

# Density profiles

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## 1 Useful Quantities and definitions

In this section some common quantities useful for describe the denstites profile are defined and explained.

### 1.1 Critical density of the Universe:

### 1.2 Virialization

A dark matter halo is virialized when its in equilibrium, such an equilibrium occurs after the dark matter have collapsed and the force of gravity equals the **relaxtion** processes (Binney & Tremaine pag 380).

**How is related the virialization with the radius, since what redshift you can define a  $r_{vir}$**

A halo can be characterized using and overdensity  $\Delta_{vir}$  defined as ratio of the density of a virialized halo over the critical density of the Universe  $\Delta_{vir} = \frac{\rho_{vir}}{\rho_c}$ . For a cosmology with  $(\Omega_m + \Omega_\Lambda = 1)$

$$\Delta_{vir} = (18\pi^2 + 82x - 39x^2)/\Omega(z) \quad (1)$$

(Bryan & Norman 1998) it's a good approximation, here  $x = \Omega(z) - 1$ . For the present time ( $z = 0$ )  $\Delta_{vir} = 360$ .

<http://arxiv.org/pdf/astro-ph/9710107v1.pdf>

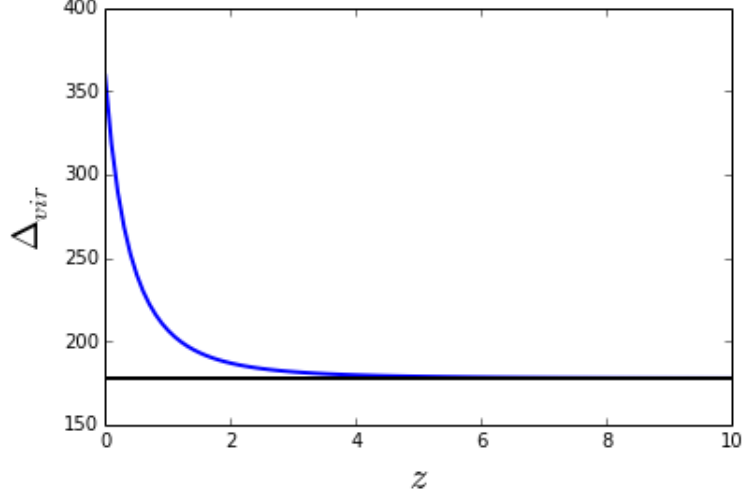
<http://arxiv.org/pdf/astro-ph/9601088v1.pdf>

This overdensity is enclosed in a volume which can be charcterized with a radius  $r_{vir}$  which correspond to a  $M_{vir}$  given  $\Delta_{vir}$

$$\rho_{vir} = \frac{3M_{vir}}{4\pi r_{vir}^3} = \Delta_{vir}\Omega_m\rho_{crit} \quad (2)$$

$$r_{vir} = \left( \frac{3M_{vir}}{4\pi\Delta_{vir}\Omega_m\rho_{crit}} \right)^{1/3} \quad (3)$$

For example for a halo of mass  $M = 1 \times 10^{12}M_\odot$  the corresponding radius is  $r_{vir} = 262.4$  Kpc



### 1.3 $r_{200}$ & $M_{200}$

$$M_{200} = 200\rho_c \frac{4}{3}\pi r_{200}^3 \quad (4)$$

$$M_{vir} = \Delta_{vir}\Omega_m\rho_c \frac{4}{3}\pi r_{vir}^3 \quad (5)$$

Matching  $\rho_c$  for the above two equations we get.

$$\frac{M_{200}}{M_{vir}} = \left( \frac{200}{\Delta_{vir}\Omega_m} \right) \left( \frac{r_{200}}{r_{vir}} \right)^3 \quad (6)$$

Here is common to call  $q = \left( \frac{200}{\Delta_{vir}\Omega_m} \right)$  at  $z = 0$   $q = 2.053$

$$\frac{M_{200}}{M_{vir}} = q \left( \frac{r_{200}}{r_{vir}} \right)^3 \quad (7)$$

## 2 Densities profiles

### 2.1 Plummer

The plumer density profile is one of the simplest models which describes a constant density near the center and falls at large radius.

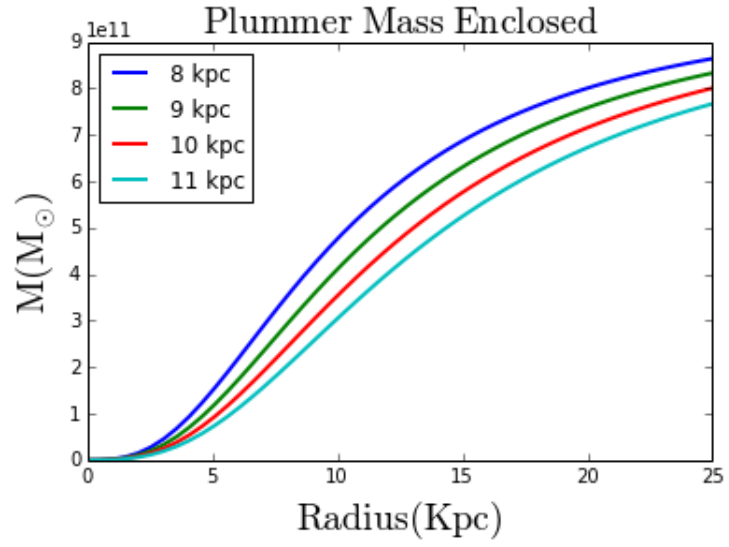
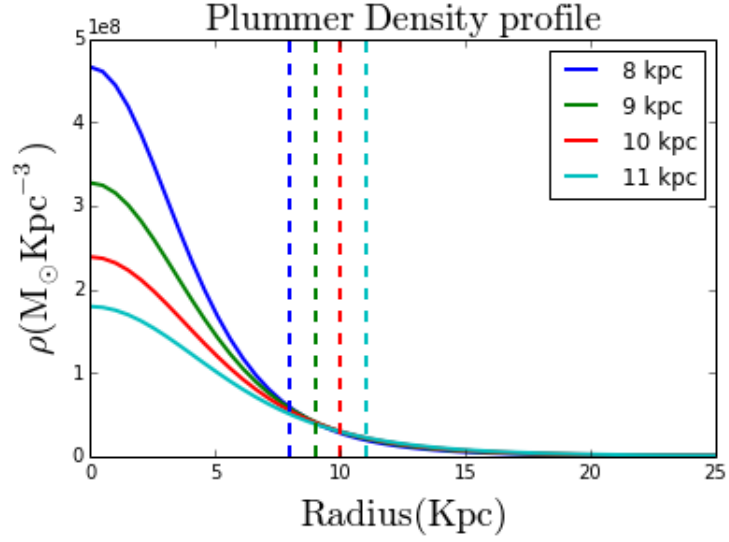
$$\rho_P(r) = \frac{3M}{4\pi a^3} \left( 1 + \frac{r^2}{a^2} \right)^{-5/2} \quad (8)$$

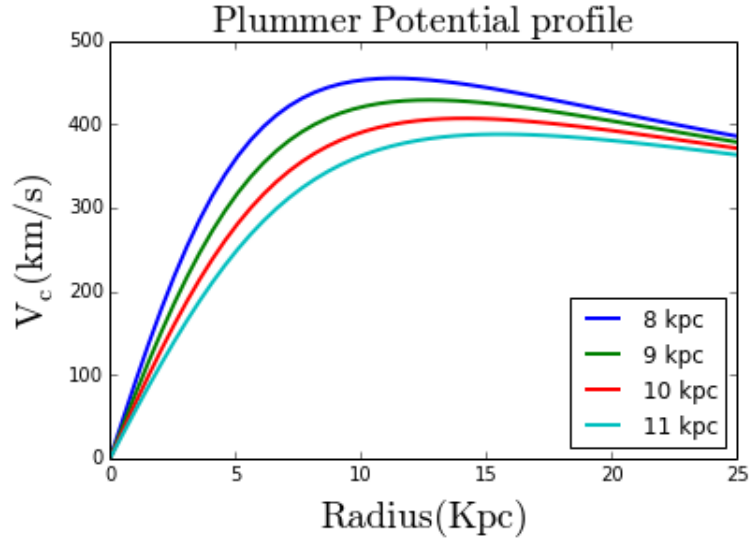
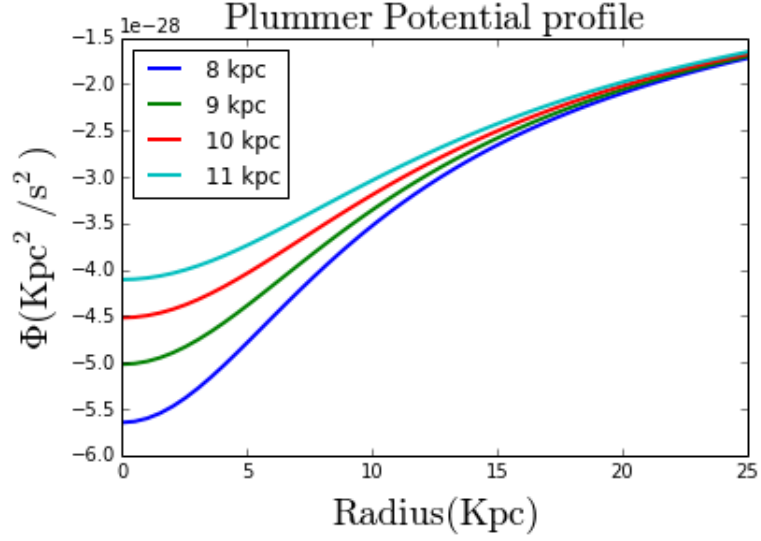
Where  $a$  is call the scale length. The scale length set the length  $a$  in which the majority of the density is enclosed. Note that if  $a$  is zero the plummer potential would be exactly as the potential of a point mass. In the other hand if  $a$  goes to infity the potential is rewpresenting a very extended mass source. In other words the scale length set up the size of the volume in which the mass  $M$  is enclosed.

The enclosed mass can be derived from the density by integrating over a volume.

$$M_P(< r) = 4\pi \int_0^r r'^2 \frac{3M}{4\pi a^3} \left(1 + \frac{r'^2}{a^2}\right)^{-5/2} dr' = \frac{3M}{a^3} \left( \frac{a^4 r^3 \sqrt{r^2/a^2 + 1}}{3(r^2 + a^2)^2} \right) \quad (9)$$

$$M_P(< r) = M \frac{r^3}{(a^2 + r^2)^{3/2}} \quad (10)$$





## 2.2 Hernquist

The Hernquist profile is derived in such a way that it follows the

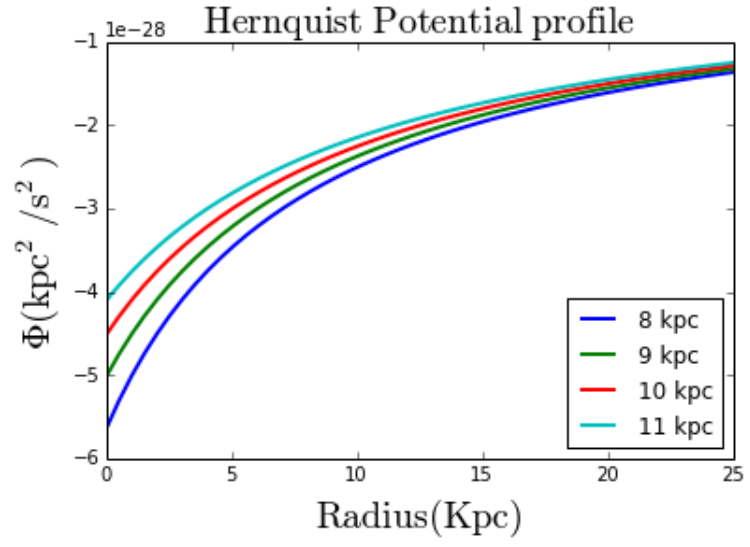
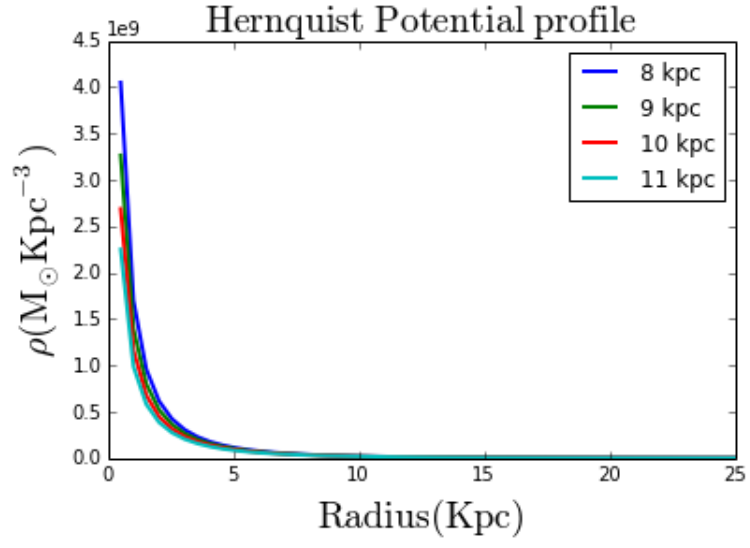
$$\rho_{\text{Hernquist}}(r) = \frac{M}{2\pi} \frac{a}{r(r+a)^3} \quad (11)$$

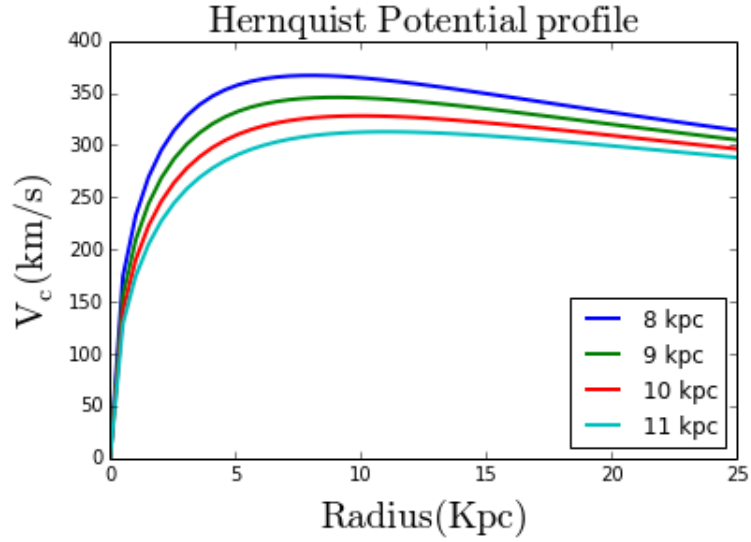
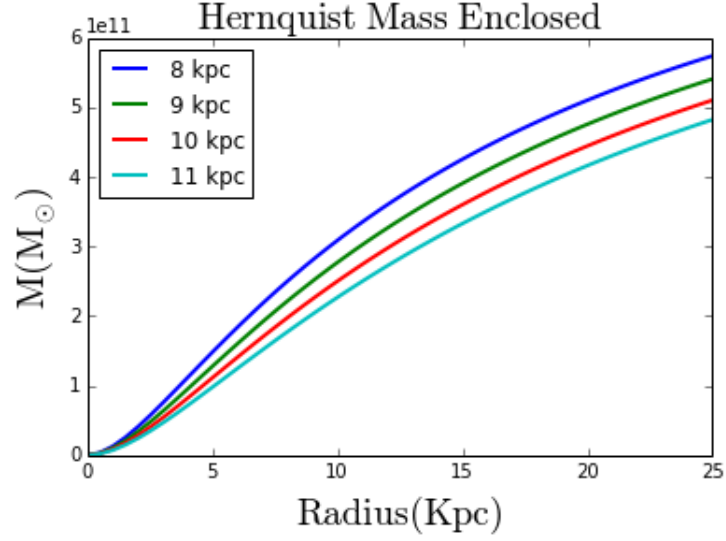
$$M_{\text{Hernquist}}(< r) = 2aM \int \frac{r}{(r+a)^3} dr \quad (12)$$

$$M_{\text{Hernquist}}(< r) = M \frac{r^2}{(r+a)^2} \quad (13)$$

$$\Phi = -\frac{GM}{r+a} \quad (14)$$

$$v_c(r) = GM \frac{r}{(r+a)^2} \quad (15)$$





### 2.3 Singular Isothermal Sphere

The Singular Isothermal Sphere (SIS) describes a system in which the particles follow a Maxwellian density distribution. With this distribution and the Poisson equation the follow density profiles could be derived.

$$\rho_{iso}(r) = \frac{\sigma^2}{2\pi G r^2} \quad (16)$$

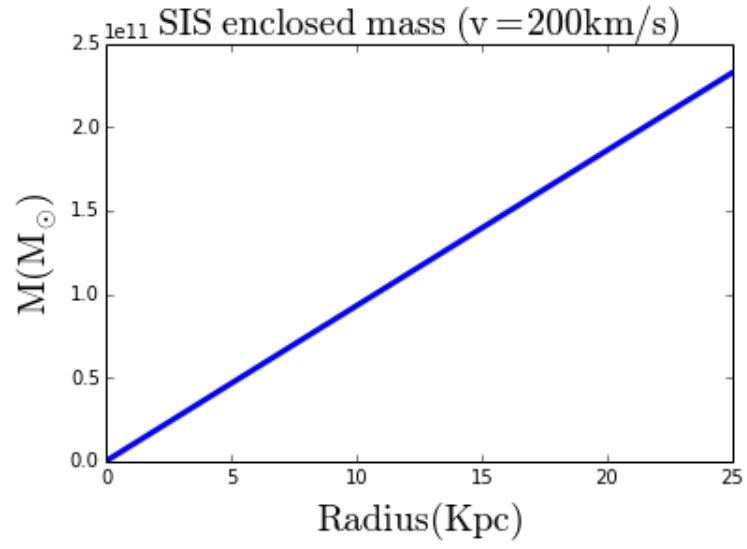
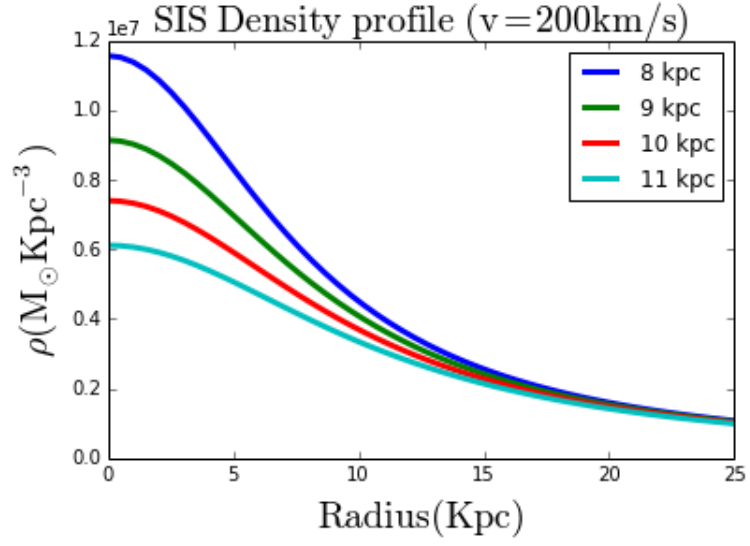
Following the same procedure as with the previous profiles we find  $M$ ,  $\Phi$  and  $v_c$ .

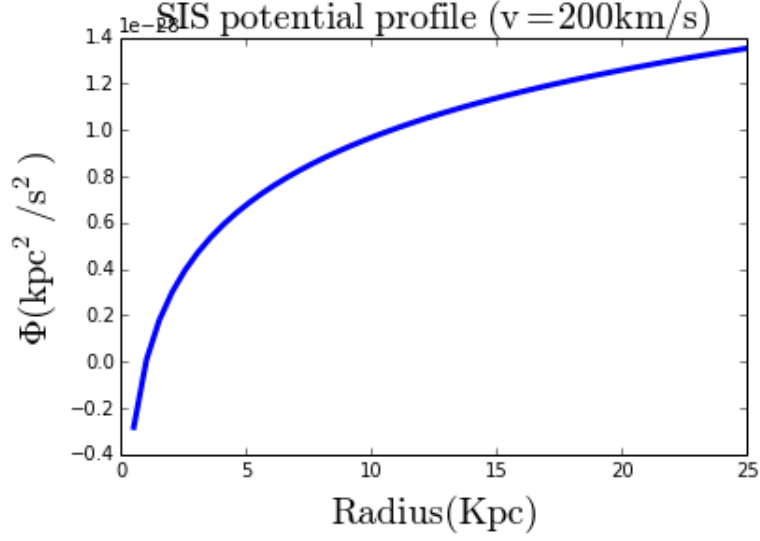
$$M_{iso}(< r) = \frac{2\sigma r}{G} \quad (17)$$

$$\Phi_{iso}(r) = 2\sigma^2 \ln(r) + const. \quad (18)$$

$$v_c(r) = \sqrt{2}\sigma \quad (19)$$

This profile is quite different to the previous ones due to the fact that here the input is the velocity instead of the total Mass.





## 2.4 NFW

$$\rho_{NFW}(r) = \frac{M}{2\pi a^3 (r/a)(1+r/a)^2} \quad (20)$$

$$M_{NFW}(r) = M \left( \ln(1+x) - \frac{x}{1+x} \right) \quad (21)$$

Where  $x = r/a$ , is useful to define the function  $f(x)$  as:

$$f(x) = \ln(1+x) - \frac{x}{1+x} \quad (22)$$

Then 21 can be expressed as:

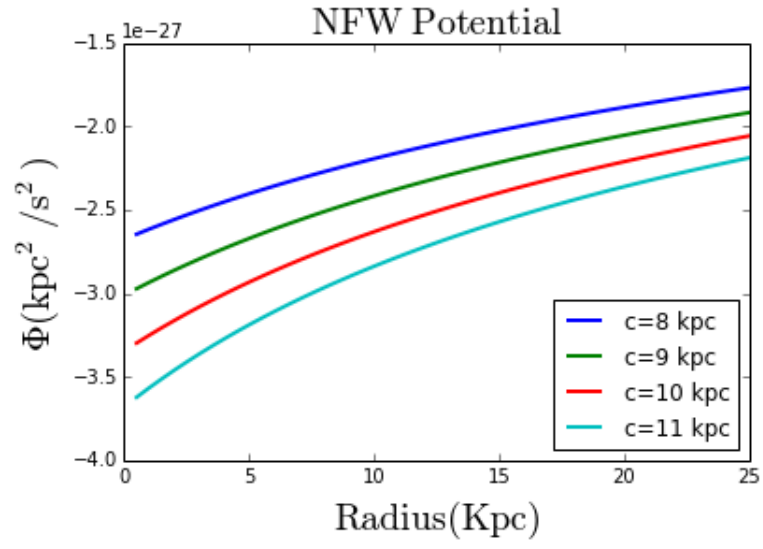
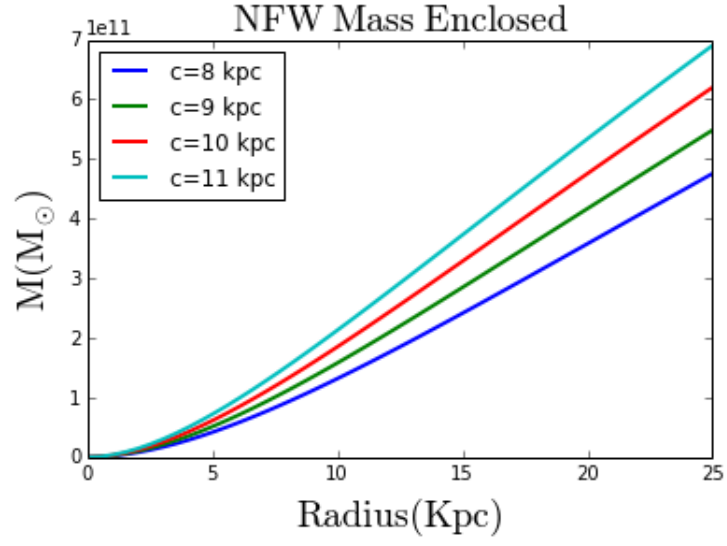
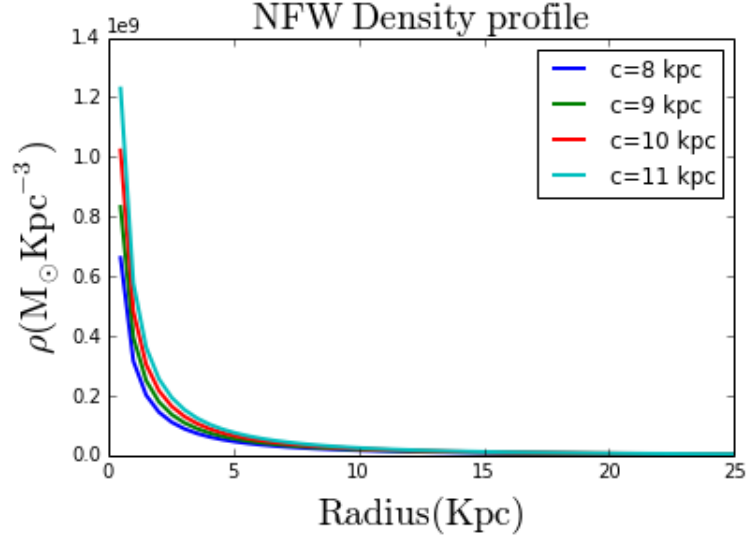
$$M_{NFW} = 4\pi\rho_a a^3 f(x) \quad (23)$$

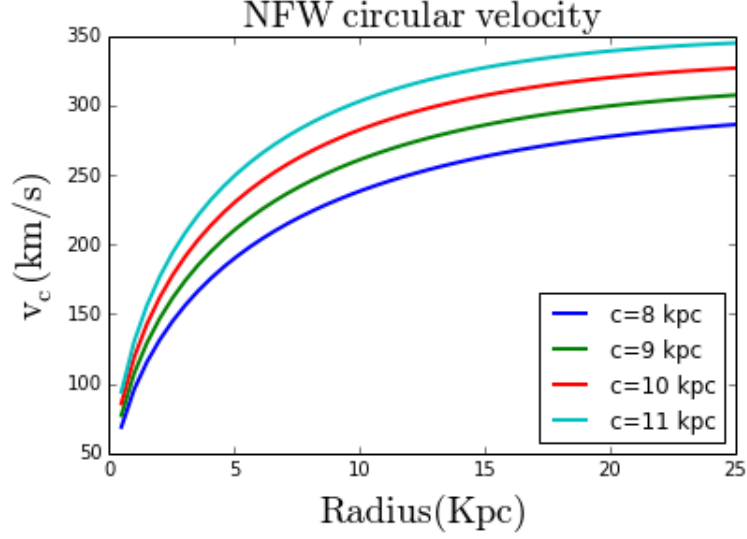
$$\Phi_{NFW} = -4\pi G M \frac{\ln(1+r/a)}{r} \quad (24)$$

$$c(M_{vir}) = 9.60 \left( \frac{M_{vir}}{10^{12} h^{-1} M_{\odot}} \right)^{-0.075} \quad (25)$$

$$v_c(r) = \sqrt{\left( \frac{M(r)G}{r} \right)} = \sqrt{\left( \frac{2M \left( \ln(1+c) - \frac{c}{1+c} \right)}{r} \right)} \quad (26)$$







### 3 Conversion from NFW to the Hernquist profile

The average density of the NFW distribution can be expressed as:

$$\bar{\rho}_{NFW}(r) = \frac{3M_{NFW}(r)}{4\pi r^3} \quad (27)$$

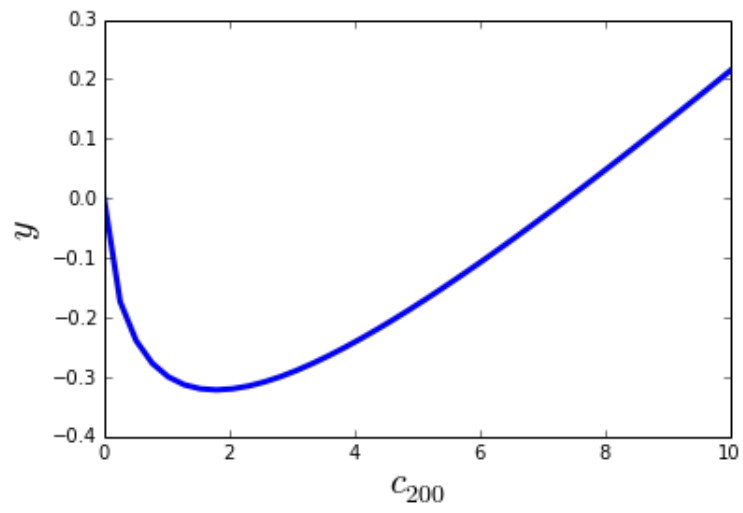
And with eq.23 the  $\rho_{NFW}(r)$  takes the form:

$$\bar{\rho}_{NFW}(r) = 3\rho_a \left(\frac{a}{r}\right)^3 f(x) \quad (28)$$

Now if we want to find the relationship between  $r_{200}$  and  $r_{vir}$  for the NFW profile we have to apply eq7.

$$q = \frac{3\rho_a \frac{a}{r_{200}} f(c_{200})}{3\rho_a \frac{a}{r_{vir}} f(c)} = \frac{c_{200}^3 f(c_{200})}{c_{vir}^3 f(c_{vir})} \quad (29)$$

$$\frac{c_{200}}{c_{vir}} = \left( \frac{f(c_{200})}{q f(c_{vir})} \right)^{1/3} \quad (30)$$



## 4 Miyamoto-Nagai Disk