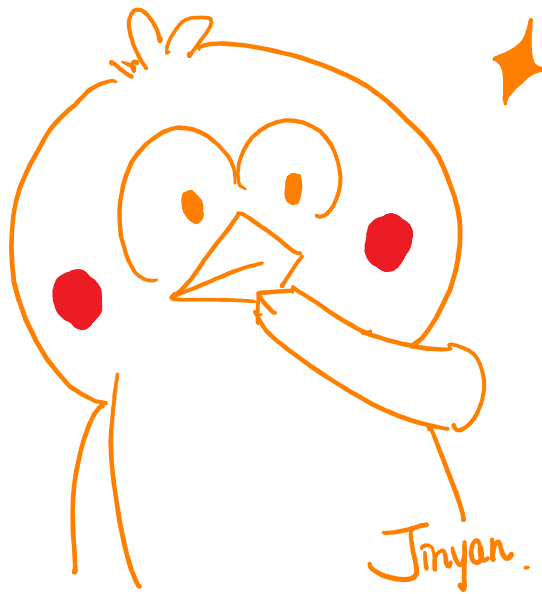


# Homework 9

Physics 542 - Quantum Optics  
Professor Alex Kuzmich



**Jinyan Miao**

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# 1

Here we consider

$$S(N) = |A(N, t)|^2$$

where we have

$$A(N, t) = \sum_{j=1}^N e^{-i\omega t + i\phi_j}$$

with  $\phi_j \in \mathbb{R}$  being random phase. Here we will compute, assume taking the time average,

$$\langle A \rangle = \left\langle e^{-i\omega t} \sum_{j=1}^N e^{i\phi_j} \right\rangle = 0 \cdot \sum_{j=1}^N e^{i\phi_j} = 0.$$

$$\langle S \rangle = \left\langle \sum_{j=1}^N e^{i\omega t - i\phi_j} \sum_{k=1}^N e^{-i\omega t + i\phi_k} \right\rangle = \left\langle \sum_{j,k=1}^N e^{i(\phi_k - \phi_j)} \right\rangle = \sum_{j,k=1}^N e^{i(\phi_k - \phi_j)} = \sum_{j \leq k}^N 2 \cos(\phi_j - \phi_k).$$

$$\langle S^2 \rangle = \left\langle \left( \sum_{j,k=1}^N e^{i(\phi_k - \phi_j)} \right) \left( \sum_{j',k'=1}^N e^{i(\phi_{k'} - \phi_{j'})} \right) \right\rangle = \left( \sum_{j \leq k}^N 2 \cos(\phi_j - \phi_k) \right) \left( \sum_{j' \leq k'}^N 2 \cos(\phi_{j'} - \phi_{k'}) \right).$$

$$(\Delta S)^2 = \langle S^2 \rangle - \langle S \rangle^2 = 0.$$

Now compute, assuming taking the ensemble average (over  $\phi_j$ ),

$$\langle A \rangle = \left\langle e^{-i\omega t} \sum_{j=1}^N e^{i\phi_j} \right\rangle = e^{-i\omega t} \cdot 0 = 0.$$

$$\langle S \rangle = \left\langle \sum_{j=1}^N e^{i\omega t - i\phi_j} \sum_{k=1}^N e^{-i\omega t + i\phi_k} \right\rangle = \left\langle \sum_{j,k=1}^N e^{i(\phi_k - \phi_j)} \right\rangle = N.$$

$$\langle S^2 \rangle = \left\langle \left( \sum_{j,k=1}^N e^{i(\phi_k - \phi_j)} \right) \left( \sum_{j',k'=1}^N e^{i(\phi_{k'} - \phi_{j'})} \right) \right\rangle = \left\langle \sum_{j,k,j',k'=1}^N e^{i(\phi_k - \phi_j + \phi_{k'} - \phi_{j'})} \right\rangle = 2N^2 - N.$$

$$(\Delta S)^2 = \langle S^2 \rangle - \langle S \rangle^2 = 2(N^2 - N).$$

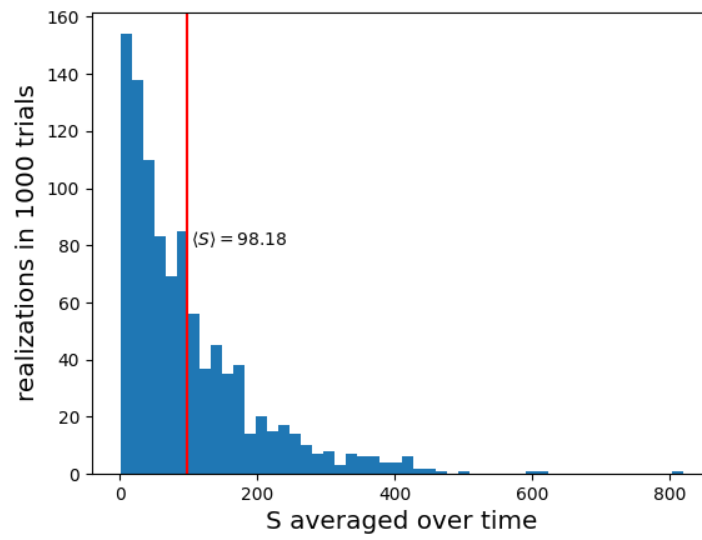
Quantities are computed with the help of Eq. (13.61) and (13.64) on Berman's text.

## 2

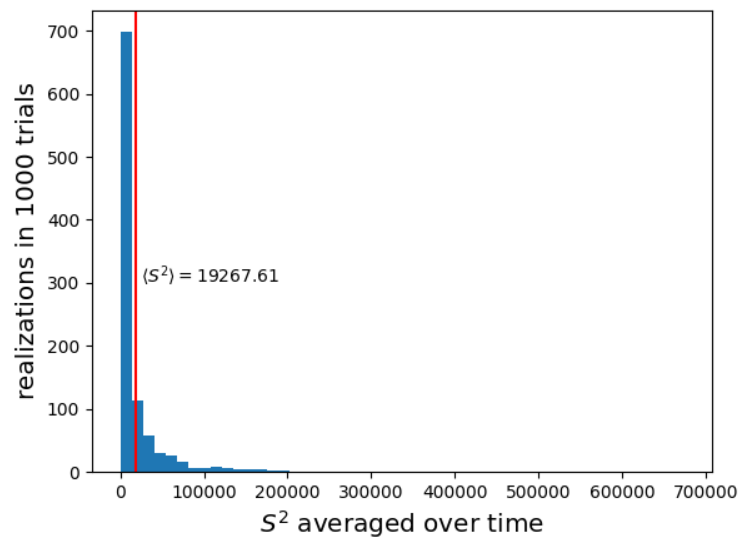
We evaluate

$$S = \left| \sum_{j=1}^{100} e^{-i\omega t + i\phi_j} \right|$$

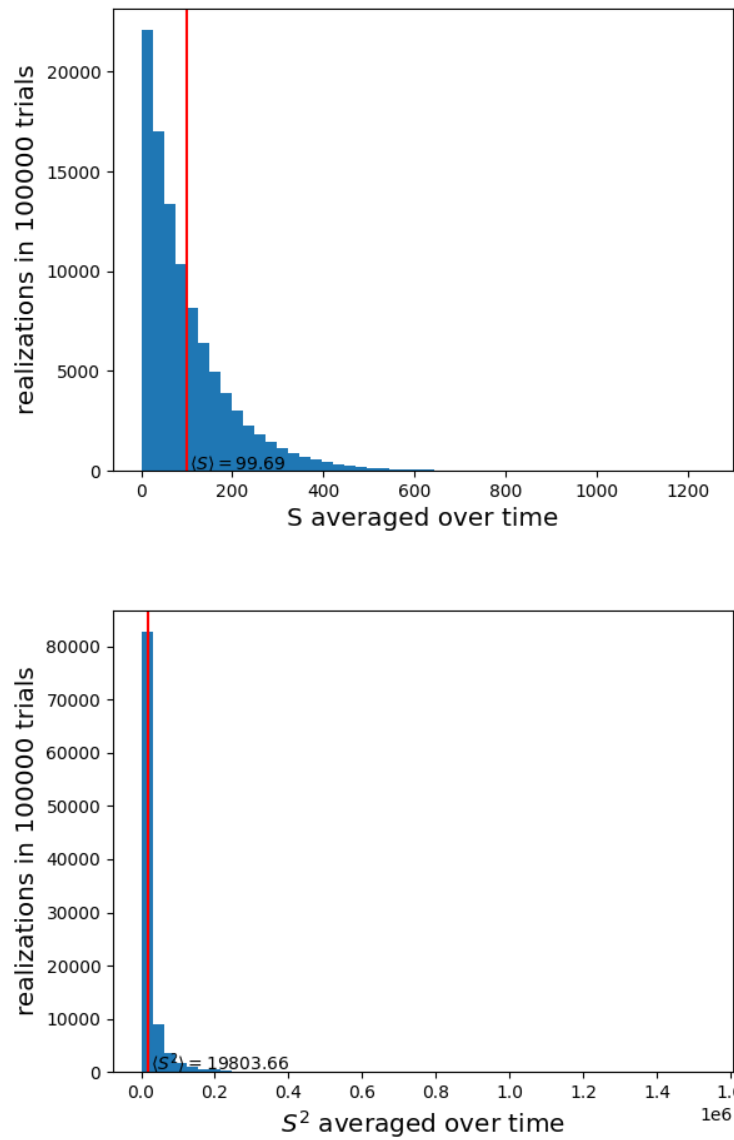
with random phases  $\phi_j$  picked via random number generator. The three  $S$ 's that we obtain are  $S = 3.59$ ,  $S = 114.65$ , and  $S = 94.02$ . Then we repeat the calculation of  $S$  for a 1000 times and take the average,  $\langle S \rangle = 98.18$  in this case, the result is shown on next page.



We do the same thing for  $S^2$ .



Now we evaluate them 100000 times, the numbers that we get here are very close to what we have calculated in problem 1.



### 3

First we consider

$$I(t) = \sum_{n=-\infty}^{\infty} \Theta(t + 0.1 - n) - \Theta(t - 0.1 - n)$$

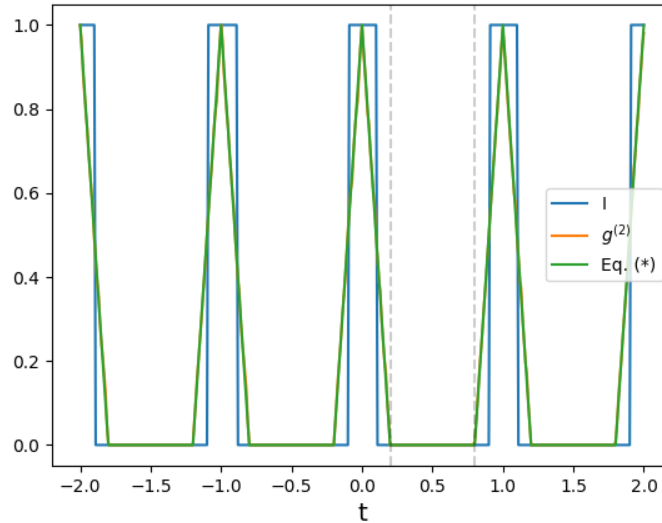
where  $\Theta$  is the Heaviside step function. It is not hard to see that

$$\bar{I} = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} I(t) dt = \lim_{T \rightarrow \infty} \frac{1}{T} (\lfloor T \rfloor \cdot 0.2 + \mu) = 0.2,$$

where  $|\mu| < 1$ . To compute  $\langle I(t) I(t + \tau) \rangle$ , as  $I$  is periodic with periods 1, it is not hard to see that  $\langle I(t + n) I(t) \rangle = 0.2$  for all  $n \in \mathbb{N}$  by symmetry. Furthermore, for  $0.2 \leq \tau \leq 0.8$ , we have either (1)  $I(t + \tau) = 0$  with  $I(t) = 1$ , or (2)  $I(t + \tau) = 1$  with  $I(t) = 0$ , thus their product must be  $\langle I(t) I(t + \tau) \rangle = 0$  when  $0.2 \leq \tau \leq 0.8$ . Lastly, for  $0 < \tau < 0.2$ , by the periodicity of  $I$ , it is not hard to see that the time average of the product  $\langle I(t) I(t + \tau) \rangle = 0$  decreases linearly from 0.2 to 0 as a function of  $\tau$ . Then using periodicity and symmetry, we conclude

$$g^{(2)}(\tau) = \begin{cases} 1 - 5 \cdot |\tau - \text{round}(\tau)| & |\tau - \text{round}(\tau)| \leq 0.2 \\ 0 & |\tau - \text{round}(\tau)| > 0.2 \end{cases}, \quad (*)$$

where  $\text{round}(\cdot)$  is the function that rounds to the nearest integer. We calculate  $g^{(2)}$  from the definition using  $I$ , and compare that with the result we obtain in Eq. (\*).



The gray dashed lines are  $t = 0.2$  and  $t = 0.8$ , respectively. Here  $g^{(2)}(\tau)$  does not approach 0 as  $\tau \rightarrow \infty$  as  $I$  is a periodic function of time, which implies  $I(t)$  and  $I(t + \tau)$  are for sure correlated.

Here we consider

$$E^+(t) = \frac{e^{-i\omega t}}{\pi^{1/4}} e^{-t^2/2}.$$

Thus we can compute

$$E^-(t) = (E^+)^* = \frac{e^{i\omega t}}{\pi^{1/4}} e^{-t^2/2},$$

then we have

$$\begin{aligned} g^{(1)}(\tau) &= \frac{\langle E^-(t) E^+(t+\tau) \rangle}{\langle E^-(t) E^+(t) \rangle} = \frac{\lim_{T \rightarrow \infty} \frac{1}{T\pi^{1/2}} \int_{-T/2}^{T/2} e^{-(t^2+(t+\tau)^2)/2} e^{-i\omega\tau} dt}{\lim_{T \rightarrow \infty} \frac{1}{T\pi^{1/2}} \int_{-T/2}^{T/2} e^{-t^2} dt} \\ &= \frac{e^{-i\omega\tau} \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{-(t^2+(t+\tau)^2)/2} dt}{\lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{-t^2} dt} \\ &= \frac{e^{-i\omega\tau}}{\pi^{1/2}} \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{-(t^2+(t+\tau)^2)/2} dt \\ &= e^{-i\omega\tau - \tau^2/4}. \end{aligned}$$

While on the other hand,

$$\begin{aligned} g^{(2)}(\tau) &= \frac{\langle E^-(t) E^-(t+\tau) E^+(t+\tau) E^+(t) \rangle}{\langle E^-(t) E^+(t) \rangle^2} \\ &= \frac{\langle e^{i\omega t - t^2/2} e^{i\omega(t+\tau) - (t+\tau)^2/2} e^{-i\omega t - t^2/2} e^{-i\omega(t+\tau) - (t+\tau)^2/2} \rangle}{\langle e^{-t^2} \rangle^2} \\ &= e^{-\tau^2} \frac{\lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} e^{-2t^2 - 2\tau t} dt}{\lim_{T \rightarrow \infty} \frac{\pi}{T^2}} \\ &= \frac{\pi^{1/2} e^{\tau^2/2} e^{-\tau^2}}{2^{1/2} \pi} \frac{\lim_{T \rightarrow \infty} \frac{1}{T}}{\lim_{T \rightarrow \infty} \frac{1}{T^2}} \\ &= \frac{e^{-\tau^2/2}}{(2\pi)^{1/2}} \lim_{T \rightarrow \infty} T \end{aligned}$$

which is an ill-defined quantity and obviously  $g^{(2)}(0) = \infty$ .



Diagram illustrating the geometry of two point sources  $R_1$  and  $R_2$  separated by distance  $d$ , with a point  $R_j$  at a distance  $D$  from the midpoint. The diagram shows the distances  $d/2$  and  $D$ , and the angle  $\theta \approx \frac{d}{2D}$ . The position vector  $R_j$  is given by  $R_j = D\hat{z} + x_j\hat{x}$ .

Equations for the distances  $R_j - R_a$ :

$$|R_j - R_a| = (D^2 + (x_j \pm d/2)^2)^{1/2}$$

$$= D \left(1 + \frac{(x_j \pm d/2)^2}{D^2}\right)^{1/2}$$

$$\sim D \left(1 + \frac{(x_j \pm d/2)^2}{2D^2}\right)$$

$$= D + \frac{(x_j \pm d/2)^2}{2D}$$

Equation for the distance  $d$  in terms of wavelength  $\lambda$  and angle  $\theta$ :

$$d \approx \frac{\lambda}{\Delta\theta} = \frac{\lambda D}{a}$$

Calculate  $g^{(2)}(R_1, t; R_2, t) = \frac{\langle E_1^- E_2^- E_2^+ E_1^+ \rangle}{\langle E_1^- E_1^+ \rangle \langle E_2^- E_2^+ \rangle}$

at  $R_a$

$$E_a^+(R_1, t) = \frac{1}{2} \sum_j E_0 e^{-i\omega t_j + i\phi_j(t-D/c)}$$

$$E_a^-(R_1, t) = \frac{1}{2} \sum_j E_0 e^{i\omega t_j - i\phi_j(t-D/c)}$$

With  $t_j = t - |R_j - R_a|/c$   
 $r_j = |R_j - R_a|/c$

for simplicity. assume  $c=1 \Rightarrow \omega \sim \frac{1}{\lambda}$

$\Rightarrow r_j \sim D + \frac{(x_j \pm d/2)^2}{2D}$

$R_2$  is -  
 $R_1$  is +

$$\begin{aligned}
\langle E_a^- E_a^+ \rangle &= \frac{E_0^2}{4} \left\langle \sum_j e^{i\omega t_j - i\phi_j(t-D/c)} \sum_j e^{-i\omega t_j + i\phi_j(t-D/c)} \right\rangle \\
&= \frac{E_0^2}{4} \left\langle \sum_{j,k} e^{i\omega(t-r_j-t+r_k) - i(\phi_j(t-D/c) - \phi_k(t-D/c))} \right\rangle \\
&= \frac{E_0^2}{4} \left\langle \sum_{j,k} e^{i\omega(r_k-r_j) - i(\phi_j(t-D/c) - \phi_k(t-D/c))} \right\rangle \\
&= \frac{E_0^2 N}{4}
\end{aligned}$$

$$\begin{aligned}
&\langle E_1^- E_2^- E_2^+ E_1^+ \rangle \\
&= \frac{E_0^4}{16} \left\langle \sum_j e^{-i\omega t_j + i\phi_j^1(t-D/c)} \sum_j e^{-i\omega t_j^2 + i\phi_j^2(t-D/c)} \right. \\
&\quad \left. \sum_k e^{i\omega t_k^1 - i\phi_k^1(t-D/c)} \sum_{k'} e^{i\omega t_k^2 - i\phi_k^2(t-D/c)} \right\rangle \\
&= \frac{E_0^4}{16} \left\langle \sum_{j,j',k,k'} \exp(-i\omega(t-r_j^1+t-r_{j'}^2-t+r_k^1-t+r_k^2) + i(\phi_j^1+\phi_{j'}^2-\phi_k^1-\phi_k^2)) \right\rangle \\
&= \frac{E_0^4}{16} \left\langle \sum_{j,j',k,k'} \exp(-i\omega(r_k^1+r_k^2-r_j^1-r_{j'}^2) + i(\phi_j^1+\phi_{j'}^2-\phi_k^1-\phi_k^2)) \right\rangle
\end{aligned}$$

count # of  $e^0$   
 $j=j'=k=k' \Rightarrow N$   
 $j=k \neq j'=k' \Rightarrow N(N-1)$   
 $j=k' \neq j'=k \Rightarrow N(N-1)$

$$\begin{aligned}
r_k^1+r_k^2-r_j^1-r_{j'}^2 &= \frac{(x_k+d/2)^2}{2D} + \frac{(x_{k'}-d/2)^2}{2D} - \frac{(x_j+d/2)^2}{2D} - \frac{(x_{j'}-d/2)^2}{2D} \\
&= (2D)^{-1} (x_k^2 - x_j^2 + x_{k'}^2 - x_{j'}^2 + d(x_k - x_{k'} - x_j + x_{j'}))
\end{aligned}$$