Scattering n-p interaction

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The nucleon-nucleon potential is assume as:

$$V_{ij} = (V_R + \frac{1}{2}(1 + P_{ij}^{\sigma})V_t + \frac{1}{2}(1 - P_{ij}^{\sigma})V_s)(\frac{1}{2}u + \frac{1}{2}(2 - u)P_{ij}^{\tau}) + \frac{1}{4}(1 + \tau_{iz})(1 + \tau_{jz})\frac{e^2}{r_{ij}}$$
(1)

take $u \approx 1$

$$V_{ij} = \frac{1}{2} \left[V_R + \frac{1}{2} (1 + P_{ij}^{\sigma}) V_t + \frac{1}{2} (1 - P_{ij}^{\sigma}) V_s \right] (1 + P_{ij}^r) + \frac{1}{4} (1 + \tau_{iz}) (1 + \tau_{jz}) \frac{e^2}{r_{ij}}$$
(2)

in which the operator come defined as $P_{ij}^{\sigma}=(-1)^{s+1}$ and $P_{ij}^{r}=(-1)^{l}$. For n-p interaction we have $\tau_{iz}=1$ and $\tau_{jz}=-1$. We take account all possible cases to s and l. In each case, the quantum number is in 0,1 set hence there are 4 possible combination:

Case I: l = 0 y s = 0

$$V_{ij} = \frac{1}{2}(V_R + \frac{1}{2}(1 + (-1))V_t + \frac{1}{2}(1 - (-1))V_s)(1 + 1) + \frac{1}{4}(1 + 1)(1 + (-1))\frac{e^2}{r_{ij}}$$
(3)

$$V_{ij} = V_R + V_s \tag{4}$$

Case II: l = 0 y s = 1

$$V_{ij} = \frac{1}{2} \left(V_R + \frac{1}{2} (1+1) V_t + \frac{1}{2} (1-1) V_s \right) (1+1) + \frac{1}{4} (1+1) (1+(-1)) \frac{e^2}{r_{ij}}$$
 (5)

$$V_{ij} = V_R + V_t \tag{6}$$

Case III: l = 1 y s = 0

$$V_{ij} = \frac{1}{2}(V_R + \frac{1}{2}(1 + (-1))V_t + \frac{1}{2}(1 - (-1))V_s)(1 + (-1)) + \frac{1}{4}(1 + 1)(1 + (-1))\frac{e^2}{r_{ij}}$$
(7)

$$V_{ij} = 0 (8)$$

Case IV: l = 0 y s = 1

$$V_{ij} = \frac{1}{2}(V_R + \frac{1}{2}(1+1)V_t + \frac{1}{2}(1-1)V_s)(1+(-1)) + \frac{1}{4}(1+1)(1+(-1))\frac{e^2}{r_{ij}}$$
(9)

$$V_{ij} = 0 (10)$$

In this way is possible take account only two firsts cases

$$V(r) = \begin{cases} V_R + V_s & \text{if } l = 0 \text{ and } s = 0\\ V_R + V_t & \text{if } l = 0 \text{ and } s = 1\\ 0 & \text{other case.} \end{cases}$$
 (11)

with

$$V_R = V_{0R} \exp(-\kappa_R r^2),$$

$$V_t = -V_{0t} \exp(-\kappa_t r^2),$$

$$V_s = -V_{0s} \exp(-\kappa_s r^2),$$
with

$$V_{0R} = 200.0 \text{ MeV}, \quad \kappa_R = 1.487 \text{ fm}^{-2},$$

 $V_{0t} = 178.0 \text{ MeV}, \quad \kappa_t = 0.639 \text{ fm}^{-2},$
 $V_{0s} = 91.85 \text{ MeV}, \quad \kappa_s = 0.465 \text{ fm}^{-2}.$

to find $f(\theta)$ put potential in the next integral

$$f(\theta) = \int_0^\infty r \sin(qr) V(r) dr \tag{12}$$

for all non-vanish cases

Case I:

$$f(\theta) = \int_0^\infty r \sin(qr)(V_R(r) + V_s(r))dr \tag{13}$$

$$f(\theta) = \int_0^\infty r \sin(qr) \left(V_{0R} \exp(-\kappa_R r^2) - V_{0s} \exp(-\kappa_s r^2) \right) dr$$
(14)

Case II:

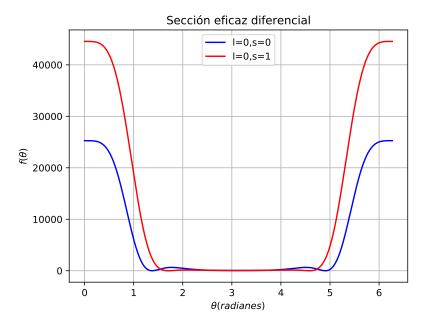
$$f(\theta) = \int_0^\infty r \sin(qr)(V_R(r) + V_t(r))dr \tag{15}$$

$$f(\theta) = \int_0^\infty r \sin(qr) \left(V_{0R} \exp(-\kappa_R r^2) - V_{0t} \exp(-\kappa_t r^2) \right) dr$$
(16)

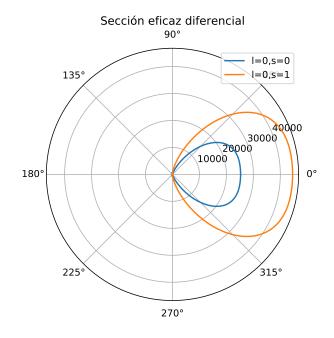
solving integral numerically we get $f(\theta)$, with this

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2 \tag{17}$$

the results are shown in Figure 1



(a) Plot in cartesian coordinates.



(b) Plot in polar coordinates.

Figure 1: Plot of numerical results for differential cross-section in cartesian and polar coordinates.