# Ab initio methods in solid state physics

XIII. Lattice Thermodynamics

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# Thermodynamics of lattices

- This is not a thermodynamics lecture
- Basic thermodynamics of crystals
  - Harmonic phonons
  - Anharmonic perturbation
- Thermodynamic properties
  - Thermodynamic functions
  - Thermal displacements
  - Heat capacity
  - Thermal expansion
  - Heat conductance
  - Phase transitions

# Thermodynamics of crystals

The thermodynamics in the statistical mechanics formulation is founded on probability distribution and partition function Z:

$$Z = \sum e^{-E/k_BT}$$

The remining thermodynamic functions are defined in relation to Z. For example Helmholtz free energy F:

$$F = -k_B T \ln Z,$$

also entropy, pressure, stresses, heat capacity are defined in terms of various functions of Z.

Calculation of Z requires knowledge of the energy levels of the system E. The energy can be expanded in the small parameter  $\epsilon = \sqrt{\langle u^2 \rangle}/\langle R \rangle$  where u is a displacement and R is an interatomic distance. Thus, terms in Hamiltonian H can be ordered in the powers of  $\epsilon$ :

$$H = H_0 + H_2 + H_3 + H_4 + \dots$$

### Harmonic phonons

If we keep just two first terms in the Hamiltonian, the energy levels written in the *occupation number* formulation will take the form:

$$E_2(\dots,n_k,\dots) = \sum_k \hbar \omega_k \left(n_k + \frac{1}{2}\right)$$

which leads to the harmonic partition function  $Z_H$ :

$$Z_H = \sum_{(\dots,n_k,\dots)} e^{-\sum_k \hbar \omega_k (n_k + \frac{1}{2})/k_B T} = \prod_k \frac{e^{-\hbar \omega_k/2k_B T}}{1 - e^{-\hbar \omega_k/k_B T}}$$

giving harmonic free energy  $F_H$ :

$$F_{H} = \sum_{k} \left[ \frac{1}{2} \hbar \omega_{k} + k_{B} T \ln \left( 1 - e^{-\hbar \omega_{k}/k_{B}T} \right) \right] \label{eq:fh}$$

## Anharmonic perturbation

Sustaining fourth order term in the Hamiltonian ( $H=H_0+H_2+H_4$ ) leads to the anharmonic partition function Z:

$$Z = e^{-E_0/k_BT} \sum e^{-(E_2 + E_4)/k_BT} = e^{-E_0/k_BT} Z_H \left[ 1 - Z_H^{-1} \sum \frac{E_4}{k_BT} e^{-E_2/k_BT} \right]$$
 
$$Z = e^{-E_0/k_BT} Z_H \left[ 1 - \frac{\langle E_4 \rangle_H}{k_BT} \right]$$

giving free energy F:

$$F = E_0 - k_B T \ln Z_H + \langle E_4 \rangle_H = E_0 + F_H + F_A$$

# Anharmonic free energy

After rather lenghty calculations the anharmonic free energy term becomes:

$$\begin{split} F_A = & 12 \sum_{k,k'} \Phi_{k,-k,k',-k'}(n_k + 1/2)(n_{k'} + 1/2) \\ & - \frac{18}{\hbar} \sum_{k,k',k''} \left\{ |\Phi_{kk'k''}|^2 \left[ \frac{n_k n_{k'} + n_k + 1/3}{(\omega_k + \omega_{k'} + \omega_{k''})_p} + \frac{2n_k n_{k''} - n_k n_{k'} + n_{k''}}{(\omega_k + \omega_{k'} - \omega_{k''})_p} \right] \\ & + & 2\Phi_{k,-k,k''} \Phi_{k',-k',k''} \frac{n_k n_{k'} + n_k + 1/4}{(\omega_{k''})_p} \right\} \end{split}$$

which can be used to calculate further termodynamic functions (entropy, heat capacity, forces, stresses, etc.). The formulas for harmonic parts are reasonably compact, e.g.  $C_V$ :

$$C_V = \left(\frac{\partial U}{\partial T}\right)_V = -T\frac{\partial^2 F}{\partial T^2} = \sum_{\mathbf{q}\nu} k_{\rm B} \left(\frac{\hbar\omega(\mathbf{q}\nu)}{k_{\rm B}T}\right)^2 \frac{\exp(\hbar\omega(\mathbf{q}\nu)/k_{\rm B}T)}{[\exp(\hbar\omega(\mathbf{q}\nu)/k_{\rm B}T)-1]^2}$$

Expectation value of the squared atomic displacement:

$$\left\langle |u^{\alpha}(jl,t)|^2 \right\rangle = \frac{\hbar}{2Nm_j} \sum_{\mathbf{q},\nu} \omega_{\nu}(\mathbf{q})^{-1} (1 + 2n_{\nu}(\mathbf{q},T)) |e^{\alpha}_{\nu}(j,\mathbf{q})|^2,$$

### Thermal expansion

$$F_{H} = \sum_{k} \left[ \frac{1}{2} \hbar \omega_{k} + k_{B} T \ln \left( 1 - e^{-\hbar \omega_{k}/k_{B}T} \right) \right] \label{eq:fh}$$

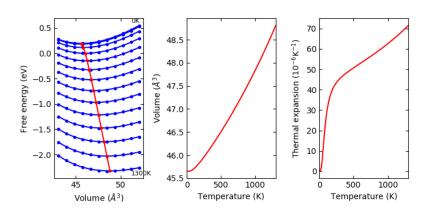
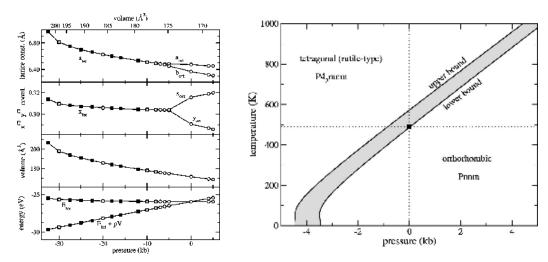


Figure 1: Figure from documentation of Phonopy package

#### Phase transitions



J. Łażewski, P. T. Jochym, P. Piekarz, and K. Parlinski; Phys. Rev. B 70 (2004) 104109