On the Instantaneous Phase and Frequency Estimation of a Non-stationary Signal. The JADE Algorithm

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Abstract—Many real-life signals are strongly non-stationary, like gravitational wave measurements, biomedical signals, or geophysical data. Many methods are now available to decompose them into mono-component signals which contain only one active frequency over time. The problem now is how to compute, in an accurate and stable way, the instantaneous frequency and phase of such mono-component signals. Many approaches have been developed so far, but they are unstable to noise and have difficulty capturing quick instantaneous changes in frequency. In this work, we present an alternative approach, called the JADE method, which is based on the Dynamic Time Warping algorithm. We test its robustness to noise and run comparisons with classical methods used for the instantaneous frequency and phase estimation tasks.

Index Terms—Instantaneous Frequency, Instantaneous Phase, Non-stationary Signal

I. INTRODUCTION

REAL life signals are non-stationary in general. They are generated by complex and nonlinear systems. In recent years many innovative algorithms have been developed that are able to decompose real-life signals into simple oscillatory components. We can think, for instance, of Empirical Mode Decomposition [1], Iterative Filtering [2], Ensemble empirical mode decomposition [3], sparse time-frequency representation [4], Geometric Mode Decomposition [5], Variational Mode Decomposition [6], Empirical wavelet transform [7], and Resampled Iterative Filtering [8], just to mention a few.

Once a signal has been decomposed into simple oscillatory and non-stationary components, there is a need to study their frequency content to characterize them. Many papers have been published in the past introducing the concept of instantaneous frequency and reviewing the methods developed for its computation [9]–[19].

From all these works, it emerges that the methods mainly used in the literature, like Hilbert transform (HT), Wigner–Ville distribution (WVD) and its variations, such as the generalized zero-crossing (GZC) can only capture up to interwave modulations in frequency, i.e. changes in frequency that take place from one period to the next one. However, the non-stationary frequency can vary even inside each period of

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the oscillatory components, this is what we call an intrawave modulated signal [15]. As an example of such a signal, we can consider, for instance, a biomedical signal like the electrocardiogram of a patient [20], or an astrophysical signal like a gravitational wave produced by the collision of two black holes in deep space [21], or a geophysical signal as the Earth's magnetic field measured through a satellite orbiting around the globe [22]. All these signals present rapid changes in their frequency content. To be properly analyzed they require an algorithm able to capture the instantaneous frequency contained in them. In this work, we propose an innovative approach that is based on the so-called Dynamic Time Warping algorithm, and we show its robustness against noise compared to traditional methods available in the literature.

The rest of this work is organized as follows: in Section II we review a few methods for the computation of the instantaneous frequency of a monocomponent signal and then we propose our approach. In Section III, results relative to a few artificial examples are presented, including the test of the robustness of the proposed method and previously published techniques to noise. Section IV concerns the application of the proposed method to a few real-life applications. In the last section, we derive conclusions and highlight future directions of research.

II. METHODS

A. Instantaneous Frequency Estimation Methods

Non-stationary signals are characterized by their instantaneous frequency (IF), which is defined as the local frequency of the signal at a specific time. We describe several methods to compute the IF of a monocomponent signal in this section. IF is commonly computed through the analytic signal produced by the HT. For an input signal x(t), this approach constructs the complex-valued analytic signal $Z(t) = x(t) + j\hat{x}(t)$, where

$$\hat{x}(t) = \frac{1}{\pi} P.V. \int_{-\infty}^{+\infty} \frac{x(\tau)}{t - \tau} d\tau \tag{1}$$

is the Hilbert Transform of x(t), and the phase $\phi(t)$ is the angle of Z(t) [1]. The IF of the signal, $\omega(t)$, is then obtained from:

$$\omega(t) = \frac{1}{2\pi} \frac{\mathrm{d}\phi(t)}{\mathrm{d}t} \tag{2}$$

Huang et al. [15] determined that several conditions, including those proposed by Bedrosian [23] and Nuttall [24], need to be satisfied to obtain an accurate IF estimate using the HT

1

analytic signal approach. To satisfy these conditions, Huang et al. [15] introduced the Normalized Hilbert Transform (NHT) approach, in which an empirical method separates amplitude modulation (AM) from frequency modulation (FM). In this method, a cubic spline curve $e_1(t)$ is fit to the envelope of the signal, and this curve is used to normalize the input x(t):

$$y_1(t) = \frac{x(t)}{e_1(t)},$$
 (3)

where $y_1(t)$ is the normalized signal. This normalization scheme is applied iteratively until all signal values are less than or equal to unity. After the nth iteration, the normalized signal

$$y_n(t) = \frac{y_{n-1}(t)}{e_n(t)},\tag{4}$$

is only frequency modulated, with the AM component removed. The IF of the normalized signal $y_n(t)$ is then computed using the HT analytic signal approach.

Another IF estimation method proposed by Huang et al. [15] is Direct Quadrature (DQ), which also uses the previously described normalization scheme. Here, assuming the original signal $x(t) = A(t)y_n(t)$, and $y_n(t) = cos(\phi(t))$, the quadrature of $y_n(t)$ is computed as $\sqrt{1-y_n(t)^2}$. The phase is obtained as:

$$\phi(t) = \arctan \frac{y_n(t)}{\sqrt{1 - y_n(t)^2}} \tag{5}$$

The IF is then computed from the derivative of Eq. 5. The advantage of the DQ approach is that it bypasses the HT integral and solely computes the phase from differentiation.

B. Dynamic Time Warping

Dynamic Time Warping (DTW) is a widely used algorithm to compare time series data by performing a temporal alignment. DTW is capable of determining the optimal alignment despite differing lengths and time shifts of input signals. Suppose there are two time series $\mathbf{X} = [x_1,...,x_n]$ and $\mathbf{Y} = [y_1,...,y_m]$ of lengths n and m respectively. The two series can be aligned to form a cost matrix $\mathbf{D} \in \mathbb{R}^{n \times m}$. The cost matrix \mathbf{D} is initialized under two constraints:

- $\mathbf{D}_{i,0} = \infty$ for $i \in [1,n)$ and $\mathbf{D}_{0,j} = \infty$ for $j \in [1,m)$
- $\mathbf{D}_{0,0} = 0$

The local distance measure $d(\cdot,\cdot)$ used in this study is the Euclidean distance. The matrix ${\bf D}$ is then populated according to the formula:

$$\mathbf{D}_{i,j} = d(\mathbf{X}_i, \mathbf{Y}_j) + \min \left\{ \begin{matrix} \mathbf{D}_{i-1,j-1} \\ \mathbf{D}_{i-1,j} \\ \mathbf{D}_{i,j-1} \end{matrix} \right\}$$
(6)

A warping path **P** is a set containing the indices of the aligned elements: $\mathbf{P} = [(i_1, j_1), ..., (i_L, j_L)]$. **P** is defined as a path through the cost matrix **D** that satisfies the following conditions:

- $\mathbf{P}_1 = (i_0, j_0) = (0, 0)$
- $\mathbf{P}_L = (i_L, j_L) = (n-1, m-1)$
- $i_{k-1} \le i_k \le (i_{k-1} + 1)$ for $k \in (1, L)$
- $j_{k-1} \le j_k \le (j_{k-1} + 1)$ for $k \in (1, L)$

The cost C_P of a given path $\mathbf{P} = [(i_1, j_1), ..., (i_L, j_L)]$ is defined as:

$$C_P = \sum_{k=1}^{L} \mathbf{D}_{i_k, j_k} \tag{7}$$

With the set of all possible paths **P** defined as $\mathbb{P}_{n,m}$, the optimal path Ω is:

$$\Omega = \operatorname*{arg\,min}_{\mathbf{P}} \left\{ C_P | \mathbf{P} \in \mathbb{P}_{n,m} \right\} \tag{8}$$

This path Ω can be determined by tracing back through the cost matrix from $\mathbf{D}_{n-1,m-1}$ to $\mathbf{D}_{0,0}$. The cost $C_{\Omega} = DTW(\mathbf{X},\mathbf{Y})$ of Ω is often used as a metric of similarity between the inputs \mathbf{X} and \mathbf{Y} .

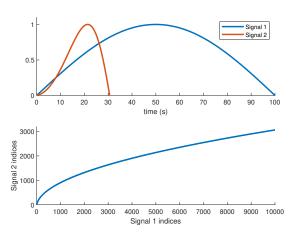


Fig. 1. Example of DTW alignment between two signals. The original signals (top) are aligned according to the warping path (bottom).

C. JADE algorithm

In this study, we use DTW as a tool to estimate the instantaneous phase and frequency of an input signal. We call this newly developed approach the JADE algorithm, from the initials of the authors' names and the word estimator. Consider a sampled signal $x(n) = cos(\phi(n))$, where T_s is the sampling period. Assume x(n) is time-warped using DTW to align with a discrete-time template signal $x_T(k) = cos(\omega_T kT_s)$, where ω_T is the known template frequency. x(n) can be expressed as $cos(\omega_T \psi(n))$ where $\psi(n) = \frac{\phi(n)}{\omega_T}$. After the DTW alignment, the optimal warping path $\Omega(k)$ approximates the following relation:

$$x(\mathbf{\Omega}(k)) = \cos(\omega_T \psi(\mathbf{\Omega}(k))) \approx \cos(\omega_T k T_s)$$
 (9)

Therefore, $\psi(\mathbf{\Omega}(k)) = T_s k$, and:

$$\mathbf{\Omega}^{-1}(k) = \frac{1}{T_s}\psi(k) = \frac{1}{\omega_T T_s}\phi(k) \tag{10}$$

By time-warping x(n) into a known template signal using DTW, the phase $\phi(n)$ can be recovered using the warping path $\Omega(k)$. This study leverages the relationship in (10) to estimate the phase of input signals. Then, by applying (2) we compute the corresponding instantaneous frequency.

The discrete-time signal model we use to represent a zeromean oscillatory signal x(n) is $x(n) = cos(\phi(n)) + w(n)$, where w(n) is Gaussian noise, and the function $\phi(n)$ is to be estimated. The noise can cause several problems with finding $\phi(n)$. First, we have found that when x(n) is nonmonotonic, the noise can have impact on $\phi(n)$ in particular around the inflection point. Accordingly, we may elect to break the signal into monotonic segments for further analysis. Alternatively, we can break the signal into segments at the zero crossings. We denote the k zero-crossings of x(n) as $\mathbf{z} = \{z_1, ..., z_k\}$. Because the noise w(n) can interfere with zero-crossing detection, a smoothed version of x(n), $x_s(n)$, is created by applying a moving average window with a fixed window length. A heuristic estimates the window length that attenuates 25% of the energy of the signal x(n). The zerocrossings z are then determined:

- $i \in \mathbf{z} \text{ if } x_s(i-1) < 0 < x_s(i), i \in n$
- $i \in \mathbf{z}$ if $x_s(i) < 0 < x_s(i-1), i \in n$

Next, x(n) is split into sections $\mathbf{x}_k = \{x_1(n), ..., x_k(n)\}$ at zero-crossing indices $\mathbf{z} = \{z_1, ..., z_k\}$. Each section $x_i(n) \in \mathbf{x}_k$ is defined by:

$$x_i(n) = \begin{cases} x(n+z_i) & \text{if } 0 \le n \le (z_{i+1} - z_i) \\ 0 & \text{if } n > (z_{i+1} - z_i) \end{cases}$$
 (11)

Each section $x_i(n) \in \mathbf{x}_k$ resembles a half-period sinusoid, which we denote the template sinusoid $x_{i_T}(n)$. The template sinusoid has amplitude A_{i_T} and frequency ω_{i_T} . The frequency of each template, ω_{i_T} , is an approximation of the frequency of the corresponding section:

$$\omega_{i_T} = \frac{\pi}{z_{i+1} - z_i} \tag{12}$$

Then, the template sinusoid $x_{i\tau}(n)$ can be expressed as:

$$x_{i_T}(n) = \begin{cases} \pm A_{i_T} sin(\omega_{i_T} n) & \text{if } 0 \le n \le (z_{i+1} - z_i) \\ 0 & \text{if } n > (z_{i+1} - z_i) \end{cases}$$
 (13)

where $x_{i_T}(n)$ matches the sign of $x_i(n)$. If we split the signal at local extrema to create monotonic segments, the template sinusoid will have the form $A_{i_T} sin(\omega_{i_T} n \pm \frac{\pi}{2})$.

For zero-mean artificial signals, it is usually sufficient to set the template amplitude A_{i_T} to the maximum or minimum value of the current segment. However, for low-SNR signals and multicomponent examples, we choose the A_{i_T} that corresponds to the lowest-cost DTW alignment between the template and the current segment:

$$A_{i_T} = \underset{A}{\operatorname{arg\,min}} \left\{ C_{\Omega} = DTW(x_{i_T}(n), x_i(n)) \right\}, \quad (14)$$

where C_{Ω} is the DTW alignment cost defined in (7), and $x_{i_T}(n) = A \cdot sin(\omega_{i_T} n)$.

Because $x_i(n)$ and $x_{i_T}(n)$ are similar, DTW can now find a low-cost alignment between these two signals, which can help create a smoother warping path for phase retrieval. Equation (10) requires the inverse of the warping path, which is determined by inverting the axes of the warping path of $DTW(x_{i_T}(n), x_i(n))$.

In order to further smoothen aberrations in the warping path due to noise, we can elect to fit a polynomial curve to the inverted warping path to obtain an expression for $\Omega_i^{-1}(k)$. The polynomial fit coefficients are denoted $\{\Omega_{i_0}^{-1},...,\Omega_{i_N}^{-1}\}$. The order of the polynomial N is left as a hyperparameter for the algorithm; however, for the examples discussed in this study, N was usually set to 3 or 4.

Assuming the original continuous-time phase $\phi_i(t)$ of the section $x_i(n)$ has the form $\phi_{i_0} + \phi_{i_1}t + \ldots + \phi_{i_N}t^N$, the sampled phase is $\phi_i(k) = \phi_{i_0} + \phi_{i_1}kT_s + \ldots + \phi_{i_N}k^NT_s^N$. Using (10),

$$\Omega_i^{-1}(k) = \frac{1}{\omega_{i_T} T_s} \phi_i(k) = \sum_{j=0}^N \frac{\phi_{i_j} k^j T_s^{j-1}}{\omega_{i_T}}$$
(15)

Since $\Omega_i^{-1}(k) = \sum_{j=0}^N \Omega_{i_j}^{-1} k^j$, the predicted phase of the section $x_i(n)$, $\hat{\phi}_i(k)$, is calculated by:

$$\hat{\phi}_i(k) = \sum_{j=0}^{N} \frac{\omega_{i_T}}{T_s^{j-1}} \mathbf{\Omega}_{i_j}^{-1} k^j$$
 (16)

If polynomial fitting is not desired, the phase estimate can be directly computed from the inverse of the DTW warping path using (5):

$$\hat{\phi}_i(k) = \omega_{i\pi} T_s \mathbf{\Omega}_i^{-1}(k) \tag{17}$$

Lastly, the estimates $\hat{\phi}_i(k)$ are concatenated to obtain the entire phase estimate $\hat{\phi}(k)$ for the input signal x(n). To obtain the IF estimate, we compute $\hat{\omega}_i(k) = \frac{1}{2\pi} \frac{\mathrm{d}\hat{\phi}_i(k)}{\mathrm{d}k}$ for each section, and concatenate the results to obtain $\hat{\omega}(k)$. Taking the derivative of each section helps avoid issues of discontinuity at boundary points. The overall approach is summarized in Algorithm 1.

Algorithm 1 JADE($x(n), T_s, N$).

1: $\mathbf{z} \leftarrow \text{Zero-crossings of } x(n)$

Inputs: Target signal x(n), sampling period T_s , polynomial fit order N (optional)

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Output: Total phase estimate \hat{\phi}(n), IF estimate \hat{\omega}(n)
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2: \mathbf{x}_k = \{x_1(n),...,x_k(n)\} \leftarrow x(n) split at crossings \mathbf{z} 3: \hat{\phi}(n) \leftarrow \{\}
4: k \leftarrow length of \mathbf{z}
5: for i=1,2,...,k do
6: x_{i_T}(n) \leftarrow template signal from Eqs.(12, 13, 14)
7: \Omega_i \leftarrow warping path of DTW(x_{i_T}(n),x_i(n))
8: if polynomial fitting desired then
9: Perform N-order polynomial fit to \Omega_i^{-1}
10: \{\Omega_{i_0}^{-1},...,\Omega_{i_N}^{-1}\} \leftarrow polynomial fit coefficients
11: \hat{\phi}_i(n) \leftarrow phase computed from Eq.(16)
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 $\hat{\phi}_i(n) \leftarrow \text{phase computed from Eq.}(17)$

14: **end if** 15: $\hat{\phi}(n) \leftarrow \left\{ \hat{\phi}(n), \hat{\phi}_i(n) \right\}$

16: **end for**

12:

13:

17: **return** $\hat{\phi}(n), \hat{\omega}(n)$

III. ANALYSIS OF ARTIFICIAL SIGNALS

To demonstrate our approach, we first consider its performance on artificial signals. A simple case is the linear chirp signal with quadratic phase and amplitude 1: $x(n) = cos(\alpha n^2 + \beta n) + \gamma \xi_n$, where ξ_n are Gaussian random numbers with zero mean and standard deviation 1, and γ is a scaling factor. In Fig. 2, phase estimation is shown on a quadratic phase signal with $\gamma = 0.05$, $\alpha = \frac{\sqrt{2}}{300}$ and $\beta = \frac{\sqrt{5}}{1000}$. The Signal to Noise Ratio, computed as SNR = $20 \cdot log(\|signal\|_2 / \|noise\|_2)$, of this signal is 23.18 dB. In Fig. 2, we show the result of the JADE phase estimate with and without polynomial fitting, in the bottom and middle panels respectively. In both cases, it can be seen that the phase estimated using JADE matches the ground truth phase function of the signal well.

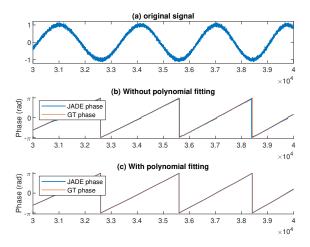


Fig. 2. Phase estimation of a quadratic-phase signal with added Gaussian noise (top). The middle panel shows the JADE-estimated phase (red) and the ground truth (blue) without polynomial fitting. Polynomial fitting can remove slight aberrations in the JADE phase estimate, as seen in the bottom panel.

Different values of SNR were tested to determine the algorithm's resilience to noise. We define the relative error ϵ to be:

$$\epsilon = \frac{\left\|\phi(n) - \hat{\phi}(n)\right\|_2}{\left\|\phi(n)\right\|_2},\tag{18}$$

where the ground truth phase $\phi(n) = \alpha n^2 + \beta n$, and $\hat{\phi}(n)$ is the JADE phase estimate. When testing performance over a range of values of SNR, ground truth zero crossings are used. This ensures that the performance of the JADE algorithm is evaluated independently of zero-crossing detection. The assumption of ground truth zero crossings is close to reality up to considerable levels of noise. Table 1 shows measured values of ϵ for different SNR.

The approach is next evaluated on a more complicated artificial signal with several components:

$$x(n) = [A(n) \cdot \cos(\phi(n))] + \gamma \xi_n, \tag{19}$$

where $A(n) = A_1 cos(\omega_1 n)$, $\phi(n) = n + A_2 cos(\omega_2 n)$, and $\gamma \xi_n$ are Gaussian random numbers with amplitude γ . This signal is amplitude and phase-modulated, as seen in the upper panel of Fig. 3. It is difficult to obtain an expression for the ground truth phase of x(n) from equation 19. In order to obtain the ground truth phase, we assume that multi-component signals can be

TABLE I RELATIVE ERROR IN PHASE ESTIMATE VERSUS NOISE LEVEL

SNR (dB)	ϵ
25.55	0.00015
13.62	0.00055
9.19	0.00089
4.11	0.0012
-1.45	0.0019
-6.36	0.0027
-10.86	0.0040

considered single-component if the amplitude and frequency components of the signal can be separated. Adapting definition 3.1 from Daubechies et al. [25], we assume the components of a function $f(n) = A(n) \cdot cos(\phi(n))$ can be separated if:

$$|A'(n)|, |\phi''(n)| \le \epsilon |\phi'(n)|, \forall n \in \mathbb{Z}$$
 (20)

up to accuracy ϵ . For x(n) in (19), choosing $\epsilon=0.01$, $A_1=0.3$, $\omega_1=1/35$, $A_2=1$, $\omega_2=1/100$, and $\gamma=0.05$, the conditions in (20) are satisfied. Therefore for this signal, we assume that the ground truth phase is $\phi(n)=n+\cos(n/100)$. In the bottom panel of Fig. 3 the JADE phase estimate is shown in yellow, along with the ground truth phase $\phi(n)$ in red, wrapped to the range $[-\pi,\pi]$. We also compare the result to the phase estimate based on the Hilbert Transform (HT). In Fig. 3, the HT phase output is plotted in blue. It can be seen that as the amplitude of the signal decreases, the noise amplitude dominates. In this section, the HT phase estimate becomes erratic, while the JADE phase estimate accurately represents the ground truth phase. The relative error ϵ as defined in (18) between the JADE estimate and the ground truth phase is 0.055.

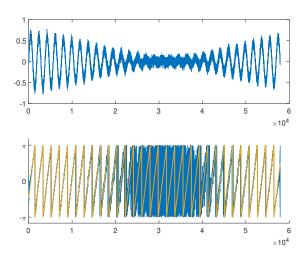


Fig. 3. Original AM+FM signal from (19) shown in the top panel, with SNR of 28.17 dB. In the bottom panel, the HT phase is shown in blue. The JADE phase estimate (yellow) aligns with the ground truth phase (red).

In Fig. 4, we also show the results of the NHT and DQ methods in estimating the instantaneous phase of the signal in (19) at 28.17 dB SNR. The removal of the AM component causes a visible improvement in the NHT phase compared

to the HT phase in Fig. 3. However, both the DQ and NHT phase estimates are erratic in noise. Because the normalization scheme common to both methods involves spline-fitting, the addition of noise can impact this method, and the form of the normalized signal is significantly altered. This can contribute to the inaccuracy in phase estimation seen in Fig. 4.

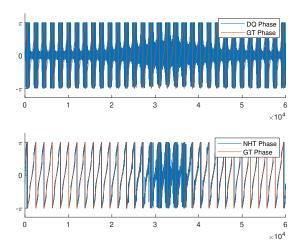


Fig. 4. Outputs of DQ (top) and NHT (bottom) phase estimation algorithms for signal in (19) at 28.17 dB SNR. The ground truth phase is plotted in red in both panels.

We also consider the phase estimation of the solution to the Duffing oscillator, a non-linear second order differential equation used to model driven oscillators. The generalized Duffing equation has the form:

$$\frac{\partial^2 x}{\partial t^2} + x + \varepsilon x^3 = \gamma \cos(\omega t) \tag{21}$$

where ε is a small parameter, γ is the coefficient modeling the driving force to the oscillations, which has frequency ω . The frequency response of the forced Duffing oscillator varies with the choice of nonlinearity. For the cubic nonlinearity in Eq. 21, the homotopy analysis method can be used to obtain the frequency response equation:

$$\left[\left(\omega^2 - 1 - \frac{3}{4} \varepsilon z^2 \right)^2 + \omega^2 \right] z^2 = \gamma^2 \tag{22}$$

We consider the case where $\varepsilon=1,\ \gamma=0.1,$ and $\omega=1.$ The particular solution is obtained using the Euler method, a numerical procedure used to solve ordinary differential equations (ODEs). It can be seen from Eq. 22 that the solution of the Duffing equation is multicomponent and non-stationary. We therefore decompose the Duffing equation solution into intrinsic mode functions (IMFs) using the Fast Iterative Filtering (FIF) algorithm. The first IMF is shown in the top panels of Fig. 5 and 6. In these figures, we show a comparison of the IF estimation outputs from STFT, CWT, SST, HT, and JADE methods for the first IMF of the Duffing equation solution.

IV. ANALYSIS OF REAL WORLD EXAMPLES

To demonstrate the abilities of the JADE algorithm, we apply it to analyze electrocardiogram (ECG) signals, gravi-

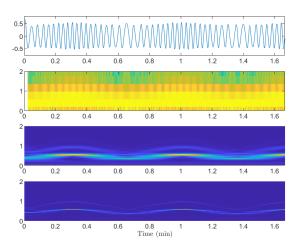


Fig. 5. First IMF of Duffing equation solution (top panel), and corresponding IF estimates from STFT (second), CWT (third), and SST (bottom) in Hertz.

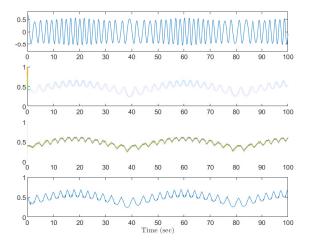


Fig. 6. First IMF of Duffing equation solution (top panel), and corresponding IF estimates from HT (second), IMFogram (third), and JADE (bottom) in Hertz.

tational wave (GW) signals, and geomagnetic response waveforms.

A. Gravitational Wave Analysis

In this example, we consider the phase estimation of a simulated inspiral gravitational wave (GW) signal. Inspiral GW signals are generated during the end-of-life merging of neutron star or black hole binary systems. There exist many approximate analytical models to simulate inspiral GW waveforms; we use SEOBNRv4 [26], an effective-one-body (EOB) model employed by the LIGO-VIRGO collaboration. Using the PyCBC software package [27], we simulate an inspiral GW signal using SEOBNRv4, with a ringdown time step of 1/8192 and a lower frequency of 40Hz. This simulated waveform is shown in the upper panel of Fig. 7. The ground truth phase and IF are obtained directly from PyCBC. The middle panel shows that the JADE phase estimate matches

the ground truth phase, and the bottom panel shows the IF estimate matches the ground truth as well.

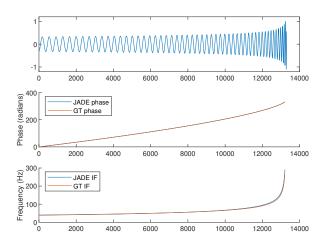


Fig. 7. Simulated inspiral gravitational waveform is shown in the top panel. In the middle panel, the JADE phase estimate (blue) aligns with the ground truth phase from PyCBC (red). Similarly in the bottom panel, the JADE IF (blue) aligns with the ground truth frequency from PyCBC (red).

Because the ground truth phase and IF are available from the PyCBC module, we test the JADE estimate performance in different noise conditions. We add Gaussian noise to the original waveform to create a range of SNR signals and measure the relative error between the JADE phase estimate and ground truth phase, defined in (18), for each case. The results are summarized in Table 2.

TABLE II
RELATIVE ERROR IN PHASE ESTIMATE VERSUS NOISE LEVEL FOR
GRAVITATIONAL WAVEFORM

SNR (dB)	ϵ
6.02	0.00072
3.63	0.00055
1.59	0.0011
-1.31	0.0014
-3.41	0.0015
-6.58	0.0018

The -6.58 dB case is shown in Fig. 8. We zoom into the latter portion of the signal to show the effect of the noise on the gravitational waveform.

B. Electrocardiogram Analysis

We next consider the phase estimation and reconstruction of electrocardiogram (ECG) signals from the MIT-BIH Arrhythmia database publicly available at *PhysioNet* [28]. This database consists of 48 two-channel ECG recordings of approximately 30 minutes in length, from 47 different patients. The recordings were sampled at 360Hz. From a particular recording (serial number 105), we isolate a singular ECG pulse of a duration of 0.7 seconds. We apply a moving average filter in order to remove the high-frequency component from the

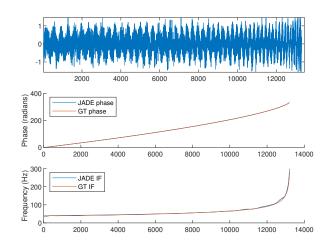


Fig. 8. Simulated inspiral gravitational waveform with added Gaussian noise (-6.58 dB SNR) is shown in the top panel. The phase and IF estimates compared with ground truths are shown in the middle and bottom panels respectively.

signal. The resulting signal is shown in the top panel of Fig. 9 in red.

Because the ground truth phase of the ECG signal is unknown, the performance of the JADE phase estimate cannot be directly evaluated. Instead, we perform a reconstruction of the original signal using the JADE phase estimate. We aggregate the sinusoidal template amplitudes into an amplitude function $\hat{A}(n)$, and also construct a function $\hat{\mu}(n)$ that contains the mean values of each section which were removed to perform the phase estimation. Then, to reconstruct the signal from the JADE phase estimate $\hat{\phi}(n)$, we evaluate:

$$\widehat{ECG} = [\hat{A}(n) \cdot \cos(\hat{\phi}(n))] + \hat{\mu}(n)$$
 (23)

The reconstruction \widehat{ECG} is shown in the top panel of Fig. 9 in blue. It can be seen that the reconstruction based on the JADE phase estimate accurately represents the original ECG signal. In the bottom panel, we show the JADE IF estimate along with the HT IF for comparison. However, our approach only provides a singular frequency estimate per time index, while the ECG is clearly a multi-component signal.

C. Magnetic Field Response Analysis

In this example, we consider phase estimation of magnetospheric and geomagnetic response to solar wind. Collision of the solar wind with the magnetic field causes a magnetospheric response characterized by sudden enhancement of the magnetic field intensity [29]. We analyze one particular event detected by GEOS (Geostationary Operational Environmental Satellites) spacecraft; this signal is shown in the upper panel of Fig. 10. First, the signal is decomposed into IMFs using the Fast Iterative Filtering (FIF) method. We estimate the phase of the second and third IMFs using JADE and compute the signal reconstruction according to (23) for both IMFs. The results are shown in the middle and bottom panels of Fig. 10.

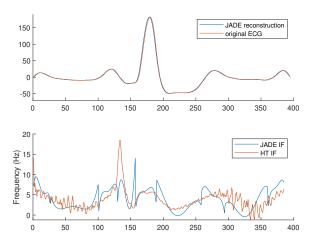


Fig. 9. Top panel shows ECG pulse from MIT-BIH Arrhythmia database (red) and reconstruction of the signal from JADE phase estimate (blue) according to Eq. 23. The bottom panel shows the JADE IF estimate (blue) and HT IF (red).

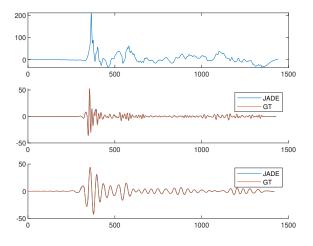


Fig. 10. Top panel shows the original geomagnetic response waveform. The middle and bottom panels show the second and third IMFs respectively (red), and the corresponding reconstructions using JADE phase estimates (blue).

V. CONCLUSION

Many real-life applications require the analysis of non-stationary signals whose frequencies vary rapidly over time. In recent years many innovative and nonlinear approaches have been proposed for the decomposition of such signals into mono-component signals. Once the signals have been decomposed into mono-component ones, we need to study their phase and frequency content over time with, possibly, high accuracy. In this work, we propose a new approach, called JADE, based on the Dynamic Time Warping method, for the estimation of the instantaneous phase and frequency of a mono-component signal. In contrast with traditional phase estimation methods such as HT, CWT, and STFT, JADE is basis-independent and is resilient in low SNR. This benefit is proven to be useful in the analysis of intrinsic mode functions of real-life signals, where decomposition algorithms such as

EMD and FIF suffer from challenges related to mode-mixing. Because the JADE method is able to ignore low-amplitude white and colored noise, it is able to perform favorably in the IF and phase estimation of IMFs. In this paper, we test this method on both synthetic and real-life signals comparing results with other algorithms proposed so far in the literature. The results show clearly that JADE outperforms any other method developed so far in the literature and proves to be extremely stable, even in the presence of heavy noise.

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