

Analysis of Equations of State for Neutron Star Modeling

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 - ▶ QHD-I

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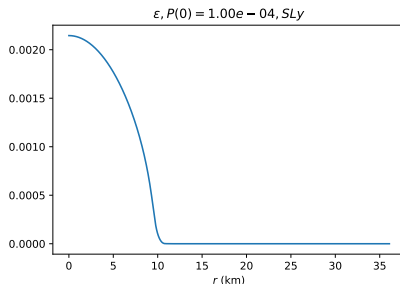
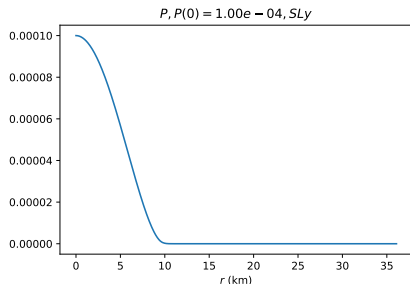


Figure: Example static solution for $P_0 = 10^{-4}$ for an EoS called “SLy.”

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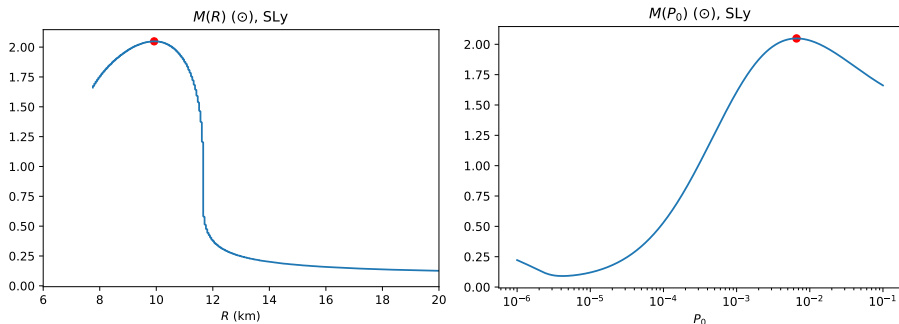
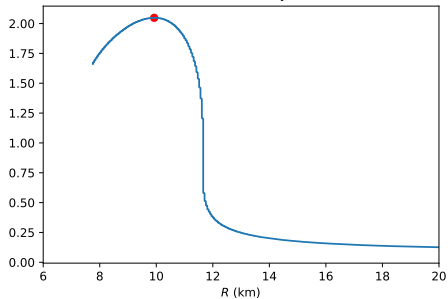


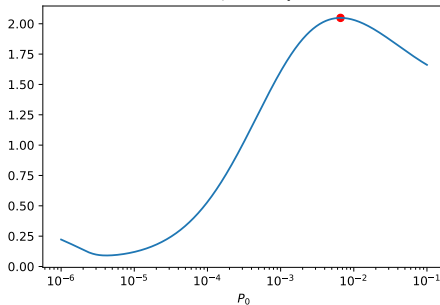
Figure: Example curves for EoS “SLy.” $1 \odot = 1.989 \times 10^{30}$ kg (solar mass)

Critical Values of P , R , and M

$M(R)$ (\odot), SLy

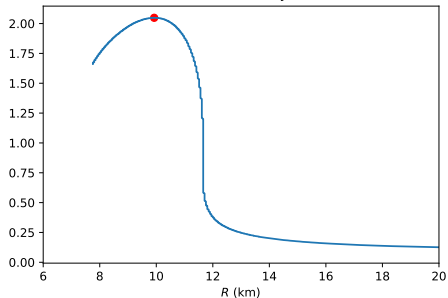


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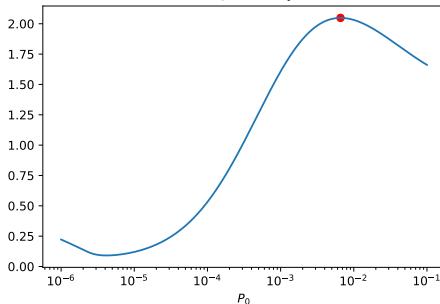


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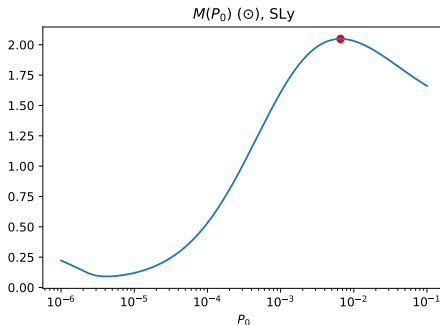
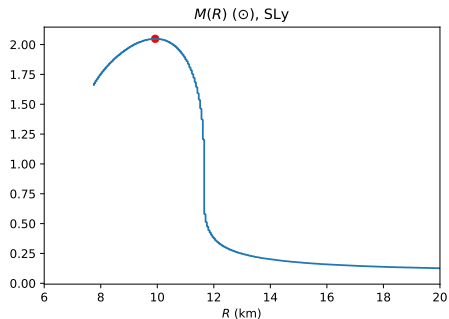


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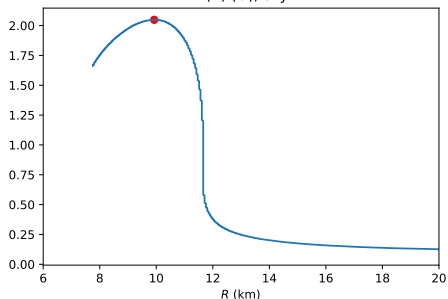
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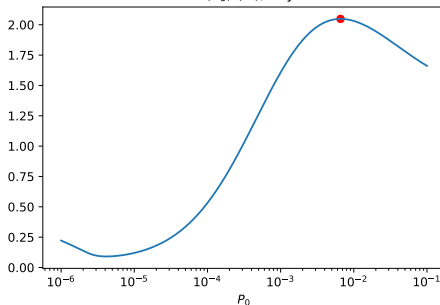
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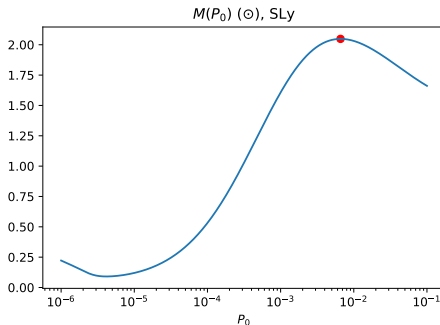
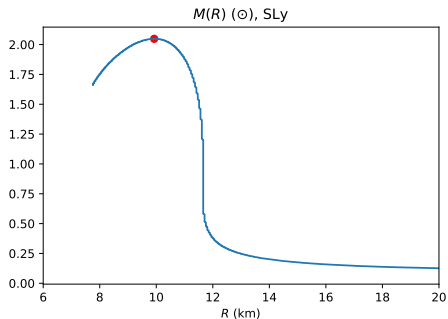
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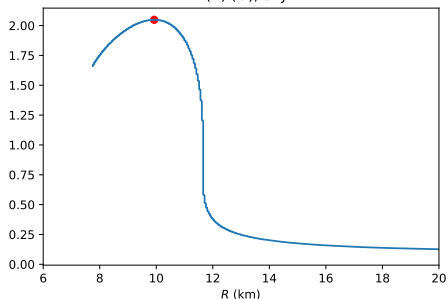


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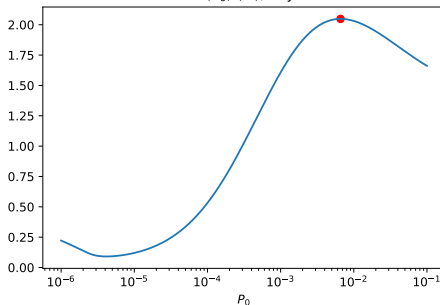
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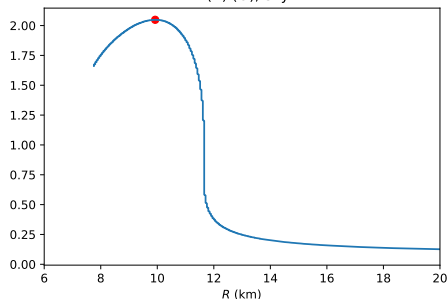


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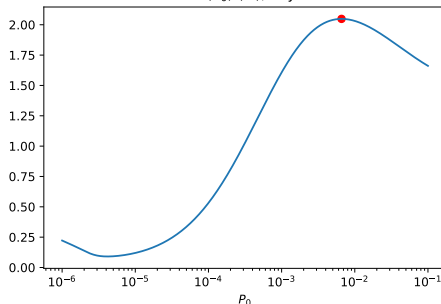
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- For SLy, $M_{\max} = 2.05 \odot$, $R_{\max} = 9.93$ km, and $P_{\text{crit}} = 6.59 \times 10^{-3}$.

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From \mathcal{L} , we can determine ϵ and P , the EoS we desire.

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$$\partial_\nu \left(\frac{\partial \mathcal{L}}{\partial(\partial_\nu \varphi_\alpha)} \right) - \frac{\partial \mathcal{L}}{\partial \varphi_\alpha} = 0, \quad T^{\mu\nu} = \frac{\partial \mathcal{L}}{\partial(\partial_\mu \varphi_\alpha)} \partial^\nu \varphi_\alpha - \mathcal{L} \eta^{\mu\nu}.$$

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$$\phi \rightarrow \langle \phi \rangle = \phi_0, \quad V_\mu \rightarrow \langle V_\mu \rangle = V_0,$$

where ϕ_0 and V_0 are constants. This allows us to simplify \mathcal{L} considerably:

$$\mathcal{L}_{\text{RMF}} = \bar{\psi}[i\gamma_\mu \partial^\mu - g_v \gamma_0 V_0 - (M - g_s \phi_0)]\psi - \frac{1}{2}m_s^2 \phi_0^2 + \frac{1}{2}m_\omega^2 V_0^2,$$

Determining ϕ_0 , V_0 , ϵ , and P : $\varphi_\alpha \in \{\phi_0, V_0, \psi\}$

$$\partial_\nu \left(\frac{\partial \mathcal{L}}{\partial(\partial_\nu \varphi_\alpha)} \right) - \frac{\partial \mathcal{L}}{\partial \varphi_\alpha} = 0, \quad T^{\mu\nu} = \frac{\partial \mathcal{L}}{\partial(\partial_\mu \varphi_\alpha)} \partial^\nu \varphi_\alpha - \mathcal{L} \eta^{\mu\nu}.$$

$$\epsilon = \langle T^{00} \rangle, \quad P = \langle T^{ii} \rangle$$

Resulting Equations

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$$\phi_0 = \frac{g_\phi}{m_\phi^2} \frac{1}{\pi^2} \int_0^{k_f} dk \frac{k^2 m^*}{\sqrt{k^2 + m^{*2}}},$$

$$V_0 = \frac{g_v}{m_v^2} \frac{k_f^3}{3\pi^2},$$

$$\epsilon = \frac{1}{2} m_\phi^2 \phi_0^2 + \frac{1}{2} m_v^2 V_0^2 + \frac{1}{\pi^2} \int_0^{k_f} dk k^2 \sqrt{k^2 + m^{*2}},$$

$$P = -\frac{1}{2} m_\phi^2 \phi_0^2 + \frac{1}{2} m_v^2 V_0^2 + \frac{1}{3} \left(\frac{1}{\pi^2} \int_0^{k_f} dk \frac{k^4}{\sqrt{k^2 + m^{*2}}} \right).$$

where $m^* = (M - g_\phi \phi)$, the *reduced mass*.

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$M(R)$ and $M(P_0)$ Curves for QHD-I

We use the tabulated values of P and ϵ to create an *interpolated function*

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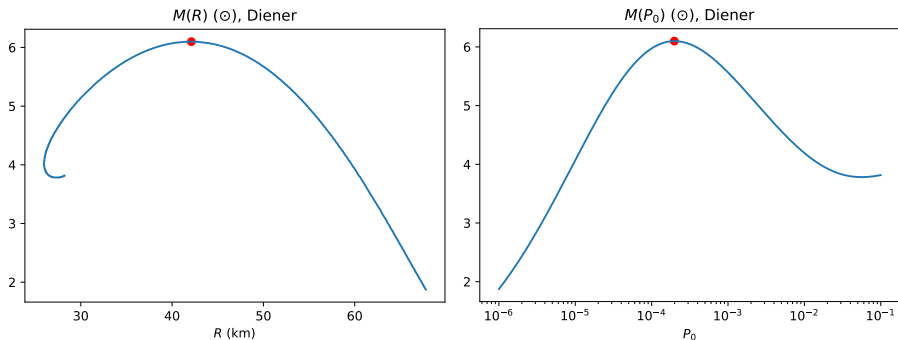


Figure: $M(R)$ and $M(P_0)$ curves for QHD-I EoS.

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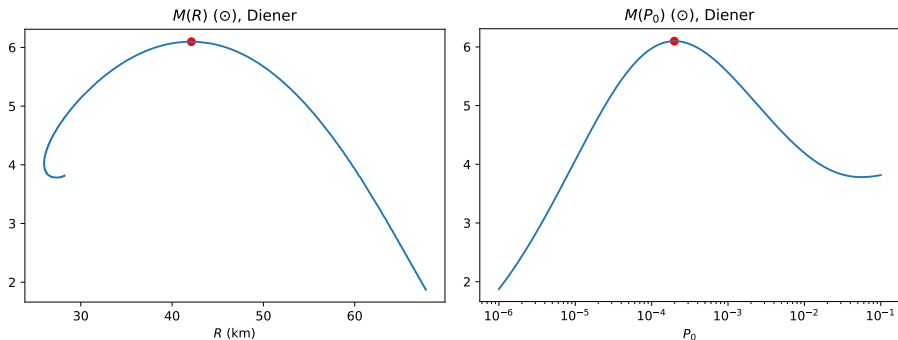


Figure: $M(R)$ and $M(P_0)$ curves for QHD-I EoS.

These curves give

$$M_{\max} = 6.1 \odot, \quad R_{\max} = 42.1 \text{ km}, \quad P_{\text{crit}} = 1.98 \times 10^{-4}.$$

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- We use the TOV equations to predict the maximum mass and radius that a given EoS will produce
- We use the QHD-I parameter set and RMF simplifications to solve a system of equations and generate an equation of state

Thanks

- Prof. Ben Kain

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- Holy Cross Physics Department
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