

Physics 562 - Computational Physics

Assignment 3: Chaotic Double Pendulum

Josh Fernandes
Department of Physics & Astronomy
California State University Long Beach

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Abstract

This paper examine emergent chaos by analyzing a double pendulum. A double pendulum is a chaotic system that is nonlinear, deterministic, and highly sensitive to initial conditions. There are three states for the double pendulum to be in: no chaos, emerging chaos, and full chaos. These states correspond to a Lyapunov exponent that is negative, zero, and positive, respectively

1 Single Pendulum

A single pendulum is a very predictable system that undergoes simple harmonic motion. The kinetic energy of a pendulum is given by

$$T = \frac{1}{2}mv^2 = \frac{1}{2}m(l\dot{\theta})^2 \quad (1)$$

and the potential energy is given by

$$V = mgh = mgl(1 - \cos(\theta)) \quad (2)$$

where l is the length of the pendulum, and m is the mass of the bob. The lagrangian is given by

$$\mathcal{L} = T - V = \frac{1}{2}m(l\dot{\theta})^2 - mgl(1 - \cos(\theta)). \quad (3)$$

The euler-lagrange equations can be applied to the lagrangian in order to get the equations of motions, but instead we will use the hamiltonian. The hamiltonian is given by

$$\mathcal{H} = \Sigma p_{q_i} \dot{q}_i - \mathcal{L} = p_{\theta} \dot{\theta} - \frac{1}{2}m(l\dot{\theta})^2 + mgl(1 - \cos(\theta)). \quad (4)$$

Furthurmore we know that

$$p_{\theta} = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = ml^2 \dot{\theta} \quad (5)$$

which can be rewritten

$$\dot{\theta} = \frac{p_{\theta}}{ml^2} \quad (6)$$

so we can substitute into the hamiltonian to get it of the form

$$\mathcal{H} = \frac{p_{\theta}^2}{ml^2} - \frac{1}{2}ml^2 \left(\frac{p_{\theta}}{ml^2} \right)^2 + mgl(1 - \cos(\theta)). \quad (7)$$

$$\mathcal{H} = \frac{p_{\theta}^2}{ml^2} - \frac{1}{2} \frac{p_{\theta}^2}{ml^2} + mgl(1 - \cos(\theta)). \quad (8)$$

$$\mathcal{H} = \frac{1}{2} \frac{p_{\theta}^2}{ml^2} + mgl(1 - \cos(\theta)). \quad (9)$$

We can then solve the hamiltonian equations of motion which are

$$\dot{\theta} = \frac{\partial \mathcal{H}}{\partial p_{\theta}} = \frac{p_{\theta}}{ml^2} \quad (10)$$

$$\dot{p}_{\theta} = -\frac{\partial \mathcal{H}}{\partial \theta} = -mgl \sin \theta \quad (11)$$

2 Double Pendulum

The double pendulum is more complex. The equations of motion for the double pendulum is given by

$$\dot{\theta}_1 = \frac{\partial \mathcal{H}}{\partial p_1} = \frac{lp_1 - lp_2 \cos(\theta_1 - \theta_2)}{l^3[m_1 + m_2 \sin^2(\theta_1 - \theta_2)]} \quad (12)$$

$$\dot{\theta}_2 = \frac{\partial \mathcal{H}}{\partial p_2} = \frac{l(m_1 + m_2)p_2 - lm_2 p_1 \cos(\theta_1 - \theta_2)}{l^3 m_2 [m_1 + m_2 \sin^2(\theta_1 - \theta_2)]} \quad (13)$$

$$\dot{p}_1 = -\frac{\partial \mathcal{H}}{\partial \theta_1} = -(m_1 + m_2)gl \sin(\theta_1) - C_1 + C_2 \quad (14)$$

$$\dot{p}_2 = -\frac{\partial \mathcal{H}}{\partial \theta_2} = -m_2 gl \sin(\theta_2) + C_1 - C_2 \quad (15)$$

where

$$C_1 \equiv \frac{p_1 p_2 \sin(\theta_1 - \theta_2)}{l^2 [m_1 + m_2 \sin^2(\theta_1 - \theta_2)]} \quad (16)$$

$$C_2 \equiv \frac{l^2 m_2 p_1^2 + l^2 (m_1 + m_2) p_2^2 - l^2 m_2 p_1 p_2 \cos(\theta_1 - \theta_2)}{2l^4 [m_1 + m_2 \sin^2(\theta_1 - \theta_2)]^2} \sin[2(\theta_1 - \theta_2)] \quad (17)$$

In order to understand how the double pendulum behaves, the momentum will be plotted against the angle for each pendulum. In addition, the lyapunov exponent will be calculated for three different cases. These cases are

Listing 1: Module Cases

```

1
2      !!Low Energy - No Chaos
3      y(1)=1._dp*(pi/180) ! initial angle for 1st pendulum
4      y(2)=0._dp          ! initial momentum for 1st pendulum
5      y(3)=1._dp*(pi/180) ! initial angle for 2nd pendulum
6      y(4)=0._dp          ! initial momentum for 2nd pendulum
7
8      !!Medium Energy - Emerging Chaos But Stable Orbit
9      y(1)=1._dp*(pi/180) ! initial angle for 1st pendulum
10     y(2)=1._dp           ! initial momentum for 1st pendulum
11     y(3)=1._dp*(pi/180) ! initial angle for 2nd pendulum
12     y(4)=1.01_dp        ! initial momentum for 2nd pendulum

```

```

13
14      !!High Energy - Full Chaos
15      y(1)=1._dp*(pi/180) ! initial angle for 1st pendulum
16      y(2)=1._dp          ! initial momentum for 1st pendulum
17      y(3)=1._dp*(pi/180) ! initial angle for 2nd pendulum
18      y(4)=1.3_dp         ! initial momentum for 2nd pendulum

```

The mass for the first bob is 1 kg and the mass of the second bob is 2 kg. The length of both pendulums is 1 meter. The lyapunov exponents are given by

$$|p_1 - p_2| = e^{\lambda \cdot t} \quad (18)$$

where P_1 and p_2 are the momenta for the two pendulums, and λ is the lyapunov exponent.

3 The Fortran95 code

The code solves the equation of motion using the Runge-Kutta method. First a module called `NumType` is created to store all my global parameters.

Listing 2: Module `NumType`

```

1
2  module NumType
3
4      save
5      integer, parameter :: dp = kind(1.d0)
6      real(dp), parameter :: pi = 4*atan(1._dp), &
7      e = exp(1._dp)
8      complex(dp), parameter :: iic = (0._dp,1._dp)
9
10 end module NumType

```

Listing 3: `rkf45.f95`

```

1
2  subroutine rkf45step(t,y,h) !4-th order Runge-Kutta step
3
4      use setup, only : dp, n_eq
5      implicit none

```

```

6      real(dp), intent(inout) :: t, h
7      real(dp), dimension(n_eq), intent(inout) :: y
8      real(dp), dimension(n_eq) :: k1, k2, k3, &
9          k4, k5, k6, y1, y2
10     real(dp), parameter :: epsilon = 1.e-6_dp, &
11         tiny = 1.e-20_dp
12     real(dp) :: rr, delta
13
14     call deriv(t,      h,  y,      k1)
15     call deriv(t+h/4,  h,  y+ k1/4,  k2)
16     call deriv(t+3*h/8, h,  y+ (3*k1+9*k2)/32, k3)
17     call deriv(t+12*h/13, h, y+ (1932*k1-7200*k2+7296*k3) &
18         2197 , k4)
19     call deriv(t+h,h,  y+ (439*k1/216-8*k2+3680*k3/513 &
20         -845*k4/4104), k5)
21     call deriv(t+h/2,  h,  y+ (-8*k1/27 +2*k2-3544*k3/2565 + &
22         1859*k4/4104 -11*k5/40), k6)
23
24     y1 = y + 25*k1/216 + 1408*k3/2565 + 2197*k4/4104 &
25         - k5/5
26     y2 = y + 16*k1/135 + 6656*k3/12825 + 28561*k4/ &
27         56430 - 9*k5/50 + 2*k6/55
28
29     rr = sqrt(dot_product(y1-y2,y1-y2))/h + tiny
30
31     if ( rr < epsilon ) then
32         t = t + h
33         y = y1
34         delta = 0.92_dp * (epsilon/rr)**(0.2_dp)
35         h = delta*h
36         write (unit = 3,fmt='(3f20.10)') t, y(1)
37         write (unit = 4,fmt='(3f20.10)') t, y(2)
38         write (unit = 5,fmt='(3f20.10)') t, y(3)
39         write (unit = 6,fmt='(3f20.10)') t, y(4)
40         write (unit = 7,fmt='(3f20.10)') y(1), y(2)
41         write (unit = 8,fmt='(3f20.10)') y(3), y(4)
42         write (unit = 9,fmt='(3f20.10)') -sin(y(1)), -cos(y(1))
43         write (unit = 10,fmt='(3f20.10)') -sin(y(1))-sin(y(3)) &
44             , -cos(y(1))-cos(y(3))
45     else

```

```

46         delta = 0.92_dp * (epsilon/rr)**(0.25_dp)
47         h = delta*h
48     end if
49
50
51     contains
52
53         subroutine deriv(t,h,y,k)      ! derivative
54
55             use setup, only : dp, n_eq, g, length, mass1, mass2
56             implicit none
57             real(dp), intent(in) :: t, h
58             real(dp), dimension(n_eq), intent(in) :: y
59             real(dp), dimension(n_eq) :: f, k
60             real(dp) :: c1, c2
61
62             c1 = (y(2)*y(4)*sin(y(1)-y(3)))/ &
63                 (length*length*(mass1+mass2*(sin(y(1)-y(3)))**2))
64             c2 = (length**2*mass2*y(2)**2 + &
65                 length**2*(mass1+mass2)*y(4)**2- &
66                 length*length*mass2*y(2)*y(4)*cos(y(1)-y(3)))/ &
67                 (2*length**4*(mass1+mass2*(sin(y(1)-y(3)))**2)**2)* &
68                 sin(2*(y(1)-y(3)))
69
70             f(1) = (length*y(2) - length*y(4)*cos(y(1)-y(2)))/ &
71                 (length**3*(mass1+mass2*(sin(y(1)-y(2)))**2))
72             f(2) = -(mass1+mass2)*g*length*sin(y(1)) - c1 + c2
73             f(3) = (length*(mass1+mass2)*y(4) - length*&
74                 mass2*y(2)*cos(y(1)-y(2)))/ &
75                 (length**3*mass2*(mass1+mass2*(sin(y(1)-y(2)))**2))
76             f(4) = -mass2*g*length*sin(y(3)) + c1 - c2
77
78             k(1:n_eq) = h*f(1:n_eq)
79
80         end subroutine deriv

```

The main program is `adpend` and it begins with its own module.

Listing 4: `adpend.f95`

```

1
2

```

```

3  module setup
4
5      use NumType
6      implicit none
7      integer, parameter :: n_eq = 4
8      real(dp), parameter:: g = 10.0_dp, length = 1.0_dp, &
9          mass1 = 1.0_dp, mass2 = 2.0_dp
10     real(dp) :: t, tmax, dt, lambda
11     real(dp), dimension(n_eq) :: y
12
13 end module setup
14
15 program pendulum
16
17     use setup
18     implicit none
19
20     !!initial conditions!!
21
22     t = 0._dp                ! time to start
23     tmax = 100._dp           ! time to exit
24     dt = 0.001_dp           ! time step
25     lambda = 0._dp           !intiate a value for lyapanov exponent
26
27
28     !!Below are four cases - Make sure only one set of y is uncommented!!
29
30     !!Low Energy - No Chaos
31     y(1) = 1._dp*(pi/180)    ! initial angle 1
32     y(2) = 0._dp            ! initial momentum 1
33     y(3) = 1._dp*(pi/180)    ! initial angle 2
34     y(4) = 0._dp            ! initial momentum 2
35
36     !!Medium Energy - Emerging Chaos But Stable Orbit
37     ! y(1) = 1._dp*(pi/180)    ! initial angle 1
38     ! y(2) = 1._dp            ! initial momentum 1
39     ! y(3) = 1._dp*(pi/180)    ! initial angle 2
40     ! y(4) = 1.01_dp          ! initial momentum 2
41
42     !!High Energy - Full Chaos

```

```

43 !   y(1) = 1._dp*(pi/180)           ! initial angle 1
44 !   y(2) = 1._dp                     ! initial momentum 1
45 !   y(3) = 1._dp*(pi/180)           ! initial angle 2
46 !   y(4) = 1.3_dp                     ! initial momentum 2
47
48   !!High Energy - Full Chaos
49 !   y(1) = 60._dp*(pi/180)           ! initial angle 1
50 !   y(2) = 1._dp                     ! initial momentum 1
51 !   y(3) = 60._dp*(pi/180)           ! initial angle 2
52 !   y(4) = 2._dp                     ! initial momentum 2
53
54
55   !!open all the files that data will be written to!!
56
57   open(unit = 3, file = 'pend_angle1.data', &
58         action = 'write', status = 'replace')
59   open(unit = 4, file = 'pend_mom1.data', &
60         action = 'write', status = 'replace')
61   open(unit = 5, file = 'pend_angle2.data', &
62         action = 'write', status = 'replace')
63   open(unit = 6, file = 'pend_mom2.data', &
64         action = 'write', status = 'replace')
65   open(unit = 7, file = 'angle_mom1.data', &
66         action = 'write', status = 'replace')
67   open(unit = 8, file = 'angle_mom2.data', &
68         action = 'write', status = 'replace')
69   open(unit = 9, file = 'pend_xy1.data', &
70         action = 'write', status = 'replace')
71   open(unit = 10, file = 'pend_xy2.data', &
72         action = 'write', status = 'replace')
73   open(unit = 11, file = 'lambda.data', &
74         action = 'write', status = 'replace')
75
76   !!calculate the momenta and angle of the pendulums!!
77
78   do while ( t < tmax )
79       if ( t + dt > tmax) dt = tmax -t
80       call rkf45step(t,y,dt)
81       !use the runga kutta method to calculate for theta and momentum
82       lambda = lambda + log(abs(y(4)-y(2)))/t

```



```

83         !update the summation of the lyapanov exponent
84         write (unit = 11,fmt='(3f20.10)') lambda
85         !write lambda to file
86     end do
87
88 end program pendulum

```

The code is run by typing `./pend.` Various data sets are plotted to different files for easy graphing.

4 Results

A single pendulum is very ordered and predictable.

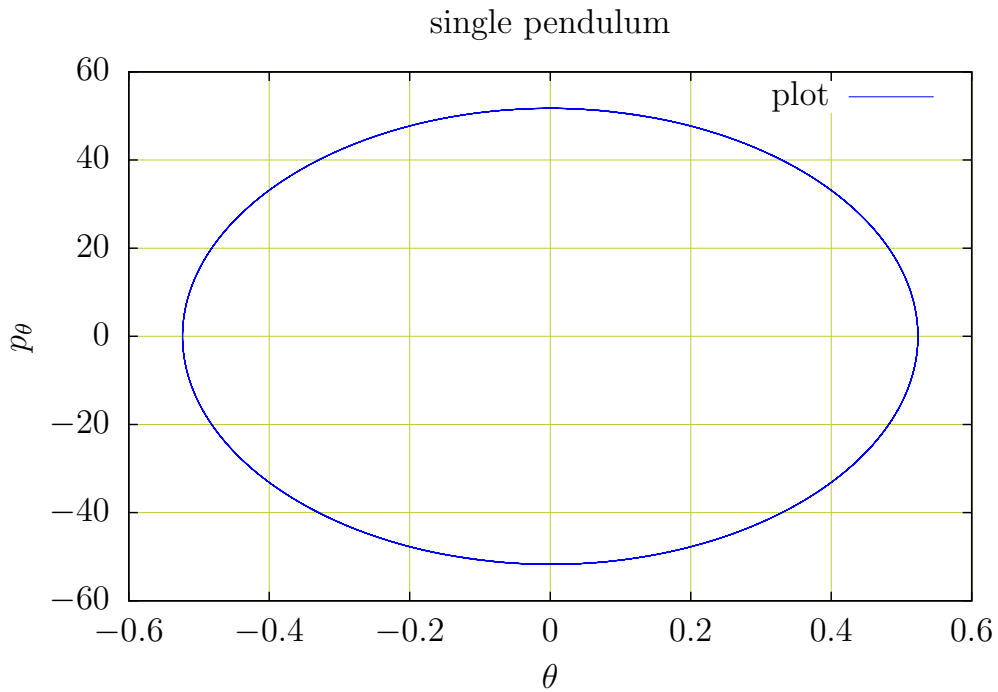


Figure 1: A simple pendulum.

Figure 1 shows that you can predict the values of the momentum for any time t . The momentum is maximum when the pendulum is at the bottom of

the arc and zero at the ends of the swing. Because of this, there is no chaos in the pendulums movement.

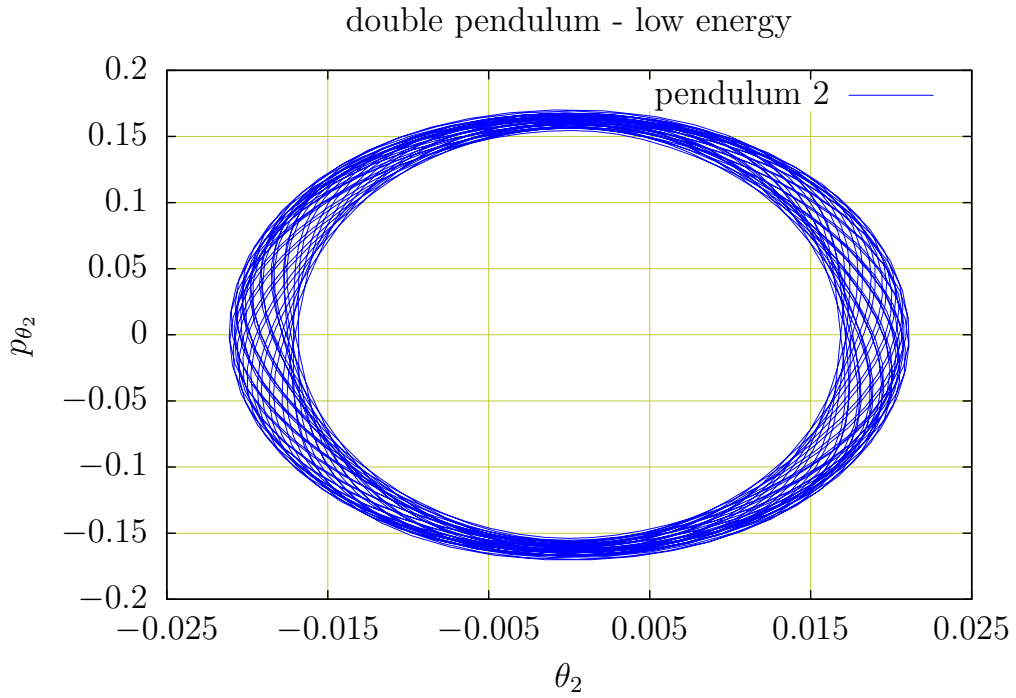


Figure 2: Plot for the second bob.

Figure 2 is similar to the single pendulum. The orbit has slight variation but it is a stable orbit. The calculated lyapanov exponent is a negative number, which confirms that the orbit is closed.

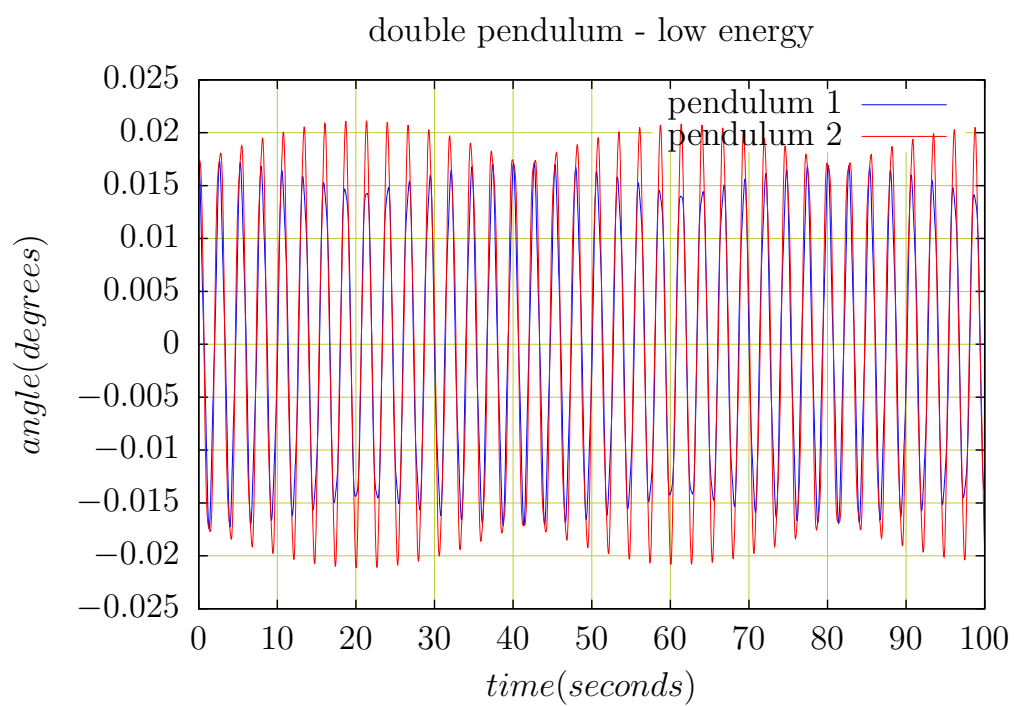


Figure 3: The angle of the two pendulums plotted over time

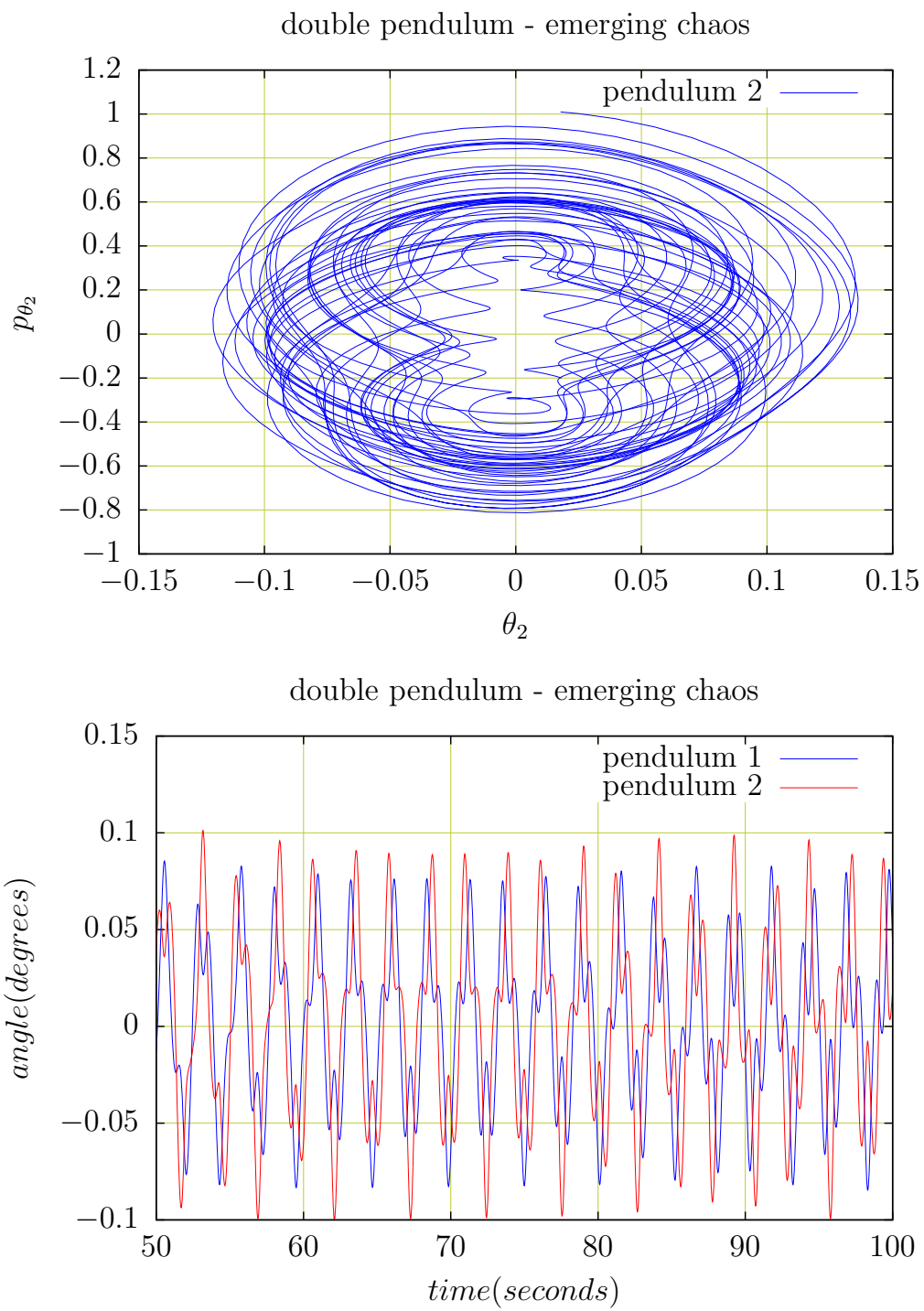


Figure 4: The second bob begins to show chaotic motion, but is in a stable orbit

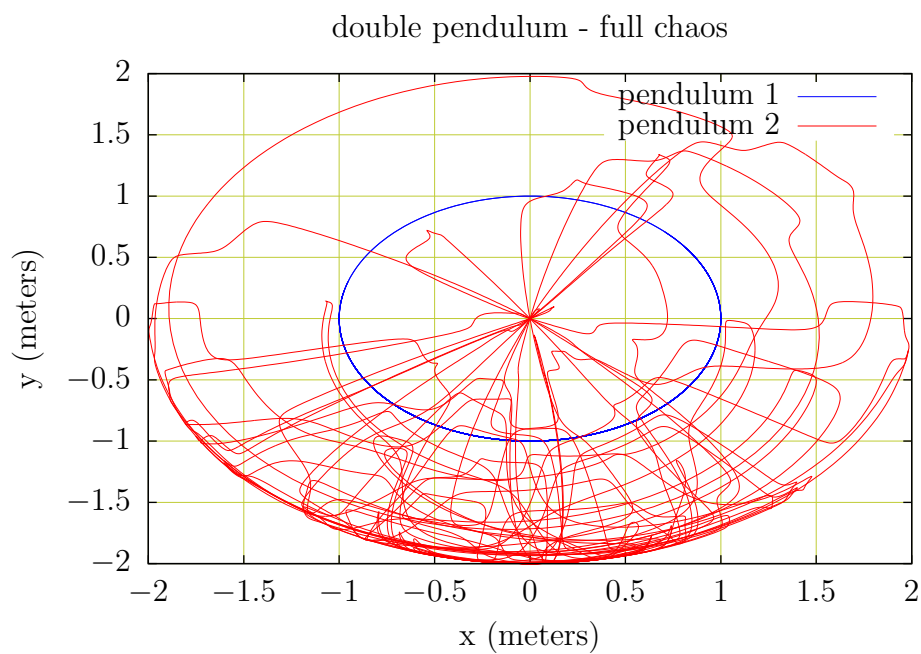
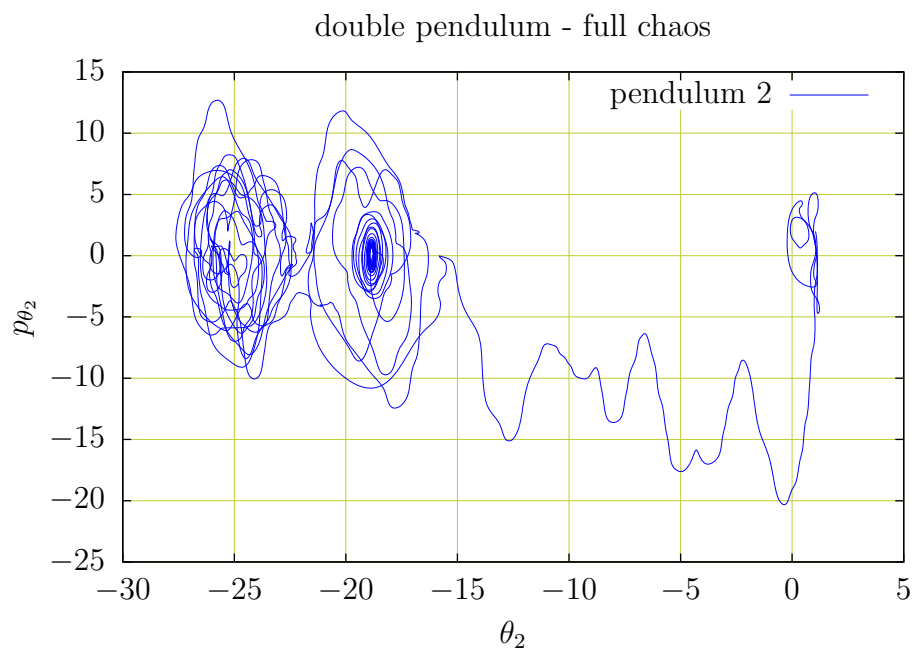


Figure 5: The second bob does not have a stable orbit. A graph of the x and y positions of the bobs are plotted so the chaotic motion can be visualized.

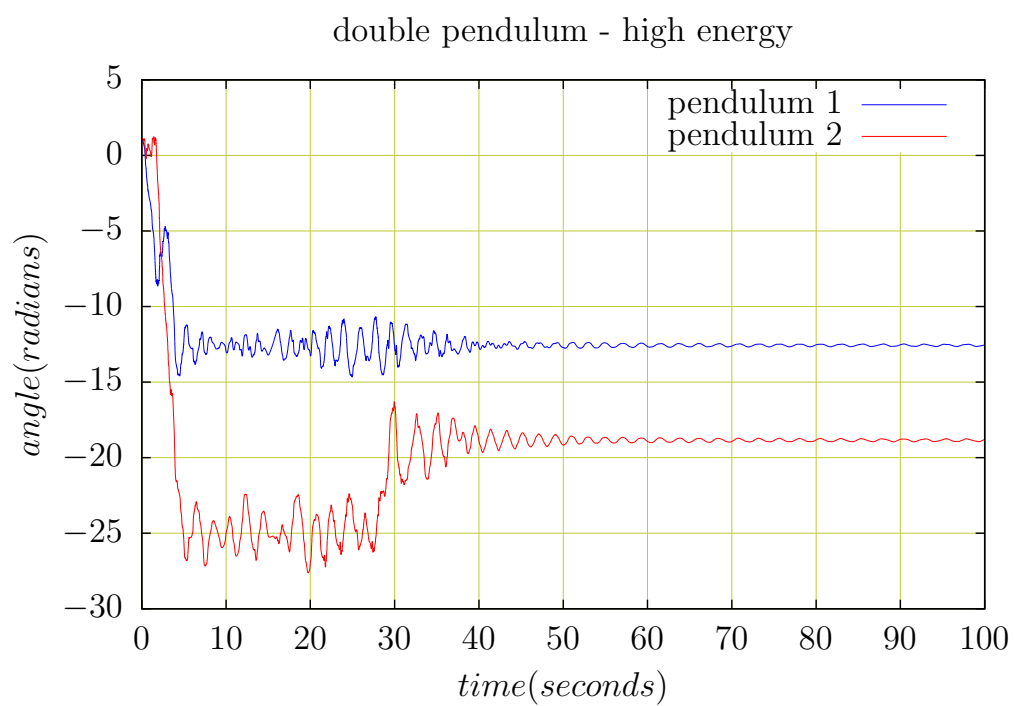


Figure 6: The chaotic behavior appears to die out after 40 seconds.

5 Summary and conclusions

A single pendulum is a non-chaotic, predictable system. A double pendulum is an inherently chaotic system, where the chaotic nature becomes more apparent as the difference between the initial conditions gets larger.

References

- [1] M. Metcalf, J. Reid and M. Cohen, *Fortran 95/2003 explained*. Oxford University Press, 2004.