

# Hight $p_T$ jets in Run II of the ATLAS Experiment

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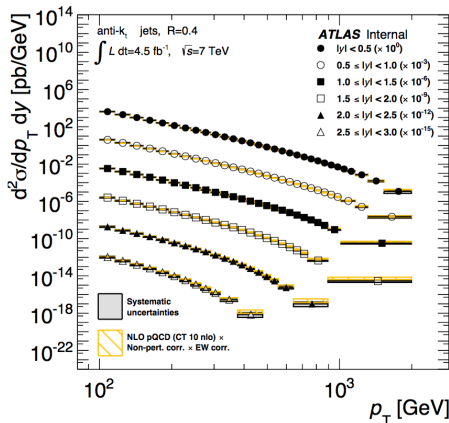
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ATLAS Meeting

May 14, 2015

# My Analysis

Inclusive jet double differential cross section in  $p_T$  and rapidity  $y$  (inclusive means  $pp \rightarrow \text{jet} + \text{"anything"}$ ) in Run II of the ATLAS Experiment ( $\sqrt{s} = 13 \text{ TeV}$ )



## Why Inclusive Jets?

- ▶ They Cover a *wide range of momentum transfers* ( $\sim 1 \text{ GeV} - 1 \text{ TeV}$  on the LHC)  $\rightarrow$  predictions sensitive to the properties of the running coupling constant  $\alpha_s$
- ▶ They probe the structure of proton at *small distance scales*

$$\lambda \sim 1/p_T \sim \text{TeV}^{-1} \sim 10^{-19} \text{ m}$$

- ▶ They contribute to our understanding of PDFs
- ▶ They *appreciate the increase in the center-of-mass energy* as no other physics process observed on hadron colliders

# My Analysis

## Data

Monte Carlo generated events of  $pp$  collisions at  $\sqrt{s} = 13$  TeV.

- ▶ collisions generated with PYTHIA8 (particle level)
- ▶ ATLAS detector response obtained with GEANT4 full simulation (detector level)

## Jet Corrections

Cross section corrected from the detector to the particle level in two steps

- ▶ Calibration
- ▶ Unfolding

## LO vs. NLO QCD

Particle level cross section from PYTHIA8 (LO QCD) compared with the parton level NLO QCD cross section prediction.

# Why Do We Need Jets?

**Gluon radiation cross section:**

**Divergences:**

- ▶ *Infrared* ( $E_k = 0$ )
- ▶ *Collinear* ( $\theta = 0$ )

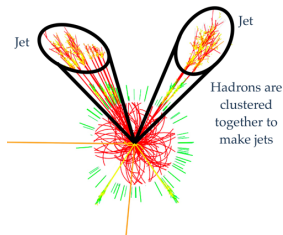
**Jet:** A group of collimated particles

**Jet algorithm:** A prescription, how particles (or other objects) are clustered into separate jets. It should fulfill

- ▶ *Infrared safety:* The presence of an additional soft particle should not affect the recombination of particles into a jet.
- ▶ *Collinear safety:* Jet reconstruction should not depend on the fact, if the energy is carried by one particle, or if the particle is split into more collinear particles.

**Summary:**  $q$  or  $g$  **CANNOT** be observed. Jets **CAN**.

$$\sigma_{q \rightarrow qg} \sim \frac{d\theta}{|\sin \theta|} \frac{dE_k}{E_k}$$



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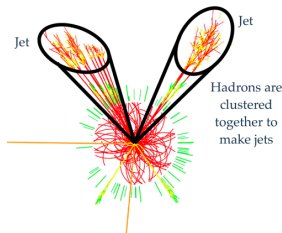
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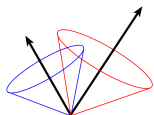


# Cone Jet Algorithms

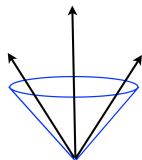
- ▶ The **most illustrative** jet algorithms. Different modifications.
- ▶ Used in Tevatron. Not used in ATLAS.

## Algorithm:

1. Take a particle with the highest  $p_T > p_T^{cutoff}$
2. Recombine all particles within the fixed cone
3. Update the cone direction
4. If the direction has changed, go to 2, else you have a jet
5. Go to 1 until there is no particle left with  $p_T > p_T^{cutoff}$



infrared unsafety



collinear unsafety

# $k_t$ Jet Algorithms

## Algorithm:

1. For each input object  $i$  and all pairs of input objects  $(i, j)$  calculate

$$d_i = p_{T,i}^{2p} \quad , \quad d_{ij} = \min \left( p_{T,i}^{2p}, p_{T,j}^{2p} \right) \frac{\Delta y^2 + \Delta \phi^2}{R^2} \quad (R = 0.4)$$

- ▶  $p = 1$   $k_t$  algorithm,
- ▶  $p = 0$  Cambridge/Aachen algorithm,
- ▶  $p = -1$  **anti- $k_t$  algorithm**.

2. Find minimum  $d_{min}$  between all  $d_{ij}$  and  $d_i$

- ▶  $d_{min}$  is between  $d_{ij}$ 's.  
Recombine  $i, j$  into a new object  $k$ . Remove  $i, j$  from the list, add  $k$  to the list.
- ▶  $d_{min}$  is between  $d_i$ 's.  
Object  $i$  is a jet. Remove  $i$  from the list.

3. Go to 1 until all input objects are part of a jet.

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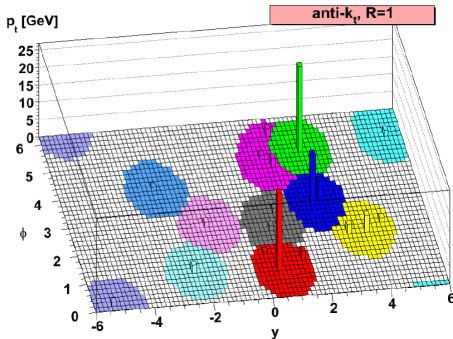
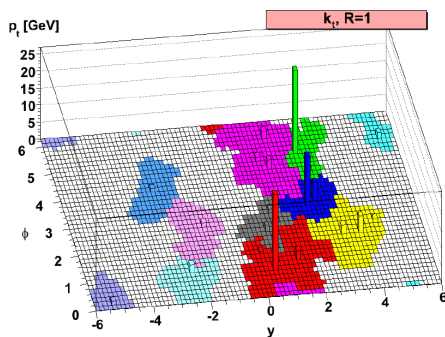
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# $k_t$ Jet Algorithms

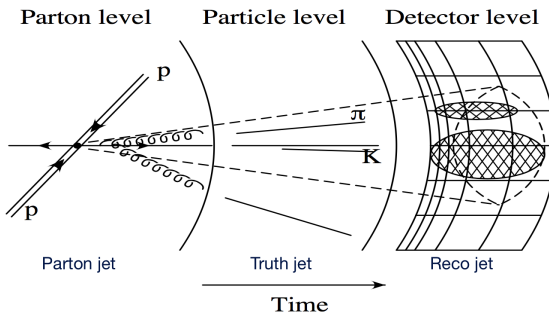


- ▶  $k_t$  jet algorithms are both infrared and collinear safe
- ▶ ATLAS uses anti- $k_t$  jet algorithm

# Jet Reconstruction

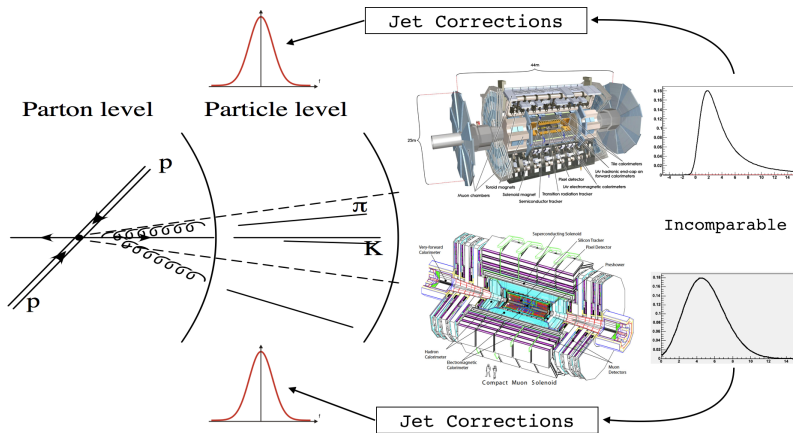
Jet can be defined on a *three different levels of collision*:

- ▶ **Parton level** - quarks, gluons and other particles created just after the collision. Directly connected to the QCD processes.
- ▶ **Particle level** - particles created by the hadronization.
- ▶ **Detector level** - recorded signal. Detector imperfections cause a *distortion of observables*.



# Jet Corrections

- ▶ Correct observables derived from detector to particle level by removing the detector effects
- ▶ Two main procedures - *Calibration* and *Unfolding*



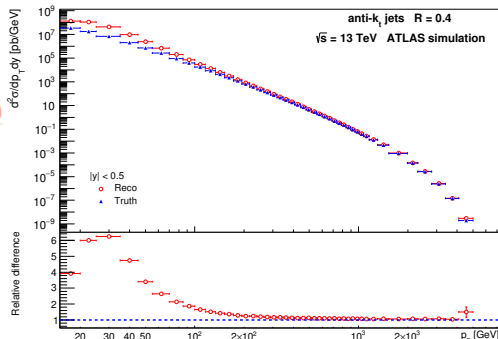
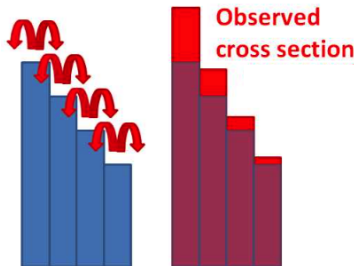
# Calibration

- Modifies the kinematic properties of *individual jets* - the most important correction: Energy
- Tries to minimize the calorimeter non-compensation, noise, losses in dead material and cracks, longitudinal leakage and particle deflection in magnetic field.
- *Universal* for each jet analysis. Uses the standard `APPLYJETCALIBRATION` library.

Jet index	0	1	2	3	4
pT (Before calibration) [GeV]	112.647	74.6027	69.601	39.5936	24.4818
pT (After calibration) [GeV]	150.576	115.018	103.515	54.6169	32.5758
E (Before calibration) [GeV]	120.755	163.125	147.191	327.249	105.504
E (After calibration) [GeV]	161.763	254.623	223.325	460.422	142.304
Eta (Before calibration)	0.348982	-1.41703	1.37864	-2.80151	2.1387
Eta (After calibration)	0.355368	-1.4309	1.40124	-2.82141	2.15266

# Unfolding

- Corrects the observables from detector level, to observables on particle level.
- Tries to minimize the effects of detector *finite resolution*.
- *Analysis dependent*.



# Unfolding - Mathematical Formulation

- ▶ **I want:**  $f(p_T)$  (distribution of inclusive jet  $p_T$  for  $p_T \in \langle a, b \rangle$ )
- ▶ From detector level **I have:**  $g(x)$  (distribution of unphysical variable  $x$ )

$$g(x) = \int_a^b A(x, p_T) f(p_T) dp_T$$

- ▶ Detector smearing described by  $A(x, p_T)$
- ▶ Complicated *integral equation* for  $f(p_T)$
- ▶ Luckily  $g(x)$  and  $f(p_T)$  are for practical purpose discretized and in analysis, I assume  $x \in \langle a, b \rangle$ ,  $N(i) \subset \langle a, b \rangle$

$$g_i = \int_{N(i)} g(x) dx \quad , \quad f_i = \int_{N(i)} f(p_T) dp_T$$

- ▶ So the response of the detector is described by a "simple" *matrix equation*, with  $A$  being called *Transfer Matrix*

$$g = Af$$

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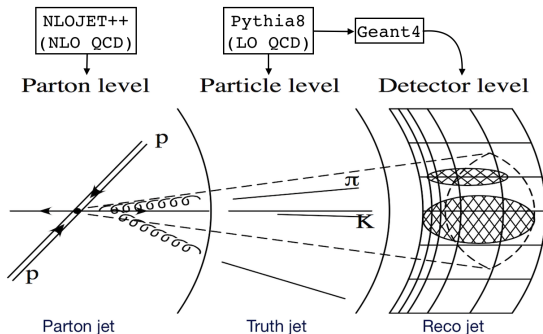
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# Data Characteristics

- ▶  $pp$  collisions at  $\sqrt{s} = 13$  TeV, anti- $k_t$  jet algorithm with  $R = 0.4$ , CT10 PDFs, AU2
- ▶ Measuring of inclusive jet double differential cross section in  $p_T$  and rapidity  $y$



- ▶ **Parton level** - cross section prediction calculated with NLOJET++ program (NLO QCD)
- ▶ **Particle level** - events generated by PYTHIA8 (LO QCD)
- ▶ **Detector level** - detector response on PYTHIA8 events obtained by GEANT4 full detector simulation.



# PYTHIA8 Data Characteristics

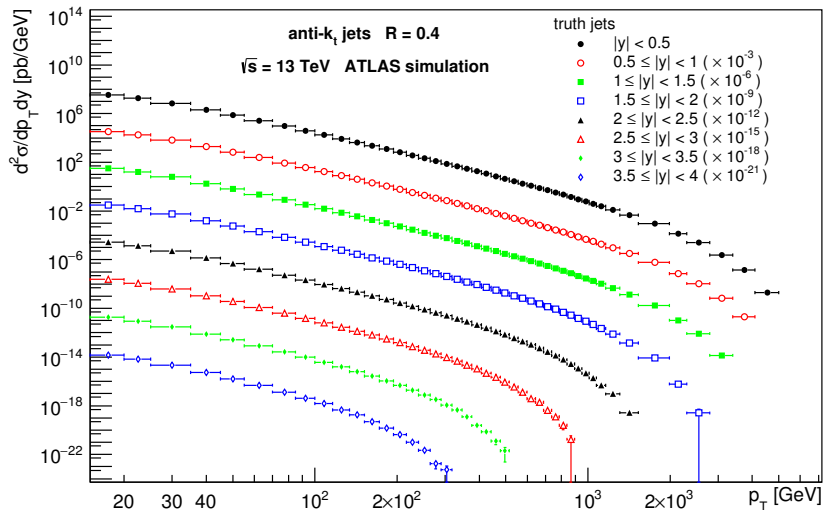
- ▶ Events were generated in a slices according to the leading truth jet  $p_T$ .
- ▶ Slices differ in event weight which is for all event calculated as

$$\text{weight} = \frac{(\text{Cross-section}) \cdot (\text{Filter Efficiency}) \cdot w_0}{(\# \text{ events})}$$

- ▶  $w_0$  is additional weight factor stored in `EventInfoAux` container

JZ	$p_T$ range (GeV)			Cross-section (fb)	Filter Efficiency	# events
JZ0W	0	-	20	7.8420e+13	9.7193e-01	3498000
JZ1W	20	-	80	7.8420e+13	2.7903e-04	2998000
JZ2W	80	-	200	5.7312e+10	5.2261e-03	500000
JZ3W	200	-	500	1.4478e+09	1.8068e-03	499500
JZ4W	500	-	1000	2.3093e+07	1.3276e-03	477000
JZ5W	1000	-	1500	2.3793e+05	5.0449e-03	499000
JZ6W	1500	-	2000	5.4279e+03	1.3886e-02	493500
JZ7W	2000	+		9.4172e+02	6.7141e-02	497000

# $p_T$ spectra of Truth Jets



# Event Selection

**Remove the jets** with  $p_T$  and rapidity  $y$  out of used binning

►  **$p_T$  Cut**

Reco and truth jets with  $p_T > 15 \text{ GeV}$  were kept.

►  **$y$  Cut**

Reco and truth jets with  $|y| < 4$  were kept.

**Remove the events** badly reconstructed by the detector

► **Zero Jet (0-jet) Cut**

Events with at least one reco and one truth jet, after the  $p_T$  and  $y$  cuts, are considered.

► **Leading Ratio (LR) Cut**

If  $0.6 < LR < 1.4$  the event is considered

$$LR = p_{T,\text{leading}}^{\text{reco}} / p_{T,\text{leading}}^{\text{truth}}$$

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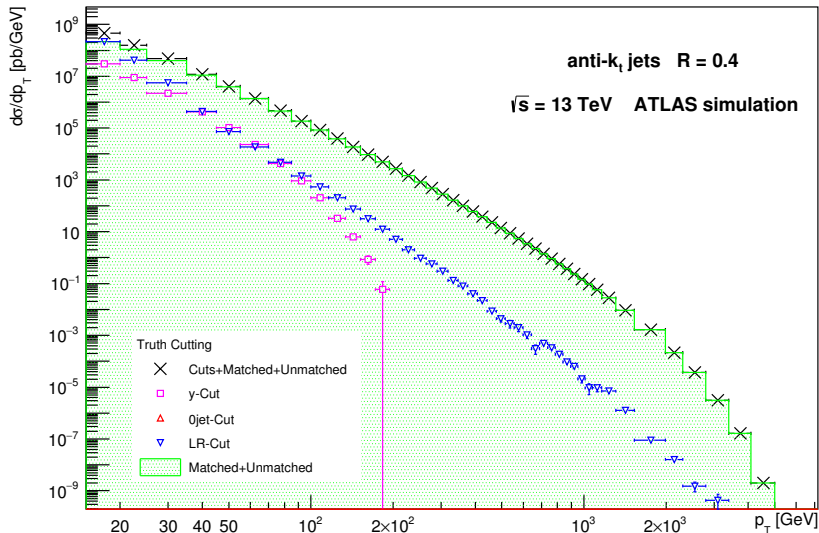
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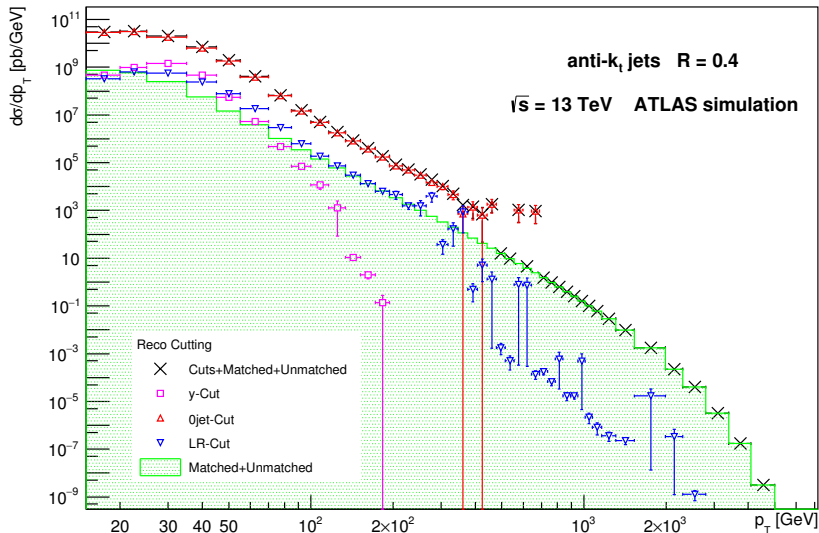
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# Event Selection - Truth Jets



# Event Selection - Reco Jets



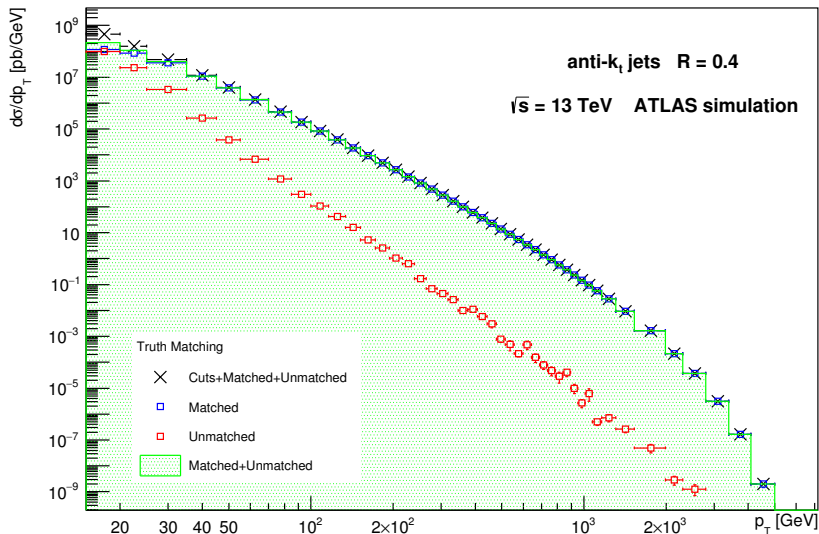
# Jet Matching

- ▶ In each event, for each truth jet, the corresponding reco jet has to be found.
- ▶ I have used *angular matching*
  1. For each pair  $(i,j)$  of reco and truth jets

$$dR_{ij} = \sqrt{d\phi_{ij}^2 + dy_{ij}^2}$$

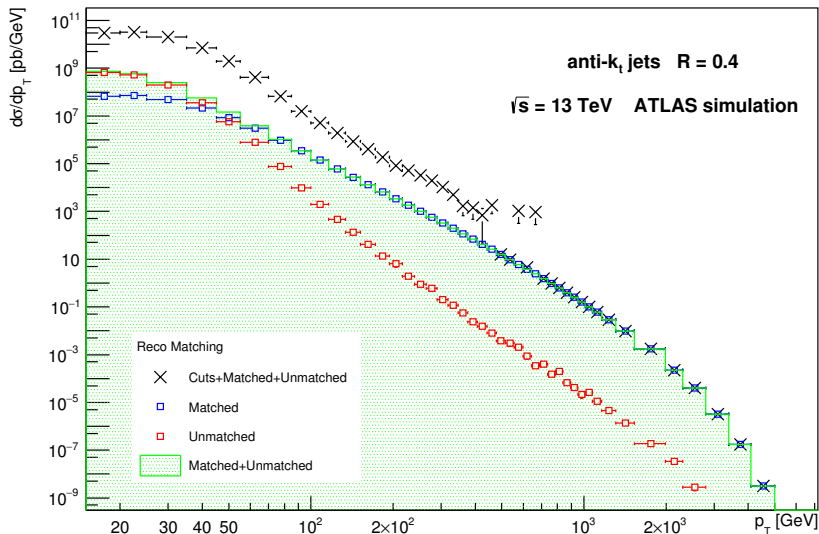
2. If  $\min(dR_{ij}) = dR_{pq} < dR^{cutoff} = 0.2$  the jets  $(p,q)$  were matched and further not assumed
3. Matching was done, when  $\min(dR_{ij}) < dR^{cutoff}$  was not satisfied or all of the reco or truth jets were matched.

# Jet Matching - Truth Jets





# Jet Matching - Reco Jets



# Inputs for Unfolding

Unfolding(calibrated reco spectrum)  $\approx$  truth spectrum

- ▶ Inputs for unfolding procedure
  - ▶ **Matching efficiencies** - describing the ratio of matched jets to all jets
  - ▶ **Transfer matrix**  $A_{ij}$  - containing the number of reco jets in bin  $i$  with a matched truth jets generated in bin  $j$
- ▶ I test two approaches to the unfolding, allowing a dealing with the double binning (in  $p_T$  and  $y$ )

## 1. Simple unfolding

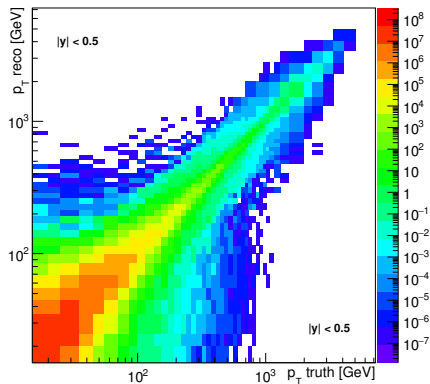
Matching jets within different rapidity bins is not allowed. There are **8 independent** 46x46 transfer matrices, one for each rapidity bin (46 = number of  $p_T$  bins)

## 2. 2D unfolding

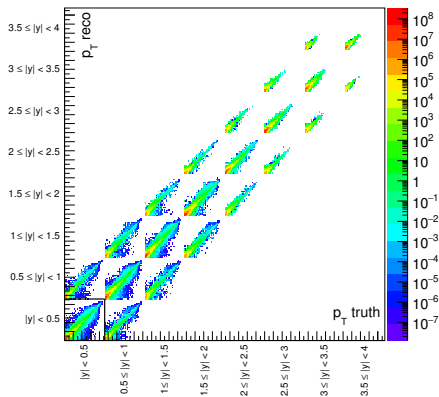
Matching within different rapidity bins allowed. **Only one** 368x368 transfer matrix ( $368 = 8 \times 48$ )

# Transfer Matrices

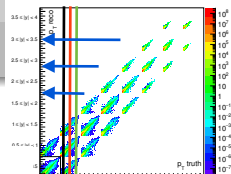
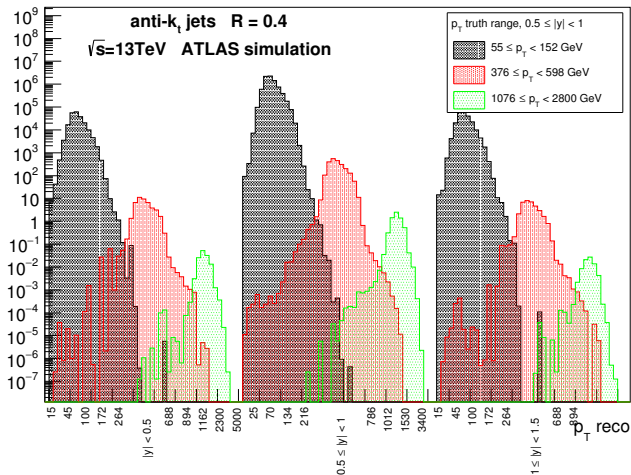
Simple unfolding



2D unfolding

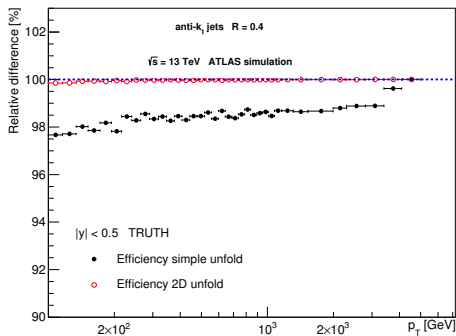


# Slices in Transfer Matrix of 2D Unfolding

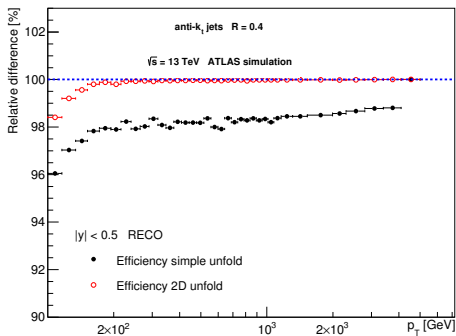


# Matching Efficiencies

Truth jets



Reco jets



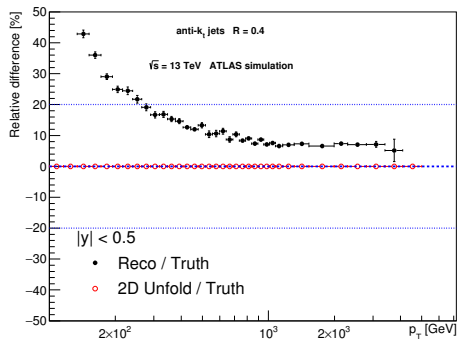
# Steps of Unfolding

## Three main steps of the unfolding procedure

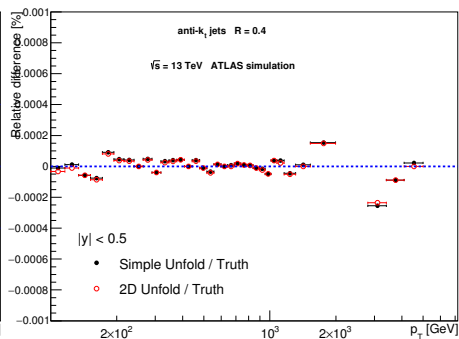
1. Input data are multiplied by the matching efficiencies of reco jets
2. Transfer matrix is used to correct data spectrum for detector effects. I use the *Iterative Dynamical Stabilized* unfolding method with one iteration
3. The spectrum obtained by the step 2 is divided by the matching efficiencies of truth jets, in order to correct resulting spectrum for the unmatched truth jets

# Unfolding Results

## Reco and Unfolded vs. Truth Spectrum



## Simple and 2D unfolded vs. Truth Spectrum



# NLO QCD Prediction

- ▶ NLO QCD predictions on parton level for  $\sqrt{s} = 8 \text{ TeV}$  and  $\sqrt{s} = 13 \text{ TeV}$
- ▶ *Theoretical uncertainties* which are taken into account
  - ▶ **Scale uncertainty**  
Choice of renormalization and factorization scales, including neglecting the higher order terms beyond the NLO
  - ▶  $\alpha_S$  **uncertainty**  
Because of experimental measurements of  $\alpha_S$ .
  - ▶ **PDF uncertainty**  
Prediction depends on the concrete choice of a PDF
- ▶ *Other corrections* (not so significant)
  - ▶ **Nonperturbative corrections**  
Hadronization and Underlying Event corrections.
  - ▶ **Electroweak corrections**  
Next to the QCD processes, the electroweak processes should be assumed.

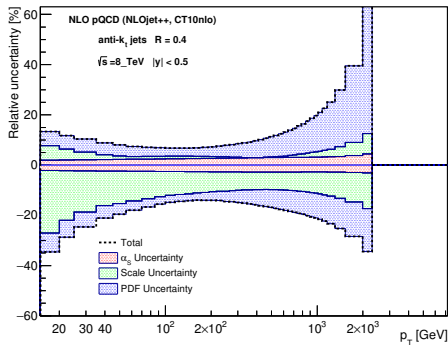


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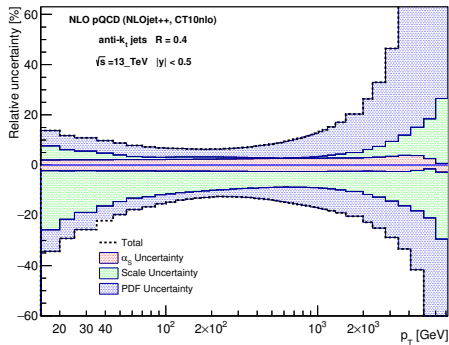
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# NLO Systematic Errors

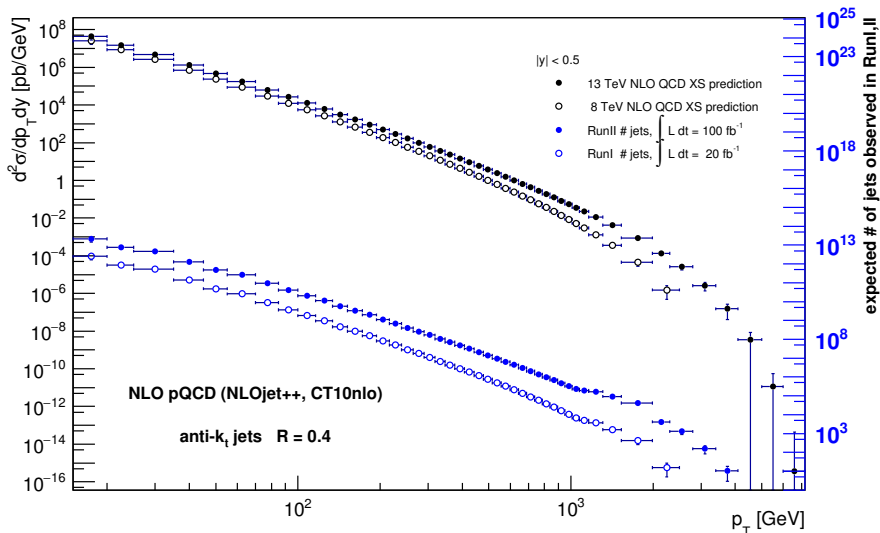
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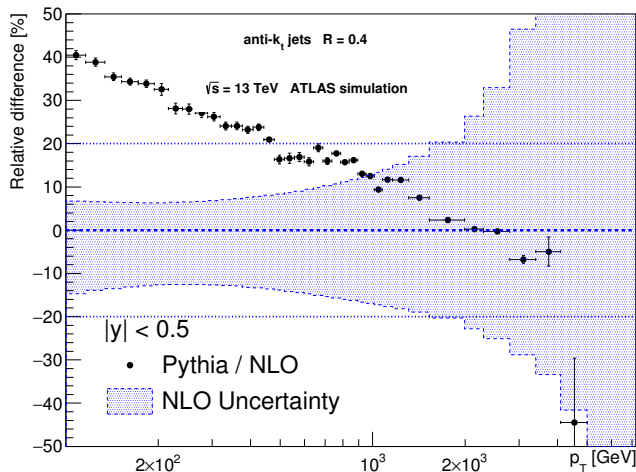
$$\sqrt{s} = 13 \text{ TeV}$$



# Comparison of NLO QCD Predictions



# Comparison of LO and NLO QCD



# Thesis Conclusions

## Unfolding

Two approaches were probed.

No significant differences between these two approaches imply, for the real analysis, the **Simple Unfolding approach should be used** for its simpler implementation.

Agreement of the unfolded  $p_T$  spectra with the truth  $p_T$  spectra up to systematic error  $< 10^{-3} \%$ .

## LO and NLO QCD

Significant differences showing the influence of the NLO QCD processes on physical observables.