Ab initio methods in solid state physics

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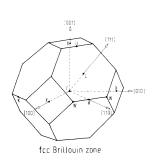
March 30, 2023

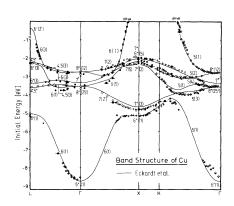
Table of Contents

- Electron interactions theoretical background
- 2 Density functional theory
- 3 Electronic states in crystals
- Band structure calculations
- 5 Functionals and potentials dependent on orbitals
- 6 Strongly correlated and magnetic systems
- Modeling of nanostructures

Band structure of Cu - comparison with the experiment

Band structure of copper obtained from the angle resolved photoemission spectroscopy (ARPES) and DFT calculations.

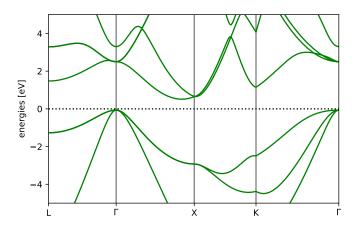




- R. Courths and S. Hüfner, Phys. Rep. 112, 53 (1984)
- H. Eckardt, L. Fritsche, and J. Noffke, J. Phys. F: Met. Phys. 14, 97 (1984)

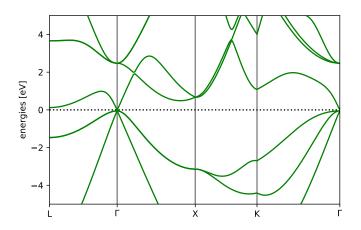
Band structure of Si

Calculations for silicon gives the energy gap $E_{\rm g}=0.58$ eV (GGA-PBE), which is underestimated by $\sim 50\%$ comparing to the experimental value $E_{\rm g}^{\rm exp}=1.14$ eV.

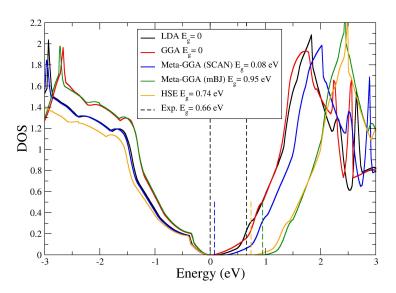


Band structure of Ge

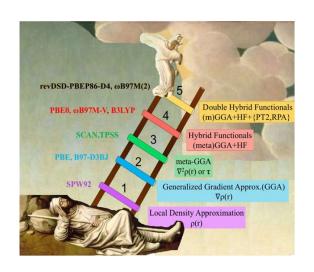
Calculations for germanium, which is a semiconductor with the energy gap $E_{\rm g}^{\rm exp}=0.66$ eV, gives the metallic state (in LDA or GGA).



Electronic density of states of Ge



Jacob's ladder of DFT



J. M. L. Martin and G. Santra, IJC 60, 787 (2020)

Energy gap problem

The fundamental energy gap can be obtained as the difference between the first ionization energy I and the first electron affinity A of the neutral solid

$$I = E(N-1) - E(N), \tag{1}$$

$$A = E(N) - E(N+1), \tag{2}$$

$$E_{g} = I - A = [E(N-1) - E(N)] - [E(N) - E(N+1)], \tag{3}$$

$$E_{g} = E(N+1) - E(N) - [E(N) - E(N-1)] = \varepsilon_{N+1}(N+1) - \varepsilon_{N}(N)$$

$$= \varepsilon_{N+1}(N) - \varepsilon_{N}(N) + \varepsilon_{N+1}(N+1) - \varepsilon_{N+1}(N) = \Delta_{KS} + \Delta_{xc}.$$
(4)

Kohn-Sham energy gap

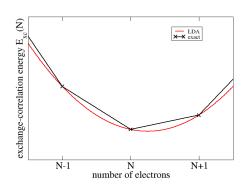
$$\Delta_{\mathsf{KS}} = \varepsilon_{\mathsf{N}+1}(\mathsf{N}) - \varepsilon_{\mathsf{N}}(\mathsf{N}), \tag{5}$$

 $\varepsilon_{N+1}(N)$ – lowest unoccupied KS state, $\varepsilon_N(N)$ – highest occupied KS state.

Energy gap problem

Derivative of the exchange-correlation functional is discontinuous with the change of number of electrons – this effect is not described by analytical functionals.

$$\Delta_{xc} = \frac{\delta E_{xc}[n]}{\delta n(\mathbf{r})}|_{N+\delta} - \frac{\delta E_{xc}[n]}{\delta n(\mathbf{r})}|_{N-\delta} = 0 \quad (LDA, GGA)$$
 (6)



Meta-GGA potentials

Taylor expansion of the exact spherically averaged exchange charge density

$$\rho_{x\sigma}(\mathbf{r},s) = n_{\sigma}(\mathbf{r}) + \frac{1}{6} \left[\nabla^2 n_{\sigma}(\mathbf{r}) - 2\tau_{\sigma}(\mathbf{r}) + \frac{1}{2} \frac{(\nabla n_{\sigma}(\mathbf{r}))^2}{n_{\sigma}(\mathbf{r})} \right] s^2 + ...,$$
(7)

$$\tau_{\sigma}(\mathbf{r}) = \sum_{i} |\nabla \psi_{i\sigma}(\mathbf{r})|^{2}.$$
 (8)

A. D. Becke and E. R. Johnson, J. Chem. Phys. 124, 221101 (2006)

$$V_{x\sigma}^{\mathsf{BJ}}(\mathbf{r}) = V_{x\sigma}^{\mathsf{Slater}}(\mathbf{r}) + \frac{1}{\pi} \sqrt{\frac{5}{12}} \sqrt{\frac{2\tau_{\sigma}(\mathbf{r})}{n_{\sigma}(\mathbf{r})}},\tag{9}$$

$$V_{x\sigma}^{\text{Slater}}(\mathbf{r}) = -\frac{1}{2} \int d\mathbf{r}' \frac{\rho_{\sigma}(\mathbf{r}, \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}.$$
 (10)

F. Tran and P. Blaha, Phys. Rev. Lett. 102, 226401 (2009)

modified Becke-Johnson potential



Meta-GGA functionals

The exchange-correlation functional in meta-GGA is written in a general form

$$E_{\rm xc}[n_{\uparrow}, n_{\downarrow}] = \int d\mathbf{r} n(\mathbf{r}) \varepsilon_{\rm xc}(n_{\uparrow}, n_{\downarrow}, \nabla n_{\uparrow}, \nabla n_{\downarrow}, \tau_{\uparrow}, \tau_{\downarrow}). \tag{11}$$

J. Sun, A. Ruzsinszky, and J. P. Perdew, Phys. Rev. Lett. 115, 036402 (2015)

Strongly constrained and appropriately normed (SCAN) functional obeys all 17 known exact constraints. The exchange part can be written in a form

$$E_{x}[n] = \int d\mathbf{r} n(\mathbf{r}) \varepsilon_{x}(n) F_{x}(s, \alpha), \qquad (12)$$

where $\varepsilon_{\rm x}(n)$ is the exchange energy for uniform gas per one electron and s is the dimensionless gradient of density

$$s = \frac{|\nabla n|}{2(3\pi^2)^{1/3} n^{4/3}}. (13)$$

Electron localization function (ELF)

Meta-GGA functionals use the dimensionless parameter

$$\alpha = \frac{\tau - \tau^{w}}{\tau^{u}},\tag{14}$$

where $\tau^{\rm w}$ is the Weizsäcker density of kinetic energy $\tau^{\rm w}=|\nabla n|^2/8n$, which defines τ for a single orbital, and $\tau^{\rm u}$ is the value for a uniform gas

$$\tau^{\mathsf{u}} = \frac{3}{10} \left(\frac{3}{\pi^2}\right)^{\frac{2}{3}} n^{\frac{5}{3}}.\tag{15}$$

Parameter $lpha \in [0,1]$ is connected with the electron localization function

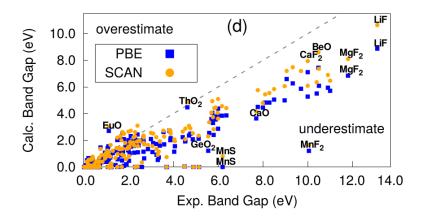
$$ELF = \frac{1}{1 + \alpha^2},\tag{16}$$

which is used to characterize the distribution of electron density

- **1** slowly changing density, typical for metals, $\alpha \approx 1$, ELF $\approx \frac{1}{2}$.
- ② covalent bonds between two orbitals, $\alpha = 0$, ELF = 1.
- ullet weak noncovalent bonds between closed atomic shells, $\alpha \to \infty$, ELF $\to 0$.

Meta-GGA functionals

E. B. Isaacs and C. Wolverton, Phys. Rev. Mater. 2, 063801 (2018)



Self-interaction problem

LDA or GGA

$$E_{\rm H}$$
 - exact, $E_{\rm xc}$ - approximate $E_{\rm H} + E_{\rm xc} \neq 0$ (self-interaction)

Hartree-Fock

$$E_{x} = -\int d\mathbf{r}_{1} d\mathbf{r}_{2} \frac{\psi_{i}^{\sigma*}(\mathbf{r}_{1})\psi_{i}^{\sigma*}(\mathbf{r}_{2})\psi_{i}^{\sigma}(\mathbf{r}_{2})\psi_{i}^{\sigma}(\mathbf{r}_{1})}{|\mathbf{r}_{1} - \mathbf{r}_{2}|}$$

$$= -\frac{1}{2} \int d\mathbf{r}_{1} d\mathbf{r}_{2} \frac{|\psi_{i}^{\sigma}(\mathbf{r}_{1})|^{2}|\psi_{i}^{\sigma}(\mathbf{r}_{2})|^{2}}{|\mathbf{r}_{1} - \mathbf{r}_{2}|}$$

$$= -\frac{1}{2} \int d\mathbf{r}_{1} d\mathbf{r}_{2} \frac{n_{i}(\mathbf{r}_{1})n_{i}(\mathbf{r}_{2})}{|\mathbf{r}_{1} - \mathbf{r}_{2}|} = -E_{H}$$
(17)

 $\rightarrow E_{H} + E_{x} = 0$ (no self-interaction)



Hybrid functionals

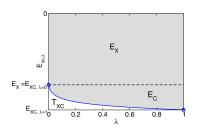
A. D. Becke, J. Chem. Phys. 98, 1372 (1993)

$$H_{\lambda} = E_{k} + \lambda V$$
 $\frac{dE_{\lambda}}{d\lambda} = \langle \psi_{\lambda} | V | \psi_{\lambda} \rangle.$ (18)

Adiabatic connection formula

$$E[n] = E_{k}[n] + \int_{0}^{1} d\lambda \langle \psi_{\lambda} | V | \psi_{\lambda} \rangle, \tag{19}$$

$$E_{xc}[n] = \int_0^1 d\lambda \langle \psi_{\lambda} | V | \psi_{\lambda} \rangle - E_{H}[n] = \int_0^1 d\lambda E_{xc,\lambda}[n], \qquad (20)$$



$$E_{\mathsf{xc},\lambda}[n] = (1 - \lambda)E_{\mathsf{x}}^{\mathsf{HF}}[n] + \lambda E_{\mathsf{xc}}^{\mathsf{DFT}}[n], \quad (21)$$

$$E_{\text{xc}}^{\text{hyb}}[n] = \frac{1}{2}E_{\text{x}}^{\text{HF}}[n] + \frac{1}{2}E_{\text{xc}}^{\text{DFT}}[n].$$
 (22)

Hybrid functionals

J. P. Perdew, M. Ernzerhof, and K. Burke, J. Chem. Phys. 105, 9982 (1996)

$$E_{xc}^{hyb}[n] = E_{xc}^{DFT}[n] + \alpha (E_{x}^{HF}[n] - E_{x}^{DFT}[n]),$$

$$\alpha = 0.25 \text{ (PBE0)},$$

$$\alpha = 0.5 \text{ (Fock-0.5)}.$$

J. Heyd, G. E. Scuseria, and M. Ernzerhof, J. Chem. Phys. 118, 8207 (2003) (HSE)

$$E_{\text{xc}}^{\text{hyb}}(\mu) = (1 - \alpha)E_{\text{x}}^{\text{GGA,SR}}(\mu) + \alpha E_{\text{x}}^{\text{HF,SR}}(\mu) + E_{\text{x}}^{\text{GGA,LR}} + E_{\text{c}}^{\text{GGA}}.$$
 (24)

Meta-GGA potential mBJ vs functional HSE

F. Tran and P. Blaha, Phys. Rev. Lett. **102**, 226401 (2009)

Solid	LDA	MBJLDA	HSE	G_0W_0	GW	Expt.
Ne (A1)	11.42	22.72		19.59 ^e	22.1 ^g	21.70
Ar (A1)	8.16	13.91	10.34 ^a	13.28e	14.9 ^g	14.20
Kr (A1)	6.76	10.83				11.6
Xe (A1)	5.78	8.52				9.8
C (A4)	4.11	4.93	5.49 ^b	5.50 ^e	6.18^{g}	5.48
Si (A4)	0.47	1.17	1.28 ^b	1.12 ^e	1.41 ^g	1.17
Ge (A4)	0.00	0.85	0.83^{b}	0.66^{f}	0.95^{g}	0.74
LiF (B1)	8.94	12.94		13.27e	15.9 ^g	14.20
LiCl (B1)	6.06	8.64				9.4
MgO (B1)	4.70	7.17	6.67 ^b	7.25 ^e	9.16 ^g	7.83
ScN (B1)	-0.14	0.90		$0.95^{\rm f}$	1.4 ^h	~0.9
MnO (B1)	0.76	2.95	2.8°		3.5^{i}	3.9 ± 0.4
FeO (B1)	-0.35	1.82	2.2°			2.4
NiO (B1)	0.42	4.16	4.2°	1.1^{f}	4.8^{i}	4.0, 4.3
SiC (B3)	1.35	2.28	2.40^{b}	2.27 ^c	2.88g	2.40
BN (B3)	4.39	5.85	5.99 ^b	6.10 ^e	7.14 ^g	~ 6.25
GaN (B3)	1.63	2.81	3.14 ^b	2.80^{e}	3.82^{g}	3.20
GaAs (B3)	0.30	1.64	1.12 ^b	1.30 ^e	1.85^{g}	1.52
AlP (B3)	1.46	2.32	2.51 ^b	2.44 ^e	2.90g	2.45
ZnS (B3)	1.84	3.66	3.49 ^b	3.29 ^e	$4.15^{\rm g}$	3.91
CdS (B3)	0.86	2.66	2.25 ^b	2.06 ^e	2.87^{g}	2.42
AlN (B4)	4.17	5.55	5.81 ^b	5.83 ^f		6.28
ZnO (B4)	0.75	2.68	2.49^{d}	2.51 ^f	3.8g	3.44

