect the conditions. L9 has a Q of 250 and  $r_d = 5k\Omega$ , now what is the predicted  $V$ .
3. Find the loss ratio  $|\frac{V}{V_i}|$  and calculate what it should be.
4. Measure the temperature dependence of the VFO
5. How much does the temperature have to change to cause a 100Hz shift?
6. What is the oscillation change if we remove one turn of the inductor
7. What is the RIT tuning range?

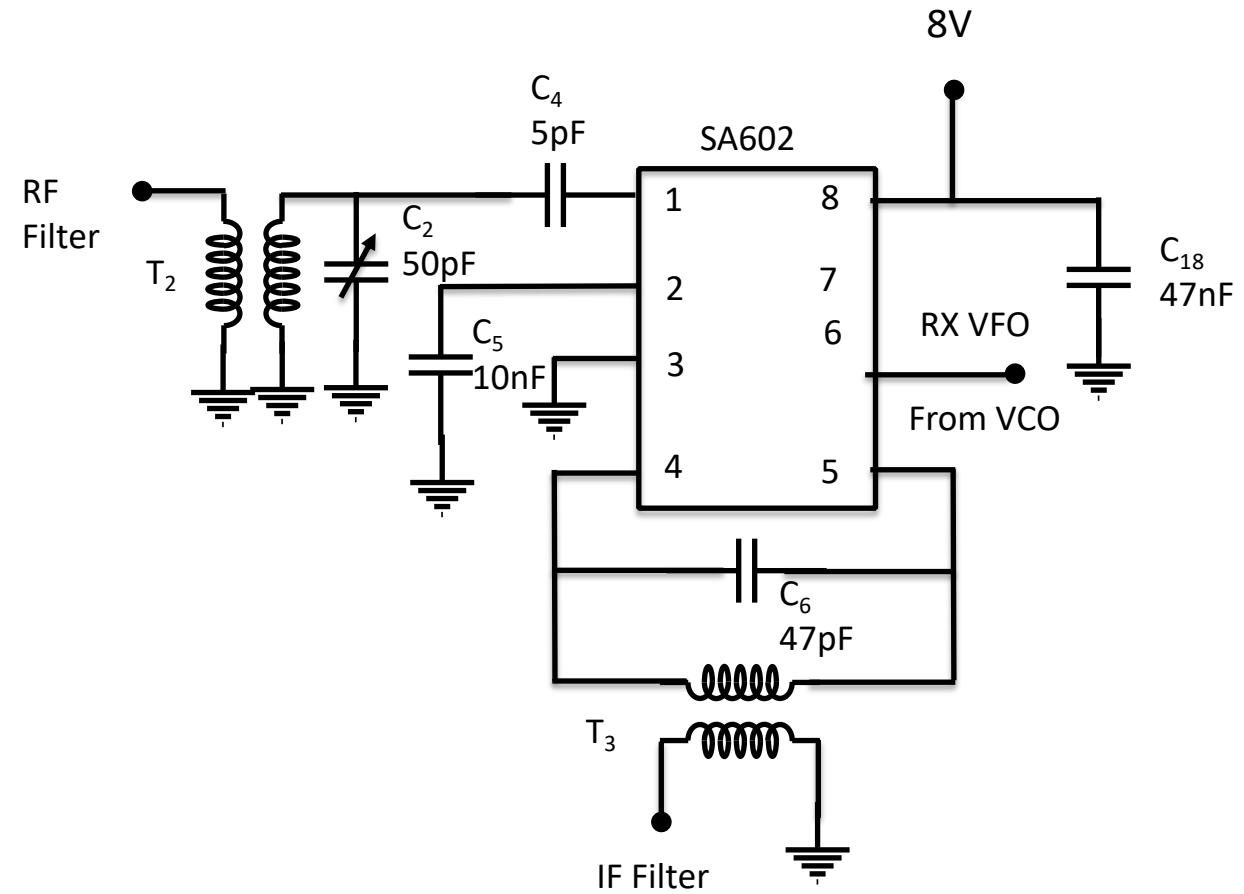
# Mixers

- $V_{lo}(t)$  is a square wave with period  $\omega_{lo}$ . Expanding this in a Fourier series, we get:
- $V_{lo}(t) = \frac{4}{\pi}(\cos(\omega_{lo}t) - \frac{\cos(3\omega_{lo}t)}{3} + \frac{\cos(5\omega_{lo}t)}{5} \dots), V_{rf}(t) = V_{rf}\cos(\omega_{rf}t)$
- $V_{lo}(t)V_{rf}(t) = \frac{2V_{rf}}{\pi}(\cos(\omega_-t) - \frac{\cos(3\omega_-t)}{3} + \frac{\cos(5\omega_-t)}{5} \dots) + \frac{2V_{rf}}{\pi}(\cos(\omega_+t) - \frac{\cos(3\omega_+t)}{3} + \frac{\cos(5\omega_+t)}{5} \dots)$
- $\omega_+ = \omega_{lo} + \omega_{rf}$  and  $\omega_- = |\omega_{lo} - \omega_{rf}|$
- We define  $\omega_{k+} = (k\omega_{lo} + \omega_{rf})$  and  $\omega_{k-} = |k\omega_{lo} - \omega_{rf}|$  and  $V_{k+}(t) = \frac{2V_{rf}}{k\pi}\cos(\omega_{k+}t)$  and  $V_{k-}(t) = \frac{2V_{rf}}{k\pi}\cos(\omega_{k-}t)$
- $\omega_i = \omega_{if} - \omega_{lo}$  and  $\omega_{if} = \omega_{if} + \omega_i$ ,  $\omega_i$  is a spurious signal.  $\omega_{k+}$  and  $\omega_{k-}$  are the spurs from the  $k$ th harmonic



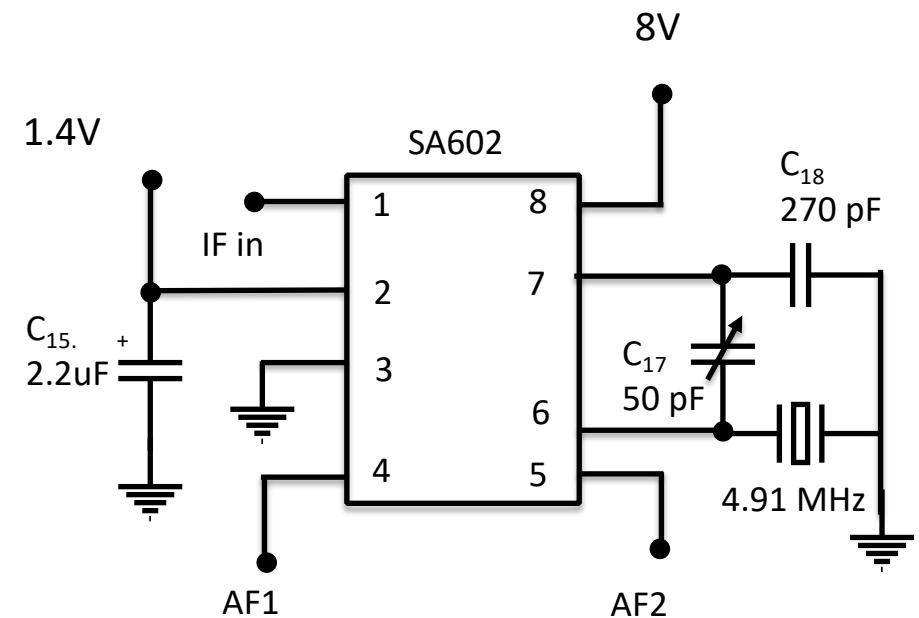
# Exercise 28: Norcal RF Mixer

1. Measure conversion gain of the Mixer.
2. How much attenuation is provided by pot?
3. By how many dB is the image response suppressed



# Exercise 29: Norcal Product Detector

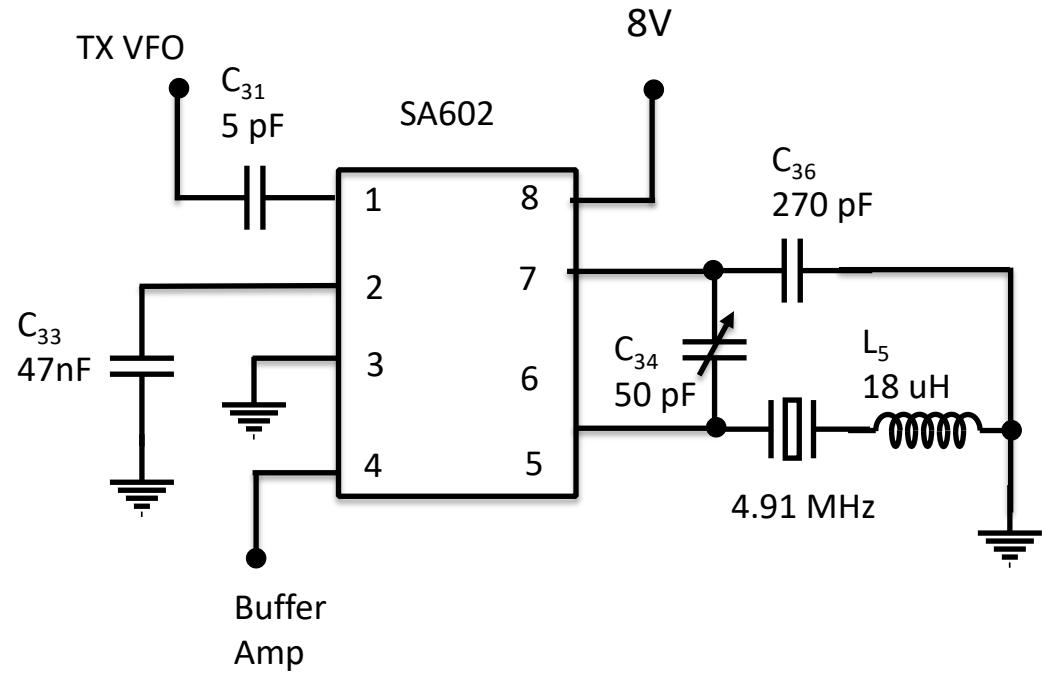
1. Adjust C17 for minimum oscillation frequency and record it
2. Calculate the minimum oscillation frequency you'd expect
3. Measure the temperature coefficient for the BFO
4. Measure the gain through the receiver from the antenna through the product detector
5. Find the f5 spur calculate the expected f3
6. By how much is the if spur suppressed



- 620 Hz output through AF1 and AF2

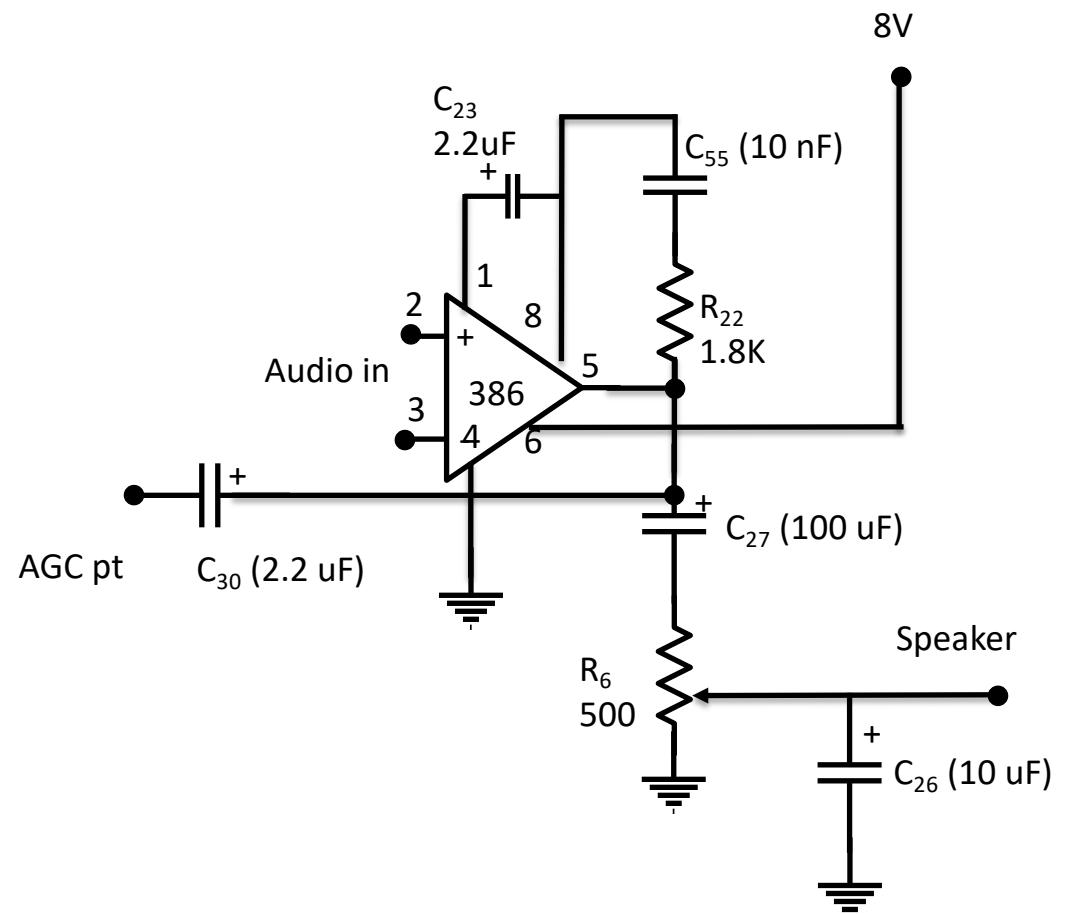
# Exercise 30: Norcal transmit mixer and oscillator

1. How much would you expect the inductor to lower the oscillation frequency
2. Use the TX VFO and the voltage attenuation to calculate the input power from the transmit mixer. Calculate the gain through the entire chain
3. Measure the rise and fall time of keying response
4. There is a spurious  $f_{mn} = mf_{vfo} + nf_{to}$ .



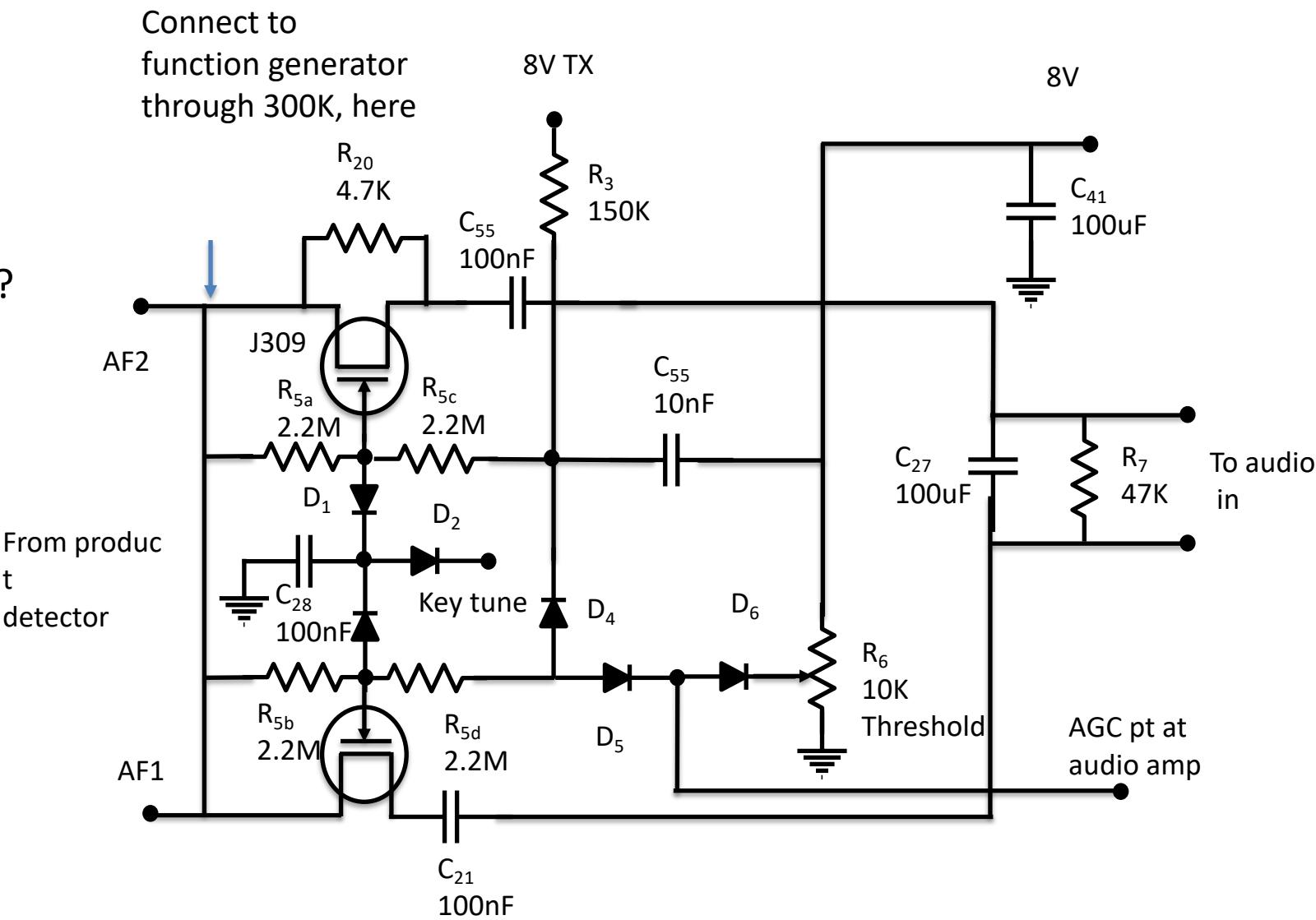
# Exercise 31: Norcal Audio Amp

1. Calculate input  $V_i$  assuming very high input impedance
2. Measure the voltage gain  $G_v$  at high frequency and 3dB rolloff

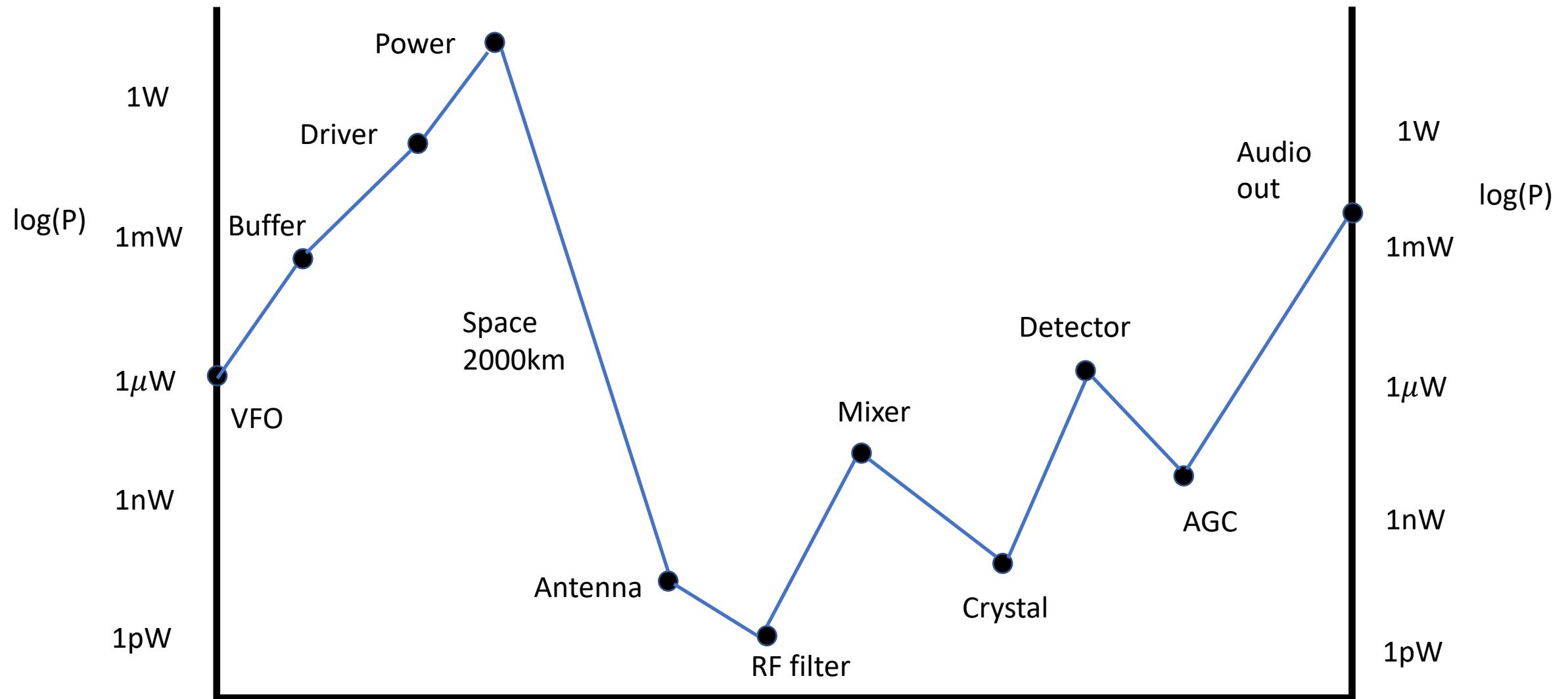


# Exercise 32: Norcal AGC

1. Plot audio output v dc control
2. What is the maximum control voltage we can measure? Infer cutoff voltage  $V_c$
3. What is the minimum control voltage?

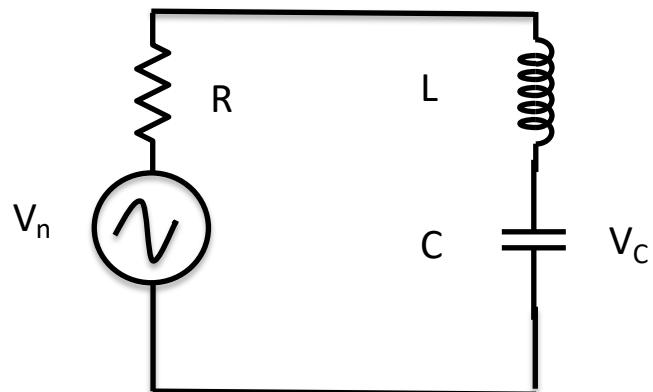


# NorCal power levels



# Noise

- $V_{n(rms)} = \sqrt{\frac{1}{\tau} \int_0^\tau V(t)^2 dt}$
- $P_n = \frac{V_{n(rms)}^2}{R}$ ,  $R$  is load resistance
- $SNR = \frac{P}{P_n}$
- $MDS = \frac{P_n}{G}$
- $P_n = NB$ ,  $N$  is noise power density,  $B$  is bandwidth
- $NEP = \frac{N}{G}$

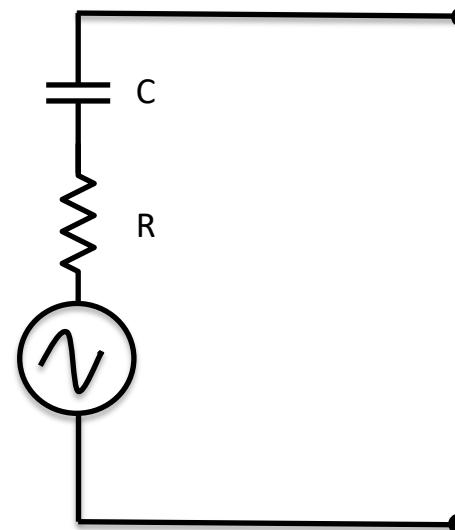


- Nyquist
  - $V_C = \frac{1}{j\omega C} \frac{V_n}{R+j\omega L+\frac{1}{j\omega C}}$
  - $|V_c|^2 = \frac{|V_n|^2}{|1-\omega^2 LC+j\omega RC|^2}$
  - Expected energy at resonance is  $kT = \frac{c}{2} \int_0^\infty |V_c|^2 df$ , by equipartition theorem
  - $\int_0^\infty \frac{1}{|1-\omega^2 LC+j\omega RC|^2} df = \frac{1}{4RC}$
  - So,  $|V_n|^2 = 8kTR$
  - $N = kT = \frac{|\frac{V_n|^2}{2R}|}{2R}$
  - $T_c = \frac{N}{k}$ ,  $T_n = \frac{NEP}{k}$ ,  $V_{rms} = \sqrt{4kTR}$

# Antennas

- From Maxwell, for a plane wave ( $E$  in  $x$  direction,  $H$  in  $y$  direction), wave is of form  $\exp(j\omega t - j\beta z)$
- $\nabla \times E = -j\mu_0 \omega H$
- $\nabla \times B = j\epsilon_0 \omega E$
- $\beta \hat{z} \times E = \mu_0 \omega H, \beta E_x \hat{y} = \mu_0 \omega H$
- Substituting and taking the restricted cross products, we get:  $\beta E_x = \omega \mu_0 \frac{\omega \epsilon_0}{\beta}$ , so  $\beta = \omega \sqrt{\mu_0 \epsilon_0}$
- Power density:  $S = \operatorname{Re} \left( \frac{E_x H_y}{2} \right) = \frac{(|E_x|)^2}{2\eta_0}$
- $\eta_0 = \frac{E_x}{H_y} = \sqrt{\frac{\mu_0}{\epsilon_0}} = 377 \Omega$
- Impedance:  $P_t = \frac{R|I|^2}{2}$ ,  $R$  is real part of  $Z$ ,  $R = R_r + R_l$ ,  $\eta = \frac{R_r}{R}$
- Power density for isotropic antenna:  $S_i = \frac{P_t}{4\pi r^2}$
- Define  $G(\theta, \phi) = \frac{S(\theta, \phi)}{S_r}$ .  $S(\theta, \phi)$  is just the Poynting vector

Receiving antenna Thevenin



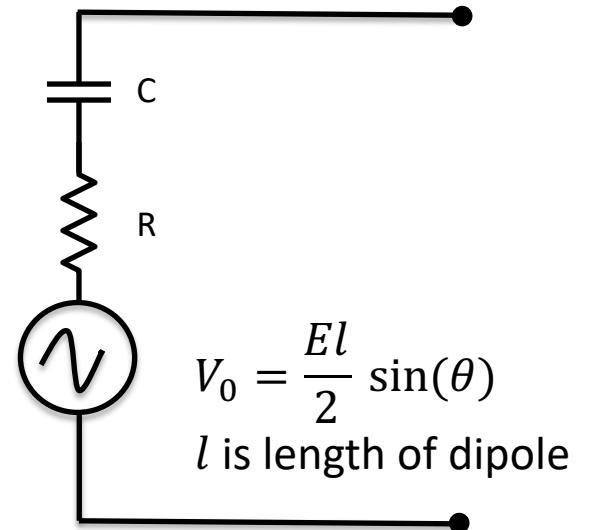
# Transmitting Antenna

- Define  $G(\theta, \phi) = \frac{S(\theta, \phi)}{S_r}$ .  $S(\theta, \phi)$  is just the Poynting vector
- For isotropic reference:  $S_i = \frac{P_t}{4\pi r^2}$ ,  $G = \frac{4\pi r^2 S}{P_t}$
- $\int G d\Omega = 4\pi$

# Receiving Antenna

- $V_0 = hE$ ,  $h$  is effective antenna length ( $h = \frac{l}{2}$  for short antenna)
- For dipole:  $V_0 = \frac{l}{2} E \sin(\theta)$
- $A(\theta, \phi) = \frac{P_r}{S(\theta, \phi)}$ . This is the definition of the effective area,  $A$ .
- By reciprocity,  $A(\theta, \phi) = \frac{\lambda^2}{4\pi} G(\theta, \phi)$
- $P_r = \frac{|V_0|^2}{8R_a} = \frac{|hE|^2}{8R_a}$ , so
- $P_r = \frac{h^2 S \eta_0}{4R}$
- $A = \frac{h^2 \eta_0}{4R}$

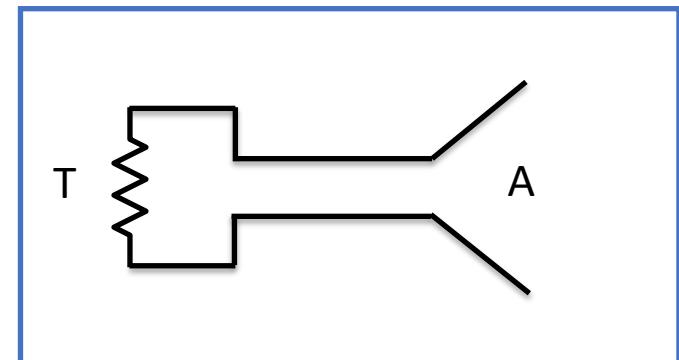
Dipole Thevenin equivalent circuit



# Friis, blackbody and Antenna Theorem

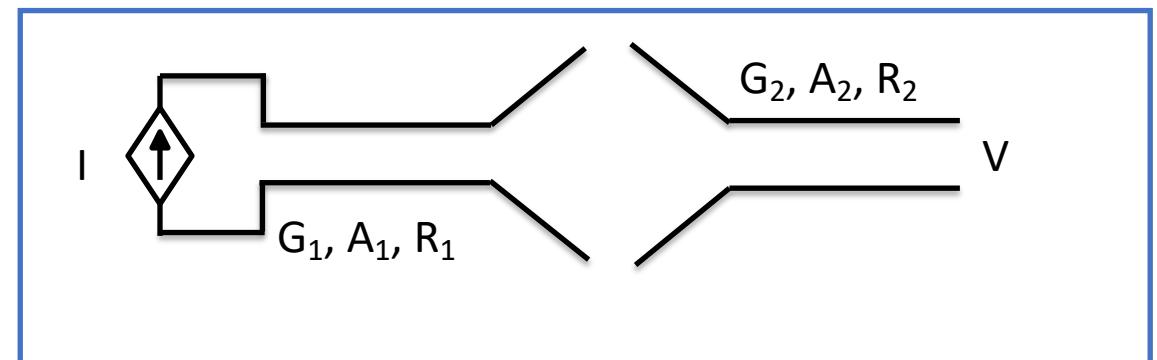
- For transmitting/receiving antenna pairs:  $G_1 A_2 = \frac{|V|^2 \pi r^2}{|I|^2 R_1 R_2} = G_2 A_1$ . So  $\frac{G_1}{A_1} = \frac{G_2}{A_2} = \frac{4\pi}{\lambda^2}$
- $S = \frac{P_t G}{4\pi r^2}$
- $P_r = SA = \frac{P_t G A}{4\pi r^2}$ . --- Friis radiation formula
- For us,  $G = 1, A = 150 \text{ m}^2, r = 2000 \text{ km}, P_t = 2 \text{ W}$
- $P_r = 6 \text{ pW}$
- Antenna theorem:  $\oint A d\Omega = \lambda^2$
- For cavity on right, T is constant at thermodynamic equilibrium and the same power is transmitted and emitted, the Johnson noise is  $kT$ . The energy received is
  - $E = \frac{4\pi kT}{c\lambda^2}$ .
  - Set  $B = \frac{kT}{\lambda^2}$ .
  - $kT = \oint BA d\Omega = \oint A \frac{kT}{\lambda^2} d\Omega$ , which gives the antenna theorem

Insulated cavity



# Reciprocity

- *Reciprocity:* The position of an ideal voltmeter and ideal current source can be interchanged without changing the voltmeter reading.
- $\frac{G}{A} = \frac{4\pi}{\lambda^2}$
- $\frac{G_1}{A_1} = \frac{G_2}{A_2}$



# Reciprocity and dipoles

- For dipole (Length:  $l = \frac{\lambda}{2}$ )
- $\lambda^2 = \int A d\Omega = \int \frac{h^2 \eta_0}{4R_r} d\Omega$ , so
- $R_r = \frac{l^2 \eta_0}{16\lambda^2} \int \sin^2(\theta) d\Omega = \eta_0 \frac{\pi}{6} (\frac{l}{\lambda})^2$
- $A = \frac{3\lambda^2}{8\pi} \sin^2(\theta)$  and  $G = 1.5 \sin^2(\theta)$ . Note we used  
 $h = \frac{l}{2} \sin(\theta)$
- $\frac{|V|^2}{8R_2} = \frac{|I|^2 R_1 G_1 A_2}{8\pi r^2}, G_1 A_2 = G_2 A_1$
- $P_t = \frac{|I|^2 R_1}{4\pi r^2}, P_t = \frac{|V|^2}{8R_2}$
- $P_r = \frac{P_t G_1 A_2}{4\pi r^2}$

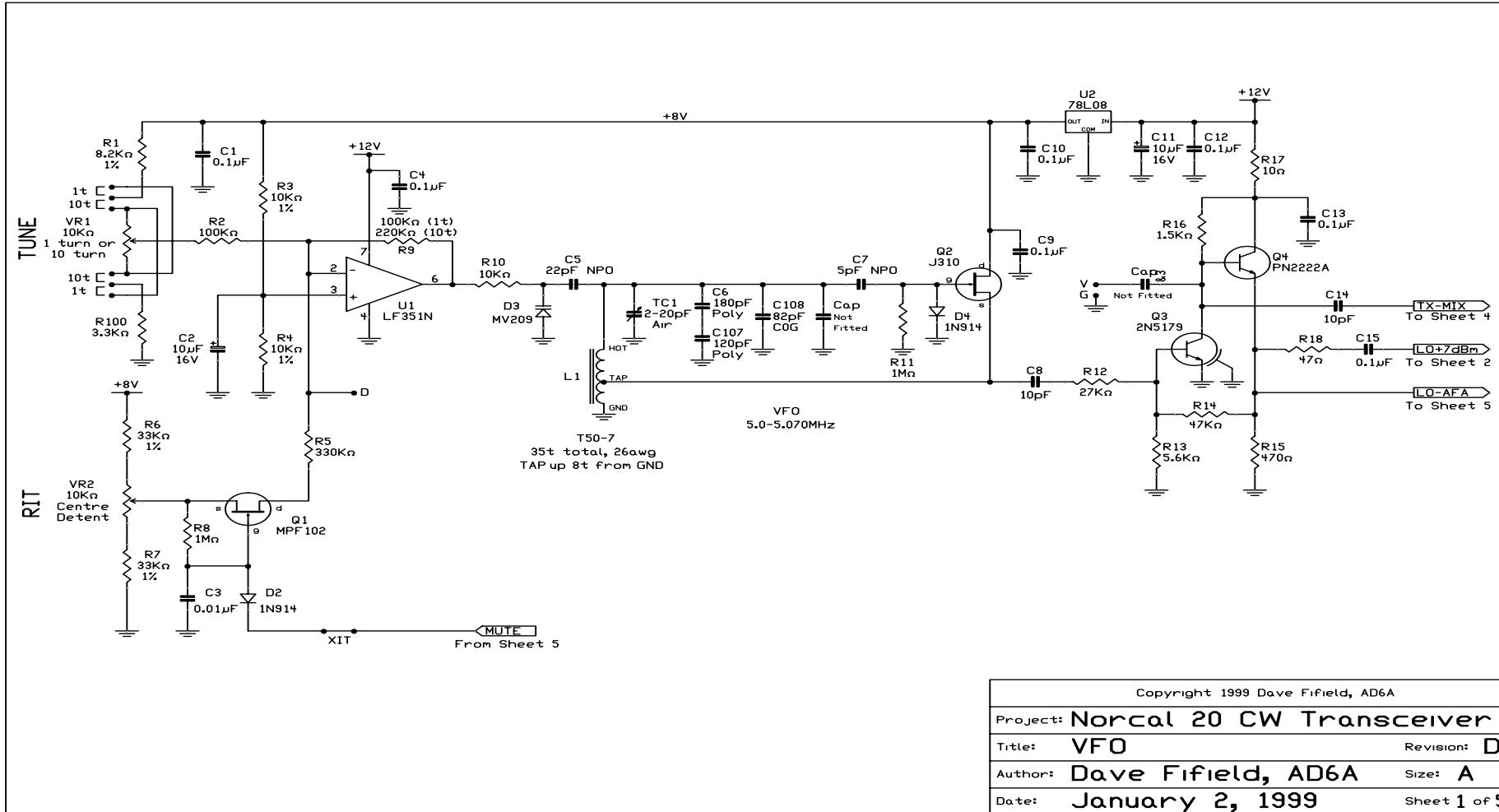
# Exercise 35: Intermodulation

- Only  $f_{3\uparrow} = 2f_1 - f_2$ ,  $f_{3\downarrow} = 2f_2 - f_1$ ,  $f_{5\uparrow} = 3f_1 - 2f_2$  and  $f_{5\downarrow} = 3f_2 + 2f_1$  are close enough to the rf frequency to matter for intermodulation
1. Find coefficients and frequencies for  $[\cos(\omega_1 t) + [\cos(\omega_2 t)]^5$
  2. Find  $f_{3\uparrow}$ ,  $f_{3\downarrow}$ ,  $f_{5\uparrow}$  and  $f_1$
  3. Find the MDS and the antenna limited MDR

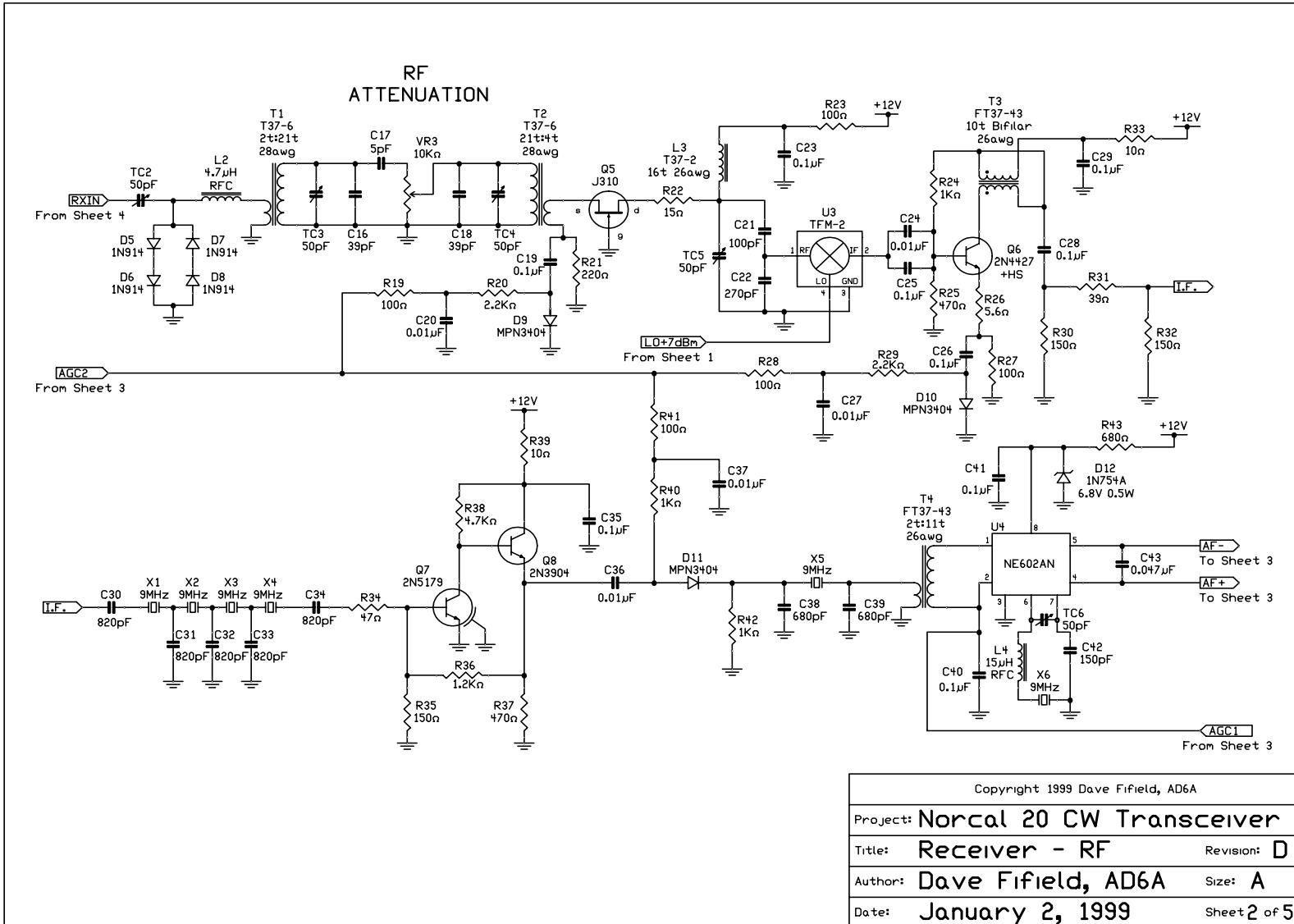
# Exercise 37: Antennas

1. Use the relation between gain and effective area to rewrite the Friis transmission formula in terms of gain only. Consider UHF for airplanes. If the frequency makes the quarter length stub antenna have gain 2, find the maximum possible LOS at 10km height. Required receiver power is  $-90$  dBm. Find the minimum transmission power.
2. Find the inductance to resonate with a 3m whip. Assuming the Q of the coil is 200, find the turns ratio required to give a transceiver a 50 ohm load. What is the radiation efficiency?
3. Repeat 2 with a capacitive end loading, assuming the capacitance doubles.

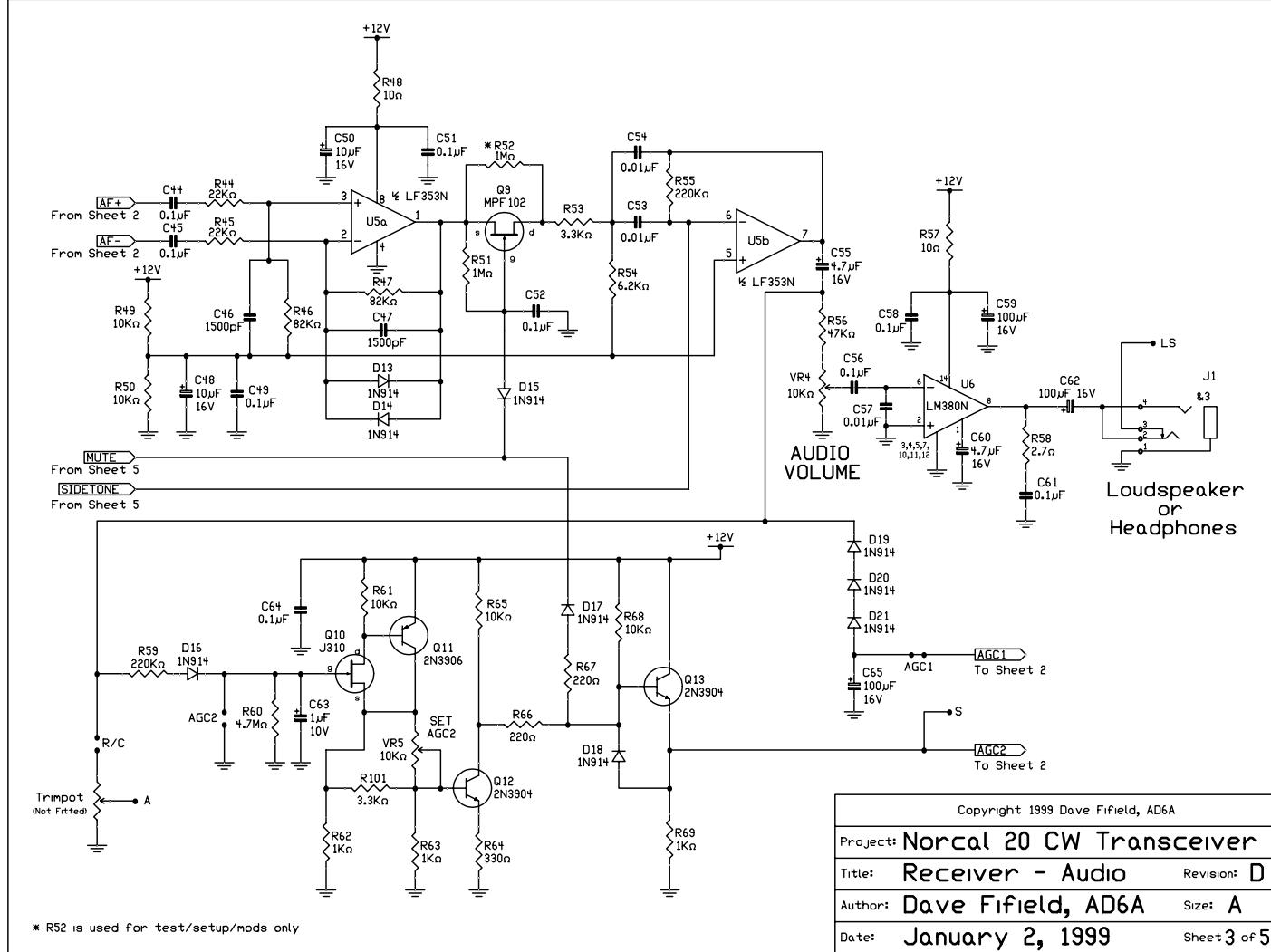
# Norcal circuit diagram, 1



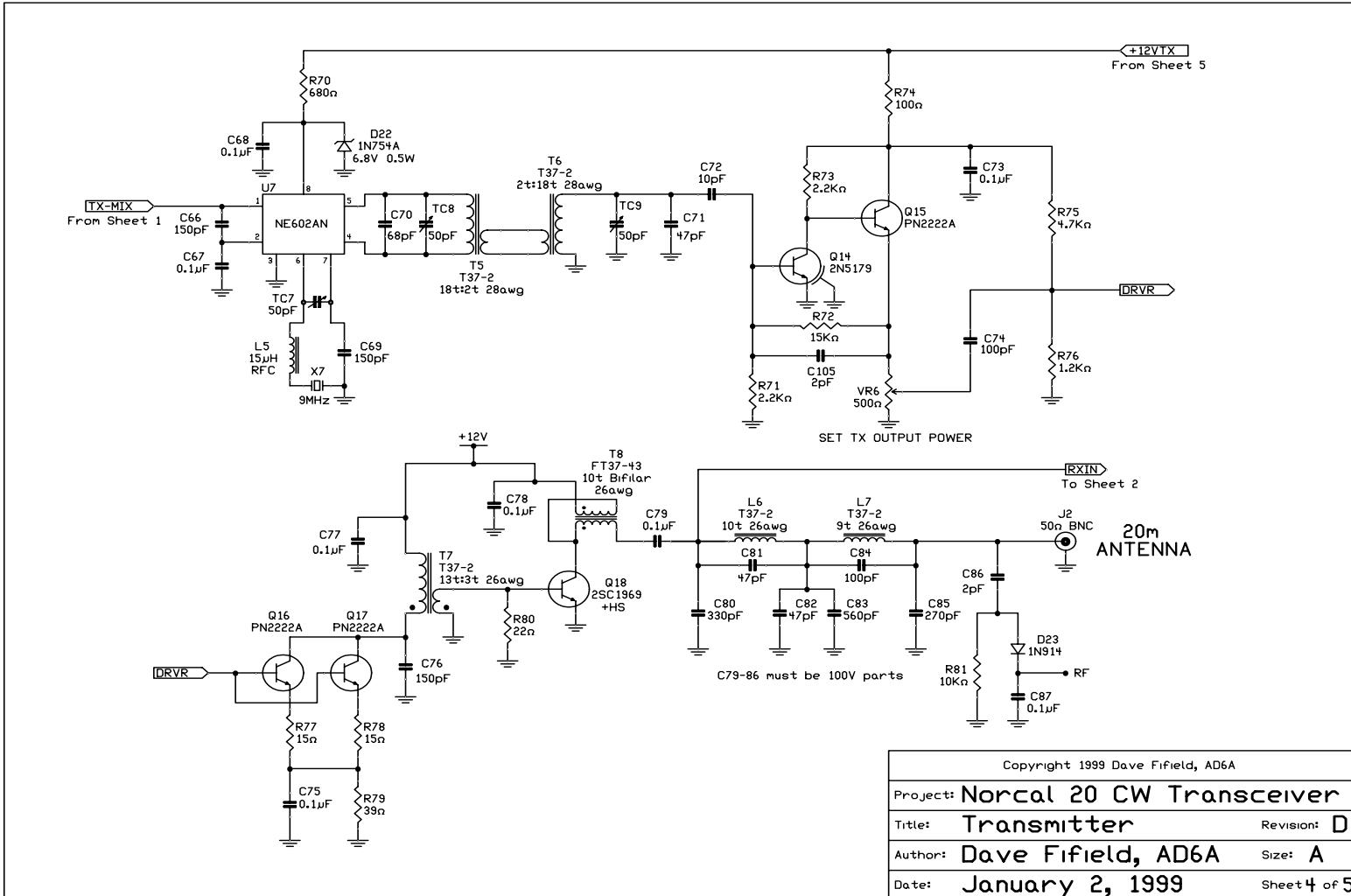
# Norcal circuit diagram, 2



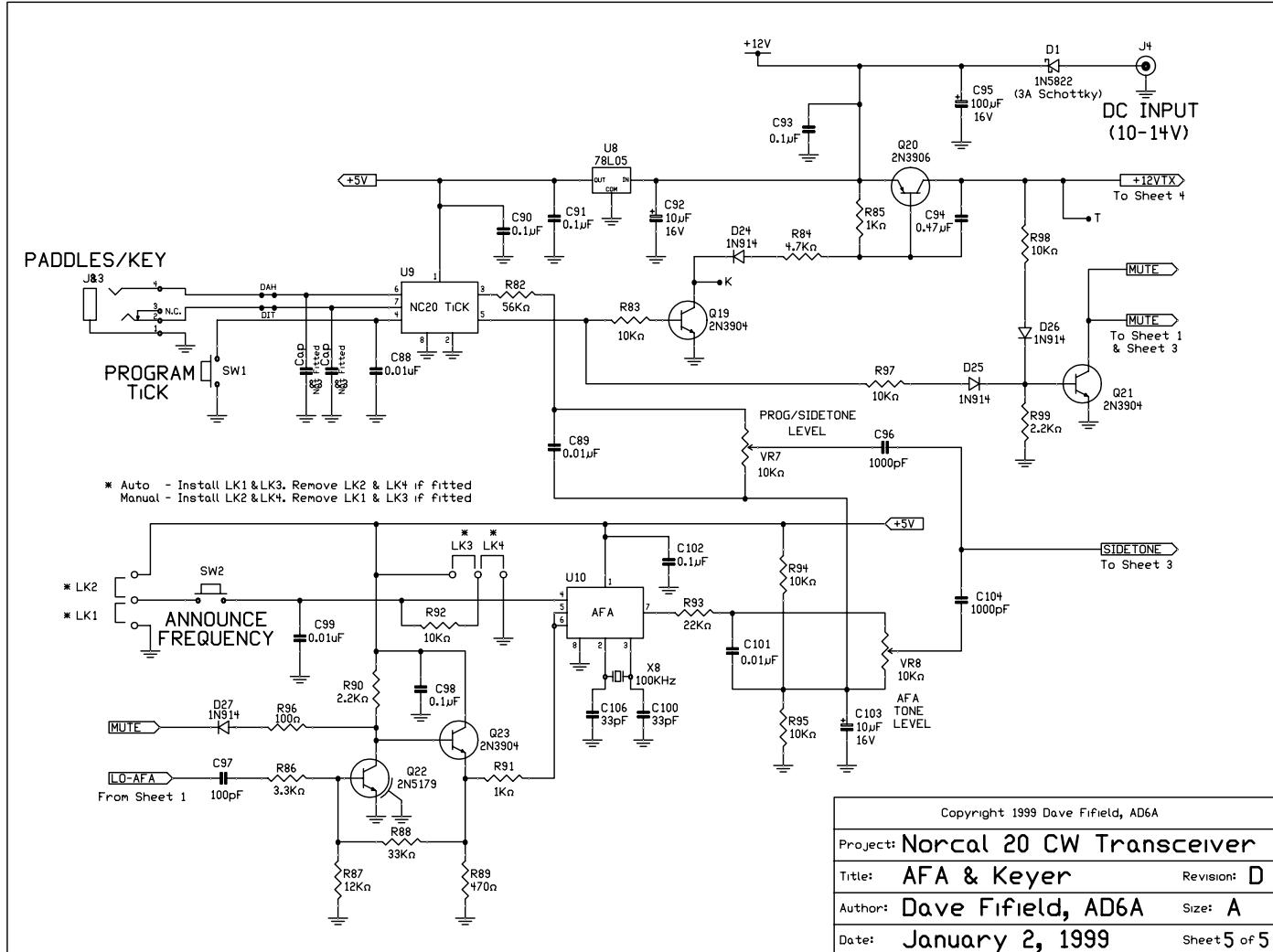
# Norcal circuit diagram, 3



# Norcal circuit diagram, 4



# Norcal circuit diagram, 5

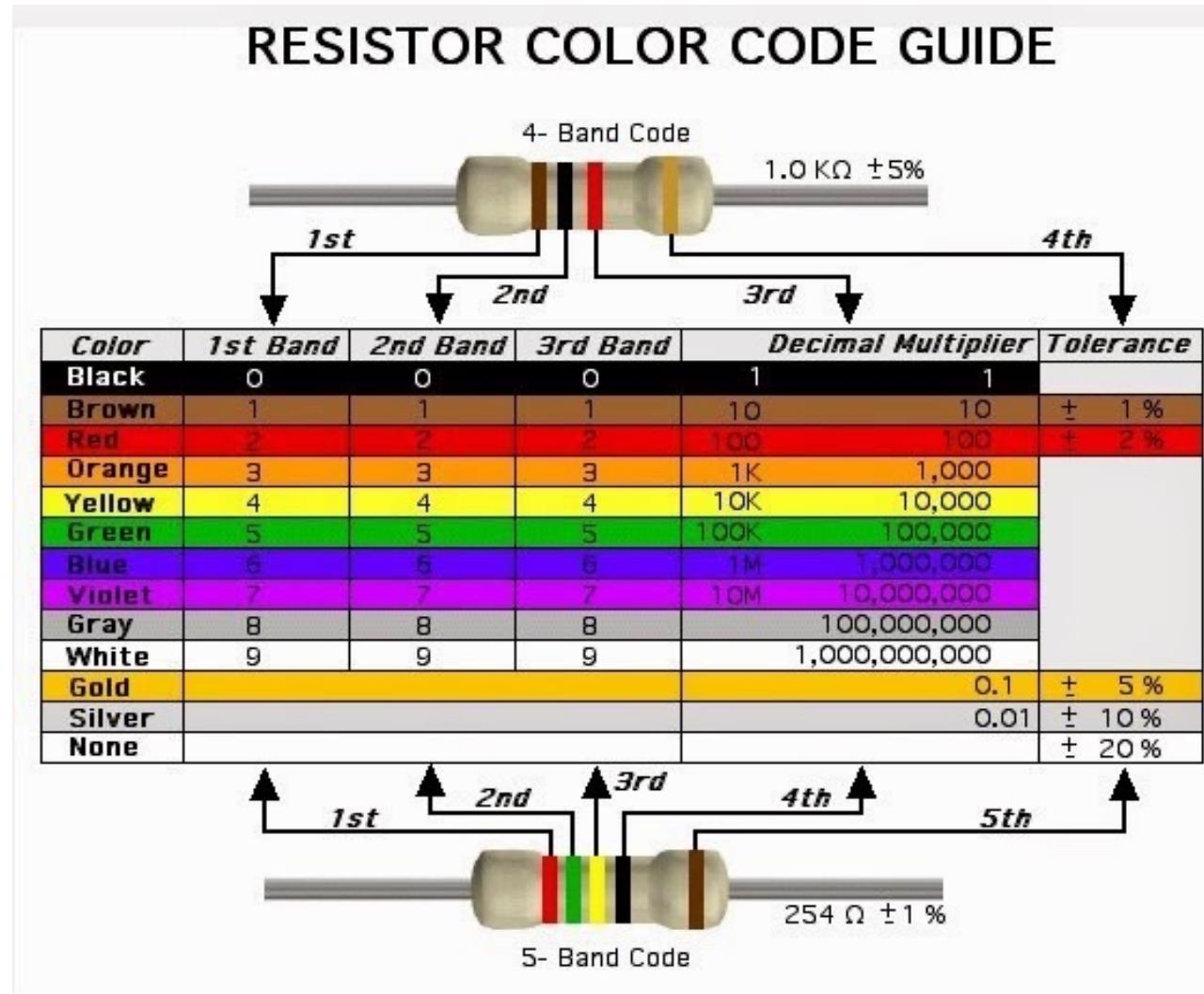


# Appendix

# Morse

Symbol	Code	Symbol	Code	Symbol	Code
a	.—	m	—	y	—.—
b	—...	n	—.	z	—..
c	—.-	o	---	0	----
d	—..	p	.—.	1	.----
e	.	q	—.-	2	..----
f	..—.	r	—.	3	...--
g	—.	s	...	4	....-
h	....	t	—	5	.....
i	..	u	..—	6	—....
j	.—	v	...—	7	—...
k	—.-	w	—.	8	—..
l	—..	x	—..—	9	-----.

# Color codes



- Resistors markings in ohms
- Capacitors markings in picoFarads
- Inductors markings in microHenries

# Component data



Core Size	26	3	15	1	2	6	10	12/17	0
T-12-( )	*	60	50	48	20	17	12	7.5	2.4
T-16-( )	145	61	55	44	22	19	13	8	3
T-20-( )	185	76	65	52	25	22	16	10	3.5
T-25-( )	245	100	85	70	34	27	19	12	4.5
T-30-( )	335	140	93	85	43	36	25	16	6
T-37-( )	285	120	90	80	40	30	25	15	4.9
T-44-( )	370	180	160	105	52	42	33	18.5	6.5
T-50-( )	330	175	135	100	49	40	31	18	6.4
T-68-( )	435	195	180	115	57	47	32	21	7.5
T-80-( )	460	180	170	115	55	45	32	22	8.5
T-94-( )	600	248	200	160	84	70	58	*	10.6
T-106-( )	930	450	345	325	135	116	*	*	19
T-130-( )	810	350	250	200	110	96	*	*	15
T-157-( )	1000	420	*	320	140	115	*	*	*
T-184-( )	1690	720	*	500	240	195	*	*	*
T-200-( )	920	425	*	250	120	100	*	*	*
T-200A-( )	1600	*	*	*	218	*	*	*	*

## IRON POWDER TOROIDS - A<sub>i</sub> Values \*\*

\* size not available in this material

\*\* L =  $\mu\text{H}/100 \text{ turns}$

# In the beginning...

- The laws of EM according to Clerk Maxwell are:

$$1. \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$2. \nabla \cdot \mathbf{B} = 0$$

$$3. \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$4. c^2 \nabla \times \mathbf{B} = \frac{\mathbf{j}}{\epsilon_0} + \frac{\partial \mathbf{E}}{\partial t}, \epsilon_0 = 8.854 \times 10^{-12} \frac{C^2}{N \cdot m^2}, \frac{1}{c^2} = \epsilon_0 \mu_0$$

$$5. \nabla \cdot \mathbf{j} = -\frac{\partial \rho}{\partial t}$$

- In non-dispersive matter,  $B = \mu H = \mu_0(H + M)$ ,  $\mu = \mu_0(1 + \chi_m)$ ,  $D = \epsilon E = \epsilon_0 E + P$ ,  $\epsilon = \kappa \epsilon_0$ .

- (1) becomes  $\nabla \cdot \mathbf{D} = \rho_f$ , (4) becomes  $\nabla \times \mathbf{H} = \mathbf{j}_f + \frac{\partial \mathbf{D}}{\partial t}$

- Here  $\mathbf{E}$  is the electric field,  $\mathbf{B}$  is the magnetic field,  $\mathbf{j}$  is the current density through a closed surface,  $c$  is the speed of light and  $\rho$  is the charge density at a point. The rest of classical physics, including special relativity, is:

- Newton-Einstein:  $\mathbf{p} = m\mathbf{v}$ ,  $m = \frac{m_0}{\sqrt{1 - (\frac{v}{c})^2}}$ ,  $\mathbf{F} = m \frac{d\mathbf{p}}{dt}$ .

- Gravity:  $\mathbf{F} = -\frac{Gm_1 m_2}{r^2} \mathbf{u}_r$  where  $\mathbf{u}_r$  is the unit vector from  $m_1$  to  $m_2$  and  $\mathbf{F}$  is the force on  $m_2$ .

# Solutions to the wave equation

- The solution of  $\nabla^2\psi - \frac{1}{c^2}\frac{\partial^2\psi}{\partial t^2} = -s$  is  $\psi(x, y, z, t) = \frac{s(t-\frac{r}{c})}{4\pi r}$  where  $S = \int_V s dV$
- Later, we will use this to find the "general" solution to Maxwell's equations
  - $\phi(r_1, t) = \int_{V_2} \frac{\rho(r_2, t-\frac{|r_1-r_2|}{c})}{4\pi\epsilon_0|r_1-r_2|} dV_2$  and  $\mathbf{A}(r_1, t) = \int_{V_2} \frac{\mathbf{j}(r_2, t-\frac{|r_1-r_2|}{c})}{4\pi\epsilon_0 c^2|r_1-r_2|} dV_2$ , where
  - $B = \nabla \times A$ ,  $E = -\nabla\phi - \frac{\partial A}{\partial t}$ , and,  $c^2\nabla \cdot A = -\frac{\partial \phi}{\partial t}$
- You are not expected to have guessed this answer
- To do this, we'll need the "BAC-CAB" identity:  $A \times (B \times C) = B(A \cdot C) - C(A \cdot B)$
- When we apply this to  $\nabla \times (\nabla \times \mathbf{A})$ , we get  $\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$

# General solution to Maxwell's equations

- Returning to the general Maxwell equations, from  $\nabla \cdot B = 0$ , we get  $B = \nabla \times A$
- Substituting into  $c^2 \nabla \times \mathbf{B} = \frac{\mathbf{j}}{\epsilon_0} + \frac{\partial \mathbf{E}}{\partial t}$ , we get  $c^2 \nabla \times (\nabla \times A) = \frac{\mathbf{j}}{\epsilon_0} + \frac{\partial \mathbf{E}}{\partial t}$
- Applying “BAC-CAB”, we get  $\nabla(\nabla \cdot A) - \nabla^2 A = \frac{\mathbf{j}}{c^2 \epsilon_0} + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$  (Equation 1)
- Now,  $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ , so substituting for  $\mathbf{B}$ , we get  $\nabla \times \left( \mathbf{E} + \frac{\partial \mathbf{A}}{\partial t} \right) = 0$  and so  $\mathbf{E} = -\nabla \phi - \frac{\partial \mathbf{A}}{\partial t}$
- Substituting into equation 1,  $\nabla(\nabla \cdot A) - \nabla^2 A = \frac{\mathbf{j}}{c^2 \epsilon_0} + \frac{1}{c^2} \frac{\partial}{\partial t} \left( -\nabla \phi - \frac{\partial \mathbf{A}}{\partial t} \right)$ , or
- $\nabla(\nabla \cdot A) - \nabla^2 A = \frac{\mathbf{j}}{c^2 \epsilon_0} - \frac{1}{c^2} \nabla \frac{\partial \phi}{\partial t} - \frac{1}{c^2} \frac{\partial^2 A}{\partial t^2}$ .
- $\nabla^2 A - \frac{1}{c^2} \frac{\partial^2 A}{\partial t^2} = -\frac{\mathbf{j}}{c^2 \epsilon_0} + \nabla \left[ \frac{1}{c^2} \frac{\partial \phi}{\partial t} + (\nabla \cdot A) \right]$
- Now if  $A$  and  $\phi$  give  $B = \nabla \times A$  and  $E = -\nabla \phi - \frac{\partial A}{\partial t}$ , then  $A' = A + \nabla \varphi$  and  $\phi' = \phi - \frac{\partial \varphi}{\partial t}$  give  $B = \nabla \times A'$  and  $E = -\nabla \phi' - \frac{\partial A'}{\partial t}$ , for any function  $\varphi$

# General solution to Maxwell's equations

- Thus, we can pick a solution  $(A, \phi)$  with  $\nabla \cdot A = -\frac{1}{c^2} \frac{\partial \phi}{\partial t}$ . Then we get
- $\nabla^2 A - \frac{1}{c^2} \frac{\partial^2 A}{\partial t^2} = -\frac{j}{c^2 \epsilon_0}$
- Substituting for  $E = -\nabla \phi - \frac{\partial A}{\partial t}$  into  $\nabla \cdot E = \frac{\rho}{\epsilon_0}$ , we get
- $\nabla \cdot (\nabla \phi) + \frac{\partial \nabla \cdot A}{\partial t} = -\frac{\rho}{\epsilon_0}$ , or  $\nabla \cdot (\nabla \phi) - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = -\frac{\rho}{\epsilon_0}$ , or  $\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = -\frac{\rho}{\epsilon_0}$
- The solutions are  $\phi(r_1, t) = \int_{V_2} \frac{\rho(r_2, t - \frac{|r_1 - r_2|}{c})}{4\pi\epsilon_0 |r_1 - r_2|} dV_2$  and
- $\mathbf{A}(r_1, t) = \int_{V_2} \frac{\mathbf{j}(r_2, t - \frac{|r_1 - r_2|}{c})}{4\pi\epsilon_0 c^2 |r_1 - r_2|} dV_2$  with  $\nabla \cdot A = -\frac{1}{c^2} \frac{\partial \phi}{\partial t}$ , with  $\nabla \cdot A = -\frac{1}{c^2} \frac{\partial \phi}{\partial t}$ .

# Solution to Maxwell's equations in free space

- Free space is defined by  $\rho = 0$  and  $j = 0$ , so our potentials satisfy
- $\nabla^2 A - \frac{1}{c^2} \frac{\partial^2 A}{\partial t^2} = 0$  and  $\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 0$
- These have the usual wave equation solutions  $\phi(x, y, z, t) = f(k \cdot r - \omega t)$ , etc
- Thus, in free space  $\phi$  and  $\mathbf{A}$  and hence  $\mathbf{E}$  and  $\mathbf{B}$  propagate as waves.

# Solution to Maxwell's equations in conductors

- In conductors,  $\mathbf{j} = \sigma \mathbf{E}$ 
  - $c^2 \nabla \times \mathbf{B} = \frac{\mathbf{j}}{\epsilon} + \frac{\partial \mathbf{E}}{\partial t} = \frac{\sigma}{\epsilon} \mathbf{E} + \frac{\partial \mathbf{E}}{\partial t}$
  - This becomes  $c^2 \frac{\partial(\nabla \times \mathbf{B})}{\partial t} = \frac{\sigma}{\epsilon} \frac{\partial \mathbf{E}}{\partial t} + \frac{\partial^2 \mathbf{E}}{\partial t^2}$
  - $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ , so we get  $c^2 \nabla \times \frac{\partial \mathbf{B}}{\partial t} = -c^2 \nabla \times (\nabla \times \mathbf{E}) = \frac{\sigma}{\epsilon} \frac{\partial \mathbf{E}}{\partial t} + \frac{\partial^2 \mathbf{E}}{\partial t^2} = -c^2 [\nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}] = c^2 \nabla^2 \mathbf{E}$   
(since  $\rho = 0$  in a conductor)
  - Applying the trial solution  $\mathbf{E} = E_0 \exp(i\omega t - kr)$ , we get  $-k^2 - i\omega\mu\sigma + \omega^2\mu\epsilon = 0$ .
  - Putting  $k = \alpha - \beta i$ ,  $\alpha = \frac{\omega}{2} \sqrt{\mu\epsilon} \left(1 + \sqrt{1 + \frac{\sigma^2}{\omega^2\epsilon^2}}\right)$  and  $\beta = \frac{\omega\mu\sigma}{2\alpha}$ .
  - For copper,  $\sigma = 5.78 \times 10^7 \Omega^{-1}\text{-m}$ . This explains the “skin effect” in conductors.

# Radiation, antennas

- Accelerating charges radiate energy in the form of electromagnetic waves (companion E and B fields).
- The radiation from accelerating charge  $q$  is  $\mathbf{E}_{rad} = -\frac{1}{4\pi\epsilon_0 c^2} \frac{q}{r} \mathbf{a}_\perp (t - \frac{r_{12}}{c})$ .
- Here,  $\mathbf{a}_\perp$  is the acceleration  $\perp$  to the line from  $\mathbf{r}_1$  to  $\mathbf{r}_2$ .
- For example, applying a time varying potential  $V_0 \sin(\omega t)$  to an antenna will cause the antenna to radiate power since the voltage and hence charges affected accelerate within the antenna, that is, their positions have a non-zero second derivative. That's how a transmitter "couples" to the antenna of a receiver. In the receiver, the radiated wave accelerates charges in the antenna replicating the original wave (at much reduced power).
- These simple radio waves are carrier waves of frequency  $\frac{\omega}{2\pi}$ . To transfer information (voice, images, binary data), we modulate carrier waves combining them with an "information source" signal. Receivers demodulate the incoming wave and recreate the original "information source" signal.

# Maxwell's equations in a non-dispersive media

- $B = \mu H, D = \epsilon E$
- $\nabla \cdot D = \rho$
- $\nabla \cdot B = 0$
- $\nabla \times E = -\frac{\partial B}{\partial t}$
- $\nabla \times H = j + \frac{\partial E}{\partial t}$
- $\nabla \cdot j = -\frac{\partial \rho}{\partial t}$

# Radiation from a small dipole

- $A_z(r, t) = \frac{\mu_0}{4\pi} \int_{[-l/2, l/2]} \frac{I(z', t - \frac{z'}{c}k)}{|r - z' k|} dz'$

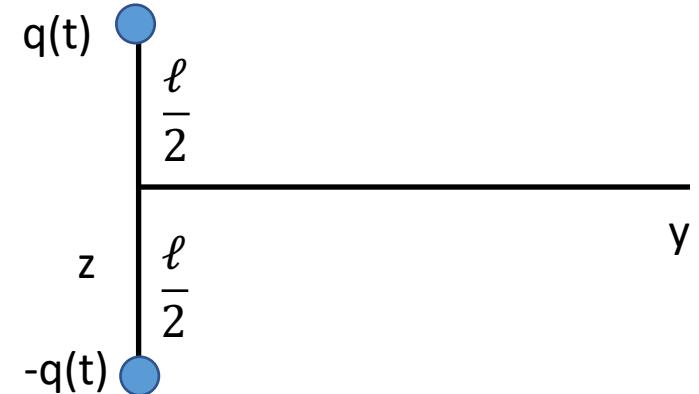
- If  $l \ll cT = \lambda$

- $A_z(r, t) = \frac{\mu_0}{4\pi} \frac{l}{r} I(z', t - \frac{r}{c})$

- Choosing gauge,  $\nabla \cdot A + \frac{1}{c^2} \frac{\partial \phi}{\partial t} = 0$

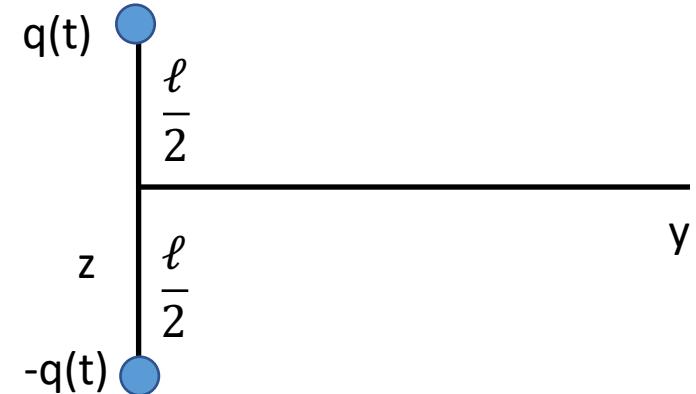
- $\frac{\partial \phi}{\partial t} = -\frac{l}{4\pi\epsilon_0} \frac{\partial}{\partial z} \left( \frac{1}{r} I\left(t - \frac{r}{c}\right) \right) = \frac{z}{r^2} \left( \frac{q(t - \frac{r}{c})}{r} - \frac{I((t - \frac{r}{c}))}{c} \right)$

- $q\left(t - \frac{r}{c}\right) = q_0 \cos\left(\omega\left[t - \frac{r}{c}\right]\right), I\left(t - \frac{r}{c}\right) = I_0 \sin\left(\omega\left[t - \frac{r}{c}\right]\right)$



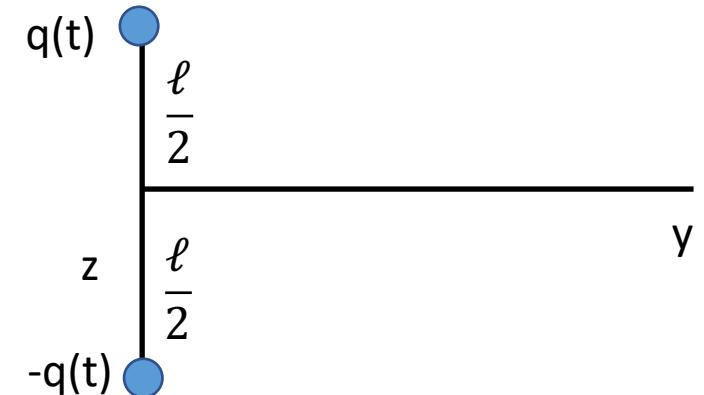
# Radiation from a small dipole

- $\nabla^2 H - \epsilon\mu \frac{\partial^2 H}{\partial t^2} - \sigma\mu \frac{\partial H}{\partial t} = 0$
- $\nabla^2 E - \epsilon\mu \frac{\partial^2 E}{\partial t^2} - \sigma\mu \frac{\partial E}{\partial t} = 0$
- $A_r = \frac{\mu_0}{4\pi} \frac{I_0 l}{r} \cos(\theta) \sin(\omega [t - \frac{r}{c}])$
- $A_\phi = 0, A_\theta = -\frac{\mu_0}{4\pi} \frac{I_0 l}{r} \cos(\theta) \sin(\omega [t - \frac{r}{c}])$
- $B_\phi = \frac{1}{r} \frac{\partial}{\partial r} (r A_\theta) - \frac{1}{r} \frac{\partial A_r}{\partial \theta} = \frac{\mu_0 I_0 l}{4\pi r} \sin(\theta) [\frac{\omega}{r} \cos(\omega [t - \frac{r}{c}]) + \frac{1}{r} \sin(\omega [t - \frac{r}{c}])]$
- $E_r = -\frac{\partial \phi}{\partial t} - \frac{\partial A_r}{\partial t} = \frac{2lI_0 \cos(\theta)}{4\pi\epsilon_0} \left[ \frac{\sin(\omega [t - \frac{r}{c}])}{r^2 c} - \frac{\cos(\omega [t - \frac{r}{c}])}{\omega r^3} \right]$
- $E_\theta = \frac{-I_0 l \sin(\theta)}{4\pi\epsilon_0} \left( \left[ \frac{1}{r^3 \omega} - \frac{\omega}{rc^2} \right] \cos(\omega [t - \frac{r}{c}]) - \frac{1}{cr^2} \sin(\omega [t - \frac{r}{c}]) \right)$
- $E_\phi = -\frac{1}{rsin(\theta)} \frac{\partial \phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} = 0$



# Radiation from a small dipole

- $\nabla \cdot (\mathbf{E} \times \mathbf{H}) = -\mathbf{H} \frac{\partial \mathbf{B}}{\partial t} - \mathbf{E} \frac{\partial \mathbf{D}}{\partial t} - \mathbf{E} \cdot \mathbf{j}, \mathbf{S} = \mathbf{E} \times \mathbf{H}, \nabla \cdot \mathbf{S} + \frac{\partial u}{\partial t} = -\mathbf{E} \cdot \mathbf{j}$  (u is energy density)
- $\int S \cdot dA = \frac{(l I_0 \omega)^2}{6\pi\epsilon_0 c^3} \cos(\omega \left[ t - \frac{r}{c} \right])^2$
- $P_{av} = \frac{(l\omega)^2}{6\pi\epsilon_0 c^3} \frac{I_0^2}{2} = \frac{2\pi}{3} \sqrt{\frac{\mu_0}{\epsilon_0}} \left(\frac{l}{\lambda}\right)^2 \frac{I_0^2}{2}$
- $R_r = \frac{2\pi}{3} \sqrt{\frac{\mu_0}{\epsilon_0}} \left(\frac{l}{\lambda}\right)^2$



# Large half wave dipole

- For large half wave, add small dipoles to produce half wave antenna.

- $dE_\theta = I_0 \frac{\sin(\theta)}{4\pi\epsilon_0 R c^2} \omega \cos(\omega) \cos\left(\frac{2\pi z'}{\lambda}\right) dz'$
- $dB_\phi = I_0 \frac{\mu_0 \omega}{4\pi R c} \omega \cos\left(\omega \left[t - \frac{r}{c}\right]\right) \cos\left(\frac{2\pi z'}{\lambda}\right) dz'$
- $K = \int_{[-\frac{\pi}{2}, \frac{\pi}{2}]} \frac{1}{R} \cos\left(t - \frac{R}{c}\right) \cos(u) du = \frac{1}{2\pi\epsilon_0 r c} \cos\left(\omega \left[t - \frac{r}{c}\right]\right) \frac{\cos\left(\frac{\pi}{2}\cos(\theta)\right)}{\sin^2(\theta)}, u = \frac{2\pi z'}{\lambda}$
- $E_\theta = I_0 \frac{1}{2\pi\epsilon_0 r c} \cos\left(\omega \left[t - \frac{r}{c}\right]\right) \frac{\cos\left(\frac{\pi}{2}\cos(\theta)\right)}{\sin(\theta)}$
- $B_\phi = I_0 \frac{\mu_0}{2\pi r} \omega \cos\left(\omega \left[t - \frac{r}{c}\right]\right) \frac{\cos\left(\frac{\pi}{2}\cos(\theta)\right)}{\sin(\theta)}$
- $P_{av} = \frac{1}{4\pi} \sqrt{\frac{\mu_0}{\epsilon_0}} I_0^2 \int_{[0, \pi]} \frac{\cos^2\left(\frac{\pi}{2}\cos(\theta)\right)}{\sin^2(\theta)} \sin(\theta) d\theta = 73.1 \Omega \frac{I_0^2}{2}$

# Radiation from an accelerating charge

- $r' + R = r, R = |\mathbf{r} - \mathbf{r}'|$
- $\varphi(r, t) = \frac{1}{4\pi\epsilon_0} \int_{V_1} \frac{\rho(r', t - \frac{R}{c})}{|r - r'|} d\nu' = \frac{1}{4\pi\epsilon_0} \left[ \frac{Q}{r} + \frac{r \cdot \mathbf{p}(t - \frac{r}{c})}{r^3} + \frac{r \cdot \frac{d\mathbf{p}}{dt}(t - \frac{r}{c})}{cr^2} \right]$
- $\mathbf{A}(r, t) = \frac{\mu_0}{4\pi} \int_{V_1} \frac{\mathbf{j}(r', t - \frac{R}{c})}{|r - r'|} d\nu' = \frac{\mu_0}{4\pi r} \frac{d}{dt} \mathbf{p}(t - \frac{r}{c})$
- $\mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t} - \nabla \varphi$
- $\mathbf{B}(r, t) = \frac{-\mu_0}{4\pi c r^2} \mathbf{r} \times \frac{d^2}{dt^2} \mathbf{p}(t - \frac{r}{c})$
- $\mathbf{E}(r, t) = -\frac{c}{r} \mathbf{r} \times \mathbf{B}(r, t)$
- $\frac{d\mathbf{p}}{dt} = q \frac{d\mathbf{r}'}{dt} = q\mathbf{v}, \frac{d^2}{dt^2} \mathbf{p}(t - \frac{r}{c}) = q \frac{d\mathbf{v}}{dt}$
- $P_R = -\frac{dW}{dt} = \frac{2}{3} \frac{q^2}{4\pi\epsilon_0 c^3} \left( \frac{d\mathbf{v}}{dt} \right)^2$

# Radiation from a single accelerating charge

- Near zone

- $\varphi(r, t) = \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{R'(1 + \frac{\nu \cdot n'}{c})} \right]$
- $R^* = R' - \frac{\nu}{c}(x_0 - x'_1)$
- $E(r, t) = \frac{1}{4\pi\epsilon_0} \frac{1}{R^{*3}} \left[ \left( R' - \frac{R' \nu'}{c} \right) \left( 1 - \frac{\nu^2}{c^2} \right) + \frac{R'}{c^2} \left( R' - \frac{R' \nu'}{c} \right) \times \frac{d\nu'}{dt} \right]$
- $B = \frac{R' \times E}{R' c}$
- $S = \frac{q^2}{16\pi^2 c^3 \epsilon_0} \frac{R' (R' \times \nu')^2}{(R')^5}$

# Radiation loss and antenna aperture

- Spreading loss:  $L_s = 32 + \log(d) + 20 \log(f)$ 
  - d in kilometers
  - F in megahertz
- $W = A_e P_e, A_e = \frac{\lambda^2}{4\pi}$

# Antenna aperture

- Deriving antenna aperture uses thermodynamic argument: black body equilibrium

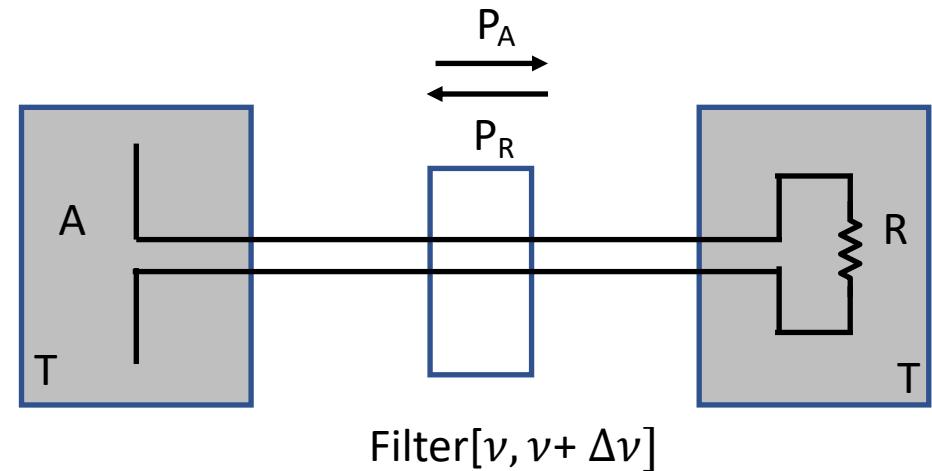
$$P_A = \frac{A_e}{2} B_v \Delta\nu \int_{[0,4\pi]} d\Omega = 2\pi A_e B_v \Delta\nu$$

$$B_v = \frac{2\nu^2 kT}{c^2} = \frac{2kT}{\lambda^2} \text{ (Rayleigh-Jeans)}$$

$$P_R = kT\Delta\nu$$

$$P_A = P_R$$

$$A_e = \frac{\lambda^2}{4\pi}$$



# Impedance matching

- L networks,  $\pi$  and T networks
- Impedance and reflection coefficients