The probability that the translational orbital n is occupied, irrespective of the state of internal motion of the molecule, is given by the ratio of the term in λ to the Gibbs sum 3:

$$f(\varepsilon_{\rm m}) = \frac{\lambda Z_{\rm int} \exp(-\varepsilon_{\rm m}/\tau)}{1 + \lambda Z_{\rm int} \exp(-\varepsilon_{\rm m}/\tau)} \simeq \lambda Z_{\rm int} \exp(-\varepsilon_{\rm m}/\tau). \tag{46}$$

The classical regime was defined earlier as $f(\varepsilon_n) \ll 1$. The result (46) is entirely analogous to (13) for the monatomic case, but $\lambda Z_{\rm int}$ now plays the role of λ . Several of the results derived for the monatomic ideal gas are different for the polyatomic ideal gas:

(a) Equation (17) for λ is replaced by

$$\lambda = n/(n_Q Z_{\rm int}) , \qquad (47)$$

with n_Q defined exactly as before. (We shall always use n_Q as defined for the monatomic ideal gas of atoms with zero spin.) Because $\lambda \equiv \exp(\mu/\tau)$ we have

$$\mu = \tau [\log(n/n_Q) - \log Z_{\rm int}]. \tag{48}$$

(b) The free energy is increased by, for N molecules,

$$F_{\rm int} = -N\tau \log Z_{\rm int}. \tag{49}$$

(c) The entropy is increased by

$$\sigma_{\rm int} = -(\partial F_{\rm int}/\partial \tau)_{V}. \tag{50}$$

The former result $U = \frac{3}{2}N\tau$ applies to the translational energy alone.

Example: Spin entropy in zero magnetic field. Consider an atom of spin I, where I may represent both electronic and nuclear spins. The internal partition function associated