

# Worksheet 9 - Solutions

## 1. Fermion Gas Mixture

As the gases are non interacting, let us consider a single gas

$$Z = Z_{tr} \times Z_s$$

↓  
full Partition function      ↓  
transport                      spin.

$$Z_s = (e^{-\beta \gamma B} + e^{\beta \gamma B})^N$$

$$= 2^N \cosh^N(\beta \gamma B)$$

And

$$\mu = -\frac{k_B T}{N} \ln \left[ \frac{1}{N!} \left( \frac{V \cosh(\beta \gamma B)}{\lambda_T} \right)^N \right]$$

↓ Stirling

$$\approx -\frac{k_B T}{N} \left[ -N \ln(N) + N + N \ln \left( \frac{V \cosh(\beta \gamma B)}{\lambda_T} \right) \right]$$

$$\approx -k_B T \ln \left[ \frac{V e \cosh(\beta \gamma B)}{N \lambda_T} \right]$$

So, equating the chemical potentials for red and green

$$\mu_{red} = \mu_{green}$$

$$-\cancel{k_B T} \ln \left[ \frac{N_c \cosh(\beta \gamma_r B)}{N_r \lambda_T} \right] = -\cancel{k_B T} \ln \left[ \frac{N_c \cosh(\beta \gamma_g B)}{N_g \lambda_T} \right]$$

$$\frac{\cosh(\beta \gamma_r B)}{N_r} = \frac{\cosh(\beta \gamma_g B)}{N_g}$$

$$\frac{N_r}{N_g} = \frac{\cosh(\beta \gamma_r B)}{\cosh(\beta \gamma_g B)}$$


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2 Molecular Oxygen.

$$1) Z = \sum_{\mu_s} e^{-\beta \epsilon(\mu_s)} = \sum_{\mu_s} e^{-\beta (\mu_i^2/2m - \mu B S_i^z)}$$

$$= \frac{1}{n!} \left( \frac{V}{\lambda^3} (e^{\beta \mu B} + 1 + e^{-\beta \mu B}) \right)^n$$

where

$$\lambda = \frac{\hbar}{\sqrt{2\pi m k_B T}}$$

2) the probabilities  $g^0$

$$P(S_z = -1) = \frac{e^{-\beta \mu B}}{2 \cosh(\beta \mu B) + 1}$$

$$P(S_z = 0) = \frac{1}{2 \cosh(\beta \mu B) + 1}$$

$$P(S=1) = \frac{e^{\beta \mu B}}{2 \cosh(\beta \mu B) + 1}$$

3) The average magnetic moment;

$$\begin{aligned} \langle M \rangle &= \frac{1}{\beta} \frac{\partial \ln Z}{\partial B} = \frac{1}{\beta} \frac{\partial (\ln (2 \cosh \beta \mu B + 1))}{\partial B} \\ &= \frac{1}{\beta} \frac{1}{2 \cosh \beta \mu B + 1} \times 2 \sinh(\beta \mu B) \times \mu \beta \\ &= \mu N \frac{2 \sinh(\beta \mu B)}{2 \cosh(\beta \mu B) + 1} \end{aligned}$$

4) And finally, the susceptibility  $\chi = \left. \frac{\partial \langle M \rangle}{\partial B} \right|_{B=0}$

$$\begin{aligned} \chi &= \mu N \left[ \frac{2 \cosh(\beta \mu B) \beta \mu}{2 \cosh(\beta \mu B) + 1} - \frac{2 \sinh(\beta \mu B) \times (2 \sinh(\beta \mu B) \beta \mu)}{(2 \cosh(\beta \mu B) + 1)^2} \right] \\ &= N \beta \mu^2 \left[ \frac{2 \cosh(\beta \mu B) (2 \cosh(\beta \mu B) + 1) - [2 \sinh(\beta \mu B)]^2}{(2 \cosh(\beta \mu B) + 1)^2} \right] \Big|_{B=0} \end{aligned}$$

$$= N \beta \mu^2 \left[ \frac{2 \times (2+1)}{(2+1)^2} \right] = \frac{2}{3} N \beta \mu^2$$

### 3. Polar Rods

1) The rotational degrees of freedom goes as

$$Z_{\text{rot}} = \frac{1}{h^2} \int_0^\pi d\theta \int_0^{2\pi} d\phi \int_{-\infty}^{\infty} dp_\theta dp_\phi \exp \left[ -\frac{\beta}{2I} \left( p_\theta^2 + \frac{p_\phi^2}{\sin^2 \theta} + \beta \mu \cos \theta \right) \right]$$

the momenta are Gaussian integrals, and taking  $x = \cos \theta$

$$\begin{aligned} Z_{\text{rot}} &= \left( \frac{2\pi I}{\beta h^2} \right) 2\pi \int_1^1 dx e^{-\frac{\beta \mu E x}{2}} \\ &= \left( \frac{8\pi^2 I}{\beta \mu h^2} \right) \frac{\sinh(\beta \mu E)}{\beta \mu E} \end{aligned}$$

2) the muon polarization.

$$P = \langle \mu \cos \theta \rangle = \frac{\partial \ln(Z_{\text{rot}})}{\partial (\beta E)} = \mu \left[ \coth(\beta \mu E) - \frac{1}{\beta \mu E} \right]$$

3) the zero field polarizability

$$\chi_T = \left. \frac{\partial P}{\partial E} \right|_{E=0}$$

$$\chi_T = \frac{\partial P}{\partial E} = \beta \mu^2 \left[ -\frac{1}{\sinh^2(\beta \mu E)} + \frac{1}{(\beta \mu E)^2} \right]$$

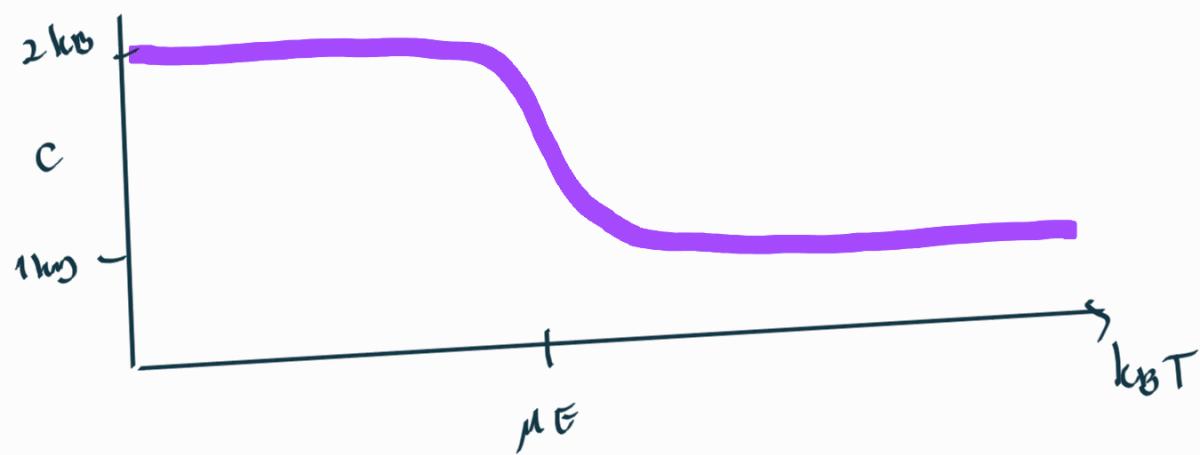
$$\left[ \frac{1}{x^2} - \frac{1}{\sinh^2(x)} \right] \Big|_{x=0} = \frac{1}{3}$$

$$\chi_T = \frac{1}{3} \mu^3 \beta$$

4) The energy stored:

$$\langle E_{\text{rot}} \rangle = - \frac{\partial \ln Z_{\text{rot}}}{\partial \beta} = 2k_B T - \mu E \text{ with } (\rho \mu E)$$

5)



4 Rotation of diatomic molecules.

a) Partition function:

$$Z = \sum_{j=0}^{\infty} (2j+1) e^{-j(j+1)\epsilon_0/\tau}$$

$$J(j+1)\epsilon_0 = \epsilon(j)$$

$(2j+1)$  = multiplicity.

$$= \sum_{j=0}^{\infty} \frac{d}{dj} \left[ e^{-(j^2+j)\epsilon_0/\tau} \right] \left( -\frac{\tau}{\epsilon_0} \right)$$

$$= -\frac{\tau}{\epsilon_0} \sum_{j=0}^{\infty} \frac{d}{dj} \left( e^{-\epsilon_0/\tau} \right)^{j^2+j}$$

2) for  $\tau \gg \frac{\epsilon_0}{\tau}$

$$Z_R(\tau) = -\frac{\tau}{\epsilon_0} \int_0^{\infty} \frac{dx}{dx} \left( e^{-\epsilon_0/\tau} \right)^{x^2+x} = -\frac{\tau}{\epsilon_0} \left( \left( e^{-\epsilon_0/\tau} \right)^{x^2+x} - \left( e^{-\epsilon_0/\tau} \right)^0 \right) \approx \frac{\tau}{\epsilon_0}$$

3) for  $\frac{\tau}{\epsilon_0} \ll 1 \rightarrow \frac{\epsilon_0}{\tau} \gg 1$ ,

$$z = (2 \times 0 + 1) e^{-0(0+1) \epsilon_0/\tau} + (2 \times 1 + 1) e^{-1(1+1) \epsilon_0/\tau} + \mathcal{O}(e^{-3})$$
$$= 1 + e^{-2 \epsilon_0/\tau}$$

4) The internal energy

$$U = \tau^2 \frac{\partial \ln z}{\partial \tau}$$

$$\gg \frac{\epsilon_0}{\tau} \quad U = \tau^2 \partial_\tau \left( \ln \frac{\tau}{\epsilon_0} \right) = \tau$$

$$1 \ll \frac{\epsilon_0}{\tau}$$

$$U = \tau^2 \left( \frac{1}{1 + 3e^{-2\epsilon_0/\tau}} \right) \times (3e^{-2\epsilon_0/\tau}) \left( \frac{2\epsilon_0}{\tau^2} \right)$$

$$= \frac{6\epsilon_0 e^{-2\epsilon_0/\tau}}{1 + 3e^{-2\epsilon_0/\tau}}$$

$$C_V = \left( \frac{\partial U}{\partial T} \right)_V \quad \xrightarrow{1 \gg \frac{\epsilon_0}{T}} \quad C_V = 1$$

$$\frac{T}{\epsilon_0} \ll 1 \quad C_V = 6 \epsilon_0 \left( \frac{e^{-2\epsilon_0/T} \left( \frac{2\epsilon_0}{T^2} \right) (1 + 3e^{-2\epsilon_0/T}) - (3e^{-2\epsilon_0/T}) \left( \frac{2\epsilon_0}{T^2} \right) e^{-2\frac{\epsilon_0}{T}}}{(1 + 3e^{-2\epsilon_0/T})^2} \right)$$

$$= 12 \epsilon_0^2 \frac{e^{-2\epsilon_0/T}}{T^2} \times \frac{1}{(1 + 3e^{-2\epsilon_0/T})^2}$$

$$\approx 12 \left( \frac{e^{-2\epsilon_0/T}}{(T/\epsilon_0)^2} \right)$$

5 zipper problem.

Consider  $N$  links.  $\epsilon_0 = 0$  one closed and  $\epsilon$  open.

$$Z = \sum_{s=0}^N \exp(-s\epsilon/\tau) \quad x = \epsilon/\tau$$
$$= e^{-0x} + e^{-x} + \dots + e^{-Nx}$$

and take  $e^{-\epsilon/\tau} Z = \sum_{s=0}^N \exp(-s\epsilon/\tau) \times e^{-\epsilon/\tau}$

$$= \sum_{s=0}^N \exp(-(s+1)\epsilon/\tau)$$

$$Z - e^{-\epsilon/\tau} Z = (1 + e^{-x} + \dots + e^{-Nx}) - (e^{-x} + e^{-2x} + \dots + e^{-(N+1)x})$$
$$= 1 - e^{-(N+1)\epsilon/\tau}$$

$$Z = \frac{1 - e^{-(N+1)\epsilon/\tau}}{1 - e^{-\epsilon/\tau}}$$

2)  $\langle s \rangle = -\frac{1}{\epsilon} \frac{\sum (-s\epsilon) e^{-s\epsilon\beta}}{Z} = -\frac{1}{\epsilon} \partial_\beta \ln(Z) = -\frac{1}{\epsilon} \frac{\partial_\beta Z}{Z}$

Let us compute  
this derivative.

$$Z = \frac{1 - e^{-(N+1)\epsilon/\tau}}{1 - e^{-\epsilon/\tau}}$$

$$= \frac{e^{\epsilon\beta} - e^{-N\beta\epsilon}}{e^{\epsilon\beta} - 1}$$

$$\partial_p z = \frac{e^{\epsilon\beta} e - e^{-N\beta\epsilon}(-N\beta\epsilon)}{(e^{\epsilon\beta} - 1)} - \frac{(e^{\epsilon\beta} - e^{-N\beta\epsilon})}{(e^{\epsilon\beta} - 1)^2} (e^{\epsilon\beta}) e$$

$$= e \frac{(e^{\epsilon\beta} + N e^{-N\beta\epsilon})(e^{\epsilon\beta} - 1) - e^{\epsilon\beta}(e^{\epsilon\beta} - e^{-N\beta\epsilon})}{(e^{\epsilon\beta} - 1)^2}$$

$$= e \frac{e^{\epsilon\beta} (e^{\epsilon\beta} + N e^{-N\beta\epsilon} - e^{\epsilon\beta} + e^{-N\beta\epsilon}) - e^{\epsilon\beta} - N e^{-N\beta\epsilon}}{(e^{\epsilon\beta} - 1)^2}$$

$$= e \frac{e^{\epsilon\beta} (N+1) e^{-N\beta\epsilon} - e^{\epsilon\beta} - N e^{-N\beta\epsilon}}{(e^{\epsilon\beta} - 1)^2}$$

$$\frac{\partial_p z}{z} = e \frac{e^{\epsilon\beta} (N+1) e^{-N\beta\epsilon} - e^{\epsilon\beta} - N e^{-N\beta\epsilon}}{(e^{\epsilon\beta} - 1)^2} \times \frac{(e^{\epsilon\beta} - 1)}{e^{\epsilon\beta} - e^{-N\beta\epsilon}}$$

$\underbrace{(e^{\epsilon\beta} - 1)(e^{\epsilon\beta} - e^{-N\beta\epsilon})}_{e^{2\epsilon\beta} - e^{\epsilon\beta} - e^{-(N-1)\epsilon\beta} + e^{-N\beta\epsilon}} \approx e^{2\epsilon\beta} + O(e^{\epsilon\beta})$

$$\frac{\partial_p z}{z} = e \frac{(N+1) e^{-(N-1)\beta\epsilon} - e^{\epsilon\beta} - N e^{-N\beta\epsilon}}{e^{2\epsilon\beta}}$$

$$\approx e (-e^{-\epsilon\beta})$$

$$\langle s \rangle = e^{-\epsilon/\nu}$$