

Modelling non-linear variability in the frequency domain

Phil Uttley

Why non-linear(non-Gaussian) modelling?

- Even if the underlying variability process is linear, many of our favourite emission mechanisms are non-linear (and in an energy-dependent way).
- Observed broadband-noise variability shows lognormal flux distributions, and this is also expected from propagating mdot-fluctuation models.
- Spectral-timing coherence is often < 1 , this can be due to independent processes in different bands, but can also be produced by non-linear transforms of variability between different energy bands.
- Conventional impulse response modelling of lags and power-spectra is simple convolution (i.e. a ***linear transform***), we need to extend this to the non-linear regime.
- And (thanks Kavitha!)..... all sources have non-zero bicoherence! This cannot be produced by linear Gaussian processes with linear transforms applied - we would like to be able to model the bispectrum of our broadband noise light curves.

Taylor series representation of the light curve

- Let's take a Gaussian-distributed light curve with mean 0, $x(t)$
- Now we transform it to create $y(t)$ with a non-linear function:

$$y(t) = f(x(t))$$

$$y(t) = f(0) + f'(0)x(t) + \frac{f''(0)}{2}x^2(t) + \frac{f'''(0)}{6}x^3(t) + \sum_{n=4}^{\infty} \frac{f_0^{(n')}}{n!}x^n(t)$$

In the Fourier domain this gives:

$$Y_\nu = f(0)\delta_\nu + f'(0)X_\nu + \frac{f''(0)}{2}(X_\nu * X_\nu) + \frac{f'''(0)}{6}(X_\nu * X_\nu * X_\nu) + \sum_{n=4}^{\infty} \frac{f_0^{(n')}}{n!}X_\nu^{(n*)}$$

constant-in-time
components shift to $\nu = 0$

(where X_ν is the Fourier transform of $x(t)$ and $*$ denotes convolution)

Calculating the power spectrum 1

Since our light curves are noise processes there is not a continuous form for X_ν , which for a given light curve and Fourier frequency has real and imaginary parts drawn from independent Gaussians with mean 0:

$$X_\nu \sim (\mathcal{N}(0,1) + i \mathcal{N}(0,1))\sqrt{P_\nu/2} \quad \text{where the power spectrum } P_\nu = \langle |X_\nu|^2 \rangle$$

(expectation value of $X_\nu X_\nu^c$ where c denotes complex conjugation)

The power spectrum of $y(t)$ is:

$$\langle |Y_\nu|^2 \rangle = \left\langle \left(\sum_{n_1=0}^{\infty} \frac{f_0^{(n_1)}}{n_1!} X_\nu^{(n_1*)} \right) \left(\sum_{n_2=0}^{\infty} \frac{f_0^{(n_2)}}{n_2!} X_\nu^{c(n_2*)} \right) \right\rangle$$

This looks like a nightmare. However....

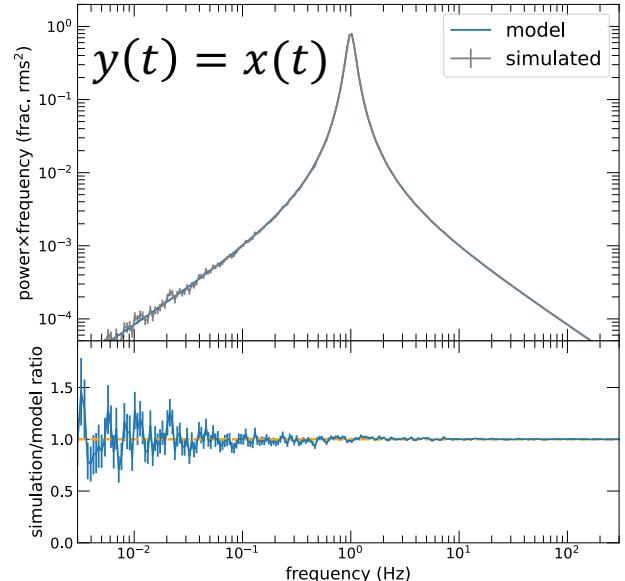
Calculating the power spectrum 2

of non-zero combinations of $\left\langle X_\nu^{(n_1*)} X_\nu^{c(n_2*)} \right\rangle$

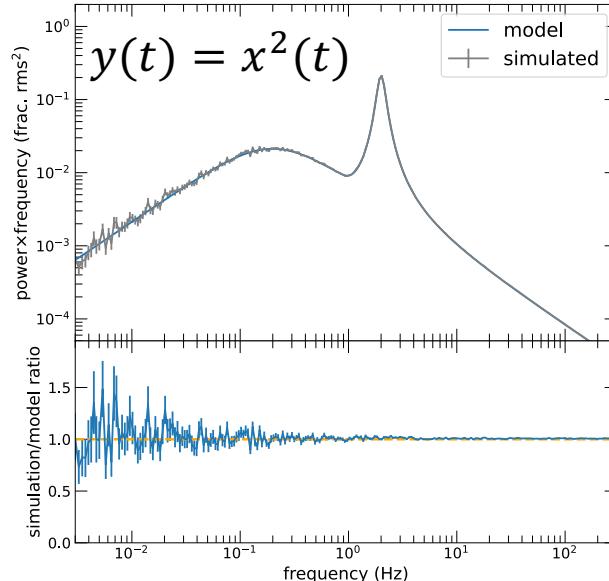
$$C_{n_1, n_2, l} = \frac{n_1! n_2! \max(0, n_1 - (l + 1))!! \max(0, n_2 - (l + 1))!!}{l! (n_1 - l)! (n_2 - l)!}$$

$$\left\langle \left(\frac{f_0^{(n_1')}}{n_1!} X_\nu^{(n_1*)} \right) \left(\frac{f_0^{(n_2')}}{n_2!} X_\nu^{c(n_2*)} \right) \right\rangle = \frac{f^{(n_1')} f^{(n_2')}}{n_1! n_2!} \sum_{\substack{l=0, \text{ even if } k \text{ even} \\ l=1, \text{ odd if } k \text{ odd}}}^{k=\min(n_1, n_2)} C_{n_1, n_2, l} \sigma^{(n_1+n_2-2l)} P_\nu^{(l*)}$$

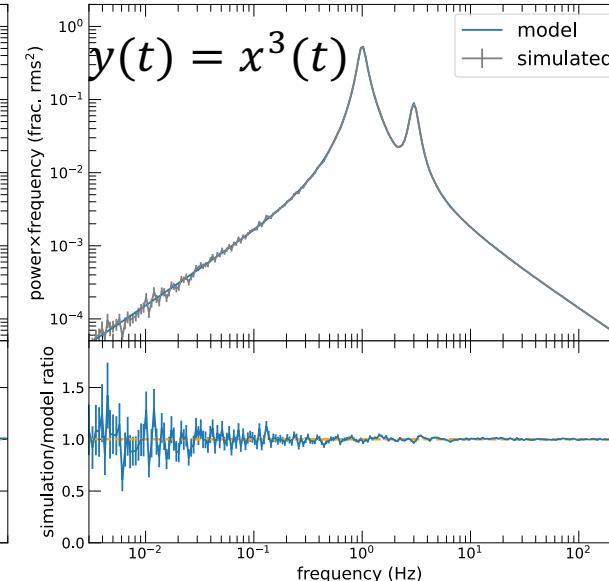
$n_1, n_2 = 1, l = 1$



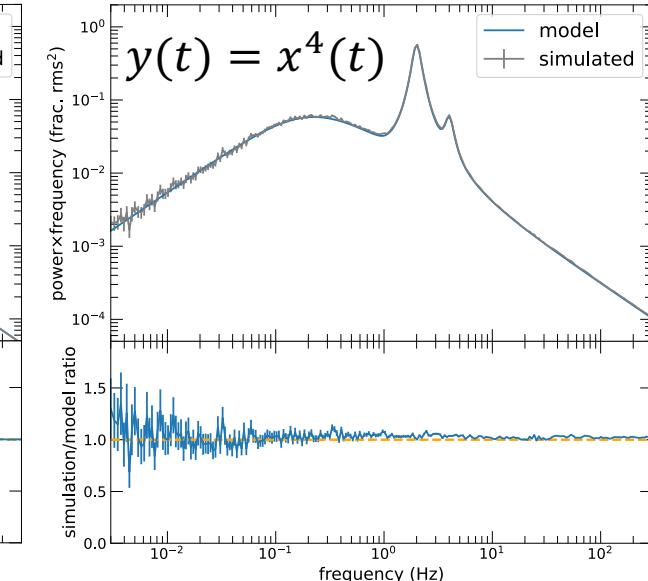
$n_1, n_2 = 2, l = 0, 2$



$n_1, n_2 = 3, l = 1, 3$



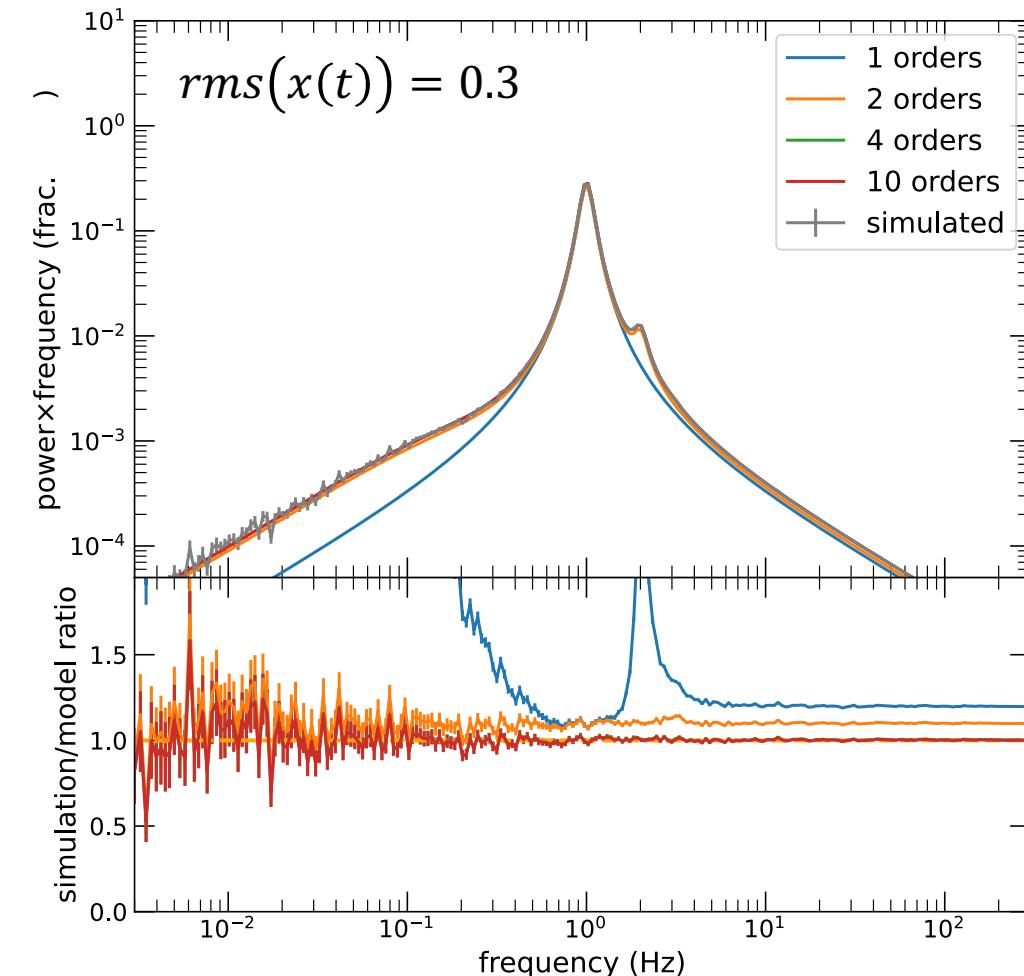
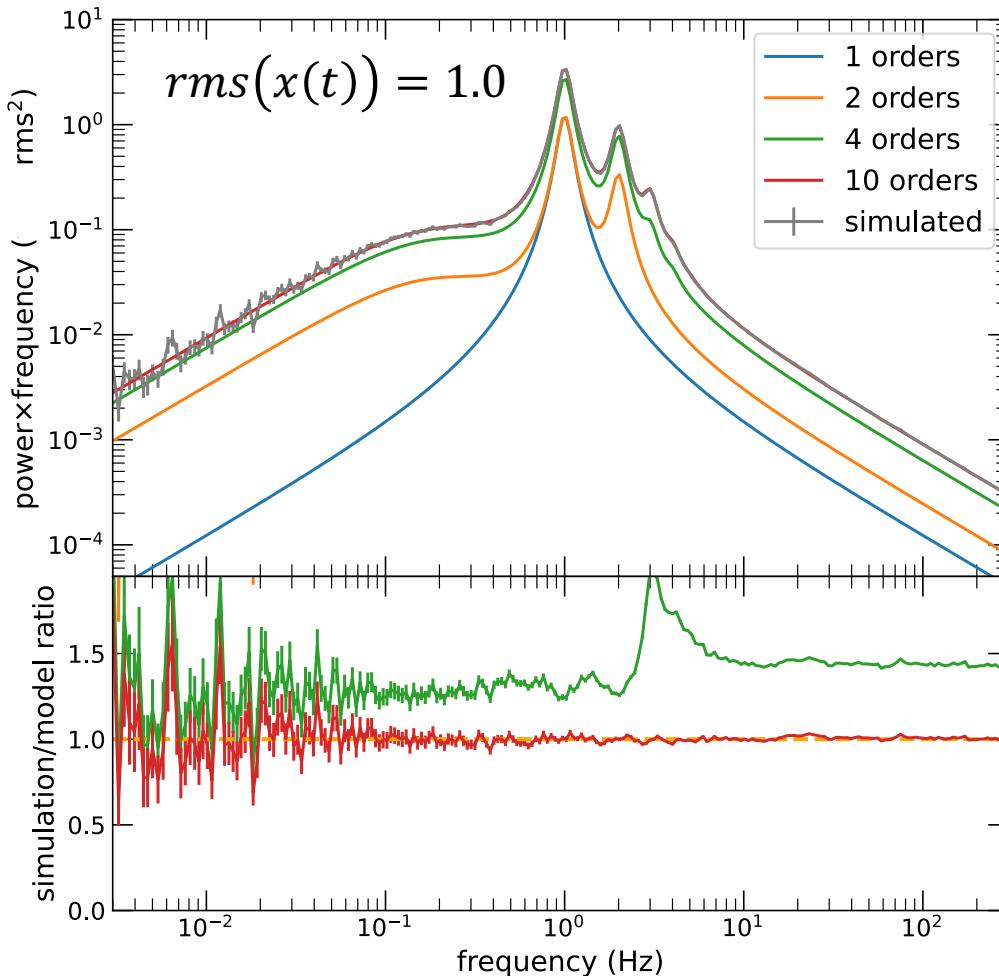
$n_1, n_2 = 4, l = 0, 2, 4$



Example: exponentiated light curve

$$y(t) = \exp(x(t)) = \sum_{n=0}^{\infty} \frac{1}{n!} x^n(t)$$

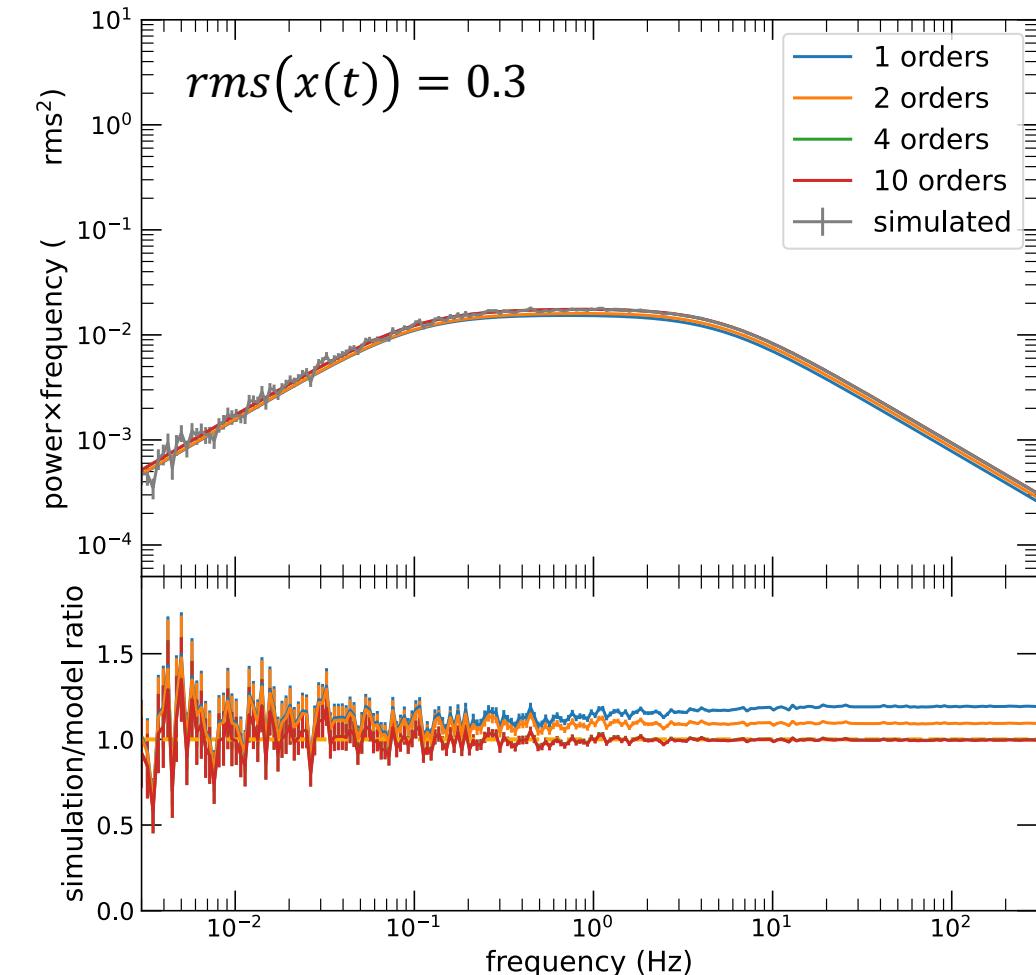
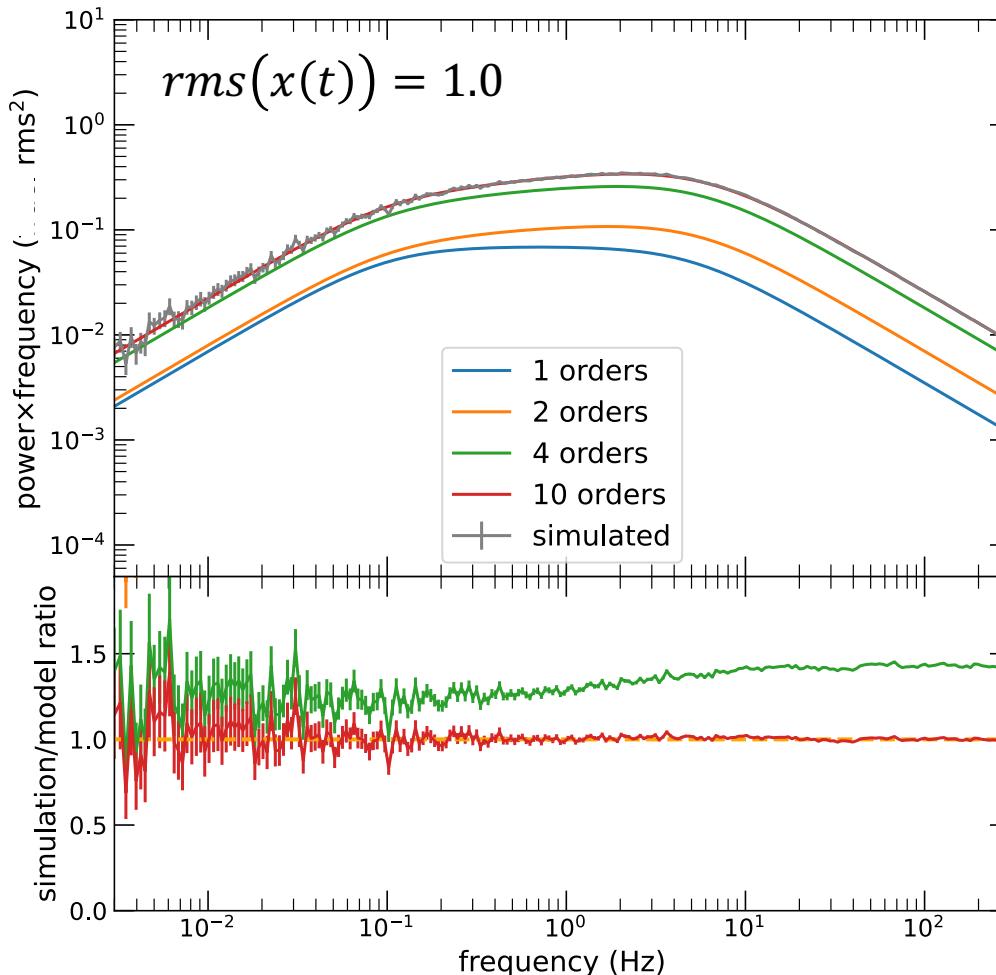
Lorentzian input power spectrum ($q=5$):



Example: exponentiated light curve

$$y(t) = \exp(x(t)) = \sum_{n=0}^{\infty} \frac{1}{n!} x^n(t)$$

Double-bending power-law input power spectrum:



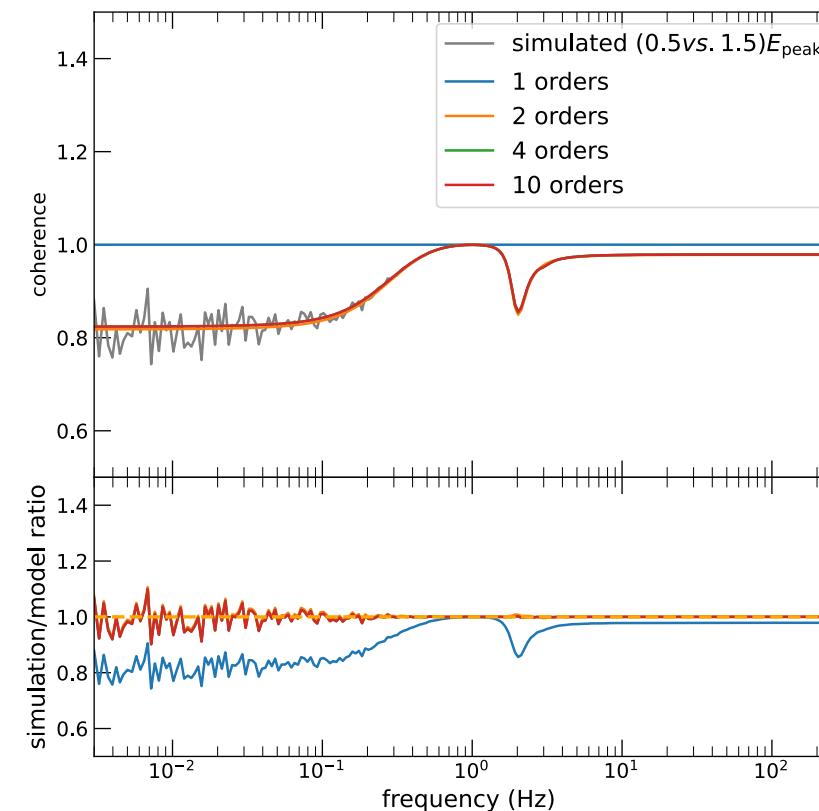
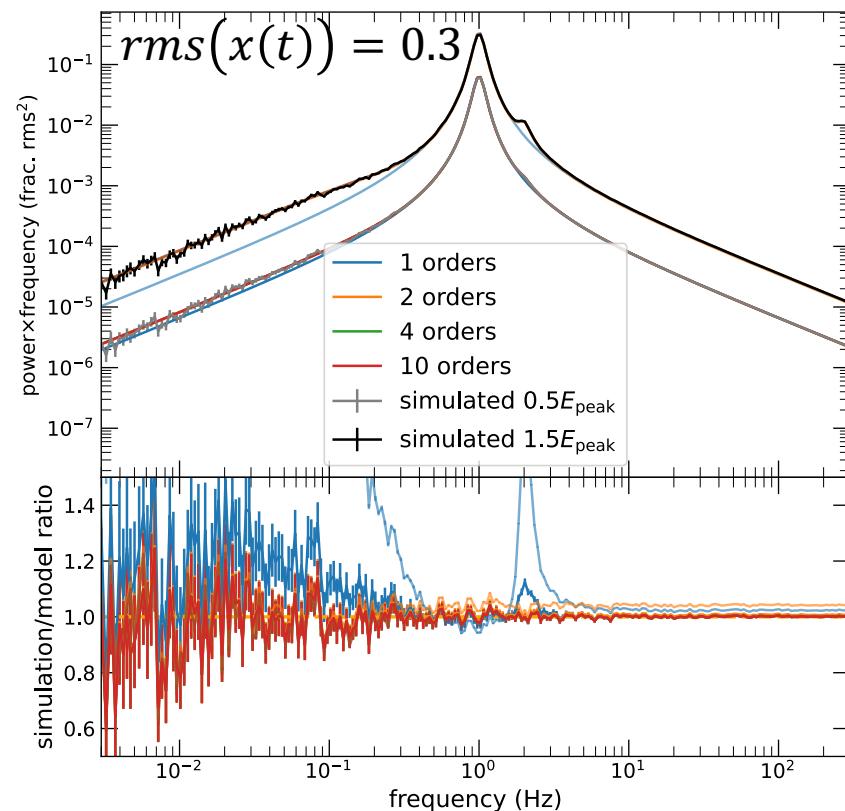
Cross-spectra and coherence (no lags yet)

$$y(t) = f(x(t)) \quad z(t) = g(x(t))$$

$f(x(t))$ and $g(x(t))$ can be different functions, or the same function with different parameters.

Cross-spectrum: $\langle Y_\nu Z_\nu^c \rangle = \langle C_\nu \rangle = \left\langle \left(\sum_{n_1=0}^{\infty} \frac{f_0^{(n_1)}}{n_1!} X_\nu^{(n_1)} \right) \left(\sum_{n_2=0}^{\infty} \frac{g_0^{(n_2)}}{n_2!} X_\nu^{c(n_2)} \right) \right\rangle$

Variable blackbody (constant area, $kT \sim (\text{exponentiated lightcurve})^{0.25}$, evaluate at $E=0.5x$ and $1.5x$ BB peak energy)



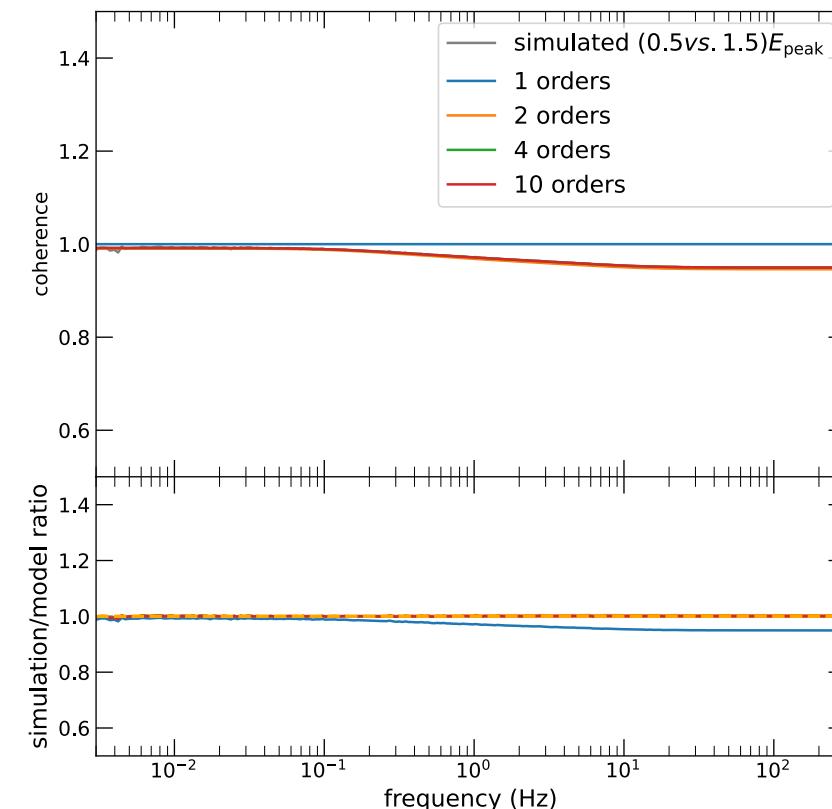
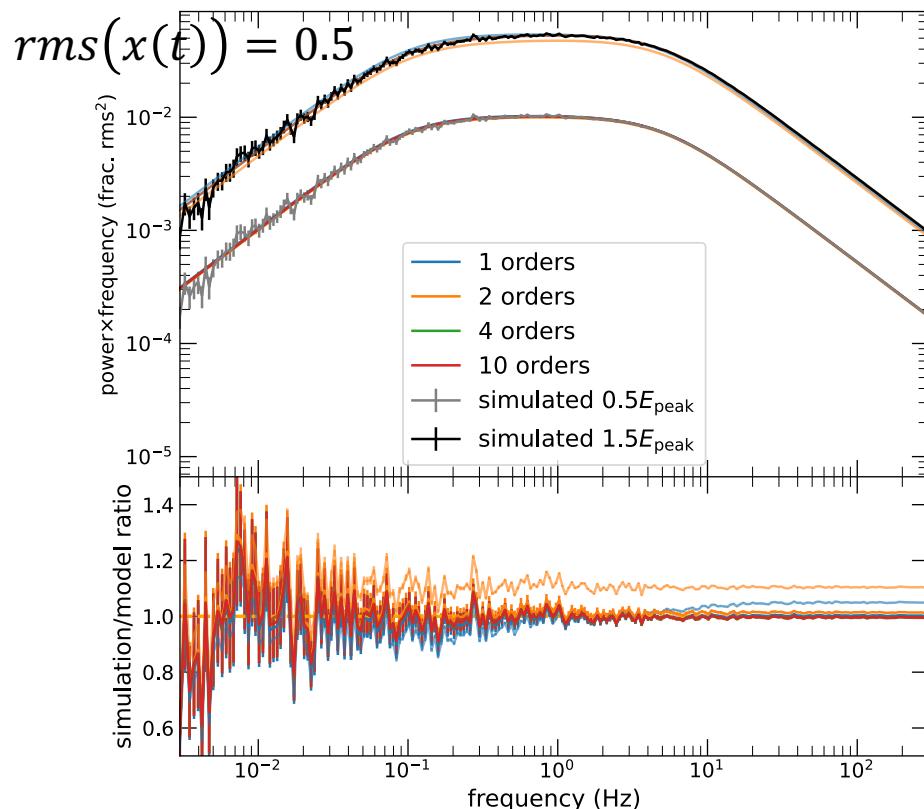
Cross-spectra and coherence (no lags yet)

$$y(t) = f(x(t)) \quad z(t) = g(x(t))$$

$f(x(t))$ and $g(x(t))$ can be different functions, or the same function with different parameters.

Cross-spectrum: $\langle Y_\nu Z_\nu^c \rangle = \langle C_\nu \rangle = \left\langle \left(\sum_{n_1=0}^{\infty} \frac{f_0^{(n_1)}}{n_1!} X_\nu^{(n_1)} \right) \left(\sum_{n_2=0}^{\infty} \frac{g_0^{(n_2)}}{n_2!} X_\nu^{c(n_2)} \right) \right\rangle$

Variable blackbody (constant area, $kT \sim (\text{exponentiated lightcurve})^{0.25}$, evaluate at $E=0.5x$ and $1.5x$ BB peak energy)



Introducing lags

Consider a single emitting region which responds to the variations with a time delay τ_i :

$$u_i(t) = \sum_{n=0}^{\infty} \frac{f_{0,i}^{(n')}}{n!} x^n(t + \tau_i) = \sum_{n=0}^{\infty} \frac{f_{0,i}^{(n')}}{n!} \delta(\tau_i) * x^n(t)$$

The delay can be included by convolving the original light curve with a δ -function

Now the total light curve of multiple regions with different delays:

$$y(t) = \sum_i u_i(t) = \sum_{n=0}^{\infty} \frac{\left(\sum_i f_{0,i}^{(n')} \delta(\tau_i) \right)}{n!} * x^n(t)$$

This is an impulse response function!

In the Fourier domain:

$$Y_\nu = \sum_{n=0}^{\infty} \frac{G_\nu^n}{n!} X_\nu^{(n*)} \quad \text{where } G_\nu^n = \text{FT} \left(\sum_i f_{0,i}^{(n')} \delta(\tau_i) \right), \text{ i.e. the transfer function for order } n$$