

**Problem 1. (Jackson 14.1)** Verify by explicit calculation that the Liénard-Wiechert expressions for *all* components of  $\mathbf{E}$  and  $\mathbf{B}$  for a particle moving with constant velocity agree with the ones obtained in the text by means of a Lorentz transformation. Follow the general method at the end of Section 14.1.

**Solution.** The Liénard-Wiechert expressions for the fields are given by Jackson (14.13–14):

$$\mathbf{B} = [\hat{\mathbf{n}} \times \mathbf{E}]_{\text{ret}}, \quad \mathbf{E}(\mathbf{x}, t) = e \left[ \frac{\hat{\mathbf{n}} - \boldsymbol{\beta}}{\gamma^2(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} + \frac{e}{c} \left[ \frac{\hat{\mathbf{n}} \times \{(\hat{\mathbf{n}} - \boldsymbol{\beta}) \times \dot{\boldsymbol{\beta}}\}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R} \right]_{\text{ret}}, \quad (1)$$

where  $\boldsymbol{\beta} = \mathbf{v}/c$  with  $\mathbf{v}$  being the particle's velocity,  $R$  is the distance from the observation point to the particle's position, and  $\hat{\mathbf{n}}$  is a unit vector defined by  $\mathbf{x} - \mathbf{r} = R\hat{\mathbf{n}}$ , where  $\mathbf{r}$  is the position of the particle.

The expressions for the components of  $\mathbf{E}$  and  $\mathbf{B}$  obtained by a Lorentz transformation are given by Jackson (11.152):

$$E_1 = -\frac{e\gamma vt}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, \quad E_2 = \frac{e\gamma b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, \quad E_3 = B_1 = B_2 = 0, \quad B_3 = \beta E_2, \quad (2)$$

where the particle is moving in the  $x_1$  direction at impact parameter  $b$  on the  $x_2$  axis, as shown in Fig. (1).

For a particle moving with constant velocity in the  $x_1$  direction with velocity  $v$  as shown in Fig. (1),  $\boldsymbol{\beta} = \beta \hat{\mathbf{x}}_1$  and  $\dot{\boldsymbol{\beta}} = 0$ . From Jackson (14.16), note that

$$(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2 R^2 = b^2 + v^2 t^2 - \beta^2 b^2 = \frac{b^2 + \gamma^2 v^2 t^2}{\gamma^2} \implies (1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2 = \frac{(b^2 + \gamma^2 v^2 t^2)^{3/2}}{R\gamma^3}.$$

This calculation comes from Fig. (2), where  $O$  is the observation point,  $P$  is the present position of the particle, and  $P'$  its retarded position. Also from Fig. 2,

$$\hat{\mathbf{n}} = \cos \theta \hat{\mathbf{x}}_1 + \sin \theta \hat{\mathbf{x}}_2 = \frac{\beta R - vt}{R} \hat{\mathbf{x}}_1 + \frac{b}{R} \hat{\mathbf{x}}_2.$$

Making these substitutions in the expression for  $\mathbf{E}(\mathbf{x}, t)$  in Eq. (1),

$$\mathbf{E}(\mathbf{x}, t) = e \left[ \frac{\hat{\mathbf{n}} - \boldsymbol{\beta}}{\gamma^2(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} = e \left[ \frac{(\beta - vt/R - \beta) \hat{\mathbf{x}}_1 + (b/R) \hat{\mathbf{x}}_2}{\gamma^2(b^2 + \gamma^2 v^2 t^2)^{3/2}} R\gamma^3 \right]_{\text{ret}} = e\gamma \frac{-vt \hat{\mathbf{x}}_1 + b \hat{\mathbf{x}}_2}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}. \quad (3)$$

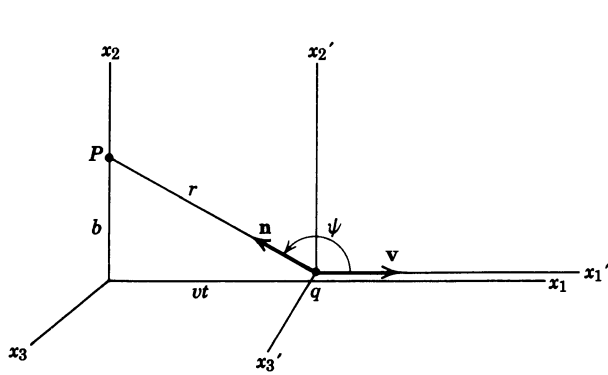


Figure 1: (Jackson Fig. 11.8) Particle of charge  $q$  moving at constant velocity  $\mathbf{v}$  passes an observation point  $P$  at impact parameter  $b$ .

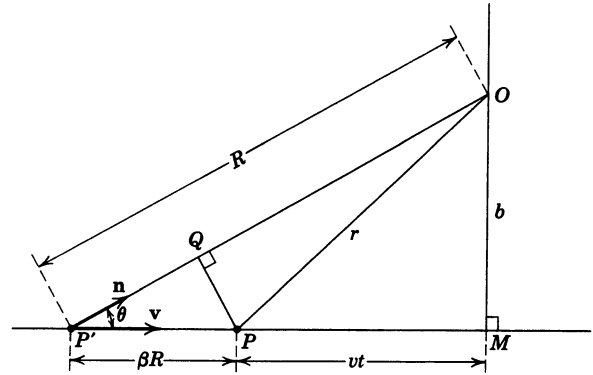


Figure 2: (Jackson Fig. 14.2) Present and retarded positions of a charge in uniform motion.

For  $\mathbf{B}(\mathbf{x}, t)$ , note that

$$\hat{\mathbf{n}} \times \mathbf{E} \propto \left( \frac{\beta R - vt}{R} \hat{\mathbf{x}}_1 + \frac{b}{R} \hat{\mathbf{x}}_2 \right) \times (-vt \hat{\mathbf{x}}_1 + b \hat{\mathbf{x}}_2) = \left( b \frac{\beta R - vt}{R} + \frac{bvt}{R} \right) \hat{\mathbf{x}}_3 = \beta b,$$

so

$$\mathbf{B}(\mathbf{x}, t) = e\gamma \frac{\beta b \hat{\mathbf{x}}_3}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}. \quad (4)$$

Writing Eqs. (3) and (4) in component notation, we find

$$\begin{aligned} E_1 &= -\frac{e\gamma vt}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, & E_2 &= \frac{e\gamma b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, & E_3 &= 0, \\ B_1 &= 0, & B_2 &= 0, & B_3 &= \frac{e\gamma \beta b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}} = \beta E_2, \end{aligned}$$

which are identical to Eq. (2) as was to be shown.  $\square$

**Problem 2. (Jackson 14.3)** The Heaviside-Feynman expression for the electric field of a particle of charge  $e$  in arbitrary motion, an alternative to the Liénard-Wiechert expression in Eq. (1), is

$$\mathbf{E} = e \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} + e \left[ \frac{R}{c} \right]_{\text{ret}} \frac{d}{dt} \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} + \frac{e}{c^2} \frac{d^2}{dt^2} [\hat{\mathbf{n}}]_{\text{ret}}, \quad (5)$$

where the time derivatives are with respect to the time at the observation point. Using the fact that the retarded time is  $t' = t - R(t')/c$  and that, as a result,

$$\frac{dt}{dt'} = 1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t'),$$

show that the form above yields the expression for  $\mathbf{E}$  in Eq. (1) when the time differentiations are performed.

**Solution.** One of the vector identities on the inside cover of Jackson is

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot \mathbf{c})\mathbf{b} - (\mathbf{a} \cdot \mathbf{b})\mathbf{c}.$$

From this,

$$\begin{aligned} \hat{\mathbf{n}} \times [(\hat{\mathbf{n}} - \boldsymbol{\beta}) \times \dot{\boldsymbol{\beta}}] &= (\hat{\mathbf{n}} \cdot \dot{\boldsymbol{\beta}})(\hat{\mathbf{n}} - \boldsymbol{\beta}) - [\hat{\mathbf{n}} \cdot (\hat{\mathbf{n}} - \boldsymbol{\beta})]\dot{\boldsymbol{\beta}} = (\hat{\mathbf{n}} \cdot \dot{\boldsymbol{\beta}})\hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \dot{\boldsymbol{\beta}})\boldsymbol{\beta} - \dot{\boldsymbol{\beta}} + (\hat{\mathbf{n}} \cdot \boldsymbol{\beta})\dot{\boldsymbol{\beta}} \\ &= \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}}}{c} - \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v}}{c^2} - \frac{\dot{\mathbf{v}}}{c} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}}}{c^2}. \end{aligned}$$

Then Eq. (1) can be written

$$\begin{aligned} \mathbf{E} &= e \left[ \frac{1}{\gamma^2(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( \hat{\mathbf{n}} - \frac{\mathbf{v}}{c} \right) \right]_{\text{ret}} + \frac{e}{c} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R} \left( \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}}}{c} - \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v}}{c^2} - \frac{\dot{\mathbf{v}}}{c} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}}}{c^2} \right) \right]_{\text{ret}} \\ &= e \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( \frac{\hat{\mathbf{n}}}{\gamma^2} - \frac{\mathbf{v}}{c\gamma^2} + \frac{R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}}}{c^2} - \frac{R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v}}{c^3} - \frac{R\dot{\mathbf{v}}}{c^2} + \frac{R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}}}{c^3} \right) \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( \frac{c^3 \hat{\mathbf{n}}}{\gamma^2} - \frac{c^2 \mathbf{v}}{\gamma^2} + cR(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} - cR\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}} \right) \right]_{\text{ret}}. \quad (6) \end{aligned}$$

Since  $R(t') = c(t - t')$ , note that

$$\left[ \frac{dR}{dt} \right]_{\text{ret}} = \frac{dR(t')}{dt'} = c \left( \frac{dt}{dt'} - 1 \right) = c(1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t') - 1) = -[\hat{\mathbf{n}} \cdot \mathbf{v}]_{\text{ret}}.$$

Note that  $[\hat{\mathbf{n}}]_{\text{ret}} = [(\mathbf{x} - \mathbf{r})/R]_{\text{ret}}$ . Then

$$\begin{aligned} \frac{d}{dt'} [\hat{\mathbf{n}}]_{\text{ret}} &= \left[ \frac{d\hat{\mathbf{n}}}{dt} \right]_{\text{ret}} = \left[ \frac{1}{R^2} \left( R \frac{d}{dt} (\mathbf{x} - \mathbf{r}) - (\mathbf{x} - \mathbf{r}) \frac{dR}{dt} \right) \right]_{\text{ret}} = \left[ \frac{1}{R^2} (-R\mathbf{v} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}}) \right]_{\text{ret}} \\ &= \left[ \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R} \right]_{\text{ret}}, \end{aligned}$$

$$\begin{aligned} \frac{d}{dt'} \left[ \frac{\hat{\mathbf{n}}}{R} \right]_{\text{ret}} &= \left[ \frac{d}{dt} \left( \frac{\hat{\mathbf{n}}}{R} \right) \right]_{\text{ret}} = \left[ \frac{1}{R^2} \left( R \frac{d\hat{\mathbf{n}}}{dt} - \hat{\mathbf{n}} \frac{dR}{dt} \right) \right]_{\text{ret}} = \left[ \frac{1}{R^2} \left\{ R \left( \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R} \right) + (\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} \right\} \right]_{\text{ret}} \\ &= \left[ \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v} + (\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} = \left[ \frac{2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R^2} \right]_{\text{ret}}, \end{aligned}$$

$$\begin{aligned} \frac{d}{dt'} \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} &= \left[ \frac{d}{dt} \left( \frac{\hat{\mathbf{n}}}{R^2} \right) \right]_{\text{ret}} = \left[ \frac{1}{R^2} \left\{ R \frac{d}{dt} \left( \frac{\hat{\mathbf{n}}}{R} \right) - \frac{\hat{\mathbf{n}}}{R} \frac{dR}{dt} \right\} \right]_{\text{ret}} = \left[ \frac{1}{R^2} \left( R \frac{2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R^2} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}}}{R} \right) \right]_{\text{ret}} \\ &= \left[ \frac{2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v} + (\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}}}{R^3} \right]_{\text{ret}} = \left[ \frac{3(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R^3} \right]_{\text{ret}}. \end{aligned}$$

For the second derivative of  $\hat{\mathbf{n}}$ , note that

$$\frac{d}{dt'} [\hat{\mathbf{n}} \cdot \mathbf{v}]_{\text{ret}} = \left[ \frac{d}{dt} (\hat{\mathbf{n}} \cdot \mathbf{v}) \right]_{\text{ret}} = \left[ \hat{\mathbf{n}} \cdot \frac{d\mathbf{v}}{dt} + \frac{d\hat{\mathbf{n}}}{dt} \cdot \mathbf{v} \right]_{\text{ret}} = \left[ \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R} \cdot \mathbf{v} \right]_{\text{ret}} = \left[ \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right]_{\text{ret}},$$

so then

$$\begin{aligned} \frac{d^2}{dt'^2} [\hat{\mathbf{n}}]_{\text{ret}} &= \left[ \frac{d^2 \hat{\mathbf{n}}}{dt^2} \right]_{\text{ret}} = \left[ \frac{d}{dt} \left( \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R} \right) \right]_{\text{ret}} = \left[ \frac{d}{dt} \left( \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}}}{R} \right) - \frac{d}{dt} \left( \frac{\mathbf{v}}{R} \right) \right]_{\text{ret}} \\ &= \left[ \frac{\hat{\mathbf{n}}}{R} \frac{d}{dt} (\hat{\mathbf{n}} \cdot \mathbf{v}) + (\hat{\mathbf{n}} \cdot \mathbf{v}) \frac{d}{dt} \left( \frac{\hat{\mathbf{n}}}{R} \right) - \frac{1}{R^2} \left( R \frac{d\mathbf{v}}{dt} - \mathbf{v} \frac{dR}{dt} \right) \right]_{\text{ret}} \\ &= \left[ \frac{\hat{\mathbf{n}}}{R} \left( \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right) + (\hat{\mathbf{n}} \cdot \mathbf{v}) \left( \frac{2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - \mathbf{v}}{R^2} \right) - \frac{1}{R^2} (R\dot{\mathbf{v}} + (\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}) \right]_{\text{ret}} \\ &= \left[ \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}}}{R} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{R^2} - \frac{\mathbf{v}^2 \hat{\mathbf{n}}}{R^2} + 2 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{R^2} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}}{R^2} - \frac{\dot{\mathbf{v}}}{R} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}}{R^2} \right]_{\text{ret}} \\ &= \left[ \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}}}{R} + 3 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{R^2} - \frac{\mathbf{v}^2 \hat{\mathbf{n}}}{R^2} - 2 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}}{R^2} - \frac{\dot{\mathbf{v}}}{R} \right]_{\text{ret}}. \end{aligned}$$

By the chain rule,

$$\frac{d}{dt} = \frac{dt'}{dt} \frac{d}{dt'} = \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \frac{d}{dt'} = \left[ \frac{1}{1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}}} \right]_{\text{ret}} \frac{d}{dt'}.$$

For the second derivative, note that

$$\begin{aligned} \frac{d}{dt'} \left[ \frac{1}{1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}}} \right]_{\text{ret}} &= \left[ \frac{d}{dt} \left( \frac{1}{\mathbf{v} \cdot \hat{\mathbf{n}}/c - 1} \right) \right]_{\text{ret}} = \left[ \frac{1}{(1 - \mathbf{v} \cdot \hat{\mathbf{n}}/c)^2} \frac{d}{dt} \left( 1 - \frac{\hat{\mathbf{n}} \cdot \mathbf{v}}{c} \right) \right]_{\text{ret}} \\ &= \frac{1}{c} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2} \left( \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right) \right]_{\text{ret}}, \end{aligned}$$

so

$$\begin{aligned} \frac{d^2}{dt^2} &= \frac{d}{dt} \left( \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \frac{d}{dt'} \right) = \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \frac{d}{dt'} \left( \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \frac{d}{dt'} \right) \\ &= \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \left\{ \frac{d}{dt'} \left( \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \right) \frac{d}{dt'} + \frac{1}{1 - \boldsymbol{\beta}(t') \cdot \hat{\mathbf{n}}(t')} \frac{d^2}{dt'^2} \right\} \\ &= \frac{1}{c} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3} \left( \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right) \right]_{\text{ret}} \frac{d}{dt'} + \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2} \right]_{\text{ret}} \frac{d^2}{dt'^2}. \end{aligned}$$

The first term of Eq. (5) can be written

$$\begin{aligned} e \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} &= e \left[ \frac{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 \hat{\mathbf{n}}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} = e \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( \hat{\mathbf{n}} - 3 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}}}{c} + 3 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{c^2} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}}}{c^3} \right) \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ c^3 \hat{\mathbf{n}} - 3c^2 (\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} + 3c (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} \} \right]_{\text{ret}}. \end{aligned} \quad (7)$$

The second term of Eq. (5) becomes

$$\begin{aligned} e \left[ \frac{R}{c} \right]_{\text{ret}} \frac{d}{dt} \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} &= \frac{e}{c} \left[ \frac{R}{1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}}} \right]_{\text{ret}} \frac{d}{dt'} \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} = \frac{e}{c} \left[ \frac{3(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - \mathbf{v}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}}) R^2} \right]_{\text{ret}} = \frac{e}{c} \left[ \frac{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2 \{ 3(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - \mathbf{v} \}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} \\ &= \frac{e}{c} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( 1 - 2 \frac{\hat{\mathbf{n}} \cdot \mathbf{v}}{c} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2}{c^2} \right) \{ 3(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - \mathbf{v} \} \right]_{\text{ret}} \\ &= \frac{e}{c} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( 3(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - \mathbf{v} - 6 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{c} + 2 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v}}{c} + 3 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}}}{c^2} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v}}{c^2} \right) \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{3c^2 (\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - c^2 \mathbf{v} - 6c (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} + 2c (\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v} + 3(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}}. \end{aligned} \quad (8)$$

The third term of Eq. (5) becomes

$$\frac{e}{c^2} \frac{d^2}{dt^2} [\hat{\mathbf{n}}]_{\text{ret}} = \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3} \left( \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right) \right]_{\text{ret}} \frac{d}{dt'} [\hat{\mathbf{n}}]_{\text{ret}} + \frac{e}{c^2} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2} \right]_{\text{ret}} \frac{d^2}{dt'^2} [\hat{\mathbf{n}}]_{\text{ret}}. \quad (9)$$

The first term of Eq. (9) is

$$\begin{aligned} &\frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3} \left( \hat{\mathbf{n}} \cdot \dot{\mathbf{v}} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 - \mathbf{v}^2}{R} \right) \frac{(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - \mathbf{v}}{R} \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3} \left( \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}}}{R} + \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}}}{R^2} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v}^2 \hat{\mathbf{n}}}{R^2} - \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}}) \mathbf{v}}{R} - \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v}}{R^2} + \frac{\mathbf{v}^2 \mathbf{v}}{R^2} \right) \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} + (\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v}^2 \hat{\mathbf{n}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}}) \mathbf{v} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + \mathbf{v}^2 \mathbf{v}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}}, \end{aligned}$$

and the second term of Eq. (9) is

$$\begin{aligned} &\frac{e}{c^2} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^2} \left( \frac{(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}}) \hat{\mathbf{n}}}{R} + 3 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}}}{R^2} - \frac{\mathbf{v}^2 \hat{\mathbf{n}}}{R^2} - 2 \frac{(\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v}}{R^2} - \frac{\dot{\mathbf{v}}}{R} \right) \right]_{\text{ret}} \\ &= \frac{e}{c^2} \left[ \frac{1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}}) \hat{\mathbf{n}} + 3(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - \mathbf{v}^2 \hat{\mathbf{n}} - 2(\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v} - R \dot{\mathbf{v}} \} \right]_{\text{ret}} \\ &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ R c (\hat{\mathbf{n}} \cdot \dot{\mathbf{v}}) \hat{\mathbf{n}} + 3c (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - c \mathbf{v}^2 \hat{\mathbf{n}} - 2c (\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v} - R c \dot{\mathbf{v}} \right. \\ &\quad \left. - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})(\hat{\mathbf{n}} \cdot \mathbf{v}) \hat{\mathbf{n}} - 3(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} + (\hat{\mathbf{n}} \cdot \mathbf{v}) \mathbf{v}^2 \hat{\mathbf{n}} + 2(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + R(\hat{\mathbf{n}} \cdot \mathbf{v}) \dot{\mathbf{v}} \} \right]_{\text{ret}}. \end{aligned}$$

Summing the two terms, we find

$$\begin{aligned}
\frac{e}{c^2} \frac{d^2}{dt^2} [\hat{\mathbf{n}}]_{\text{ret}} &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} + (\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}^2 \hat{\mathbf{n}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + \mathbf{v}^2 \mathbf{v} \right. \\
&\quad + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} + 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - c\mathbf{v}^2 \hat{\mathbf{n}} - 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} - Rc\dot{\mathbf{v}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} \\
&\quad \left. - 3(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} + (\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v}^2 \hat{\mathbf{n}} + 2(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}} \} \right]_{\text{ret}} \\
&= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ -2(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} + (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + \mathbf{v}^2 \mathbf{v} + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} + 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} \right. \\
&\quad \left. - c\mathbf{v}^2 \hat{\mathbf{n}} - 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} - Rc\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}} \} \right]_{\text{ret}}. \quad (10)
\end{aligned}$$

Summing Eqs. (7) and (8),

$$\begin{aligned}
e \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} + e \left[ \frac{R}{c} \right]_{\text{ret}} \frac{d}{dt} \left[ \frac{\hat{\mathbf{n}}}{R^2} \right]_{\text{ret}} &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ c^3 \hat{\mathbf{n}} - 3c^2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} + 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} \right. \\
&\quad \left. + 3c^2(\hat{\mathbf{n}} \cdot \mathbf{v})\hat{\mathbf{n}} - c^2 \mathbf{v} - 6c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} + 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} + 3(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} \} \right]_{\text{ret}} \\
&= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ c^3 \hat{\mathbf{n}} - 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} + 2(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - c^2 \mathbf{v} + 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} \} \right]_{\text{ret}}. \quad (11)
\end{aligned}$$

Summing Eqs. (11) and (10),

$$\begin{aligned}
\mathbf{E} &= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \{ c^3 \hat{\mathbf{n}} - 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} + 2(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - c^2 \mathbf{v} + 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} - (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} \right. \\
&\quad \left. - 2(\hat{\mathbf{n}} \cdot \mathbf{v})^3 \hat{\mathbf{n}} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} + (\hat{\mathbf{n}} \cdot \mathbf{v})^2 \mathbf{v} + \mathbf{v}^2 \mathbf{v} + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} \right. \\
&\quad \left. + 3c(\hat{\mathbf{n}} \cdot \mathbf{v})^2 \hat{\mathbf{n}} - c\mathbf{v}^2 \hat{\mathbf{n}} - 2c(\hat{\mathbf{n}} \cdot \mathbf{v})\mathbf{v} - Rc\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}} \} \right]_{\text{ret}} \\
&= \frac{e}{c^3} \left[ \frac{c^3 \hat{\mathbf{n}} - c\mathbf{v}^2 \hat{\mathbf{n}} + \mathbf{v}^2 \mathbf{v} - c^2 \mathbf{v} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} - Rc\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} \\
&= \frac{e}{c^3} \left[ \frac{c^3(1 - \beta^2)\hat{\mathbf{n}} - c^2(1 - \beta^2)\mathbf{v} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} - Rc\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}}}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \right]_{\text{ret}} \\
&= \frac{e}{c^3} \left[ \frac{1}{(1 - \boldsymbol{\beta} \cdot \hat{\mathbf{n}})^3 R^2} \left( \frac{c^3 \hat{\mathbf{n}}}{\gamma^2} - \frac{c^2 \mathbf{v}}{\gamma^2} - R(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\mathbf{v} + Rc(\hat{\mathbf{n}} \cdot \dot{\mathbf{v}})\hat{\mathbf{n}} - Rc\dot{\mathbf{v}} + R(\hat{\mathbf{n}} \cdot \mathbf{v})\dot{\mathbf{v}} \right) \right]_{\text{ret}},
\end{aligned}$$

which is identical to Eq. (6). Thus, we have proven that the Heaviside-Feynman and Liénard-Wiechert expressions for  $\mathbf{E}$  are equivalent.  $\square$

**Problem 3. (Jackson 14.4)** Using the Liénard-Wiechart fields, discuss the time-averaged power radiated per unit solid angle in nonrelativistic motion of a particle with charge  $e$ , moving as described below. Sketch the angular distribution of the radiation and determine the total power radiated in each case.

**3(a)** The particle is moving along the  $z$  axis with instantaneous position  $z(t) = \alpha \cos \omega_0 t$ .

**Solution.** For a nonrelativistic particle, the power radiated per unit solid angle is given by Jackson (14.21),

$$\frac{dP}{d\Omega} = \frac{e^2}{4\pi c^2} |\dot{\mathbf{v}}|^2 \sin^2 \Theta, \quad (12)$$

where  $\Theta$  is the angle between  $\dot{\mathbf{v}}$  and  $\hat{\mathbf{n}}$ , and  $\hat{\mathbf{n}}$  is a unit vector pointing toward the observer. The total instantaneous power radiated is given by Jackson (14.22):

$$P = \frac{2}{3} \frac{e^2}{c^3} |\dot{\mathbf{v}}|^2. \quad (13)$$

In this case, we have

$$\mathbf{x}(t) = \alpha \cos \omega_0 t \hat{\mathbf{x}}_3, \quad \mathbf{v}(t) = -\alpha \omega_0 \sin \omega_0 t \hat{\mathbf{x}}_3, \quad \dot{\mathbf{v}}(t) = -\alpha \omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_3.$$

The system is azimuthally symmetric since  $\dot{\mathbf{v}}$  always points along the  $z$  axis. Thus,  $\Theta = \theta$  where  $\theta$  is the polar angle in spherical coordinates. Equation (12) becomes

$$\frac{dP}{d\Omega} = \frac{e^2}{4\pi c^2} |-\alpha \omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_3|^2 \sin^2 \theta = \frac{e^2 \alpha^2 \omega_0^4}{4\pi c^2} \cos^2 \omega_0 t \sin^2 \theta,$$

so the time-averaged power radiated per unit solid angle is

$$\left\langle \frac{dP}{d\Omega} \right\rangle = \frac{e^2}{4\pi c^2} \alpha^2 \omega_0^4 \langle \cos^2 \omega_0 t \rangle \sin^2 \theta = \frac{e^2 \alpha^2 \omega_0^4}{8\pi c^2} \sin^2 \theta. \quad (14)$$

A plot of the angular distribution of the radiation in the  $xz$  plane is shown in Fig. 3, and in three dimensions in Fig. 4. This is the typical pattern for radiation of an accelerating, non-relativistic point charge, as we saw on p. 214 of the lecture notes. It is also the radiation pattern for a pure, single dipole oscillating parallel to the  $z$  axis ( $m = 0$ ), as is seen in Jackson Table 9.1 and Jackson Fig. 9.5.

Equation (13) becomes

$$P = \frac{2}{3} \frac{e^2}{c^3} |-\alpha \omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_3|^2 = \frac{2}{3} \frac{e^2 \alpha^2 \omega_0^4}{c^3} \cos^2 \omega_0 t,$$

so the time-averaged total power radiated is

$$\langle P \rangle = \frac{2}{3} \frac{e^2 \alpha^2 \omega_0^4}{c^3} \langle \cos^2 \omega_0 t \rangle = \frac{e^2 \alpha^2 \omega_0^4}{3c^3}.$$

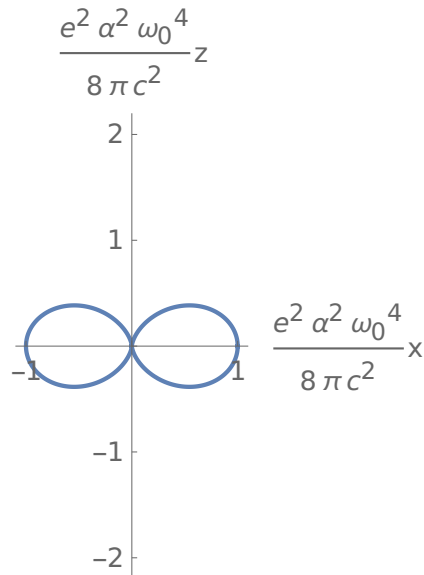


Figure 3: Plot of Eq. (14) in the  $xz$  plane.

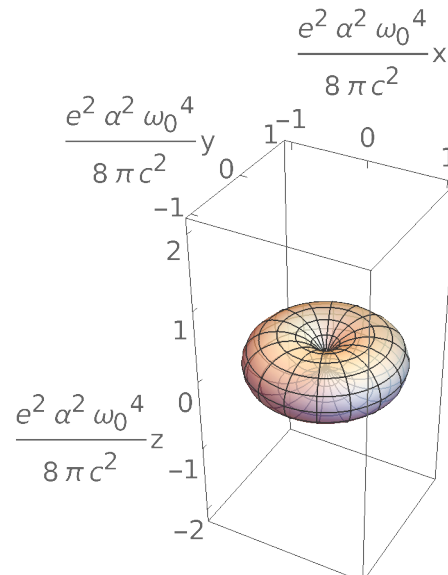


Figure 4: Three-dimensional plot of Eq. (14).

**3(b)** The particle is moving in a circle of radius  $R$  in the  $xy$  plane with constant angular frequency  $\omega_0$ .

**Solution.** For a charge moving counter-clockwise,

$$\begin{aligned}\mathbf{x}(t) &= R \cos \omega_0 t \hat{\mathbf{x}}_1 - R \sin \omega_0 t \hat{\mathbf{x}}_2, & \mathbf{v}(t) &= -R\omega_0 \sin \omega_0 t \hat{\mathbf{x}}_1 - R\omega_0 \cos \omega_0 t \hat{\mathbf{x}}_2, \\ \dot{\mathbf{v}}(t) &= -R\omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_1 + R\omega_0^2 \sin \omega_0 t \hat{\mathbf{x}}_2.\end{aligned}$$

This system is also azimuthally symmetric, so it is sufficient to restrict the position of the observer to the  $yz$  plane. In polar coordinates,  $\hat{\mathbf{n}} = \sin \theta \hat{\mathbf{x}}_2 + \cos \theta \hat{\mathbf{x}}_3$ . Then  $\sin^2 \Theta$  can be found by

$$\sin^2 \Theta = 1 - \cos^2 \Theta = 1 - \frac{(\dot{\mathbf{v}} \cdot \hat{\mathbf{n}})^2}{\dot{v}^2} = 1 - \sin^2 \theta \sin^2 \omega_0 t.$$

With these substitutions, Eq. (12) becomes

$$\begin{aligned}\frac{dP}{d\Omega} &= \frac{e^2}{4\pi c^2} |-R\omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_1 + R\omega_0^2 \sin \omega_0 t \hat{\mathbf{x}}_2|^2 (1 - \sin^2 \theta \sin^2 \omega_0 t) \\ &= \frac{e^2 R^2 \omega_0^4}{4\pi c^2} (\cos^2 \omega_0 t + \sin^2 \omega_0 t) (1 - \sin^2 \theta \sin^2 \omega_0 t) = \frac{e^2 R^2 \omega_0^4}{4\pi c^2} (1 - \sin^2 \theta \sin^2 \omega_0 t),\end{aligned}$$

giving us the time-averaged power radiated per unit solid angle:

$$\begin{aligned}\left\langle \frac{dP}{d\Omega} \right\rangle &= \frac{e^2 R^2 \omega_0^4}{4\pi c^2} (1 - \sin^2 \theta \langle \sin^2 \omega_0 t \rangle) = \frac{e^2 R^2 \omega_0^4}{4\pi c^2} \left( 1 - \frac{\sin^2 \theta}{2} \right) = \frac{e^2 R^2 \omega_0^4}{4\pi c^2} \left( 1 - \frac{1 - \cos^2 \theta}{2} \right) \\ &= \frac{e^2 R^2 \omega_0^4}{8\pi c^2} (1 + \cos^2 \theta).\end{aligned}\tag{15}$$

A plot of the angular distribution of the radiation in the  $xz$  plane is shown in Fig. 5, and in three dimensions in Fig. 6. The magnitude of the radiation on the  $xy$  plane is reduced because of the particle's motion in that plane. This is also the radiation pattern for two dipoles, one on each of the  $x$  and  $y$  axes, oscillating out of phase ( $m = \pm 1$ ), as is seen in Jackson Table 9.1 and Jackson Fig. 9.5.

From Eq. (13), we have

$$P = \frac{2e^2}{3c^3} |-R\omega_0^2 \cos \omega_0 t \hat{\mathbf{x}}_1 + R\omega_0^2 \sin \omega_0 t \hat{\mathbf{x}}_2|^2 = \frac{2e^2 R^2 \omega_0^4}{3c^3} (\cos^2 \omega_0 t + \sin^2 \omega_0 t) = \frac{2e^2 R^2 \omega_0^4}{3c^3} = \langle P \rangle.$$

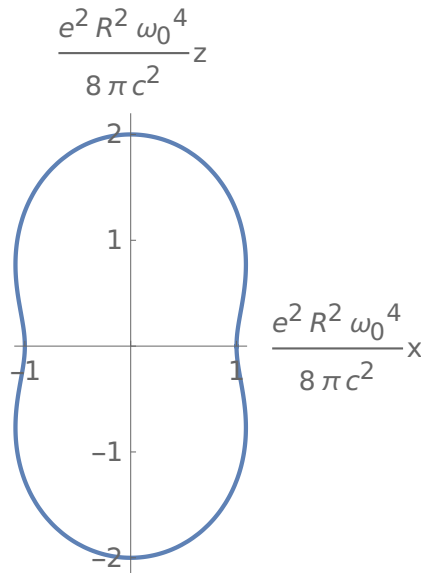


Figure 5: Plot of Eq. (15) in the  $xz$  plane.

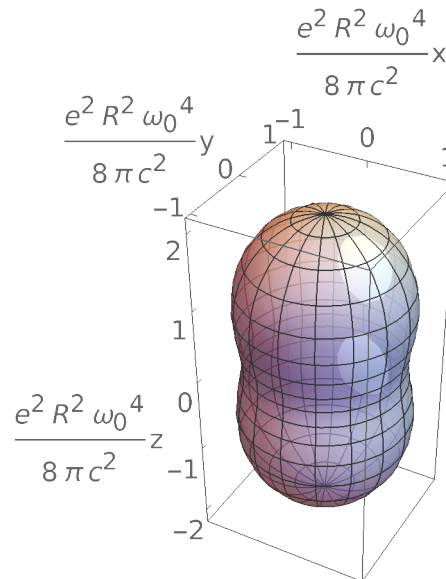


Figure 6: Three-dimensional plot of Eq. (15).