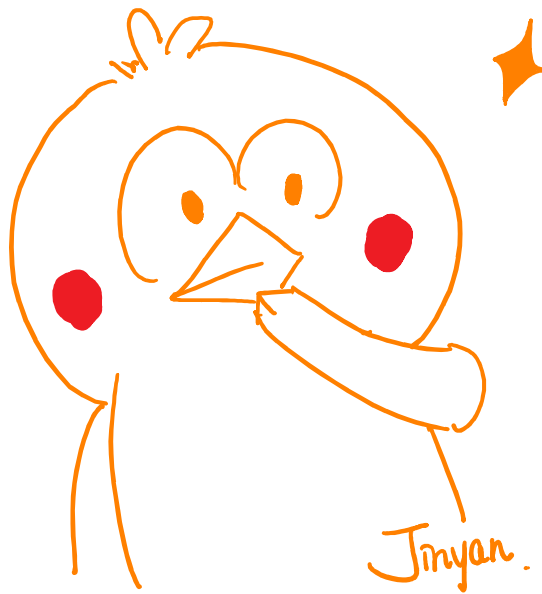


Homework 1

Physics 513 - Quantum Field Theory
Professor Ratindranath Akhoury



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Fall 2023

1

For pure rotation about the xy -axes by an angle θ , we can write

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos(\theta) & -\sin(\theta) & 0 \\ 0 & \sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

when θ is small or infinitesimal, denoted as $\delta\theta$, we can rewrite

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \left(\mathbb{I} + \delta\theta \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \right) \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

with small angle approximation $\sin(\delta\theta) = \delta\theta$, and $\cos(\delta\theta) = 1$. From which we obtain one of the rotation generator

$$L_3 = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}.$$

For pure rotation about the xz -axes by an angle θ , we can write

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with small angle approximation $\sin(\delta\theta) = \delta\theta$, and $\cos(\delta\theta) = 1$. From which we obtain one of the rotation generator

$$L_1 = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix}.$$

Now we consider a boost in x -direction, we can write

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} \gamma & -\gamma v & 0 & 0 \\ -\gamma v & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

For infinitesimal boost with speed denoted as δv , we have $\gamma = 1/\sqrt{1 - (\delta v)^2} \approx 1$. Thus we can rewrite

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \left(\mathbb{I} - \delta v \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \right) \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

and thus we obtain the one of the boost generators

$$K_1 = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}.$$

Now we consider a boost in y -direction, we can write

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} \gamma & 0 & -\gamma v & 0 \\ 0 & 1 & 0 & 0 \\ -\gamma v & 0 & \gamma & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

For infinitesimal boost with speed denoted as δv , we have $\gamma = 1/\sqrt{1 - (\delta v)^2} \approx 1$. Thus we can rewrite

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \left(\mathbb{I} - \delta v \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \right) \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

and thus we obtain the one of the boost generators

$$K_2 = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}.$$

Now we consider a boost in z -direction, we can write

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \begin{bmatrix} \gamma & 0 & 0 & -\gamma v \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\gamma v & 0 & 0 & \gamma \end{bmatrix} \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

For infinitesimal boost with speed denoted as δv , we have $\gamma = 1/\sqrt{1 - (\delta v)^2} \approx 1$. Thus we can rewrite

$$\begin{bmatrix} t' \\ x' \\ y' \\ z' \end{bmatrix} = \left(\mathbb{I} - \delta v \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \right) \begin{bmatrix} t \\ x \\ y \\ z \end{bmatrix},$$

and thus we obtain the one of the boost generators

$$K_3 = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}.$$

2

(a) Here we consider two events A and B , we can set up an inertial frame such that the spacial distance between A and B only has x -component. That is, let S denote such a frame, event A has coordinate $(x_0, x_1, 0, 0)$ in S , and event B has coordinate $(y_0, y_1, 0, 0)$ in S .

If events A and event B are spacelike related, then we can write

$$(y_0 - x_0)^2 - (y_1 - x_1)^2 < 0. \quad (2.1)$$

Here we will find an inertial frame S' in which the two event occur at the same time. Suppose S' is obtained by Loretnz boosting S in the x -direction with speed v . Then the coordinate of A in S' is given by

$$\begin{bmatrix} \gamma & -\gamma v & 0 & 0 \\ -\gamma v & \gamma & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} x_0 \\ x_1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \gamma x_0 - \gamma v x_1 \\ -\gamma v x_0 + \gamma x_1 \\ 0 \\ 0 \end{bmatrix}, \quad (2.2)$$

where $\gamma = 1/\sqrt{1-v^2}$. Similarly, the coordinate of B in S' is given by

$$\begin{bmatrix} \gamma & -\gamma v & 0 & 0 \\ -\gamma v & \gamma & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} y_0 \\ y_1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \gamma y_0 - \gamma v y_1 \\ -\gamma v y_0 + \gamma y_1 \\ 0 \\ 0 \end{bmatrix}. \quad (2.3)$$

If the two events have the same time coordinate in S' , then we require v to satisfy

$$\begin{aligned} \gamma x_0 - \gamma v x_1 &= \gamma y_0 - \gamma v y_1, \\ x_0 - v x_1 &= y_0 - v y_1, \\ v y_1 - v x_1 &= y_0 - x_0, \\ v(y_1 - x_1) &= (y_0 - x_0), \\ v &= (y_0 - x_0)/(y_1 - x_1). \end{aligned}$$

Note from (2.1) we have that

$$-1 < \frac{y_0 - x_0}{y_1 - x_1} < 1, \quad (2.4)$$

and thus $v = (y_0 - x_0)/(y_1 - x_1)$ is well-defined. If $v = (y_0 - x_0)/(y_1 - x_1)$ is satisfied, event A and B have the same time coordinate in frame S' .

If one is able to find a Lorentz frame S' such that the spatial coordinates of events A and B coincide, then, denoting the coordinates of A as (x'_0, x'_1, x'_2, x'_3) and that of B as (y'_0, y'_1, y'_2, y'_3) in frame S' , one can write

$$(y'_0 - x'_0)^2 - \sum_{i=1}^3 (y'_i - x'_i)^2 = (y'_0 - x'_0)^2 - 0 = (y'_0 - x'_0)^2 = (y_0 - x_0)^2 - (y_1 - x_1)^2 < 0,$$

but this is not possible as $(y'_0 - x'_0)^2$ is non-negative.

(b) Now suppose A and B are timelike related, then their coordinates in frame S satisfy the relation

$$(y_0 - x_0)^2 - (y_1 - x_1)^2 > 0,$$

rearranging we have

$$-1 < \frac{y_1 - x_1}{y_0 - x_0} < 1. \quad (2.5)$$

One immediately notice that S' obtained by a Lorentz boost in the x -direction with velocity v can never yield $x'_0 = x'_1$ as Eq. (2.4) gives $v > 1$. However, one can find a frame S' by a Lorentz boost in the x -direction with velocity v such that events A and B have the same spatial coordinate components. Utilizing (2.2) and (2.3), we require

$$\begin{aligned} -\gamma v x_0 + \gamma x_1 &= -\gamma v y_0 + \gamma y_1, \\ v &= \frac{y_1 - x_1}{y_0 - x_0}. \end{aligned}$$

Such a velocity is possible as guaranteed by (2.5). Thus one can obtain an inertial frame in which events A and B have the same spatial coordinates, in the case where A and B are timelike related. In this case, it is not possible to obtain a frame S' in which A and B have the same time coordinate, as otherwise we would have

$$(y'_0 - x'_0)^2 - \sum_{i=1}^3 (y'_i - x'_i)^2 = - \sum_{i=1}^3 (y'_i - x'_i)^2 = (y_0 - x_0)^2 - (y_1 - x_1)^2 > 0,$$

which is not possible as $-\sum_{i=1}^3 (y'_i - x'_i)^2$ is non-positive.

3

Let $F_{\mu\nu}$ be an antisymmetric tensor. Here we can write

$$\begin{aligned}(\partial_\nu F_\mu{}^\alpha)F^\nu{}_\alpha &= (\partial_\nu F_\mu{}^\alpha)F^{\nu\rho}g_{\rho\alpha} \\&= (\partial_\nu F_\mu{}^\alpha g_{\rho\alpha})F^{\nu\rho} \\&= (\partial_\nu F_{\mu\rho})F^{\nu\rho} \\&= -(\partial_\nu F_{\mu\rho})F^{\rho\nu} \\&= -(\partial_\beta F_{\mu\alpha})F^{\alpha\beta},\end{aligned}$$

where we have applied change of indices $\rho \rightarrow \alpha$ and $\nu \rightarrow \beta$ in the last step.

4

Consider a conserved current j^μ , and define the total charge as

$$Q = \int_M d^3x j^0(\vec{x}, t). \quad (4.1)$$

where M is the region of interest.

(a) Here we write

$$\frac{dQ}{dt} = \frac{d}{dt} \int_M d^3x j^0(\vec{x}, t) = \int_M d^3x \frac{d}{dt} j^0(\vec{x}, t) = - \int_M d^3x \nabla \cdot \vec{j},$$

where we have used the conservation law $\partial_\alpha j^\alpha(x) = 0$ in the last equality. Now applying Stokes' Theorem, with the assumption that M is large enough such that there is no net current flux on the surface ∂M , we can write

$$\frac{dQ}{dt} = - \int_M \nabla \cdot \vec{j} dV = - \int_{\partial M} \vec{j}(x) \cdot \hat{n} dS = 0.$$

Thus we see that Q is independent of time.

(b) To see that Q is a Lorentz scalar, we redefine Q as

$$Q = \int_{M'} d^4x j^a(x) \partial_a \theta(n_\beta x^\beta) \quad (4.2)$$

where M' is the entire spacetime, θ is the Heaviside step function, and n_λ (is not a tensor, but for notation here, $n_\beta x^\beta$ indicate summing same indices) is defined by $n_1 = n_2 = n_3 = 0$, with $n_0 = 1$. To see that Q (4.2) agrees with (4.1), we simply expand

$$Q = \int_{M'} d^4x j^a(x) \partial_a \theta(n_\beta x^\beta) = \int_{M'} d^4x j^a(x) \partial_a \theta(t) = \int_{M'} d^4x j^0(x) \delta t = \int_M d^3x j^0(x).$$

The effect of a Lorentz transformation on Q is then simply to change n ,

$$n'_\beta = \Lambda^\gamma_\beta n_\gamma$$

for any Lorentz transformation Λ^γ_β . Then we can write

$$\begin{aligned} Q' - Q &= \int_{M'} d^4x j^a(x) \left(\partial_a \theta(n'_\beta x^\beta) - \partial_a \theta(n_\beta x^\beta) \right) \\ &= \int_{M'} d^4x j^a(x) \partial_a \left(\theta(n'_\beta x^\beta) - \theta(n_\beta x^\beta) \right). \end{aligned} \quad (4.3)$$

From product rule, we see that, for reasonable function $k(x)$, we have

$$\int_{M'} d^4x \partial_a (j^a(x) k(x)) = \int_{M'} d^4x (k(x) \partial_a j^a(x) + j^a(x) \partial_a k(x)) = \int_{M'} d^4x j^a(x) \partial_a k(x),$$

with the help of conservation law $\partial_a j^a = 0$ everywhere. And thus (4.3) becomes

$$\begin{aligned} Q' - Q &= \int_{M'} d^4x \partial_a \left(j^a(x) \left(\theta(n'_\beta x^\beta) - \theta(n_\beta x^\beta) \right) \right) \\ &= \int_{\partial M'} d^3x j^a(x) \left(\theta(n'_\beta x^\beta) - \theta(n_\beta x^\beta) \right), \end{aligned} \quad (4.4)$$

where we have applied Stokes' Theorem. With the assumption that $j^a(x)$ vanishes for large enough $|x|$ at fixed t , and $\theta(n'_\beta x^\beta) - \theta(n_\beta x^\beta)$ vanishes for large enough $|t|$ with fixed x , we conclude that we have

$$Q - Q' = \int_{\partial M} d^3x j^a(x) \left(\theta(n'_\beta x^\beta) - \theta(n_\beta x^\beta) \right) = 0,$$

showing that Q is a Lorentz scalar.