

**Problem 1 (Jackson 11.3).** Show explicitly that two successive Lorentz transformations in the same direction are equivalent to a single Lorentz transformation with a velocity

$$v = \frac{v_1 + v_2}{1 + v_1 v_2 / c^2}.$$

This is an alternative way to derive the parallel-velocity addition law.

**Solution.** The general expression for a Lorentz transformation in the  $x_1$  direction is given by Jackson (11.16),

$$x'_0 = \gamma(x_0 - \beta x_1), \quad x'_1 = \gamma(x_1 - \beta x_0), \quad x'_2 = x_2, \quad x'_3 = x_3.$$

Define  $\beta$ ,  $\beta_1$ , and  $\beta_2$  by

$$\beta = \frac{v}{c}, \quad \beta_1 = \frac{v_1}{c}, \quad \beta_2 = \frac{v_2}{c},$$

and define  $\gamma$ ,  $\gamma_1$ , and  $\gamma_2$  correspondingly.

First applying the boost corresponding to  $\beta_1$  yields

$$x'_0 = \gamma_1(x_0 - \beta_1 x_1), \quad x'_1 = \gamma_1(x_1 - \beta_1 x_0). \quad (1)$$

Successively applying the boost corresponding to  $\beta_2$  yields

$$x''_0 = \gamma_2(x'_0 - \beta_2 x'_1), \quad x''_1 = \gamma_2(x'_1 - \beta_2 x'_0). \quad (2)$$

Substituting (1) into (2) gives us

$$\begin{aligned} x''_0 &= \gamma_1 \gamma_2 [(x_0 - \beta_1 x_1) - \beta_2 (x_1 - \beta_1 x_0)] = \gamma_1 \gamma_2 [(1 + \beta_1 \beta_2) x_0 - (\beta_1 + \beta_2) x_1] \\ x''_1 &= \gamma_1 \gamma_2 [(x_1 - \beta_1 x_0) - \beta_2 (x_0 - \beta_1 x_1)] = \gamma_1 \gamma_2 [(1 + \beta_1 \beta_2) x_1 - (\beta_1 + \beta_2) x_0], \end{aligned}$$

or

$$x''_0 = (1 + \beta_1 \beta_2) \gamma_1 \gamma_2 \left( x_0 - \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} x_1 \right), \quad x''_1 = (1 + \beta_1 \beta_2) \gamma_1 \gamma_2 \left( x_1 - \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} x_0 \right). \quad (3)$$

Note that

$$\frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} = \frac{1}{c} \frac{v_1 + v_2}{1 + v_1 v_2 / c^2} = \frac{v}{c} = \beta,$$

and that

$$\begin{aligned} (1 + \beta_1 \beta_2) \gamma_1 \gamma_2 &= \frac{1 + \beta_1 \beta_2}{\sqrt{(1 - \beta_1^2)(1 - \beta_2^2)}} = \sqrt{\frac{(1 - \beta_1^2)(1 - \beta_2^2)}{(1 + \beta_1 \beta_2)^2}}^{-1} = \sqrt{\frac{1 - \beta_1^2 - \beta_2^2 + \beta_1^2 \beta_2^2}{(1 + \beta_1 \beta_2)^2}}^{-1} \\ &= \sqrt{\frac{1 + 2\beta_1 \beta_2 + \beta_1^2 \beta_2^2 - (\beta_1^2 + 2\beta_1 \beta_2 + \beta_2^2)}{(1 + \beta_1 \beta_2)^2}}^{-1} = \sqrt{\frac{(1 + \beta_1 \beta_2)^2 - (\beta_1 + \beta_2)^2}{(1 + \beta_1 \beta_2)^2}}^{-1} \\ &= \sqrt{1 - \left( \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} \right)^2}^{-1} = \sqrt{1 - \beta^2}^{-1} = \gamma. \end{aligned}$$

Then (3) becomes

$$x''_0 = \gamma(x_0 - \beta x_1), \quad x''_1 = \gamma(x_1 - \beta x_0),$$

which is a single Lorentz transformation with velocity  $v$ , as we wanted to show.  $\square$

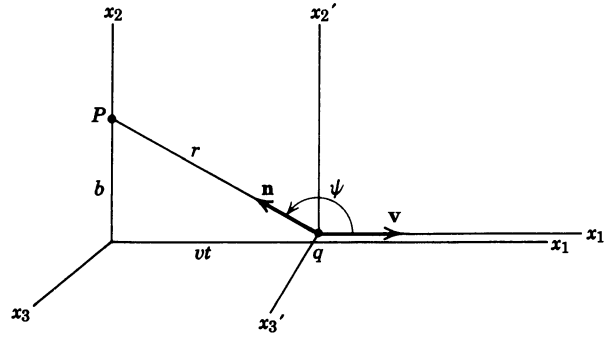


Figure 1: (Jackson 11.8) Particle of charge  $q$  moving at constant velocity  $\mathbf{v}$  passes an observation point  $P$  at impact parameter  $b$ .

**Problem 2 (Jackson 11.17).** The electric and magnetic fields (4) of a charge in uniform motion can be obtained from Coulomb's law in the charge's rest frame and the fact that the field strength  $F^{\alpha\beta}$  is an antisymmetric tensor of rank 2 without considering *explicitly* the Lorentz transformation. The idea is the following. For a charge in uniform motion the only relevant variables are the charge's 4-velocity  $U^\alpha$  and the relative coordinate  $X^\alpha = x_p^\alpha - x_q^\alpha$ , where  $x_p^\alpha$  and  $x_q^\alpha$  are the 4-vector coordinates of the observation point and the charge, respectively. The only antisymmetric tensor that can be formed is  $(X^\alpha U^\beta - X^\beta U^\alpha)$ . Thus the electromagnetic field  $F^{\alpha\beta}$  must be this tensor multiplied by some scalar function of the possible scalar products,  $X_\alpha X^\alpha$ ,  $X_\alpha U^\alpha$ ,  $U_\alpha U^\alpha$ .

**2.a** For the geometry of Fig. 1 the coordinates of  $P$  and  $q$  at a common time in  $K$  can be written  $x_p^\alpha = (ct, \mathbf{b})$ ,  $x_q^\alpha = (ct, \mathbf{v}t)$ , with  $\mathbf{b} \cdot \mathbf{v} = 0$ . By considering the general form of  $F^{\alpha\beta}$  in the rest frame of the charge, show that

$$F^{\alpha\beta} = \frac{q}{c} \frac{X^\alpha U^\beta - X^\beta U^\alpha}{[(U_\alpha X^\alpha/c)^2 - X_\alpha X^\alpha]^{3/2}}.$$

Verify that this yields the expressions

$$E_1 = E'_1 = -\frac{q\gamma vt}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, \quad E_2 = \gamma E'_2 = \frac{\gamma qb}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, \quad B_3 = \gamma \beta E'_2 = \beta E_2, \quad (4)$$

with all other components vanishing, in the inertial frame  $K$ .

**2.b** Repeat the calculation, using as the starting point the common-time coordinates in the rest frame,  $x_p'^\alpha = (ct', \mathbf{b} - \mathbf{v}t')$  and  $x_q'^\alpha = (ct', 0)$ . Show that

$$F^{\alpha\beta} = \frac{q}{c} \frac{Y^\alpha U^\beta - Y^\beta U^\alpha}{(-Y_\alpha Y^\alpha)^{3/2}},$$

where  $Y^\alpha = x_p'^\alpha - x_q'^\alpha$ . Verify that the fields are the same as in 2.a. Note that to obtain the results of (4) it is necessary to use the time  $t$  of the observation point  $P$  in  $K$  as the time parameter.

**2.c** Finally, consider the coordinate  $x_p^\alpha = (ct, \mathbf{b})$  and the "retarded-time" coordinate  $x_q^\alpha = [ct - R, \beta(ct - R)]$  where  $R$  is the distance between  $P$  and  $q$  at the retarded time. Define the difference as  $Z^\alpha = [R, \mathbf{b} - \beta(ct - R)]$ . Show that in terms of  $Z^\alpha$  and  $U^\alpha$  the field is

$$F^{\alpha\beta} = \frac{q}{c} \frac{Z^\alpha U^\beta - Z^\beta U^\alpha}{(U_\alpha Z^\alpha/c)^3}.$$

**Problem 3 (Jackson 11.18).** The electric and magnetic fields of a particle of charge  $q$  moving in a straight line with speed  $v = \beta c$ , given by (4), become more and more concentrated as  $\beta \rightarrow 1$ . Choose axes so that the charge moves along the  $z$  axis in the positive direction, passing the origin at  $t = 0$ . Let the spatial coordinates of the observation point be  $(x, y, z)$  and define the transverse vector  $\mathbf{r}_\perp$ , with components  $x$  and  $y$ . Consider the fields and the source in the limit of  $\beta = 1$ .

**3.a** Show that the fields can be written as

$$\mathbf{E} = 2q \frac{\mathbf{r}_\perp}{r_\perp^2} \delta(ct - z), \quad \mathbf{B} = 2q \frac{\hat{\mathbf{v}} \times \mathbf{r}_\perp}{r_\perp^2} \delta(ct - z),$$

where  $\hat{\mathbf{v}}$  is a unit vector in the direction of the particle's velocity.

**3.b** Show that by substitution into the Maxwell equations that these fields are consistent with a 4-vector source density,

$$J^\alpha = qc v^\alpha \delta^2(\mathbf{r}_\perp) \delta(ct - z),$$

where the 4-vector  $v^\alpha = (1, \hat{\mathbf{v}})$ .

**3.c** Show that the fields of (3.a) are derivable from either of the following 4-vector potentials,

$$A^0 = A^z = -2q\delta(ct - z) \ln(\lambda r_\perp), \quad \mathbf{A}_\perp = 0,$$

or

$$A^0 = 0 = A^z, \quad \mathbf{A}_\perp = -2q\Theta(ct - z) \nabla_\perp \ln(\lambda r_\perp),$$

where  $\lambda$  is an irrelevant parameter setting the scale of the logarithm.

Show that the two potentials differ by a gauge transformation and find the gauge function,  $\chi$ .

**Problem 4 (Jackson 11.20).** The lambda particle ( $\Lambda$ ) is a neutral baryon of mass  $M = 1115 \text{ MeV}$  that decays with a lifetime of  $\tau = 2.9 \times 10^{-10} \text{ s}$  into a nucleon of mass  $m_1 \approx 939 \text{ MeV}$  and a pi-meson of mass  $m_2 \approx 140 \text{ MeV}$ . It was first observed in flight by its charged decay mode  $\Lambda \rightarrow p + \pi^-$  in cloud chambers. The charged tracks originate from a single point and have the appearance of an inverted vee or lambda. The particles' identities and momenta can be inferred from their ranges and curvature in the magnetic field of the chamber.

**4.a** Using conservation of momentum and energy and the invariance of scalar products of 4-vectors show that, if the opening angle  $\theta$  between the two tracks is measured, the mass of the decaying particle can be found from the formula

$$M^2 = m_1^2 + m_2^2 + 2E_1 E_2 - 2p_1 p_2 \cos \theta,$$

where here  $p_1$  and  $p_2$  are the magnitudes of the 3-momenta.

**4.b** A lambda particle is created with a total energy of  $10 \text{ GeV}$  in a collision in the top plate of a cloud chamber. How far on the average will it travel in the chamber before decaying? What range of opening angles will occur for a  $10 \text{ GeV}$  lambda if the decay is more or less isotropic in the lambda's rest frame?

**Problem 5.** Show that  $\mathbf{E} \cdot \mathbf{B}$  is Lorentz invariant. You can do this either by using the Lorentz transformation laws for  $\mathbf{E}$  and  $\mathbf{B}$  derived in class, or by writing  $\mathbf{E} \cdot \mathbf{B}$  in a manifestly Lorentz invariant (and gauge invariant) form.

In addition to the course lecture notes, I consulted Jackson's *Classical Electrodynamics* while writing these solutions.