

• **Problem 3.29** The canonical partition function for N non-interacting particles is

$$Z(T, V, N) = [Z_1(T, V)]^N.$$

From it, we immediately obtain

$$U_C(T, V, N) = - \left(\frac{\partial (\ln Z)}{\partial \beta} \right)_{V, N} = NkT^2 \frac{Z'_1}{Z_1}$$

$$F_C(T, V, N) = -kT \ln Z = -NkT \ln Z_1$$

$$S_C(T, V, N) = \frac{U_C - F_C}{T} = Nk \left(T \frac{Z'_1}{Z_1} + \ln Z_1 \right)$$

where $Z'_1 = (\partial Z_1 / \partial T)_V$.

The grand canonical partition function is

$$\begin{aligned} \mathcal{Z}(T, V, \mu) &= \sum_{N=0}^{\infty} e^{\beta \mu N} Z(T, V, N) = \sum_{N=0}^{\infty} [e^{\beta \mu} Z_1(T, V)]^N \\ &= \frac{1}{1 - e^{\beta \mu} Z_1(T, V)}. \end{aligned}$$

To obtain thermodynamic quantities as functions of N , we must adjust the chemical potential μ so that

$$N = \langle N \rangle_G = \left(\frac{\partial(\ln Z)}{\partial(\beta\mu)} \right)_{T,V} = \frac{e^{\beta\mu} Z_1}{1 - e^{\beta\mu} Z_1}$$

which implies that

$$e^{\beta\mu} Z_1 = \frac{N}{N+1} \quad \mu = -kT \ln \left(\frac{(N+1)Z_1}{N} \right).$$

We then find

$$\Omega_G(T, V, N) = kT \ln Z = kT \ln(N+1)$$

$$F_G(T, V, N) = \mu N - \Omega_G = kT [N \ln N - (N+1) \ln(N+1) - N \ln Z_1]$$

$$U_G(T, V, N) = - \left(\frac{\partial(\ln Z)}{\partial\beta} \right)_{\beta\mu, V} = NkT^2 \frac{Z'_1}{Z_1}$$

$$\begin{aligned} S_G(T, V, N) &= \frac{U_G - F_G}{T} \\ &= NkT \frac{Z'_1}{Z_1} - k[N \ln N - (N+1) \ln(N+1) - N \ln Z_1]. \end{aligned}$$

Clearly, we have $U_G(T, V, N) = U_C(T, V, N)$ and

$$\begin{aligned} (S_G - S_C)/k &= - \frac{(F_G - F_C)}{kT} \\ &= (N+1) \ln(N+1) - N \ln N \\ &= \ln N + 1 + \frac{1}{2N} + \dots \end{aligned}$$

On dividing this last result by N and taking N to be large, we get

$$\frac{s_G - s_C}{k} = - \frac{f_G - f_C}{kT} \simeq \frac{\ln N}{N}.$$

(c) *Grand canonical ensemble.* This ensemble describes a system in equilibrium with a reservoir, with which it can exchange both energy and particles, so both the energy and the number of particles fluctuate. The reservoir is characterized by a definite temperature $T = 1/k\beta$ and chemical potential μ or fugacity $z = e^{\beta\mu}$. The partition function is

$$Z(T, \mu) = \sum_{N=0}^{\infty} z^N Z(T, N) = \frac{1}{1 - z(e^{\beta\varepsilon} + e^{-\beta\varepsilon})}$$

and the thermodynamic interpretation is through the grand potential $\Omega_G(T, \mu) = kT \ln[\mathcal{Z}(T, \mu)]$. The entropy will be obtained through the thermodynamic relation $\Omega = TS - U + \mu N$, where the particle number and internal energy are identified as the mean values

$$N_G = \langle N \rangle_G = z \left(\frac{\partial(\ln \mathcal{Z})}{\partial z} \right)_\beta = \frac{z(e^{\beta\epsilon} + e^{-\beta\epsilon})}{1 - z(e^{\beta\epsilon} + e^{-\beta\epsilon})}$$

$$U_G = \langle E \rangle_G = - \left(\frac{\partial(\ln \mathcal{Z})}{\partial \beta} \right)_z = -\epsilon \frac{z(e^{\beta\epsilon} - e^{-\beta\epsilon})}{1 - z(e^{\beta\epsilon} + e^{-\beta\epsilon})} = -N_G \epsilon \tanh(\beta\epsilon).$$

As above, we would like to express thermodynamic functions in terms of the parameter $x = U_G/N_G\epsilon = -\tanh(\beta\epsilon)$ and N_G , rather than T and μ , so we first solve the above equations to obtain

$$\beta\epsilon = \frac{1}{2} [\ln(1-x) - \ln(1+x)]$$

$$\beta\mu = \ln \left(\frac{N_G}{N_G + 1} \right) - \ln 2 + \frac{1}{2} [\ln(1+x) + \ln(1-x)].$$

With these results in hand, we can calculate

$$U_G = N_G \epsilon x$$

$$\Omega_G = kT \ln \mathcal{Z} = kT \ln(N_G + 1)$$

$$S_G = \frac{1}{T} (\Omega_G + U_G - \mu N_G)$$

$$= k N_G \left[\frac{\ln(N_G + 1)}{N_G} + \ln \left(1 + \frac{1}{N_G} \right) \right. \\ \left. + \ln 2 - \frac{1}{2}(1+x) \ln(1+x) - \frac{1}{2}(1-x) \ln(1-x) \right].$$

On taking the limit that N_G is very large, we find for the entropy per particle that

$$s = \lim_{N_G \rightarrow \infty} \left(\frac{S_G}{N_G} \right) = k [\ln 2 - \frac{1}{2}(1+x) \ln(1+x) - \frac{1}{2}(1-x) \ln(1-x)]$$