A Splitting Approximation for the Solution of a Self-Adjoint Quenching Problem

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Outline

- System description
- Governing equations
- Literature review on recent progress on similar subject
- Our goal
- Discretization schemes
- Positivity, monotonocity and linear stability
- Simulation results
- Conclusion and Future work

System description

- Ignition process: rate of change of the temperature blows up in a very short time interval.
- Fuel combustion chamber as an example
- A non-linear advection-diffusion equation is used to model the temperature dynamics where its first derivative tends to infinity as it reaches a certain critical point.

Advection-diffusion equation

• Let b > 0 and $\tilde{\mathcal{D}} = (0, b) \times (0, b)$ be a squared domain, with $\partial \tilde{\mathcal{D}}$ being its boundary. We consider the following semilinear advection-diffusion problem,

$$\begin{array}{rcl} u_t & = & \nabla(a\nabla u) + f(u), & (\tilde{x},\tilde{y}) \in \tilde{\mathcal{D}}, \ t > 0, \\ u(\tilde{x},\tilde{y},t) & = & 0, & (\tilde{x},\tilde{y}) \in \partial \tilde{\mathcal{D}}, \ t \geq 0, \\ u(\tilde{x},\tilde{y},0) & = & u_0(\tilde{x},\tilde{y}), & (\tilde{x},\tilde{y}) \in \tilde{\mathcal{D}}, \end{array}$$

• The nonlinear source function f(u) has following properties:

$$f(0) = f_0 > 0;$$

 $f(u)$ is strictly increasing for $0 \le u < 1;$
 $\lim_{u \to 1^-} f(u) = \infty.$



Normalized equation

• Re-scaling the space domain

$$u_t = b^{-2} \nabla(a\nabla u) + f(u), (x, y) \in \mathcal{D}, t > 0, (1.1)$$

$$u(x, y, t) = 0, \quad (x, y) \in \partial \mathcal{D}, \quad t \ge 0,$$
 (1.2)

$$u(x, y, 0) = u_0(x, y), (x, y) \in \mathcal{D},$$
 (1.3)

where $\mathcal{D} = (0, 1) \times (0, 1)$.

Semi-discretized equation

Discrization in space only:

$$(v_{t})_{k,\ell} = \frac{1}{b^{2}h^{2}} \left[a_{k-\frac{1}{2},\ell} v_{k-1,\ell} + a_{k+\frac{1}{2},\ell} v_{k+1,\ell} \right. \\ \left. - \left(a_{k-\frac{1}{2},\ell} + a_{k+\frac{1}{2},\ell} \right) v_{k,\ell} \right. \\ \left. + a_{k,\ell-\frac{1}{2}} v_{k,\ell-1} + a_{k,\ell+\frac{1}{2}} v_{k,\ell+1} \right. \\ \left. - \left(a_{k,\ell-\frac{1}{2}} + a_{k,\ell+\frac{1}{2}} \right) v_{k,\ell} \right] \\ \left. + f(v_{k,\ell}) \right.$$

for $1 \leq k, \ell \leq N$.



Semi-discrete equation matrix form

Initial value problem

$$v' = (P+Q)v + g(v), \quad t > 0,$$
 (1.4)

$$v(0) = v_0,$$
 (1.5)

in which

$$P = \frac{1}{b^2 h^2} \operatorname{diag}(P_1, P_2, \dots, P_N)$$

$$Q=rac{1}{b^2h^2}$$
tridiag $\left(Q_j^1,\,Q_j^2,\,Q_j^3
ight)$

Exact solution of the semi-discrete equation

• The exact solution of the initial value problem (1.4), (1.5) is

$$v(t) = e^{t(P+Q)}v_0 + \int_0^t e^{(t-\tau)(P+Q)}g(v)d\tau,$$

$$0 \le t \le T < \infty.$$

• The Peaceman-Rachford splitting is used to approximate the exponential, i.e $e^{t(P+Q)} = S(t) + \mathcal{O}(t^3)$ where

$$S(t) = (I - \frac{t}{2}Q)^{-1}(I - \frac{t}{2}P)^{-1}(I + \frac{t}{2}P)(I + \frac{t}{2}Q),$$

and the trapezoidal rule for approximating the integral. The following scheme is then obtained where τ is a fixed time step:



Discretized equation

Full discretized equation

$$v^{n+1} = (1.6)$$

$$(I - \frac{\tau}{2}Q)^{-1}(I - \frac{\tau}{2}P)^{-1}(I + \frac{\tau}{2}P)(I + \frac{\tau}{2}Q)\left(v^n + \frac{\tau}{2}g(v^n)\right)$$
(1.7)

$$+\frac{\tau}{2}g(w^{(n)})+\mathcal{O}(\tau^2). \tag{1.8}$$

where an Euler scheme is used for the approximation $v^{n+1} \approx w^{(n)} = v^n + \tau((Q+P)v^n + g(v^n))$



Monotonicity and positivity: lemma 1

• Lemma 1. $\max(\|P\|, \|Q\|) \leq \frac{4}{b^2h^2} \max_{i,j} \{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\}.$ Proof: We have: $\|P\| = \frac{1}{b^2h^2} \max_j \{\max\{u_{1,j} - m_{1,j}, I_{N-1,j} - m_{N,j}, \max_{i=2,...N-1}(I_{i,j} - m_{i,j} + u_{i,j+1})\}\} \leq \frac{4}{b^2h^2} \max_i \{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\}.$

Similarly the same bound can be obtained for ||Q||.

Monotonicity and positivity: lemma 2

• Lemma 2. If $\frac{\tau}{h^2} < \frac{b^2}{2\max_{i,j}\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\}}$ then the matrices $(I-\frac{\tau}{2}Q),(I-\frac{\tau}{2}P),(I+\frac{\tau}{2}P)$ and $(I+\frac{\tau}{2}Q)$ are non-singular. Also $(I-\frac{\tau}{2}R)$ and $(I-\frac{\tau}{2}P)$ are monotone and inverse positive. $(I+\frac{\tau}{2}Q)$ and $(I+\frac{\tau}{2}P)$ are nonnegative. Similarly, If $\frac{\tau}{h^2} < \frac{b^2}{8\max_{i,j}\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\}}$, then $I+\tau(P+Q)$ is nonsingular as well as nonnegative.

Proof: Lemma 2

Proof.

According to lemma 1

$$\|\frac{\tau}{2}P\| \leq \frac{2\tau}{b^2h^2} \max_{i,j} \{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\} < 1.$$

Hence $(I + \frac{\tau}{2}P)$ is nonsingular [2]. It can be

observed that $(I + \frac{\tau}{2}P)$ is also nonnegative. A similar argument shows that $(I + \frac{\tau}{2}Q)$ is nonsingular and nonnegative.

Again according to lemma (10) we have

$$\|\tau(P+Q)\| \le \tau\left(\|P\| + \|Q\|\right) \le \frac{8\tau}{h^2b^2} \max_{i,j} \{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\} < 1$$

Therefore $I + \tau (P + R)$ is nonsingular and nonnegative.

Consider the matrix $M=I-\frac{\tau}{2}P$. As $M_{i,j}\leq 0$ for $i\neq j$ and the weak row sum criterion is satisfied M is monotone and hence an inverse exists which is nonnegative [2]. Using a similar argument shows that $I-\frac{\tau}{2}Q$ is inverse-positive.



Monotonicity and positivity: lemma 3-4

• Lemma3 Let $(P+Q)v^l+g(v^l)+\frac{\tau^2}{4}PQg(v^l)>0$, where $l\geq 0$ is any begining time step. For uniform time steps τ , if $\frac{\tau}{h^2}<\frac{b^2}{2\max_{i,j}\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\}}$ then $v^{k+1}>v^k$ for all $k\geq l$. The sequence $\{v^k\}_{k=l}^\infty$ is therefore monotonically increasing.

• lemma4 If $v_0=0$, and $\frac{\tau}{h^2}<\frac{b^2}{\max_{i,j}\left\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\right\}}$, then $(P+Q)v_0+g(v_0)+\frac{\tau_0^2}{4}PQg(v_0)>0.$

Monotonicity and positivity: lemma 5 and theorem

$$\begin{array}{l} \bullet \ \ \mathsf{Lemma5} \ \ \mathsf{If} \ \frac{\tau_0}{h^2} < \frac{b^2}{2 \max_{i,j} \big\{ a_{i\pm\frac{1}{2},j}, \, a_{i,j\pm\frac{1}{2}} \big\}} \\ \\ h^2 < \frac{1}{4 M_{\frac{b^2}{2 \max_{i,j} \big\{ a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}} \big\}}} \\ \end{array}$$

then given $\mathbf{v_0} = 0$ all the components of $\mathbf{v^1} < 1$. where $M = f(\tau_0 f_0)$.

Theorem For any beginning step τ , if:

(i)
$$\frac{\tau}{h^2} < \frac{b^2}{2 \max_{i,j} \{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\}}$$

(ii) $(P+R)v_0 + g(v_0) + \frac{\tau^2}{4} PRg(v_0) > 0$

(ii)
$$(P+R)v_0 + g(v_0) + \frac{\tau^2}{4}PRg(v_0) > 0$$

then the sequence of solutions $(v^k)_{k>l}$ produced by the Peaceman-Rachford-Strang splitting increases monotonically until unity is exceed by at least component of the solution vector(i.e until quenching occurs).



Linear stability

Stability is a challenge while solving non-linear blow-up type or quenching-type problem such as the one treated in this work. Analysis by mean of linear stability is not a vigorous process, but it is still a useful and effective practice. This linear stability method freezes the non linear term. In other words the non linear term is kept constant.

Linear stability: lemma 6

 $||(I - \frac{\tau}{2}R)^{-1}|| \le 1.$

• Lemma 6 $\|(I-\frac{\tau}{2}P)^{-1}\|$, $\|(I-\frac{\tau}{2}Q)^{-1}\| \leq 1$ proof: Consider the entries of the matrix $(I-\frac{\tau_k}{2}P)$: $\alpha_j^{(i)} = 1 - \frac{\tau}{2}(m_j + u_j + l_{j-1})$ Since the matrix P is diagonal dominant with diagonal entries $m_j < 0$, then $\alpha_j^{(i)} \geq 1$. The desired result is obtained by using the Varah-bound [1]. Similarly, $\beta_i^{(i)} = 1 - \frac{\tau}{2}(m_i + u_i + l_{i-1}) \geq 1$, and as a consequence

Linear stability: lemma 7

• Lemma 7. If
$$\frac{\tau}{h^2} < \frac{b^2}{2\max_{i,j} \left\{a_{i\pm\frac{1}{2},j}, a_{i,j\pm\frac{1}{2}}\right\}}$$
 Then $\|(I+\frac{\tau}{2}P)(I+\frac{\tau}{2}R)\|=1$ Proof Since the matrices $(I+\frac{\tau}{2}P)$, and $(I+\frac{\tau}{2}R)$ are non-negative, so is the product $(I+\frac{\tau}{2}P)(I+\frac{\tau}{2}R)$. For any non-negative vector such that $\mathbf{x} \leq 1$, we have $0 \leq (I+\frac{\tau}{2}(P+R)+\frac{\tau^2}{4}PR)\mathbf{x} \leq 1$. Since $\frac{\tau}{2}(P+R)+\frac{\tau^2}{4}PR)\mathbf{x}$ has some zero entries, we can deduce that $\|(I+\frac{\tau}{2}P)(I+\frac{\tau}{2}R)\|_{\infty}=1$.

Linear stability: lemma 8

• Lemma 8. Let $\frac{\tau}{h^2}<\frac{b^2}{2\max_{i,j}\left\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\right\}}.$ Then the

Peaceman-Rachford spliting is weakly stable in the Von Neumann sense.

Proof,

Consider a perturbation at time step k, $z_{k+1} = v_k - \tilde{v_k}$, where $\tilde{v_k}$ is the computed solution. By lemmas (16)-(17) the norm is bounded.

$$\begin{aligned} \|\mathbf{z}_{k+1}\| &= \|(I - \frac{\tau}{2}R)^{-1}(I - \frac{\tau}{2}P)^{-1}(I + \frac{\tau}{2}P)(I + \frac{\tau}{2}R)\mathbf{z}_{k}\| \\ &\leq \|(I - \frac{\tau}{2}R)^{-1}(I - \frac{\tau}{2}P)^{-1}(I + \frac{\tau}{2}P)(I + \frac{\tau}{2}R)\|\|\mathbf{z}_{k}\| \leq \|\mathbf{z}_{k}\|. \end{aligned}$$

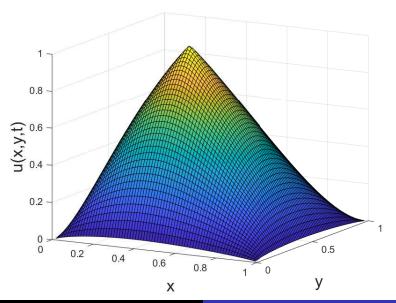
Linear stability: theorem

• Theorem: Let $\frac{\tau}{h^2} < \frac{b^2}{2\max_{i,j} \left\{a_{i\pm\frac{1}{2},j},a_{i,j\pm\frac{1}{2}}\right\}}$ for all $i \leq k$.

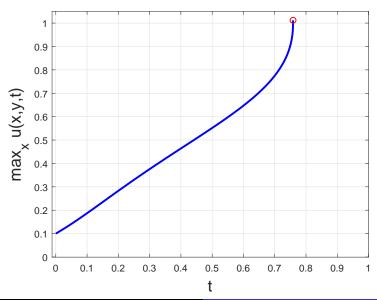
Then the linearized variable step Peace-Rachford method is weakly stable in the von neumann sense under the l_{∞} norm, i.e $\|\mathbf{z}_{k+1}\| \leq \|\mathbf{z}_0\|$.

where $\mathbf{z}_0 = \mathbf{v}_0 - \bar{\mathbf{v}}_0$ is an initial perturbation arising from the initial perturbation \mathbf{z}_0 .

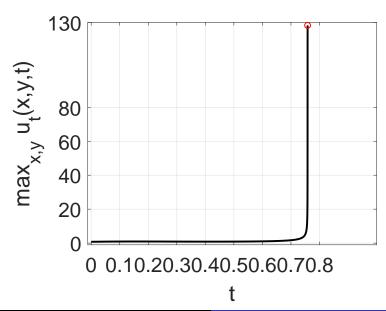
Numerical solution immediately before quenching time



Maximum temperature values

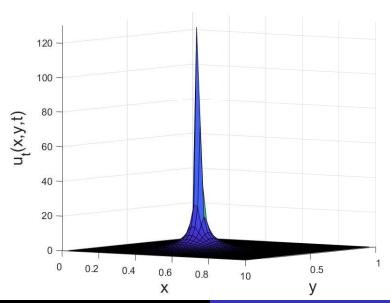


Maximum temperature derivative values





Temperature derivative values



Summary and future work

- We have derived conditions for the physical properties of the quenching equation with variable diffusion coeficient: positivity, monotonocity, and also linear stability.
- Using uniform grid and constant time step we showed some numerical results.
- One of our future goal is obtaining more accurate numerical results by using non-uniform grid and variable time step to capture the quenching location and time.
- Nonlinear stability is another important future goal for us to achieve

Thank you for your attention!



A lower bound for the smalest singular value of a matrix. *Linear Algebra Appl*, 11:3–5.



Discrete variable Methods in Ordinary Differential Equations. John Willey & Sons, 1962.