Problem 1.   A spherical shell of radius R has a total charge Q uniformly spread over the shell. The shell is now put into uniform rotation about the z axis with angular velocity ω. Find the vector potential $\mathbf{A}(\mathbf{x})$ and magnetic field $\mathbf{B}(\mathbf{x})$ everywhere, i.e., both inside and outside of the shell.

Solution.   Let $\rho(\mathbf{x})$ be the charge density everywhere in space, so

$$\rho(\mathbf{x}) = \frac{1}{4\pi} \frac{Q}{R^2} \delta(r - R).$$

The linear velocity of the moving charge everywhere is

$$\mathbf{v}(\mathbf{x}) = \omega r \, \delta(r - R) \, \hat{\boldsymbol{\varphi}}.$$

Then the current density $\mathbf{J}$ is simply the product of charge density and the linear velocity of the charge:

$$\mathbf{J}(\mathbf{x}) = \rho(\mathbf{x}) \, \mathbf{v}(\mathbf{x}) = \frac{Q\omega}{4\pi} \frac{r}{R^2} \delta(r - R) \, \hat{\boldsymbol{\varphi}}.$$

From Eq. (4.21) in the lecture notes, $\mathbf{A}(\mathbf{x})$ everywhere is given by

$$\mathbf{A}(\mathbf{x}) = \frac{1}{c} \int \frac{\mathbf{J}(\mathbf{x})'}{|\mathbf{x} - \mathbf{x}'|} \, d^3\mathbf{x}' \, .$$

The integral we need to evaluate is then

$$\mathbf{A}(\mathbf{x}) = \frac{1}{4\pi c} \frac{Q\omega}{R^2} \int \frac{r' \, \delta(r' - R)}{|\mathbf{x} - \mathbf{x}'|} \, d^3\mathbf{x}' \, .$$

The problem is azimuthally symmetric, so we will rotate our coordinate system such that $\mathbf{x}$ points along the z axis. In the new coordinate system,

$$\frac{1}{|\mathbf{x} - \mathbf{x}'|} = \frac{1}{\sqrt{\mathbf{x}^2 - 2\mathbf{x} \cdot \mathbf{x}' + \mathbf{x}'^2}} = \frac{1}{\sqrt{r^2 - 2rr' \cos\theta' + r'^2}}.$$

Let $\boldsymbol{\omega}$ be the angular velocity vector (that lay along the z axis of the original coordinate system), which we choose to lie in the xz plane. Let α be the angle between $\boldsymbol{\omega}$ and the z axis. Then the linear velocity of the moving charge is

$$\mathbf{v}(\mathbf{x}') = \boldsymbol{\omega} \times \mathbf{x}' = \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ \omega \sin\alpha & 0 & \omega \cos\alpha \\ r' \sin\theta' \cos\varphi' & r' \sin\theta' \sin\varphi' & r' \cos\theta' \end{vmatrix}$$
$$= -\omega r'(\cos\alpha \sin\theta' \sin\varphi') \, \hat{\mathbf{x}} + \omega r'(\cos\alpha \sin\theta' \cos\varphi' - \sin\alpha \cos\theta') \, \hat{\mathbf{y}} + \omega r'(\sin\alpha \sin\theta' \sin\varphi') \, \hat{\mathbf{z}},$$

so in the new coordinate system,

$$\mathbf{J}(\mathbf{x}') = \frac{Q}{4\pi} \frac{\boldsymbol{\omega} \times \mathbf{x}'}{R^2} \delta(r' - R) = \frac{Q\omega}{4\pi} \frac{r'}{R^2} (\hat{\boldsymbol{\omega}} \times \hat{\mathbf{x}}') \, \delta(r' - R),$$

where
$$\hat{\boldsymbol{\omega}} \times \hat{\mathbf{x}}' = -(\cos\alpha \sin\theta' \sin\varphi') \, \hat{\mathbf{x}} + (\cos\alpha \sin\theta' \cos\varphi' - \sin\alpha \cos\theta') \, \hat{\mathbf{y}} + (\sin\alpha \sin\theta' \sin\varphi') \, \hat{\mathbf{z}}.$$

The integral we need to evaluate becomes

$$\mathbf{A}(\mathbf{x}) = \frac{1}{4\pi c} \frac{Q\omega}{R^2} \int_0^{2\pi} \int_{-1}^1 \int_0^\infty \frac{r'^3 (\hat{\boldsymbol{\omega}} \times \hat{\mathbf{x}}') \, \delta(r' - R)}{\sqrt{r^2 - 2rr' \cos\theta' + r'^2}} \, dr' \, d(\cos\theta') \, d\varphi' \, .$$

Evaluating the radial integral, we have

$$\mathbf{A}(\mathbf{x}) = \frac{Q\omega R}{4\pi c} \int_0^{2\pi} \int_{-1}^1 \frac{\hat{\boldsymbol{\omega}} \times \hat{\mathbf{x}}'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta') \, d\varphi \,.$$

For the angular integrals, the $\hat{\mathbf{x}}$ term is

$$-\cos\alpha \, \hat{\mathbf{x}} \int_{-1}^1 \frac{\sin\theta'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta') \int_0^{2\pi} \sin\varphi' \, d\varphi \propto \left[-\cos\varphi' \right]_0^{2\pi} = 0.$$

Similarly, the $\hat{\mathbf{z}}$ term is

$$\sin\alpha \, \hat{\mathbf{z}} \int_{-1}^1 \frac{\sin\theta'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta') \int_0^{2\pi} \sin\varphi' \, d\varphi \propto \left[-\cos\varphi' \right]_0^{2\pi} = 0.$$

There are two $\hat{\mathbf{y}}$ terms. For the first,

$$\cos\alpha \, \hat{\mathbf{y}} \int_{-1}^1 \frac{\sin\theta'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta') \int_0^{2\pi} \cos\varphi' \, d\varphi \propto \left[\sin\varphi' \right]_0^{2\pi} = 0.$$

For the second,

$$-\sin\alpha \, \hat{\mathbf{y}} \int_{-1}^1 \frac{\cos\theta'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta') \int_0^{2\pi} d\varphi = -2\pi\sin\alpha \, \hat{\mathbf{y}} \int_{-1}^1 \frac{\cos\theta'}{\sqrt{r^2 - 2Rr\cos\theta' + R^2}} \, d(\cos\theta')$$

$$= -2\pi\sin\alpha \, \hat{\mathbf{y}} \left(\left[-\frac{\cos\theta'\sqrt{r^2 - 2Rr\cos\theta' + R^2}}{Rr} \right]_{-1}^1 + \frac{1}{Rr} \int_{-1}^1 \sqrt{r^2 - 2Rr\cos\theta' + R^2} \, d(\cos\theta') \right)$$

$$= -2\pi\sin\alpha \, \hat{\mathbf{y}} \left(\left[-\frac{\cos\theta'\sqrt{r^2 - 2Rr\cos\theta' + R^2}}{Rr} \right]_{-1}^1 + \frac{1}{Rr} \left[-\frac{(r^2 - 2Rr\cos\theta' + R^2)^{3/2}}{3Rr} \right]_{-1}^1 \right)$$

$$= -2\pi\sin\alpha \, \hat{\mathbf{y}} \left(-\frac{\sqrt{r^2 + 2Rr + R^2}}{Rr} + \frac{\sqrt{r^2 - 2Rr + R^2}}{Rr} - \frac{(r^2 - 2Rr + R^2)^{3/2}}{3R^2r^2} + \frac{(r^2 + 2Rr + R^2)^{3/2}}{3R^2r^2} \right)$$

$$= 2\pi\sin\alpha \frac{3Rr\sqrt{(r+R)^2} - 3Rr\sqrt{(r-R)^2} + [(r-R)^2]^{3/2} - [(r+R)^2]^{3/2}}{3R^2r^2} \hat{\mathbf{y}}$$

$$= 2\pi\sin\alpha \frac{3Rr|r+R| - 3Rr|r-R| + (r-R)^2|r-R| - (r+R)^2|r+R|}{3R^2r^2} \hat{\mathbf{y}}$$

$$= 2\pi\sin\alpha \frac{(r^2 + Rr + R^2)|r-R| - (r^2 - Rr + R^2)(r+R)}{3R^2r^2} \hat{\mathbf{y}}$$

$$= \frac{2\pi\sin\alpha \, \hat{\mathbf{y}}}{3R^2r^2} \begin{cases} (r^2 + Rr + R^2)(R-r) - (r^2 - Rr + R^2)(r+R) & r < R, \\ (r^2 + Rr + R^2)(r-R) - (r^2 - Rr + R^2)(r+R) & r > R \end{cases}$$

$$= -\frac{4}{3}\pi\sin\alpha \, \hat{\mathbf{y}} \begin{cases} \dfrac{r}{R^2} & r < R, \\ \dfrac{R}{r^2} & r > R. \end{cases}$$

Finally, in the new coordinate system we have

$$\mathbf{A}(\mathbf{x}) = -\frac{Q\omega}{3c} \sin\alpha \, \hat{\mathbf{y}} \begin{cases} \dfrac{r}{R} & r < R, \\ \dfrac{R^2}{r^2} & r > R. \end{cases}$$

Transforming back to the old coordinate system, $\sin\alpha \to -\sin\theta$. Since the original system is azimuthally symmetric, $\varphi = 0$ so $\hat{\mathbf{y}} = \sin\theta\sin\varphi\,\hat{\mathbf{r}} + \cos\theta\sin\varphi\,\hat{\boldsymbol{\theta}} + \cos\varphi\,\hat{\boldsymbol{\varphi}} = \hat{\boldsymbol{\varphi}}$. Thus we have

$$\mathbf{A}(\mathbf{x}) = \frac{Q\omega}{3c}\sin\theta\,\hat{\boldsymbol{\varphi}} \begin{cases} \dfrac{r}{R} & r < R, \\[2mm] \dfrac{R^2}{r^2} & r > R. \end{cases}$$

The magnetic field is given by Eq. (1.7),

$$\mathbf{B} = \boldsymbol{\nabla} \times \mathbf{A}. \tag{1}$$

In spherical coordinates,

$$\boldsymbol{\nabla}\times\mathbf{A} = \frac{1}{r\sin\theta}\left(\frac{\partial}{\partial\theta}(\sin\theta A_\varphi) - \frac{\partial A_\theta}{\partial\varphi}\right)\hat{\mathbf{r}} + \frac{1}{r}\left(\frac{1}{\sin\theta}\frac{\partial A_r}{\partial\varphi} - \frac{\partial}{\partial r}(rA_\varphi)\right)\hat{\boldsymbol{\theta}} + \frac{1}{r}\left(\frac{\partial}{\partial r}(rA_\theta) - \frac{\partial A_r}{\partial\theta}\right)\hat{\boldsymbol{\varphi}},$$

so

$$\mathbf{B} = \frac{1}{r\sin\theta}\frac{\partial}{\partial\theta}(\sin\theta A_\varphi)\hat{\mathbf{r}} - \frac{1}{r}\frac{\partial}{\partial r}(rA_\varphi)\hat{\boldsymbol{\theta}}.$$

For $r < R$,

$$\mathbf{B}(\mathbf{x}) = \frac{Q\omega}{3c}\frac{1}{r}\left[\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\left(\frac{r}{R}\sin^2\theta\right)\hat{\mathbf{r}} - \frac{\partial}{\partial r}\left(\frac{r^2}{R}\sin\theta\right)\hat{\boldsymbol{\theta}}\right] = \frac{Q\omega}{3c}\frac{1}{r}\left(\frac{r}{R}\frac{2\cos\theta\sin\theta}{\sin\theta}\hat{\mathbf{r}} - \frac{2r}{R}\sin\theta\,\hat{\boldsymbol{\theta}}\right)$$

$$= \frac{2}{3}\frac{Q\omega}{cR}(\cos\theta\,\hat{\mathbf{r}} - \sin\theta\,\hat{\boldsymbol{\theta}}) = \frac{2}{3}\frac{Q\omega}{cR}\,\hat{\mathbf{z}}.$$

For $r > R$,

$$\mathbf{B}(\mathbf{x}) = \frac{Q\omega}{3c}\frac{1}{r}\left[\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\left(\frac{R^2}{r^2}\sin^2\theta\right)\hat{\mathbf{r}} - \frac{\partial}{\partial r}\left(\frac{R^2}{r}\sin\theta\right)\hat{\boldsymbol{\theta}}\right] = \frac{Q\omega}{3c}\frac{1}{r}\left(\frac{R^2}{r^2}\frac{2\cos\theta\sin\theta}{\sin\theta}\hat{\mathbf{r}} + 2\frac{R^2}{r^2}\sin\theta\,\hat{\boldsymbol{\theta}}\right)$$

$$= \frac{2}{3}\frac{Q\omega}{c}\frac{R^2}{r^3}(\cos\theta\,\hat{\mathbf{r}} + \sin\theta\,\hat{\boldsymbol{\theta}}).$$

In summary,

$$\mathbf{B}(\mathbf{x}) = \frac{2}{3}\frac{Q\omega}{c}\begin{cases} \dfrac{\hat{\mathbf{z}}}{R} & r < R, \\[3mm] \dfrac{R^2}{r^3}(\cos\theta\,\hat{\mathbf{r}} + \sin\theta\,\hat{\boldsymbol{\theta}}) & r > R. \end{cases}$$

Problem 2. If an electric and magnetic field are both present, the momentum density carried by the electromagnetic field is given by Poynting's formula

$$\mathcal{P} = \frac{1}{4\pi c}(\mathbf{E}\times\mathbf{B}).$$

Consider a bounded distribution of time-independent charges and currents, i.e., $\rho(\mathbf{x})$ and $\mathbf{J}(\mathbf{x})$ are time independent and vanish when $|\mathbf{x}| > R$ for some R.

2.a Show that the total momentum can be written as

$$\mathbf{P} \equiv \int \mathcal{P}(\mathbf{x})\,d^3x = \int \phi(\mathbf{x})\,\mathbf{J}(\mathbf{x})\,d^3x.$$

Solution. Applying (1),
$$\mathbf{E} \times \mathbf{B} = -\mathbf{E} \times (\mathbf{\nabla} \times \mathbf{A}).$$

Equation (4.12) gives the vector calculus identity

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = \mathbf{b}(\mathbf{a} \cdot \mathbf{c}) - \mathbf{c}(\mathbf{a} \cdot \mathbf{b}),$$

so we find
$$-\mathbf{E} \times \mathbf{B} = \mathbf{\nabla}(\mathbf{E} \cdot \mathbf{A}) - \mathbf{A}(\mathbf{E} \cdot \mathbf{\nabla}) = \mathbf{\nabla}(\mathbf{E} \cdot \mathbf{A}) - \mathbf{A}\nabla^2\phi. \tag{2}$$

Another vector identity from the inside cover of Jackson is

$$\mathbf{\nabla}(\mathbf{a} \cdot \mathbf{b}) = (\mathbf{a} \cdot \mathbf{\nabla})\mathbf{b} + (\mathbf{b} \cdot \mathbf{\nabla})\mathbf{a} + \mathbf{a} \times (\mathbf{\nabla} \times \mathbf{b}) + \mathbf{b} \times (\mathbf{\nabla} \times \mathbf{a}),$$

so

$$\begin{aligned}
\mathbf{\nabla}(\mathbf{\nabla}\phi \cdot \mathbf{A}) &= (\mathbf{\nabla}\phi \cdot \mathbf{\nabla})\mathbf{A} + (\mathbf{A} \cdot \mathbf{\nabla})\mathbf{\nabla}\phi + \mathbf{\nabla}\phi \times (\mathbf{\nabla} \times \mathbf{A}) + \mathbf{A} \times (\mathbf{\nabla} \times \mathbf{\nabla}\phi) \\
&= \mathbf{A}\nabla^2\phi + \mathbf{\nabla}\phi(\mathbf{\nabla} \cdot \mathbf{A}) - \mathbf{E} \times \mathbf{B},
\end{aligned}$$

since $\mathbf{\nabla} \times \mathbf{\nabla}\phi = -\mathbf{\nabla} \times \mathbf{E} = 0$ in electrostatics. Substituting this into (2), we get

$$-\mathbf{E} \times \mathbf{B} = \mathbf{A}\nabla^2\phi + \mathbf{\nabla}\phi(\mathbf{\nabla} \cdot \mathbf{A}) - \mathbf{E} \times \mathbf{B} - \mathbf{A}\nabla^2\phi \implies \mathbf{\nabla} \cdot \mathbf{A} = 0,$$

which is not very helpful.

2.b Give an example of a stationary, bounded charge and current distribution for which $\mathbf{P} \neq 0$.

Problem 3. The angular momentum density of the electromagnetic field is given by

$$\mathbf{l} = \mathbf{x} \times \mathcal{P} = \frac{1}{4\pi c}\mathbf{x} \times (\mathbf{E} \times \mathbf{B}).$$

Consider a source free ($\rho = 0$, $\mathbf{J} = 0$) solution to Maxwell's equations in electrodynamics with $\mathbf{E}$ and $\mathbf{B}$ vanishing rapidly as $|\mathbf{x}| \to \infty$, so the total angular momentum

$$\mathbf{L} = \int \mathbf{l} \, d^3x$$

is well defined. Show that $\mathbf{L}$ is conserved, i.e., independent of time.

In addition to the course lecture notes, I consulted Griffiths's *Introduction to Electrodynamics*, Jackson's *Classical Electrodynamics*, and Wolfram Mathworld while writing up these solutions.