Problem 1. Consider a spin-1 particle. The unperturbed Hamiltonian is $H_0 = AS_z^2$, where A is a constant. Consider the perturbation $V = B(S_x^2 - S_y^2)$, where $|A| \gg |B|$. Note that S_i are the 3×3 spin matrices.

1.1 Calculate the first-order correction to the energies.

Solution. Firstly, note that

$$S_x = \frac{\hbar}{\sqrt{2}} \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix}, \qquad S_y = \frac{\hbar}{\sqrt{2}} \begin{bmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{bmatrix}, \qquad S_z = \hbar \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}.$$

Then

$$H_0 = A\hbar^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}, \qquad V = B\frac{\hbar^2}{2} \left(\begin{bmatrix} 1 & 0 & 1 \\ 0 & 2 & 0 \\ 1 & 0 & 1 \end{bmatrix} - \begin{bmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{bmatrix} \right) = B\hbar^2 \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}. \tag{1}$$

The eigenvalues of H_0 are

$$E_1^{(0)} = A\hbar^2, \qquad E_2^{(0)} = 0, \qquad E_3^{(0)} = A\hbar^2, \tag{2}$$

so the problem is degenerate. The eigenkets are the S_z eigenbasis kets:

$$|1^{(0)}\rangle = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = |+1\rangle, \qquad |2^{(0)}\rangle = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} = |0\rangle, \qquad |3^{(0)}\rangle = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = |-1\rangle.$$

We will begin with the correction to $E_2^{(0)}$, which is nondegenerate. From (5.1.20) and (5.1.37) in Sakurai, the first-order energy corrections in the unperturbed case are given by

$$\Delta_n^{(1)} \equiv E_n^{(1)} - E_n^{(0)} = \langle n^{(0)}|V|n^{(0)}\rangle.$$

This gives us

$$\Delta_2^{(1)} = \langle 2^{(0)}|V|2^{(0)}\rangle = \langle 2|V|2\rangle = 0.$$

For $E_1^{(0)}$ and $E_2^{(0)}$, consider the degenerate subspace spanned by $\{|+1\rangle, |-1\rangle\}$. Let P_0 be a projection onto this subspace, and let

$$V_0 = P_0 V P_0 = B\hbar^2 \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = B\hbar^2 \sigma_x,$$

where σ_x is the Pauli matrix. Therefore, we know that V_0 has eigenvalues $v_\pm = \pm B\hbar^2$. These eigenvalues are equivalent to the corresponding energy shifts.

In summary, we have

$$\Delta_1^{(1)} = B\hbar^2, \qquad \Delta_2^{(1)} = 0, \qquad \Delta_3^{(1)} = -B\hbar^2.$$

1.2 Solve the problem exactly, and compare your result to the perturbation theory result.

Solution. From (1), the perturbed Hamiltonian is given by

$$H = H_0 + \lambda V = \hbar^2 \begin{bmatrix} A & 0 & \lambda B \\ 0 & 0 & 0 \\ \lambda B & 0 & A \end{bmatrix}.$$

Let $E_i = \hbar^2 \mu_i$ denote the eigenvalues of H, where μ are the roots of the equation

$$\begin{aligned} 0 = \det\left(\frac{H}{\hbar^2} - \mu I\right) &= \begin{vmatrix} A-\mu & 0 & \lambda B \\ 0 & -\mu & 0 \\ \lambda B & 0 & A-\mu \end{vmatrix} = -\mu(A-\mu)^2 + \mu(\lambda B)^2 = \mu(A^2 - 2A\mu + \mu^2 + \lambda^2 B^2 \\ &= \mu(\mu - A + \lambda B)(\mu - A - \lambda B). \end{aligned}$$

The roots are $\mu = 0$ and $\mu = A \pm \lambda B$, which give us the eigenvalues

$$E_1 = A + \lambda B, \qquad E_2 = 0, \qquad E_3 = A - \lambda B.$$

Taking the difference $\Delta_n^{(1)} = E_n^{(1)} - E_n^{(0)}$ for $E_i^{(0)}$ given by (2), the energy shifts to first order in λ are

$$\Delta_1^{(1)} = B\hbar^2, \qquad \Delta_2^{(1)} = 0, \qquad \Delta_3^{(1)} = -B\hbar^2,$$

which are the same as those found in 1.1.

Problem 2. Consider the Stark effect for the $n = 3$ states of hydrogen. There are initially nine degenerate states $|3, l, m\rangle$ (neglect spin), and an electric field E is turned on in the z direction.

2.1 Construct the 9×9 matrix representing the perturbed Hamiltonian in this case. Show your work when deriving the nonzero matrix elements, and provide an explanation as to why the other elements are zero.

Solution. The perturbation operator for the $\mathbf{E}$ field is given by (5.2.17) in Sakurai:

$$V = -eZ|\mathbf{E}|.$$

V is a dipole interaction because the hydrogen atom has a nonzero dipole moment. Therefore V obeys the dipole selection rule, which is given by (17.2.21) in Shankar:

$$\langle nlm \big| Z \big| n'l'm' \rangle = 0 \quad \text{unless} \quad \begin{cases} l' = l \pm 1, \\ m' = m. \end{cases}$$

The dipole selection rule is a combination of the angular momentum and parity selection rules. The angular momentum selection rule stipulates that $\langle nlm|Z|n'l'm'\rangle = 0$ unless $l' = l, l \pm 1$ and $m' = m + q$ where $q = 0$ is the magnetic quantum number of the tensor operator Z. The parity selection rule eliminates $l = l'$ because Z is parity odd, so $\langle nlm|Z|n'l'm'\rangle = 0$ unless l and l' have opposite parity.

For the nonzero elements, the hydrogen atom wave functions are given by (A.6.3) in Sakurai:

$$\langle \mathbf{r}|nlm\rangle = \psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi),$$

where

$$R_{nl}(r) = -\sqrt{\left(\frac{2}{na_0}\right)^3 \frac{(n-l-1)!}{2n(n+l)!^3}} e^{-\rho/2} \rho^l L_{n+l}^{2l+1}(\rho) \quad \text{where} \quad \rho = \frac{2r}{na_0}. \tag{3}$$

The associated Laguerre polynomials L_p^q are given by (A.6.4) and (A.6.5),

$$L_p^q(\rho) = \frac{d^q L_p(\rho)}{d\rho^q} \quad \text{where} \quad L_p(\rho) = e^\rho \frac{d^p}{d\rho^p}. \tag{4}$$

The spherical harmonics Y_l^m are given by (3.6.37) and (3.6.38),

$$Y_l^m(\theta,\phi) = \frac{(-1)^l}{2^l l!}\sqrt{\frac{2l+1}{4\pi}\frac{(l+m)!}{(l-m)!}}e^{im\phi}\frac{1}{\sin^m\theta}\frac{d^{l-m}}{d(\cos\theta)^{l-m}}{}^{2l}, \qquad Y_l^{-m}(\theta,\phi) = (-1)^m Y_l^{m*}(\theta,\phi) \tag{5}$$

for $m \geq 0$.

The nonzero elements all have $l \in \{0,1,2\}$ and $m \in \{-1,0,1\}$. Substituting into (3), the relevant R_{nl} are

$$R_{30}(r) = -\sqrt{\left(\frac{2}{3a_0}\right)^3 \frac{(3-1)!}{2(3)3!^3}}e^{-\rho/2}L_3^1(\rho) = -\sqrt{\frac{2^3}{3^3a_0^3}\frac{2}{2^4 3^4}}e^{-\rho/2}L_3^1(\rho) = -\sqrt{\frac{e^{-\rho}}{3^7a_0^3}}L_3^1(\rho),$$

$$R_{31}(r) = -\sqrt{\left(\frac{2}{3a_0}\right)^3 \frac{(3-1-1)!}{2(3)(3+1)!^3}}e^{-\rho/2}\rho L_{3+1}^{2+1}(\rho) = -\sqrt{\frac{2^3}{3^3a_0^3}\frac{1}{2^{10}3^4}}e^{-\rho/2}\rho L_4^3(\rho) = -\sqrt{\frac{e^{-\rho}}{2^7 3^7a_0^3}}\rho L_4^3(\rho),$$

$$R_{32}(r) = -\sqrt{\left(\frac{2}{3a_0}\right)^3 \frac{(3-2-1)!}{2(3)(3+2)!^3}}e^{-\rho/2}\rho^2 L_{3+2}^{4+1}(\rho) = -\sqrt{\frac{2^3}{3^3a_0^3}\frac{1}{2^{10}3^4 5^3}}e^{-\rho/2}\rho^2 L_5^5(\rho) = -\sqrt{\frac{e^{-\rho}}{2^7 3^7 5^3 a_0^3}}\rho^2 L_5^5(\rho).$$

From (4), the relevant L_p are

$$\begin{aligned}
L_3(\rho) &= e^\rho\frac{d^3}{d\rho^3} = e^\rho\frac{d^2}{d\rho^2} = e^\rho\frac{d}{d\rho} = 6 - 18\rho + 9\rho^2 - \rho^3,\\
L_4(\rho) &= e^\rho\frac{d^4}{d\rho^4} = e^\rho\frac{d^3}{d\rho^3} = e^\rho\frac{d^2}{d\rho^2}\\
&= e^\rho\frac{d}{d\rho} = 24 - 96\rho + 72\rho^2 - 16\rho^3 + \rho^4,\\
L_5(\rho) &= e^\rho\frac{d^5}{d\rho^5} = e^\rho\frac{d^4}{d\rho^4} = e^\rho\frac{d^3}{d\rho^3}\\
&= e^\rho\frac{d^2}{d\rho^2}\\
&= e^\rho\frac{d}{d\rho} = 120 - 600\rho + 600\rho^2 - 200\rho^3 + 25\rho^4 - \rho^5
\end{aligned}$$

and then the relevant L_p^q are

$$\begin{aligned}
L_3^1(\rho) &= \frac{dL_3(\rho)}{d\rho} = -18 + 18\rho - 3\rho^2 = -3(6 - 6\rho + \rho^2),\\
L_4^3(\rho) &= \frac{d^3L_4(\rho)}{d\rho^3} = -(3!)16 + \left(\frac{4!}{1!}\right)\rho = 24(-4+\rho) = 2^3 3(-4+\rho),\\
L_5^5(\rho) &= \frac{d^5L_5(\rho)}{d\rho^5} = -5! = -120 = -2^3 3^1 5.
\end{aligned}$$

Substituting into (5), the relevant Y_l^m are

$$\begin{aligned}
Y_0^0(\theta,\phi) &= \sqrt{\frac{1}{2^2\pi}}, \\
Y_1^0(\theta,\phi) &= \sqrt{\frac{3}{2^2\pi}}\cos\theta, & Y_1^{\pm1}(\theta,\phi) &= \mp\sqrt{\frac{3}{2^3\pi}}e^{\pm i\phi}\sin\theta, \\
Y_2^0(\theta,\phi) &= \sqrt{\frac{5}{2^4\pi}}(3\cos^2\theta-1), & Y_2^{\pm1}(\theta,\phi) &= \mp\sqrt{\frac{3^15}{2^3\pi}}e^{\pm i\phi}\cos\theta\sin\theta.
\end{aligned}$$

Note that $Z = r\cos\theta$ in polar coordinates. In general, the nonzero matrix elements are then

$$\begin{aligned}
\big\langle 3lm\big|V\big|3l'm'\big\rangle &= -e|\mathbf{E}|\int_0^{2\pi}\int_0^{\pi}\int_0^{\infty}\psi_{3lm}^*(r,\theta,\phi)r\cos\theta\psi_{3l'm'}(r,\theta,\phi)r^2\sin\theta\,dr\,d\theta\,d\phi \\
&= -e|\mathbf{E}|\left(\frac{3a_0}{2}\right)^4\int_0^{2\pi}\int_{-1}^{1}\int_0^{\infty}\psi_{3lm}^*(r,\theta,\phi)\psi_{3l'm'}(r,\theta,\phi)\rho^3\cos\theta\,d\rho\,d(\cos\theta)\,d\phi \\
&= -\frac{3^4a_0^4e|\mathbf{E}|}{2^4}\int_0^{2\pi}\int_{-1}^{1}Y_l^{m*}(\theta,\phi)Y_{l'}^{m'}(\theta,\phi)\cos\theta\,d(\cos\theta)\,d\phi\int_0^{\infty}R_{3l}(r)R_{3l'}(r)\rho^3\,d\rho\,.
\end{aligned}$$

Firstly,

$$\langle 310|V|300\rangle = -\frac{3^4a_0^4e|\mathbf{E}|}{2^4}\int_0^{2\pi}\int_{-1}^{1}Y_1^{0*}(\theta,\phi)Y_0^0(\theta,\phi)\cos\theta\,d(\cos\theta)\,d\phi\int_0^{\infty}R_{31}(r)R_{30}(r)\rho^3\,d\rho\,, \tag{6}$$

where

$$\begin{aligned}
&\int_0^{2\pi}\int_{-1}^{1}Y_1^{0*}(\theta,\phi)Y_0^0(\theta,\phi)\cos\theta\,d(\cos\theta)\,d\phi = \int_0^{2\pi}\int_{-1}^{1}\sqrt{\frac{3}{2^2\pi}}\cos\theta\sqrt{\frac{1}{2^2\pi}}\cos\theta\,d(\cos\theta)\,d\phi \\
&\quad= \frac{\sqrt3}{2^2\pi}\int_0^{2\pi}d\phi\int_{-1}^{1}\cos^2\theta\,d(\cos\theta) = \frac{\sqrt3}{2^2\pi}\Big[\phi\Big]_0^{2\pi}\left[\frac{\cos^3\theta}{3}\right]_{-1}^{1} = \frac{\sqrt3}{2^2\pi}(2\pi)\frac{2}{3} = \frac{1}{\sqrt3},
\end{aligned}$$

and

$$\begin{aligned}
&\int_0^{\infty}R_{31}(r)R_{30}(r)\rho^3\,d\rho = \int_0^{\infty}\sqrt{\frac{e^{-\rho}}{2^73^7a_0^3}}\rho L_4^3(\rho)\sqrt{\frac{e^{-\rho}}{3^7a_0^3}}L_3^1(\rho)\rho^3\,d\rho = \frac{1}{\sqrt{2^73^7a_0^3}}\int_0^{\infty}e^{-\rho}L_4^3(\rho)L_3^1(\rho)\rho^4\,d\rho \\
&\quad= -\frac{1}{\sqrt2 3^5a_0^3}\int_0^{\infty}e^{-\rho}(-24\rho^4+30\rho^5-10\rho^6+\rho^7)\,d\rho = -\frac{1}{\sqrt2 3^5a_0^3}(-24(4!)+30(5!)-10(6!)+7!) \\
&\quad= -\frac{2^5}{\sqrt2 3^2a_0^3},
\end{aligned}$$

where we have used

$$\int_0^{\infty}x^ne^{-x}\,dx = n!.$$

Combining these results, (6) becomes

$$\langle 310|V|300\rangle = \frac{3^4a_0^4e|\mathbf{E}|}{2^4}\frac{1}{\sqrt3}\frac{2^5}{\sqrt2 3^2a_0^3} = e|\mathbf{E}|a_0\frac{3^22}{\sqrt6} = 3\sqrt6e|\mathbf{E}|a_0 = \langle 300|V|310\rangle\,.$$

Secondly,

$$\langle 32{\pm}1|V|31{\pm}1\rangle = -\frac{3^4a_0^4e|\mathbf{E}|}{2^4}\int_0^{2\pi}\int_{-1}^{1}Y_2^{\pm1*}(\theta,\phi)Y_1^{\pm1}(\theta,\phi)\cos\theta\,d(\cos\theta)\,d\phi\int_0^{\infty}R_{32}(r)R_{31}(r)\rho^3\,d\rho\,, \tag{7}$$

where

$$\int_0^{2\pi}\int_{-1}^{1} Y_2^{\pm1*}(\theta,\phi)Y_1^{\pm1}(\theta,\phi)\cos\theta\, d(\cos\theta)\, d\phi = \int_0^{2\pi}\int_{-1}^{1}\sqrt{\frac{3^15}{2^3\pi}}e^{\mp i\phi}\cos\theta\sin\theta\sqrt{\frac{3}{2^3\pi}}e^{\pm i\phi}\sin\theta\cos\theta\, d(\cos\theta)\, d\phi$$
$$= \frac{3\sqrt{5}}{2^3\pi}\int_0^{2\pi} d\phi \int_{-1}^{1}\cos^2\theta\sin^2\theta\, d(\cos\theta) = \frac{3\sqrt{5}}{2^3\pi}\int_0^{2\pi} d\phi\int_{-1}^{1}\cos^2\theta(1-\cos^2\theta)\, d(\cos\theta)$$
$$= \frac{3\sqrt{5}}{2^3\pi}\Big[\phi\Big]_0^{2\pi}\left[\frac{\cos^3\theta}{3}-\frac{\cos^5\theta}{5}\right]_{-1}^{1} = \frac{3\sqrt{5}}{2^3\pi}(2\pi)\frac{2^2}{3^15} = \frac{1}{\sqrt{5}},$$

and

$$\int_0^{\infty} R_{32}(r)R_{31}(r)\rho^3\, d\rho = \int_0^{\infty}\sqrt{\frac{e^{-\rho}}{2^73^75^3a_0^3}}\rho^2L_5^5(\rho)\sqrt{\frac{e^{-\rho}}{2^73^7a_0^3}}\rho L_4^3(\rho)\rho^3\, d\rho = \frac{1}{2^73^7\sqrt{5^3}a_0^3}\int_0^{\infty} e^{-\rho}L_5^5(\rho)L_4^3(\rho)\rho^6\, d\rho$$
$$= -\frac{1}{2^13^5\sqrt{5}a_0^3}\int_0^{\infty} e^{-\rho}(-4+\rho)\rho^6\, d\rho = -\frac{1}{2^13^5\sqrt{5}a_0^3}\int_0^{\infty} e^{-\rho}(-4\rho^6+\rho^7)\, d\rho = -\frac{1}{2^13^5\sqrt{5}a_0^3}(-4(6!)+7!)$$
$$= -\frac{2^3\sqrt{5}}{3^2a_0^3}.$$

Then (7) becomes

$$\langle 32{\pm}1|V|31{\pm}1\rangle = \frac{3^4a_0^4e|\mathbf{E}|}{2^4}\frac{1}{\sqrt{5}}\frac{2^3\sqrt{5}}{3^2a_0^3} = \frac{3^2a_0e|\mathbf{E}|}{2} = \frac{9}{2}e|\mathbf{E}|a_0 = \langle 31{\pm}1|V|32{\pm}1\rangle\,.$$

Thirdly,

$$\langle 320|V|310\rangle = -\frac{3^4a_0^4e|\mathbf{E}|}{2^4}\int_0^{2\pi}\int_{-1}^{1} Y_2^{0*}(\theta,\phi)Y_1^0(\theta,\phi)\cos\theta\, d(\cos\theta)\, d\phi\int_0^{\infty} R_{32}(r)R_{31}(r)\rho^3\, d\rho\,, \tag{8}$$

where

$$\int_0^{2\pi}\int_{-1}^{1} Y_2^{0*}(\theta,\phi)Y_1^0(\theta,\phi)\cos\theta\, d(\cos\theta)\, d\phi = \int_0^{2\pi}\int_{-1}^{1}\sqrt{\frac{5}{2^4\pi}}(3\cos^2\theta-1)\sqrt{\frac{3}{2^2\pi}}\cos\theta\cos\theta\, d(\cos\theta)\, d\phi$$
$$= \frac{\sqrt{3}\sqrt{5}}{2^3\pi}\int_0^{2\pi} d\phi\int_{-1}^{1}(3\cos^4\theta-\cos^2\theta)\, d(\cos\theta) = \frac{\sqrt{3}\sqrt{5}}{2^3\pi}\Big[\phi\Big]_0^{2\pi}\left[\frac{3\cos^5\theta}{5}-\frac{\cos^3\theta}{3}\right]_{-1}^{1} = \frac{\sqrt{3}\sqrt{5}}{2^3\pi}(2\pi)\frac{2^3}{3^15}$$
$$= \frac{2}{\sqrt{3}\sqrt{5}},$$

and

$$\int_0^{\infty} R_{32}(r)R_{31}(r)\rho^3\, d\rho = -\frac{2^3\sqrt{5}}{3^2a_0^3}.$$

Then (8) becomes

$$\langle 320|V|310\rangle = \frac{3^4a_0^4e|\mathbf{E}|}{2^4}\frac{2}{\sqrt{3}\sqrt{5}}\frac{2^3\sqrt{5}}{3^2a_0^3} = 3\sqrt{3}e|\mathbf{E}|a_0 = \langle 310|V|320\rangle\,.$$

In summary, we have

$$V = e|\mathbf{E}|a_0 \begin{array}{c} \begin{array}{ccccccccc} 300 & 31{-}1 & 310 & 311 & 32{-}2 & 32{-}1 & 320 & 321 & 322 \end{array} \\ \begin{bmatrix} 0 & 0 & 3\sqrt{6} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 9/2 & 0 & 0 & 0 \\ 3\sqrt{6} & 0 & 0 & 0 & 0 & 0 & 3\sqrt{3} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 9/2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 9/2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 3\sqrt{3} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 9/2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix} \end{array}. \tag{9}$$

2.2 Determine the first order corrections, $\Delta^{(1)}$, to the energies due to this perturbation, and write down the degeneracies of these energies.

Solution. We have the perturbed Hamiltonian

$$H = H_0 + \lambda V,$$

where V is given by (9). For the $n = 3$ states of hydrogen, H_0 is ninefold degenerate, so we need to find the eigenvalues of the full matrix V. Let $\Delta^{(1)} = e|\mathbf{E}|a_0\mu$ denote the eigenvalues of V, where μ are the roots of the equation

$$\begin{aligned} 0 &= \det\left(\frac{V}{e|\mathbf{E}|a_0} - \mu I\right) \\ &= \begin{vmatrix} -\mu & 0 & 3\sqrt{6} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -\mu & 0 & 0 & 0 & 9/2 & 0 & 0 & 0 \\ 3\sqrt{6} & 0 & -\mu & 0 & 0 & 0 & 3\sqrt{3} & 0 & 0 \\ 0 & 0 & 0 & -\mu & 0 & 0 & 0 & 9/2 & 0 \\ 0 & 0 & 0 & 0 & -\mu & 0 & 0 & 0 & 0 \\ 0 & 9/2 & 0 & 0 & 0 & -\mu & 0 & 0 & 0 \\ 0 & 0 & 3\sqrt{3} & 0 & 0 & 0 & -\mu & 0 & 0 \\ 0 & 0 & 0 & 9/2 & 0 & 0 & 0 & -\mu & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\mu \end{vmatrix} \\ &= \begin{vmatrix} -\mu & 0 & 3\sqrt{6} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -\mu & 0 & 0 & 0 & 9/2 & 0 & 0 & 0 \\ 0 & 0 & (\mu^2-54)/\mu & 0 & 0 & 0 & 3\sqrt{3} & 0 & 0 \\ 0 & 0 & 0 & -\mu & 0 & 0 & 0 & 9/2 & 0 \\ 0 & 0 & 0 & 0 & -\mu & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & (81/4-\mu^2)/\mu & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & \mu(81-\mu^2)/(\mu^2-54) & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & (81/4-\mu^2)/\mu & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -\mu \end{vmatrix} \\ &= (-\mu)^5 \frac{\mu^2-54}{\mu}\left(\frac{81/4-\mu^2}{\mu}\right)^2 \frac{\mu(81-\mu^2)}{\mu^2-54} = -\mu^3\left(\frac{81}{4}-\mu^2\right)^2(81-\mu^2) \\ &= \mu^3\left(\frac{9}{2}-\mu\right)^2\left(\frac{9}{2}+\mu\right)^2(9-\mu)(9+\mu), \end{aligned}$$

where we have taken advantage of the determinant's invariance under elementary row addition. This gives us the energy shifts

$$\Delta^{(1)} = \begin{cases} 0 & \text{degeneracy 3,} \\ \pm\dfrac{9}{2}e|\mathbf{E}|a_0 & \text{degeneracy 2,} \\ \pm 9e|\mathbf{E}|a_0 & \text{no degeneracy.} \end{cases}$$

Problem 3. Consider the Hamiltonian H_0 acting on a three-dimensional Hilbert space spanned by the orthonormal basis $\{|1\rangle, |2\rangle, |3\rangle\}$. $H_0 = \sum_{i=3}^{3} E_i |i\rangle\langle i|$, with energy eigenvalues $E_1^{(0)}, E_2^{(0)}, E_3^{(0)}$. Assume $E_1^{(0)} = E_2^{(0)} = E_D^{(0)}$. To H_0, we add a perturbation

$$V = v_1 |1\rangle\langle 3| + v_1^* |3\rangle\langle 1| + v_2 |2\rangle\langle 3| + v_2^* |3\rangle\langle 2| .$$

Here, v_1 and v_2 are complex constants and small compared to E_3.

3.1 To second order in V, write down the explicit form of the effective Hamiltonian acting on the subspace spanned by $\{|1\rangle, |2\rangle\}$.

Solution. We have

$$H_0 = \begin{bmatrix} E_D^{(0)} & 0 & 0 \\ 0 & E_D^{(0)} & 0 \\ 0 & 0 & E_3^{(0)} \end{bmatrix}, \qquad V = \begin{bmatrix} 0 & 0 & v_1 \\ 0 & 0 & v_2 \\ v_1^* & v_2^* & 0 \end{bmatrix}, \qquad H = H_0 + \lambda V = \begin{bmatrix} E_D^{(0)} & 0 & \lambda v_1 \\ 0 & E_D^{(0)} & \lambda v_2 \\ v_1^* & v_2^* & E_3^{(0)} \end{bmatrix}.$$

From the lecture notes and (5.2.12) in Sakurai, the effective Hamiltonian is given to second order in λ by

$$H_{\text{eff}} = E_D^{(0)} + \lambda P_0 V P_0 + \lambda^2 P_0 V P_1 (E_D^{(0)} - H_0)^{-1} P_1 V P_0,$$

where P_0 is the projection onto the degenerate subspace, P_1 is the projection onto the nondegenerate subspace, and $E_D^{(0)}$ is the degenerate energy. Here, P_0 projects onto the subspace spanned by $\{|1\rangle, |2\rangle\}$ and P_1 onto that spanned by $\{|3\rangle\}$.

Note that

$$PoVP_0 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & 0 & v_1 \\ 0 & 0 & v_2 \\ v_1^* & v_2^* & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix},$$

and

$$\begin{aligned} P_0 V P_1 (E_D^{(0)} - H_0)^{-1} P_1 V P_0 &= \frac{1}{E_D^{(0)} - E_3^{(0)}} P_0 \begin{bmatrix} 0 & 0 & v_1 \\ 0 & 0 & v_2 \\ v_1^* & v_2^* & 0 \end{bmatrix} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 0 & v_1 \\ 0 & 0 & v_2 \\ v_1^* & v_2^* & 0 \end{bmatrix} P_0 \\ &= \frac{1}{E_D^{(0)} - E_3^{(0)}} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} |v_1|^2 & v_1 v_2^* & 0 \\ v_1^* v_2 & |v_1|^2 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \frac{1}{E_D^{(0)} - E_3^{(0)}} \begin{bmatrix} |v_1|^2 & v_1 v_2^* & 0 \\ v_1^* v_2 & |v_1|^2 & 0 \\ 0 & 0 & 0 \end{bmatrix}. \end{aligned}$$

In the degenerate subspace, we have

$$H_{\text{eff}} = \begin{bmatrix} E_D^{(0)} + |v_1|^2/(E_D^{(0)} - E_3^{(0)}) & v_1 v_2^*/(E_D^{(0)} - E_3^{(0)}) \\ v_1^* v_2/(E_D^{(0)} - E_3^{(0)}) & E_D^{(0)} + |v_2|^2/(E_D^{(0)} - E_3^{(0)}) \end{bmatrix}.$$

3.2 By solving the effective Hamiltonian, construct the approximate solution for the eigenvalues and eigenfunctions of $H_0 + V$. (The eigenkets only need to be constructed within the degenerate subspace.)

Solution. Let E be the eigenvalues of H_{eff}. We need to solve the characteristic equation

$$\begin{aligned}
0 = \det(H_{\text{eff}} - EI) &= \begin{vmatrix} E_D^{(0)} + |v_1|^2/(E_D^{(0)} - E_3^{(0)}) - E & v_1 v_2^*/(E_D^{(0)} - E_3^{(0)}) \\ v_1^* v_2/(E_D^{(0)} - E_3^{(0)}) & E_D^{(0)} + |v_2|^2/(E_D^{(0)} - E_3^{(0)}) - E \end{vmatrix} \\
&= \left(E_D^{(0)} + \frac{|v_1|^2}{E_D^{(0)} - E_3^{(0)}} - E \right) \left(E_D^{(0)} + \frac{|v_2|^2}{E_D^{(0)} - E_3^{(0)}} - E \right) - \frac{|v_1|^2 |v_2|^2}{(E_D^{(0)} - E_3^{(0)})^2} \\
&= {E_D^{(0)}}^2 + E_D^{(0)} \frac{|v_2|^2}{E_D^{(0)} - E_3^{(0)}} - E_D^{(0)} E + E_D^{(0)} \frac{|v_1|^2}{E_D^{(0)} - E_3^{(0)}} - E \frac{|v_1|^2}{E_D^{(0)} - E_3^{(0)}} - E_D^{(0)} E - E \frac{|v_2|^2}{E_D^{(0)} - E_3^{(0)}} + E^2 \\
&= E^2 - E_D^{(0)} E - E \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}} - E_D^{(0)} E + {E_D^{(0)}}^2 + E_D^{(0)} \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}} \\
&= (E - E_D^{(0)}) \left(E - E_D^{(0)} - \frac{|v_1|^2 + |v_2|^2}{(E_D^{(0)} - E_3^{(0)})^2} \right),
\end{aligned}$$

so the energies are

$$E_1 = E_D^{(0)}, \qquad E_2 = E_D^{(0)} + \frac{|v_1|^2 + |v_2|^2}{(E_D^{(0)} - E_3^{(0)})^2}.$$

For E_3, we can use nondegenerate perturbation theory to find the energy shift to second order. From (5.1.43) in Sakurai,

$$\Delta_n^{(1)} = \lambda V_{nn} + \lambda^2 \sum_{k \neq n} \frac{|V_{nk}|^2}{E_n^{(0)} - E_k^{(0)}} + \cdots .$$

For $n = 3$, up to second order we have

$$\Delta_3^{(1)} = \lambda V_{33} + \lambda^2 \left(\frac{|V_{31}|^2}{E_3^{(0)} - E_D^{(0)}} + \frac{|V_{32}|^2}{E_3^{(0)} - E_D^{(0)}} \right) = \lambda^2 \frac{|v_1|^2 + |v_2|^2}{E_3^{(0)} - E_D^{(0)}} = -\lambda^2 \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}}.$$

Thus, we have

$$E_3 = E_3^{(0)} - \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}}.$$

The eigenvector corresponding to E_1 can be found by

$$\begin{bmatrix} E_D^{(0)} + |v_1|^2/(E_D^{(0)} - E_3^{(0)}) & v_1 v_2^*/(E_D^{(0)} - E_3^{(0)}) \\ v_1^* v_2/(E_D^{(0)} - E_3^{(0)}) & E_D^{(0)} + |v_2|^2/(E_D^{(0)} - E_3^{(0)}) \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} = E_D^{(0)} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}$$

which is equivalent to the system of equations

$$\left(E_D^{(0)} + \frac{|v_1|^2}{E_D^{(0)} - E_3^{(0)}} \right) u_1 + \frac{v_1 v_2^*}{E_D^{(0)} - E_3^{(0)}} u_2 = E_D^{(0)} u_1, \quad \frac{v_1^* v_2}{E_D^{(0)} - E_3^{(0)}} u_1 + \left(E_D^{(0)} + \frac{|v_2|^2}{E_D^{(0)} - E_3^{(0)}} \right) u_2 = E_D^{(0)} u_2.$$

By inspection, these are satisfied when $u_1 = -v_2^*$ and $u_2 = v_1^*$.

For the eigenvector corresponding to E_2, we have

$$\begin{bmatrix} E_D^{(0)} + |v_1|^2/(E_D^{(0)} - E_3^{(0)}) & v_1 v_2^*/(E_D^{(0)} - E_3^{(0)}) \\ v_1^* v_2/(E_D^{(0)} - E_3^{(0)}) & E_D^{(0)} + |v_2|^2/(E_D^{(0)} - E_3^{(0)}) \end{bmatrix} \begin{bmatrix} w_1 \\ w_2 \end{bmatrix} = \left(E_D^{(0)} + \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}} \right) \begin{bmatrix} w_1 \\ w_2 \end{bmatrix}$$

which is equivalent to the system of equations

$$\left(E_D^{(0)} + \frac{|v_1|^2}{E_D^{(0)} - E_3^{(0)}} \right) w_1 + \frac{v_1 v_2^*}{E_D^{(0)} - E_3^{(0)}} w_2 = \left(E_D^{(0)} + \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}} \right) w_1,$$
$$\frac{v_1^* v_2}{E_D^{(0)} - E_3^{(0)}} w_1 + \left(E_D^{(0)} + \frac{|v_2|^2}{E_D^{(0)} - E_3^{(0)}} \right) w_2 = \left(E_D^{(0)} + \frac{|v_1|^2 + |v_2|^2}{E_D^{(0)} - E_3^{(0)}} \right) w_2.$$

By inspection, these are satisfied when $w_1 = v_1$ and $w_2 = w_2$. So we have the eigenvectors

$$|E_1\rangle = \begin{bmatrix} -v_2^* \\ v_1^* \end{bmatrix}, \qquad |E_2\rangle = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix}.$$

I consulted Sakurai's *Modern Quantum Mechanics*, Shankar's *Principles of Quantum Mechanics*, and Wolfram MathWorld while writing up these solutions.