Notes on Introduction to Condensed Matter Physics

Xupeng Yang

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Basic Introduction

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1	Basic Introduction	
	1. Condensed Matter: $\sim 10^{23}\ /cm^3$	
	2. Major Study: Electrons, Phonons, The interactions between	
	3. Drives: New materials & New technologies	
2	Conventional Metal Physics: Electrons and Phonons	
	1. Basic Properties of Normal Metals: • Ductile • Excellent electrical conductor • Excellent then	rmal

2. Drude Free Electron Model:

• Assumptions: • Free electrons (Ignore interaction with lattice) • Independent electrons (Ignore interactions between electrons • Electrons were treated as independent classical particles

conductor • Most are weak paramagnet, some ferromagnet • Opaque

• At low T: ρ increases with T $\chi \sim Const.$ $c_V \propto T$

- Maxwell-Boltzmann distribution
- Successes: Electrical conductivity and thermal conductivity Wiedemann-Franz law (by luck!)
 - The Hall effect and magnetoresistance AC conductivity and optical properties of metals

- Problems: Heat capacity puzzle: $c_V = \frac{3}{2}nk_B = Const.$ The susceptibility puzzle: χ does not change with temperature, non-Curie like $\chi \sim 1/T$
- 3. The Sommerfeld Model: Free Electron Gas + Schrodinger Equation + Fermi Statistics
 - The Fermi Surface: In 3D $E_F \sim n^{2/3}$ Typical Values: $E_F \sim 7~eV,~k_F \sim 1 \times 10^8~cm^{-1},~T_F \sim 8 \times 10^4~K,~v_F \sim 2 \times 10^8~cm/s$ DOS In 3D: $D(\varepsilon) \sim \varepsilon^{1/2}$ Linear T Heat Capacity at Low T
 - \bullet Pauli Paramagnetism (at low $T)\colon \chi=\mu_B^2D(E_F)=const.$
 - Successes: Realized the importance of Fermi distribution
 Established the k-space language for electrons
 Introduced Fermiology
 Resolved the Pauli susceptability puzzle
 Resolved the heat capacity puzzle
 Resolved the thermopower puzzle
 Explained the Wiedemann-Franz Law
- 4. Landau's Fermi Liquid Theory: •Quasi-particle: Same charge, spin, momentum as non-interacting electron Adiabatic Continuity, Only valid at low T and low energy Qualitatively explain the susceptibility, heat capacity Only require a Fermi sea Entropy, distribution function unchanged Energy modified by the effective mass & the Fermi interaction function Low energy excitation like single particle $\varepsilon = \frac{\hbar^2 k^2}{2m^*}$ $c_V = \frac{1}{3} \frac{m^* k_F}{\hbar^2} k_B^2 T$ $\chi = \frac{m^* k_F}{\hbar^2 \pi^2} \frac{1}{1 + F_0^a} \mu_B^2$ Landau parameter: m^*, F_0^a Wilson ration: $R_W = \frac{\pi^2 k_B^2}{3\mu_B^2} \frac{\chi}{\gamma}$ For non-interacting electron gas $R_W = 1$ T^2 Law (Experimental Signature): qp-qp scattering Scattering rate $\frac{1}{\tau} \sim k_B T \cdot k_B T \propto T^2$ But electron seems to be a wrong place to start for many novel phenomena Spinon, holon, fractional charge: Collective mode looks like a fraction of an electron
- 5. Bloch Theory: Bloch Theorem: For periodic potential, $\psi_{nk}(\mathbf{r}) = \mathrm{e}^{\mathrm{i}\mathbf{k}\cdot\mathbf{r}}u_{nk}(\mathbf{r})$, where $u_{nk}(\mathbf{r}) = u_{nk}(\mathbf{r}+\mathbf{r})$ Momentum is no longer a good quantum number Band index n
 - NFE model: Perturbation to the free electron plane wave states (maximum mixing) (highly delocalized) First order: $\varepsilon_k^1 = \overline{V}$ Second order (non-degenerate): $\varepsilon_k^2 = \sum_{g \neq 0} \frac{|V_g|^2}{\varepsilon_K^0 \varepsilon_{k-g}^0}$ Second order (degenerate): $\varepsilon_k^2 = \pm |V_n|$ Energy gap: Level repulsion, Explains metal or insulator
 - +2 Metals: Band overlap
 - Tight Binding Model: (nearly localized) such as transition metal and rare earth metal with partially filled d and f orbitals. $\psi_k(\mathbf{r}) = \sum_{\mathbf{R}} \mathrm{e}^{\mathrm{i}\mathbf{k}\cdot\mathbf{R}} \psi_a(\mathbf{r} \mathbf{R})$ Overlap integral: $t_{\mathbf{R}} = \int \psi_a^*(\mathbf{r} + \mathbf{R})(\Delta V)\psi_a(\mathbf{r})\mathrm{d}\mathbf{r} \to \mathrm{Nearest}$ neighbor approximation Useful starting point
- 6. Lattice Vibrations: Harmonic Approximation: $V(a + \delta x) = V_0 + \frac{1}{2}\beta(\delta x)^2$ $\varepsilon = \sum_k (n_k + \frac{1}{2})\hbar\omega$ • Phonons: The quantum of the lattice vibration, $n_k = \frac{1}{e^{\hbar\omega/kT}-1}$ • Mono-atomic 1D Chain: $\omega = 2\sqrt{\frac{\beta}{m}}\left|\sin\left(\frac{aq}{2}\right)\right|$ • Di-atomic 1D Chain: Acoustic & Optical phonon • Phonon specific heat: Debye model: • Assume linear dispersion • Define a cutoff in the integral: Debye frequency

- $ullet T o 0 \ c_V \sim T^3 \quad \Leftarrow ext{Blackbody radiation}$
- 7. Specific Heat: Directly related to internal energy To extract important microscopic parameters
 - To study phase transition Calorimetry: What, How, Better resolution and accuracy, Flecibility
 - Adiabatic Nernst Calorimeter: Slow, Heat leak problem, Need big sample
 - Relaxation time calorimeter: $\Delta T = \Delta T_0 \mathrm{e}^{-t/\tau}$, where $\tau = c_V l/\kappa S$ (with addenda) Advantage: Accurate, Fast, Microgram crystals, Small, Work in extreme conditions Disadvantage: The addenda
 - Membrane calorimeter: Nano-gram crystals, Measure in-situ evaporated thin films, Extreme conditions
 - Heater: Resistance stable with T Thermometer: Resistance has Linear relationship with T
- 8. Anharmonic Potential: Universal $\Leftarrow V = 0$ when $r \to \infty$ Phonon-phonon interaction: No longer independent excitations \Rightarrow Phonon heat conduction (low T, high T) Thermal Expansion:
 - $\alpha = \frac{1}{l} \frac{\partial^2 l}{\partial T \partial p} = \frac{1}{3V} \frac{\partial^2 V}{\partial T \partial p}$ Provide similar information as specific heat Bad in engineering
 - XRD \rightarrow Measure l or V, High resolution, Hard to use in extreme physical conditions
 - Capacitive dilatometer: High resolution (capacitance bridge) (0.01 Å), ultra low T and large B (compact design)
 - \bullet Negative thermal expansion: $ZrW_2O_8 \ \to {\rm Rigid}$ Unit Modes
- 9. Main Frame: Landau Fermi Liquid Theory + Band Theory

3 Transport

- 1. Basic Notions: Movement of Particles or Quantities Non-equilibrium steady state $J = L \cdot F$
 - Very informative and instructive, esp. on Novel materials and in Extreme conditions
 - Normally the first to be carried out
- 2. Fractional Quantum Hall Effect: \bullet Ultra low T, Super strong B, Very clean \bullet Strong electron correlations \bullet Most precise method to measure h
- 3. Cryogenic technology: Dilution fridge method: He-3 rich & He-3 poor phase at T < 0.87 K
 - He-3 diffuse, absorb heat Down to $\sim 10 \text{ mK}$
 - Superconducting magnet: up to 20 tesla ← critical field
 - Super High megnetic field: Florida-Bitter resistive magnet, Hybrid magnet

- 4. The Boltzmann transport equation: \bullet $\frac{\partial f_k}{\partial t}\Big|_{\text{diffusion}} + \frac{\partial f_k}{\partial t}\Big|_{\text{field}} + \frac{\partial f_k}{\partial t}\Big|_{\text{scattering}} = 0$ \bullet $\frac{\partial f_k}{\partial t}\Big|_{\text{diffusion}} = -\dot{\boldsymbol{r}} \cdot \nabla_{\boldsymbol{r}} f_k$ $\bullet \frac{\partial f_k}{\partial t}\Big|_{\text{field}} = -\dot{\boldsymbol{k}} \cdot \nabla_{\boldsymbol{k}} f_k \quad \bullet \frac{\partial f_k}{\partial t}\Big|_{\text{scattering}} = -\frac{f_k f_k^0}{\tau}$
- 5. Electrical transport: $\bullet J_e = \sigma E$ \bullet Measurements: Four-probe, Low frequency ac lock-in method
 - Drude model: $\sigma = \frac{ne^2\tau}{m}$ Semi classical: $\delta k = \frac{e\tau E}{\hbar}$ Only the surface of the FS changed !
 - Ignore the diffusion effect, Complexity of the FS
 - The Boltzmann transport equation: $\overleftrightarrow{\sigma} = \frac{1}{4\pi^3} \frac{e^2 \tau}{\hbar} \int \frac{v_k v_k \mathrm{d}S_F}{v_k}$

 - Matthiessen's rule: Different scattering mechanisms don't interfere each other $\Rightarrow \frac{1}{\tau} = \frac{1}{\tau_{imp}} + \frac{1}{\tau_{ph}} + \cdots$
 - \bullet Electron-electron scattering: $\frac{1}{\tau} \sim T^2$
 - \bullet Electron-lattice scattering: $\rho \sim T$, at high T $\rho \sim T^5$, at low T
 - Electron-impurity scattering: Roughly, Temperature-independent Residual resistivity: $\rho(T =$
 - 0) Residual resistivity ratio (RRR): $\frac{\rho_{300K}}{\rho_0}$ Higher the better
- 6. Thermal transport: $\bullet J_Q = \kappa(-\nabla T)$ \bullet Measurement: One-heater two-thermometer
 - Drude: $J_{Qe,x} = \frac{1}{2} n v_x [\varepsilon (T_{x-v\tau} T_{x+v\tau})] = \frac{1}{3} c_V v l$ $\kappa_e = \frac{\pi^2}{3} \frac{k_B^2}{e^2} T \sigma$ Phonon Thermal Conductivity: Good metals ~ 1%

4 Others

- 1. Topological Insulator: Single electron model Mainly spin-orbit interaction
- 2. Superfluidity: Liquid He-4 & He-3: Low boiling $T \Leftarrow$ Weak van der Waals force & low atomic mass He-4 (Boson) < 2.17~K (BEC), He-3 (Fermion) < 2.49~mK (BCS)

