

Search for Higgs Decay to Dark Matter and Trigger Studies

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Abstract

Here the abstract of the thesis

Declaration

This dissertation is the result of my own work, except where explicit reference is made to the work of others, and has not been submitted for another qualification to this or any other university. This dissertation does not exceed the word limit for the respective Degree Committee.

João Pela

Acknowledgements

TODO:

- Family
- Friends
- Work colleagues (include CMS collaboration)
- more

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Preface

Thesis structure and so on...

“To my grand mother”

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Chapter 1

Theory

TODO:

- Global status

1.1 Standard Model of Particle Physics

TODO:

- Very brief summary of the Standard Model.

The **Standard Model (SM)** of particle physics is the currently accepted model for describing the physics of elementary particles.

Leptons (J=1/2)				
Generation	Particle Name	Symbol	Mass (GeV/c^2)	Q/e
1^{st}	Electron	e	0.000511	1
	Electron Neutrino	ν_e	$< 3 \times 10^{-9}$	0
2^{nd}	Muon	μ	0.106	1
	Muon Neutrino	ν_μ	$< 1.9 \times 10^{-4}$	0
3^{rd}	Tau	τ	1.777	1
	Tau Neutrino	ν_τ	$< 1.82 \times 10^{-2}$	0

Table 1.1: List of leptons and their fundamental properties

Quarks ($J=1/2$)				
Generation	Particle Name	Symbol	Mass (GeV/c^2)	Q/e
1^{st}	Up	u	$1.5 - 3.3 \times 10^{-3}$	-2/3
	Down	d	$3.5 - 6 \times 10^{-3}$	1/3
2^{nd}	Charm	c	1.16-1.34	-2/3
	Strange	s	$70 - 130 \times 10^{-3}$	1/3
3^{rd}	Top	t	169-173	-2/3
	Bottom	b	$4.13 - 4.37$	1/3

Table 1.2: List of quarks and their fundamental properties

Bosons			
Particle Name	Mass (GeV/c^2)	Q/e	Spin
Photon (γ)	0	0	1
W^\pm	80.4	∓ 1	1
Z^0	91.2	0	1
Gluon (g)	0	0	1
Higgs (H^0)	> 114	0	0

Table 1.3: List of bosons and their fundamental properties

1.2 Higgs Mechanism

Summary of the Higgs Mechanism. Should include

- Motivations
- Explanation of the mechanism itself
- Consequences
- Possible decays

1.3 Higgs Invisible decays

TODO:

- Explain what are SM Higgs invisible decays.
- Go over the possibility of BSM invisible decays.

Chapter 2

Experimental Apparatus

2.1 The Large Hadron Collider

The **Large Hadron Collider (LHC)**[1] is currently the world's largest particle accelerator and is capable to produce the highest energy particle beam ever made by mankind. This gigantic machine with a total perimeter of 26.7 kilometer was built at **European Organization for Nuclear Research (CERN)** in a circular tunnel at an average depth of 100 meters below ground under the Franco-Swiss border near Geneva, Switzerland. Diagram of the LHC tunnel can be found in figure 2.1.

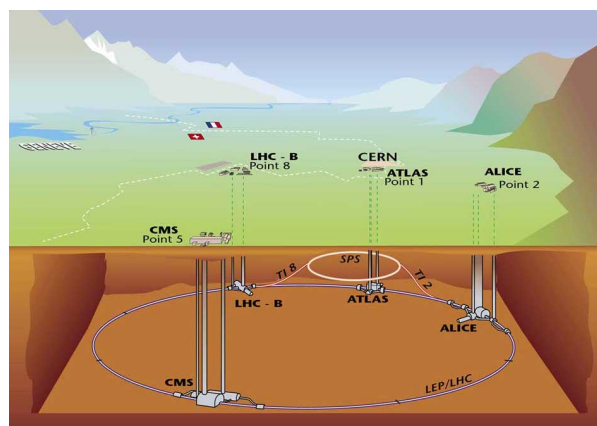


Figure 2.1: Underground diagram of the Geneva area showing the **LHC** location.

The **LHC** is a synchrotron machine with the capability to accelerate particles in two separated beam pipes with travel in opposite direction. These beams only cross and are allowed to collide in four specific points of the accelerator where huge particle detectors are installed to detect the products of such collisions. This experiments are name ATLAS[2], CMS[3], LHCb[4] and ALICE[5].

The accelerator as its name indicates can collide hadrons, more specifically proton or heavy ions. Up to now 3 modes of operation have been tried according to the particles being collided: proton-proton, proton-lead and lead-lead. Depending on the which configuration is chosen we are basically changing the quantity of nucleons available to each colliding element. The maximum design energy per proton is 7 TeV and for each lead nucleon 2.76 TeV. The design luminosity for proton-proton is of $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ and for lead-lead is of $10^{27} \text{ cm}^{-2}\text{s}^{-1}$.

The **LHC** is only the last element of a complex accelerator chain which step by step increases the energy of the particles. Protons are initially obtained by stripping the electrons of hydrogen gas. They are then accelerated at the **Linear Particle Accelerator 2 (LINAC2)** up to the energy of 50 MeV. After this initial step they are injected into the **Proton Synchrotron Booster (PSB)** and the energy ramps up to 1.4 GeV. After protons are passed to the **Proton Synchrotron (PS)** where energy further increases to 25 GeV subsequently they are injected into the **Super Proton Synchrotron (SPS)** where the particle energy level reached 450 GeV. Finally, protons pass to the **LHC** where they can be accelerated to a maximum energy of 7 TeV. A simplified diagram of the **CERN** accelerator chain can be found in figure 2.2.

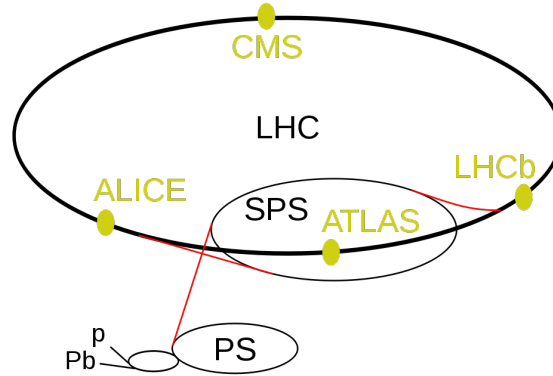


Figure 2.2: CERN Large Hadron Collider Experiment accelerator diagram.

Normal operation of the **LHC** therefore depends on the the upstream accelerators availability. The typically turn around time, the time necessary to stop the accelerator from running and restart collisions is around 2 hours. When stable beams are achieved, a single proton fill can be used to collide protons up to 24 hours, but it is common to restart more frequently to take profit of the higher collision rates possible right at the beginning of a new fill.

Some of the key parameters of the LHC proton-proton and lead-lead operation can be found in table 2.1.

		<i>pp</i>	HI	
Energy per nucleon	E	7	2.76	TeV
Dipole field at 7 TeV	B	8.33	8.33	T
Design Luminosity*	\mathcal{L}	10^{34}	10^{27}	$cm^{-2}s^{-1}$
Bunch separation		25	100	ns
No. of bunches	k_B	2808	592	
No. particles per bunch	N_p	1.15×10^{11}	7.0×10^7	
Collisions				
β -value at IP	β^*	0.55	0.5	m
RMS beam radius at IP	σ^*	16.7	15.9	μm
Luminosity lifetime	τ_L	15	6	hr
Number of collisions/crossing	n_c	≈ 20	-	

* For heavy-ion (HI) operation the design luminosity for Pb-Pb collisions is given.

Table 2.1: The machine parameters relevant for the LHC detectors.[6]

At the LHC we are looking for extremely rare processes as is can be seen in figure 2.3 the production cross section of a SM Higgs boson is more than 9 orders of magnitude smaller than the total proton-proton cross section.

To be able to record and study such rare processes we need to produce a significant amount of collisions. For this purpose the LHC was designed to operate at high instantaneous luminosity, L . This quantity is defined as,

$$L = \frac{N_b^2 n_b f_{\text{rev}} \gamma}{4\pi \epsilon_n \beta^*} F, \quad (2.1)$$

where N_b is the number of protons per bunch, n_b is the number of bunches, f_{rev} is the frequency of revolution, γ is the Lorentz factor, ϵ_n is the normalized emittance, β^* is the beta function at the collision point and F is the reduction factor due to the crossing angle.

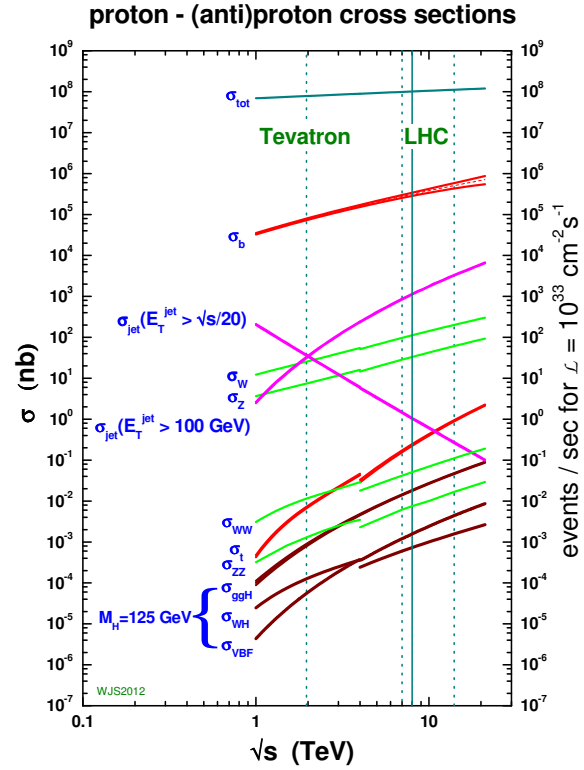


Figure 2.3: Cross sections for several processes for collisions of antiproton-proton and proton-proton as a function of the center of mass energy[3].

2.1.1 Delivered Luminosity

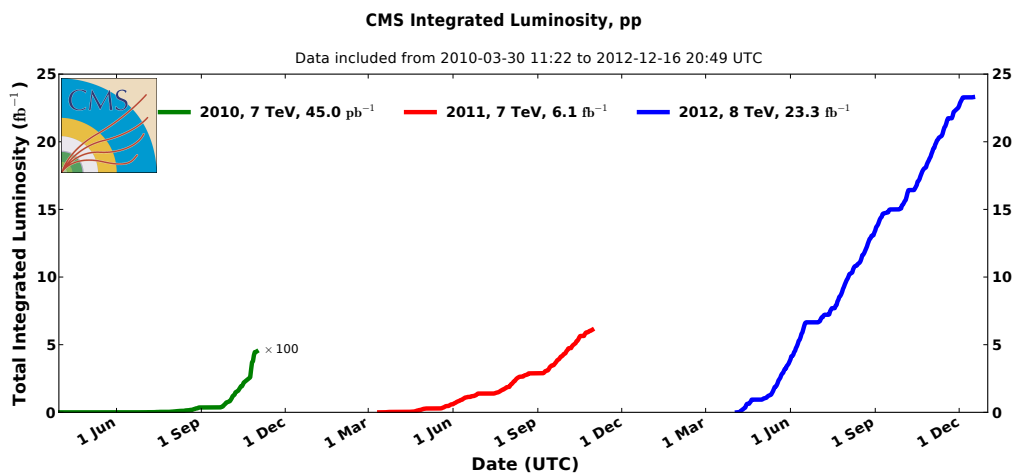


Figure 2.4: Cumulative luminosity versus day delivered to CMS during stable beams and for p-p collisions. This is shown for 2010 (green), 2011 (red) and 2012 (blue) data-taking.

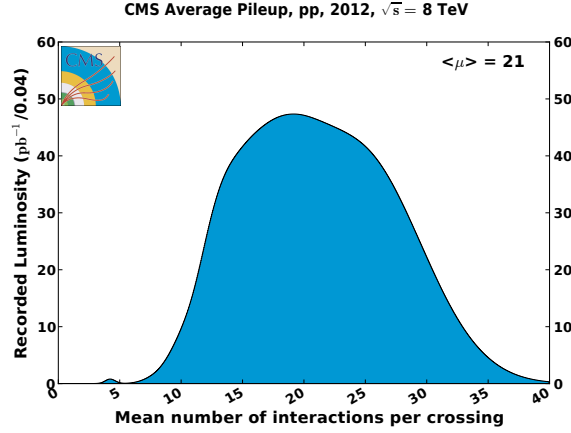


Figure 2.5: Mean number of interactions per bunch crossing at the CMS experiment during 2012.

2.2 The Compact Muon Solenoid Experiment

The **Compact Muon Solenoid (CMS)** experiment is a general purpose experiment located at point 5 of the **LHC**. It was designed to study collisions at its centre and is composed of several sub-systems in an onion shaped structure.

2.2.1 Geometry and conventions

2.2.2 Inner tracking system

The inner tracking system is the closest detector to the beam axis and the interaction region. Its function is to measure the trajectory of all charged particles, like electrons, charged hadrons and muons with momentum above 1 *GeV* being produced at each **LHC** collision. With the help of the strong magnetic field produced by the **CMS** magnet particle trajectories are bent allowing for charge and momentum determination. With the resulting tracks it is then possible to determine the primary vertex as well as vertexes secondary like other lower proton-proton collision or even displaced vertexes from the decay of long lived particles like B mesons.

Building a tracking system for an experiment at the **LHC** is very challenging. At design luminosity an average of 1000 particles will hit such system at a rate approaching 40 *MHz*, leading to high hit density at high rate. It is therefore desirable to have a fast, efficient and high granularity detector where at each layer the occupancy should be at or below 1%. On the other hand each layer should be as thin as possible in order to

not change the incoming particles trajectory or make them lose too much energy. The detector should also be radiation hard and survive for a period of at least 10 years due to its importance and location. This design requirements have lead to a tracker design entirely based on silicon detector technology.

The volume near the interaction point can be split according to the charged particle flux into three regions:

- $r < 10 \text{ cm}$: highest particle flux, up to $\approx 10^8 \text{ cm}^{-2}\text{s}^{-1}$ at $r \approx 4 \text{ cm}$, pixel detectors are used. The pixel size is $\approx 100 \times 150 \text{ }\mu\text{m}^2$, which translates into an occupancy of 10^{-4} per LHC bunch crossing.
- $20 < r < 55 \text{ cm}$: particle flux decreases enough to use silicon micro-strips with a minimum cell size of $10 \text{ cm} \times 80 \text{ }\mu\text{m}$, leading to an occupancy of $\approx 2 - 3\%$ per LHC bunch crossing.
- $50 < r < 110 \text{ cm}$: most outer region of the tracker, particle flux is low enough to use larger pitch silicon micro-strips. The maximum cell size is of $25 \text{ cm} \times 180 \text{ }\mu\text{m}$, and occupancy is of the order of $\approx 1\%$.

The CMS tracker final configuration is composed of a pixel detector with three barrel layers at radii between 4.4 cm and 10.2 cm and 2 disks on each side of the barrel. And a silicon strip tracker with 10 barrel detection layers extending up to 1.1 m with 3 plus 9 disks on each side of the barrel. A schematic of the detector module distribution can be found at figure 2.6. This detector has an acceptance covering up to pseudorapidity of $|\eta| < 2.5$ and has a total active area of about 200 m^2 making the largest silicon tracker ever built.

2.2.3 Electromagnetic Calorimeter

The Electromagnetic Calorimeter (ECAL) is an hermetic energy measurement system comprised of 61200 lead tungstate (PbWO_4) crystals mounted in the barrel and 7324 crystals in each of the 2 endcaps.

Lead tungstate has a fairly high density (8.28 g/cm^3), has a short radiation length (0.89 cm) and a small Moliere radius (2.2 cm). The crystals also have a fast scintillation decay time emitting 80% of the light yield in 25 ns (the minimal bunch crossing time at the LHC). This characteristics make it a good choice for an electromagnetic calorimeter allowing a compact design with fine granularity. However, this crystals emit a fairly

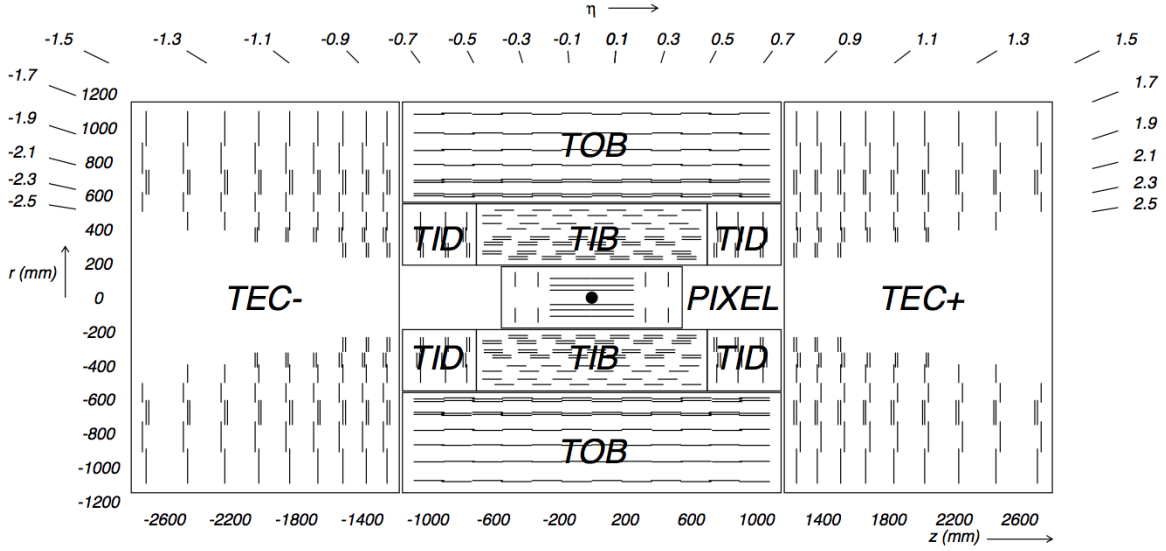


Figure 2.6: Schematic cross section of the CMS tracker. Each line represent a detector module. Double lines represent dual surface back-to-back detector modules.

low light yield ($30 \gamma/\text{MeV}$) which requires the use of photo-detectors with intrinsic gain which will perform well inside a magnetic field. In the barrel region silicon **Avalanche photo-diodes (APD)** are used and **Vacuum Photo-Triodes (VPT)** are used in the endcaps. To guarantee good response from both crystals and **APD** it is necessary to have system thermal stability, with the goal being temperature variation of less than 0.1°C .

The barrel section, the **ECAL Barrel (EB)**, has an inner radius of 129 cm and is composed of 36 identical “supermodules”, each covers the barrel length and corresponding to a pseudo-rapidity interval of $0 < |\eta| < 1.479$. The crystals are quasi-projective (the axes are tilted at 3° with respect to the line from the nominal vertex position) and cover 0.0174 (i.e. 1°) in $\Delta\phi$ and $\Delta\eta$. The crystals have a front face cross-section of $\approx 22 \times 22 \text{ mm}^2$ and a length of 230 mm , corresponding to $25.8 X_0$.

The endcap section, the **ECAL Endcap (EE)**, is at a distance of 314 cm from the vertex and covering a pseudorapidity range of $1.479 < |\eta| < 3.0$, are each structured as 2 “Dees”, consisting of semi-circular aluminium plates from which are cantilevered structural units of 5×5 crystals, known as “supercrystals”.

figure 2.7

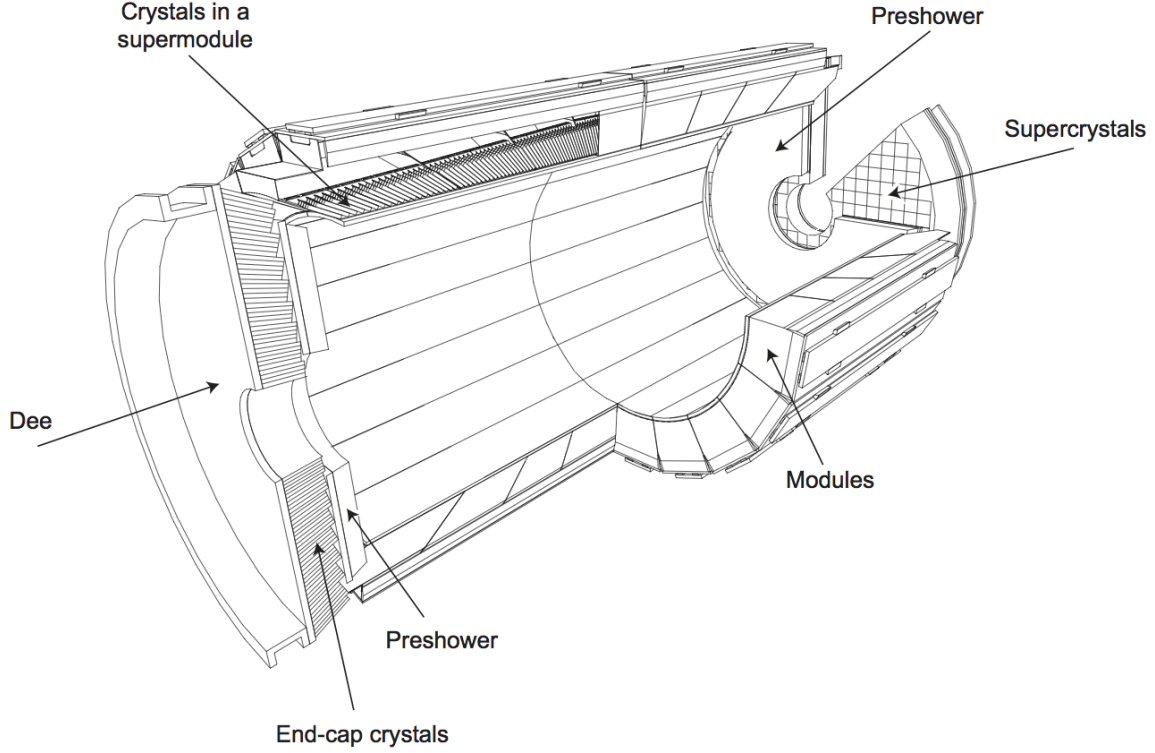


Figure 2.7

2.2.4 Hadronic Calorimeter

The **Hadronic Calorimeter (HCAL)** is a sampling calorimeter which is designed to measure the properties of hadron jets and indirectly neutrinos or other undiscovered particles that would result in apparent missing energy[7].

The design of the **HCAL** was strongly influenced by the choice of the magnet parameters since most of the calorimetry is inside of the magnet. The **HCAL Barrel (HB)** is limited from the beam side by the **ECAL** at radius $r = 1.77\text{ m}$ and outwards by the magnet at radius $r = 2.95\text{ m}$. This is a strict limitation on the amount of absorber material to be used. To improve the measurement capability, an outer calorimeter, the **HCAL Outer (HO)**, is placed outside of the magnet as a *tail catcher* outside of the solenoid magnet. Closing the barrel in each side the **HCAL Endcap (HE)** are present with acceptance up to pseudorapidity $|\eta| < 3$. Additionally, to extend acceptance to $|\eta| < 5.2$ the **HCAL Forward (HF)** is installed at 11.2 m from the interaction point providing excellent hermeticity for E_{\perp}^{miss} measurement. A diagram of the **HCAL** subsystems and their location inside **CMS** can be found in figure 2.8.

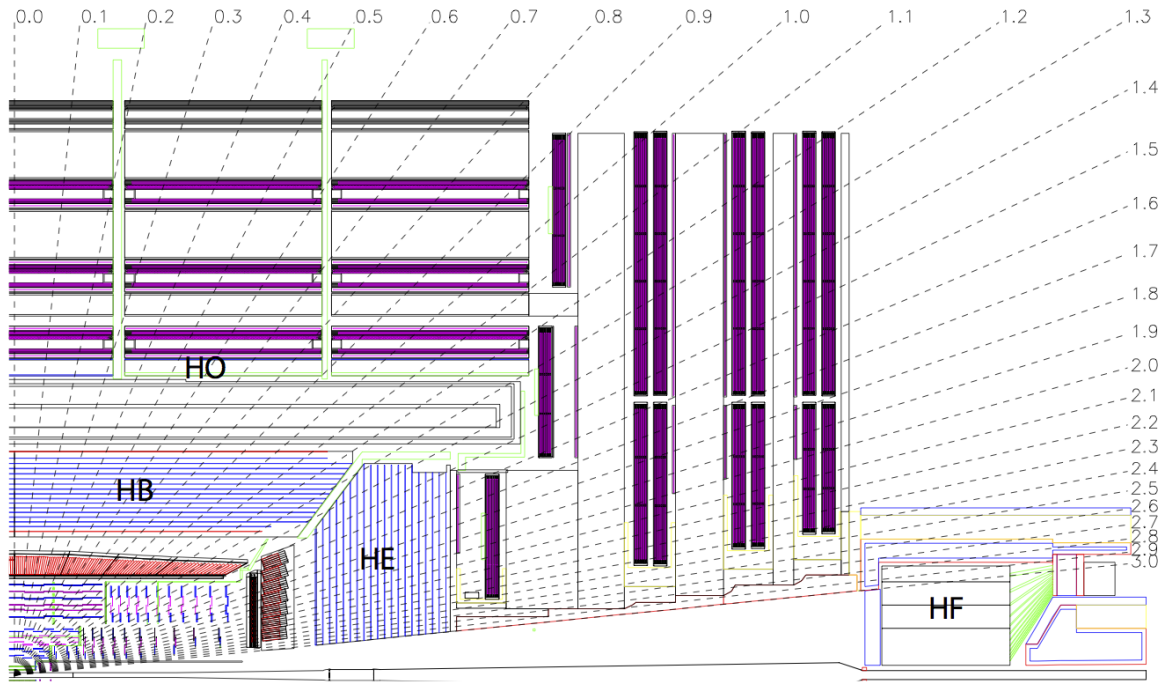


Figure 2.8: Longitudinal view of the CMS detector highlighting the location of the **HCAL** components: **HB**, **HE**, **HO** and **HF**.

The **HB** covers the region up to $|\eta| < 1.3$ and divided into towers with segmentation of $\Delta\eta \times \Delta\phi = 0.087 \times 0.087$ which corresponds to the same area of a 5×5 arrays of **ECAL** crystals.

The **HO**

The **HB** covers the range of $1.3 < |\eta| < 3.0$...

The **HF** covers the range of $3.0 < |\eta| < 5.2$...

2.2.5 Solenoid Magnet

2.2.6 Muon System

figure 2.9

Parameter	Value
Field	4 T
Inner Bore	5.9 m
Length	12.9 m
Number of turns	2168
Current	19.5 kA
Stored Energy	2.7 GJ
Hoop Stress	64 atm

Table 2.2: Parameters of the CMS superconducting solenoid

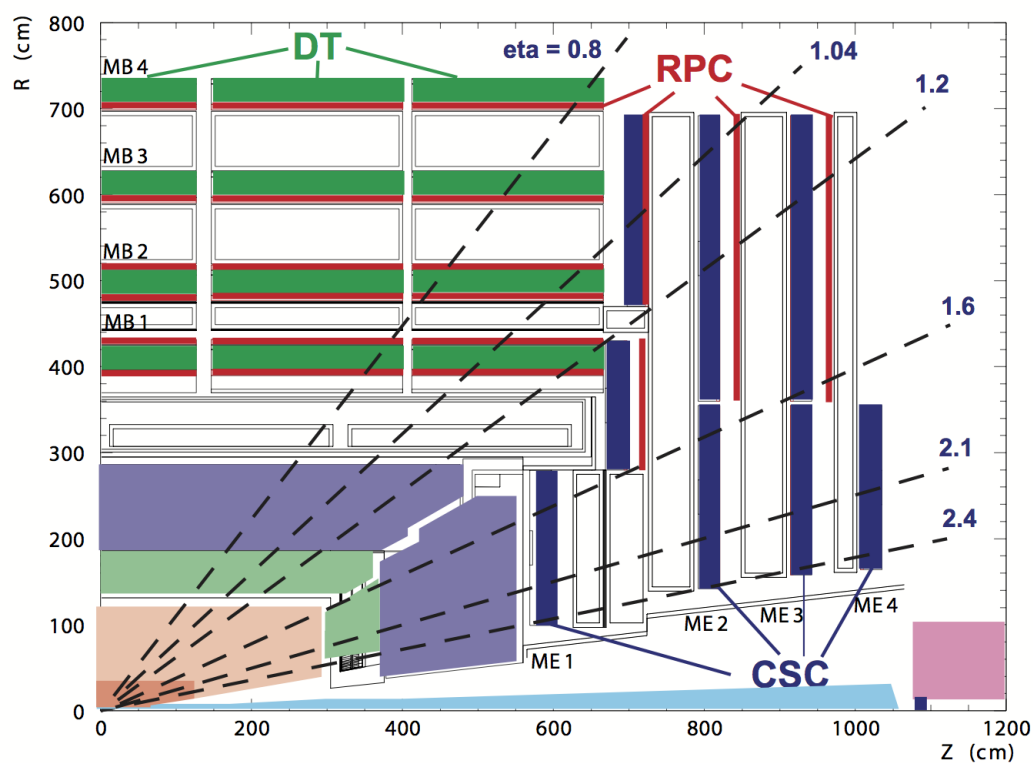


Figure 2.9: TODO

2.2.7 Data Acquisition System

The **Data Acquisition (DAQ)**

2.2.8 Trigger System

The **Level 1 Trigger (L1T)** and **High Level Trigger (HLT)**

CMS Tridas TDR[8]

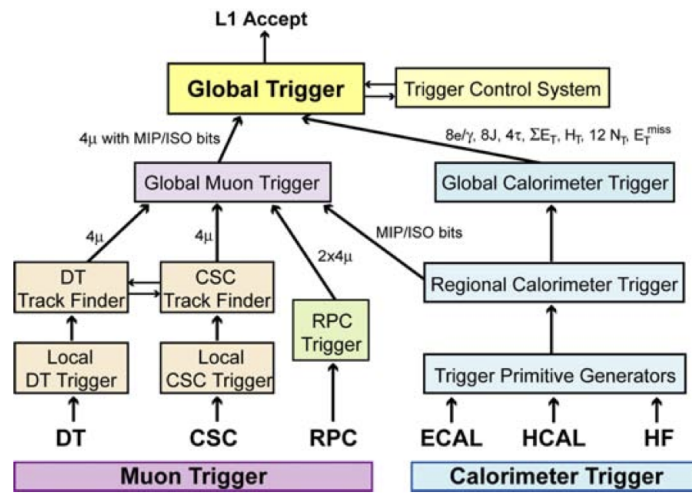


Figure 2.10: TODO

2.2.9 Computing

The **Data Quality Monitoring (DQM)**

2.2.10 Run II Upgrades

CMS L1 Trigger Upgrade TDR[9]

Chapter 3

Technical work

3.1 Level 1 Trigger Data Quality Monitoring System

Hello

Chapter 4

Physics Objects and Monte Carlo simulation

4.1 Physics objects definition

4.1.1 Electron

4.1.2 Muon

4.1.3 Tau

4.1.4 Jets

4.1.5 Missing Transverse Energy

4.2 Monte Carlo simulation

Chapter 5

Prompt Data Analysis

Chapter 6

Parked Data Analysis

Chapter 7

Run II Preparation

7.1 Run II trigger studies

7.2 Run II QCD Monte Carlo samples

Chapter 8

Conclusions

Summary of relevant results and their impact on Particle Physics

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Acronyms

APD Avalanche photo-diodes. 10

CERN European Organization for Nuclear Research. 4, 5

CMS Compact Muon Solenoid. 8, 9, 11

DAQ Data Acquisition. 13

DQM Data Quality Monitoring. 14

EB ECAL Barrel. 10

ECAL Electromagnetic Calorimeter. 9, 11, 12

EE ECAL Endcap. 10

FCT Fundação para a Ciência e a Tecnologia. vi

HB HCAL Barrel. 11, 12, 23

HCAL Hadronic Calorimeter. 11, 12, 23

HE HCAL Endcap. 11, 12, 23

HF HCAL Forward. 11, 12, 23

HLT High Level Trigger. 13

HO HCAL Outer. 11, 12, 23

L1T Level 1 Trigger. 13

LHC Large Hadron Collider. 4–6, 8, 9, 23

LINAC2 Linear Particle Accelerator 2. 5

PS Proton Synchrotron. 5

PSB Proton Synchrotron Booster. 5

SM Standard Model. 2, 6

SPS Super Proton Synchrotron. 5

VPT Vacuum Photo-Triodes. 10