

Problem 1. Show that for an arbitrary spatially bound charge-current source, the electric dipole moment \mathbf{p} satisfies

$$\frac{d\mathbf{p}}{dt} = \int \mathbf{J} d^3x.$$

Solution. The electric dipole moment \mathbf{p} is defined by Eq. (2.36),

$$\mathbf{p} = \int \mathbf{x} \rho(x) d^3x.$$

Differentiating both sides with respect to t , we find

$$\frac{d\mathbf{p}}{dt} = \frac{d}{dt} \int \mathbf{x} \rho d^3x = \int \frac{d}{dt} (\mathbf{x} \rho) d^3x = \int \mathbf{x} \frac{\partial \rho}{\partial t} d^3x, \quad (1)$$

because \mathbf{x} is simply the point at which we are evaluating the potential, and is therefore independent of time.

The charge-current conservation law is given by Eq. (5.8),

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0.$$

Multiplying by \mathbf{x} on both sides and integrating over all space, we obtain

$$\int \mathbf{x} \frac{\partial \rho}{\partial t} d^3x + \int \mathbf{x} (\nabla \cdot \mathbf{J}) d^3x = 0.$$

Applying (1), we have

$$\frac{d\mathbf{p}}{dt} = - \int \mathbf{x} (\nabla \cdot \mathbf{J}) d^3x. \quad (2)$$

It remains to be shown that the right side is equal to the integral of \mathbf{J} over all space.

Vector identity (5) in Griffiths is

$$\nabla \cdot (f\mathbf{a}) = f(\nabla \cdot \mathbf{a}) + \mathbf{a} \cdot (\nabla f).$$

Writing the right side of (2) in component notation and applying the identity gives us

$$- \int x_i (\nabla \cdot \mathbf{J}) d^3x = \int \mathbf{J} \cdot (\nabla x_i) d^3x - \int \nabla \cdot (x_i \mathbf{J}) d^3x. \quad (3)$$

Gauss's theorem is given by Eq. (2.6),

$$\int_{\mathcal{V}} \nabla \cdot \mathbf{v} d^3x = \int_S \mathbf{v} \cdot \hat{\mathbf{n}} dS,$$

Here, let \mathcal{V} be a ball of radius R , with R large enough that the entire charge-current source is enclosed. Then S is the surface of \mathcal{V} , and $\hat{\mathbf{n}} = \hat{\mathbf{r}}$. Applying Gauss's theorem to the second integral on the right side of (3), we have

$$\int \nabla \cdot (x_i \mathbf{J}) d^3x = \lim_{R \rightarrow \infty} \int_{\mathcal{V}} \nabla \cdot (x_i \mathbf{J}) d^3x = \lim_{R \rightarrow \infty} \int_S x_i \mathbf{J} \cdot \hat{\mathbf{r}} dS = 0,$$

since \mathbf{J} is bounded, and therefore \mathbf{J} evaluated on S reaches zero well before x_i becomes very large.

Returning to (3), we now have

$$-\int x_i (\nabla \cdot \mathbf{J}) d^3x = \int \mathbf{J} \cdot (\nabla x_i) d^3x = \sum_j \int J_j \partial_j x_i d^3x = \sum_j \int J_j \delta_{ij} d^3x = \int J_i d^3x,$$

where we have followed the proof in Eq. (4.24) of the course notes. Finally, (2) becomes

$$\frac{d\mathbf{p}}{dt} = \int \mathbf{J} d^3x$$

as desired. □

Problem 2. A particle of charge q_1 moves with velocity v in a circular orbit of radius R about the origin in the xy plane, such that its φ coordinate varies as $\varphi = \omega t$, with $\omega = v/R$. Assume that $v \ll c$. Another particle of charge q_2 is at rest at point \mathbf{x} , where $|\mathbf{x}| \gg R$. To order $1/|\mathbf{x}|$, find the force \mathbf{F} on the particle of charge q_2 at time t .

Problem 3. An “antenna” is a segment of conducting wire in which a current flows (driven by an external power supply). Suppose an antenna of length L is placed on the z axis between $z = 0$ and $z = L$, and suppose that the current in the antenna is

$$\mathbf{J}(t, z) = I_0 \sin\left(\frac{\pi z}{L}\right) \cos(\omega t) \delta(x) \delta(y) \hat{\mathbf{z}}.$$

3.a Find the charge density $\rho(t, z)$ in the antenna.

3.b Assume that $\omega L \ll c$. Find the electric and magnetic fields, $\mathbf{E}(t, z)$ and $\mathbf{B}(t, z)$, at large distances from the antenna (valid to order $1/|\mathbf{x}|$).

In addition to the course lecture notes, I consulted Griffiths’s *Introduction to Electrodynamics*, Jackson’s *Classical Electrodynamics*, and K. T. McDonald’s and D. K. Ghosh’s notes on electromagnetism while writing up these solutions.