

Dark Matter *Models*

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Michigan Cosmology Summer School 2023



Outline

Lecture 1: evidence and model building

- Evidence for dark matter
- Dark matter model building
 - { • What we know about DM
 - Pre-requisites for a DM model
- Mass bounds
- Landscape of models
- MOND

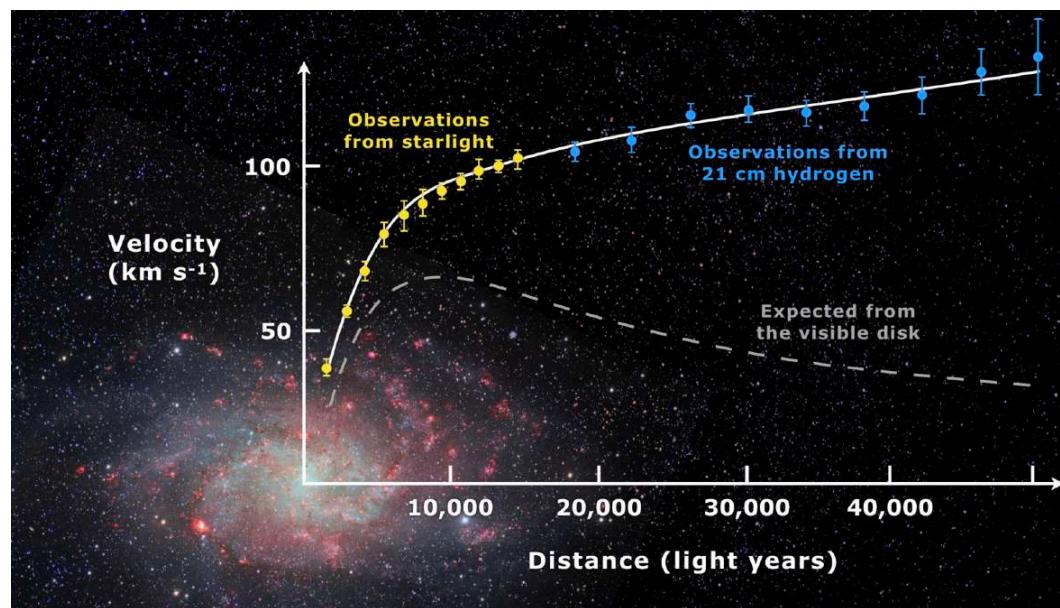
Lecture 2: DM models

- DM models
 - Particle DM: WIMPS
 - Macroscopic DM: MACHOS, Primordial BHs
 - Wave DM

Recap - lecture 1

Evidences for dark matter - properties

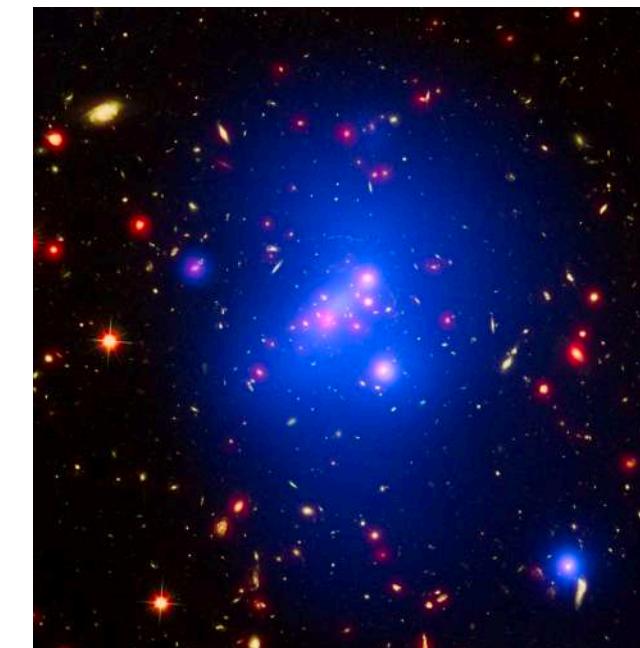
Galaxy rotation curves



Credit: Mario De Leo

- Mass fraction
- Distribution

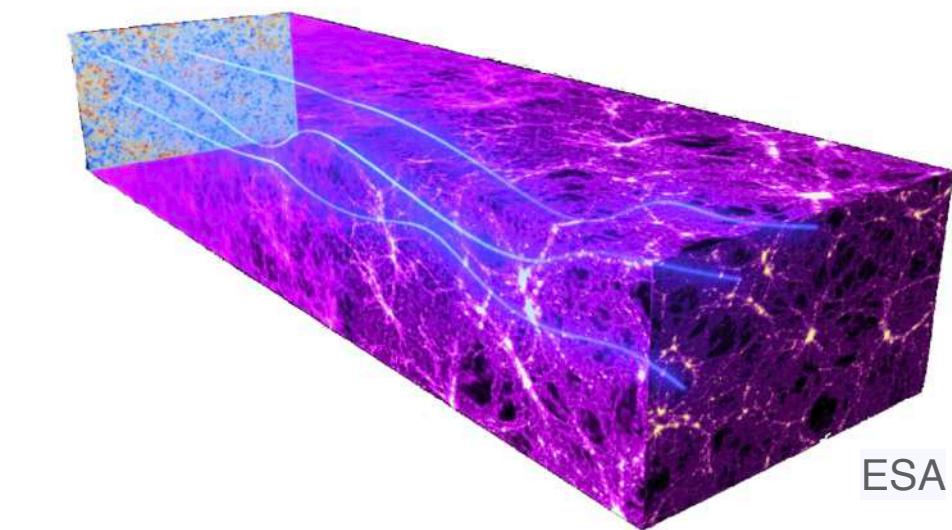
Clusters



CC BY 4.0

- Mass fraction
- Distribution

Lensing



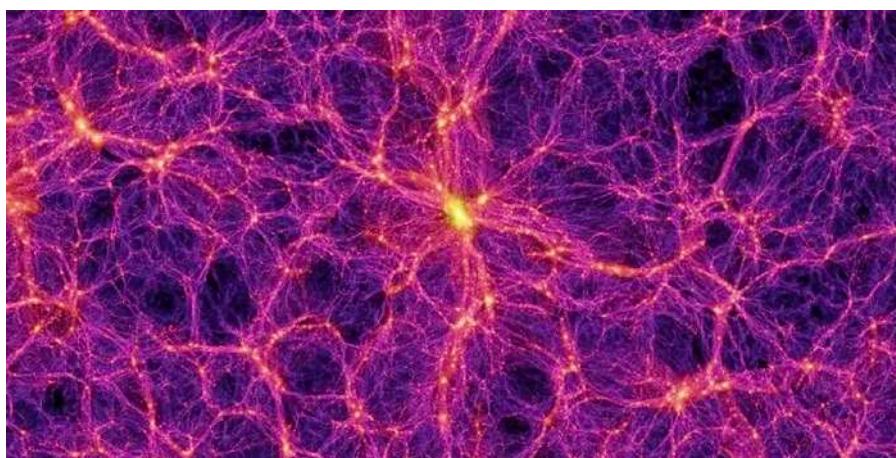
ESA

- Strong lensing
- Mass fraction
- Distribution

- Weak lensing
- Distribution
- Shape
- Structure

- Micro lensing
- Mass fraction
- Smoothness
- Structure

Large Scale Structure



Springel & others / Virgo Consortium

CMB/LSS

- Ratio of DM/collisional matter
- Thermal history

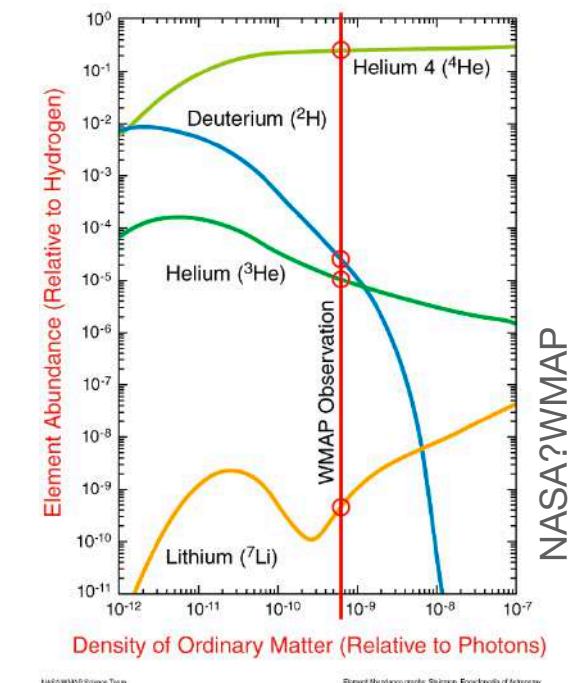
Cluster collision



NASA/CXC/CfA and NASA/STScI

- Distribution
- Separation from collisional matter
- Self-interaction

Big Bang Nucleosynthesis



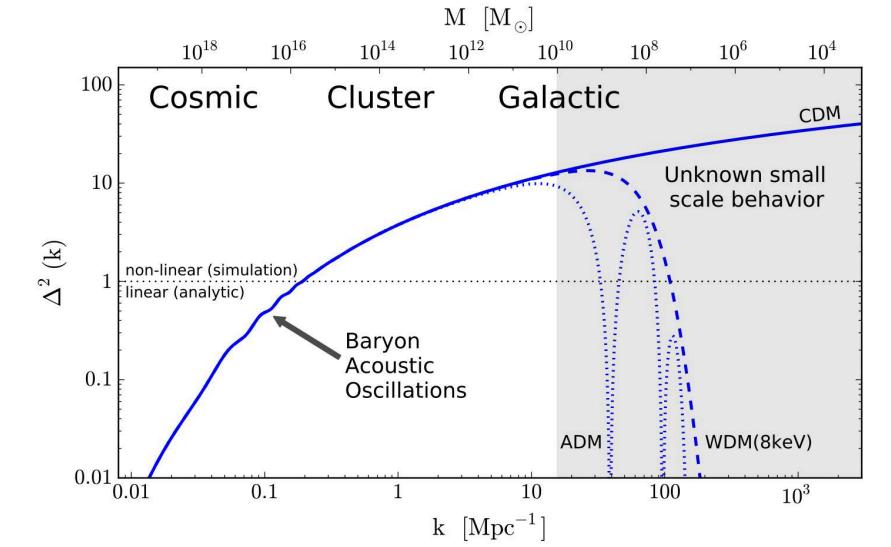
NASA/WMAP

- Amount of baryons

DM builder's guide

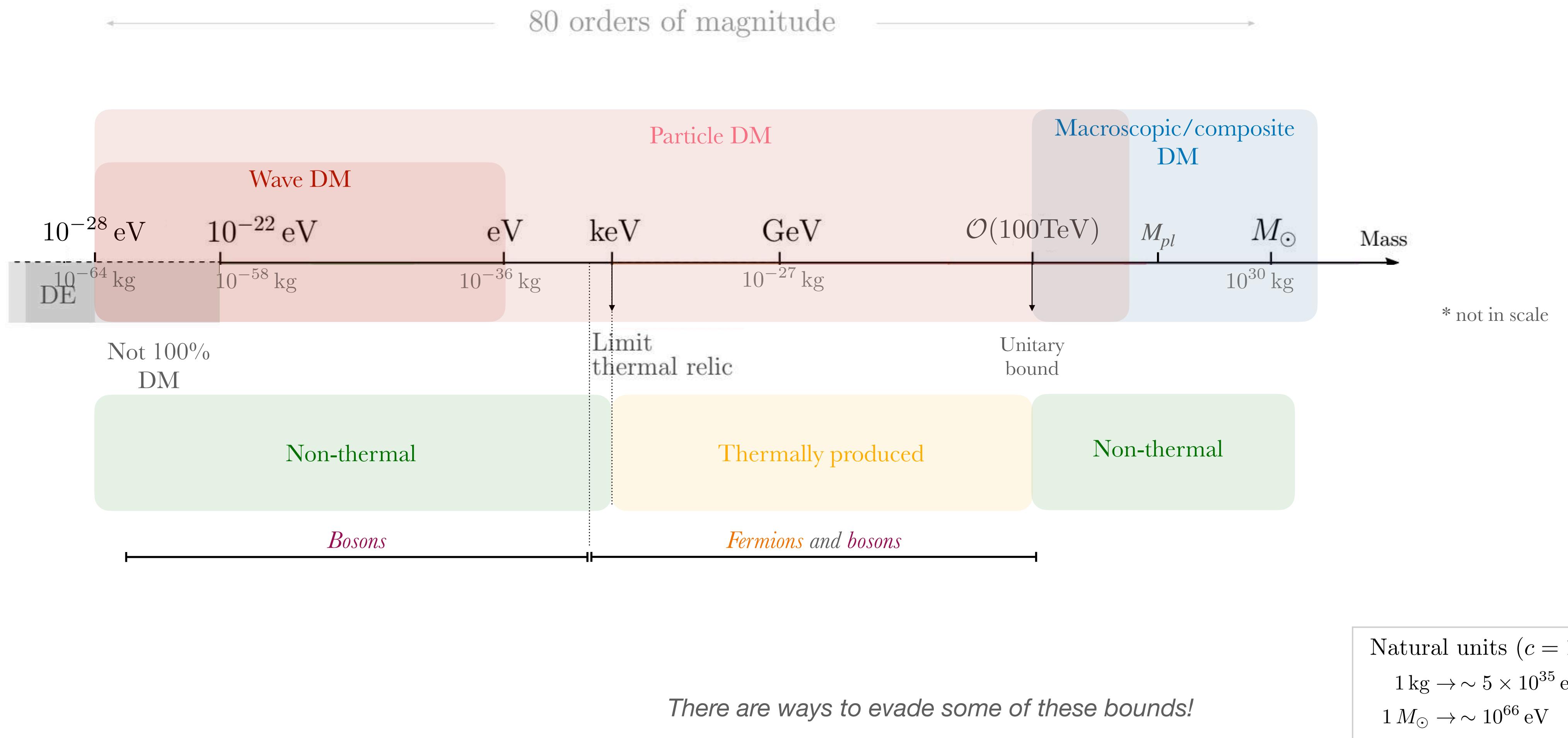
Pre-requisites for a dark matter candidate

- Cold or warm
 - Thermal candidate: $m_{dm} \geq \text{keV}$ Or produced cold by a non-thermal mechanism
 - Has to be non-relativistic at BBN
- Reproduce large and small scale distribution
 - Clusters like pressure-less fluid on large scales $k \lesssim 10 \text{ Mpc}^{-1}$
 - Clustering on scales smaller than $k \gtrsim 10 \text{ Mpc}^{-1}$ highly unconstrained
- Non-interacting or weakly interacting
 - (Dark, collisionless)
 - Can have a small electromagnetic interaction. Bound < **milicharge**
 - Can have a **self interaction**. Bounds: $\sigma/m_{dm} < 0.13 \text{ cm}^2/\text{g}$, $\sigma/m_{dm} < 0.35 \text{ cm}^2/\text{g}$
 - Can interact via the **weak force**
- Abundance $\Omega_m = 0.308 \pm 0.012$ (*Planck 2018*)
- Stable If it is a particle, it has to be stable with lifetime of DM should be much greater than the age of the universe



Mass scale of dark matter

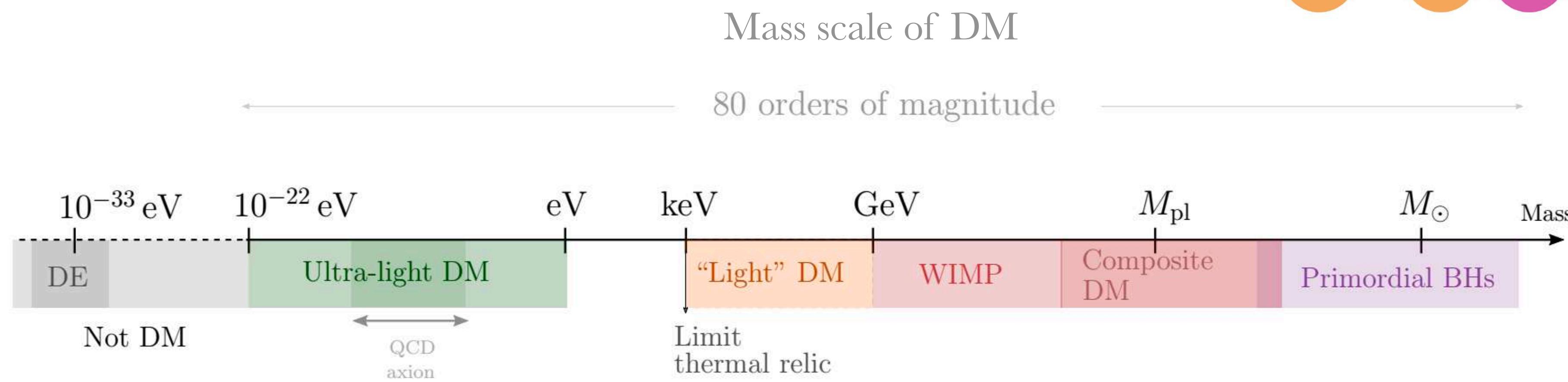
We can use observations of LSS and galaxies to put bounds in the “particle” physics properties, like mass and spin, of the DM candidate



Landscape of dark matter models

- What is DM? What is the nature of DM?

State of the “art”



MOND

Milgrom, 1983.

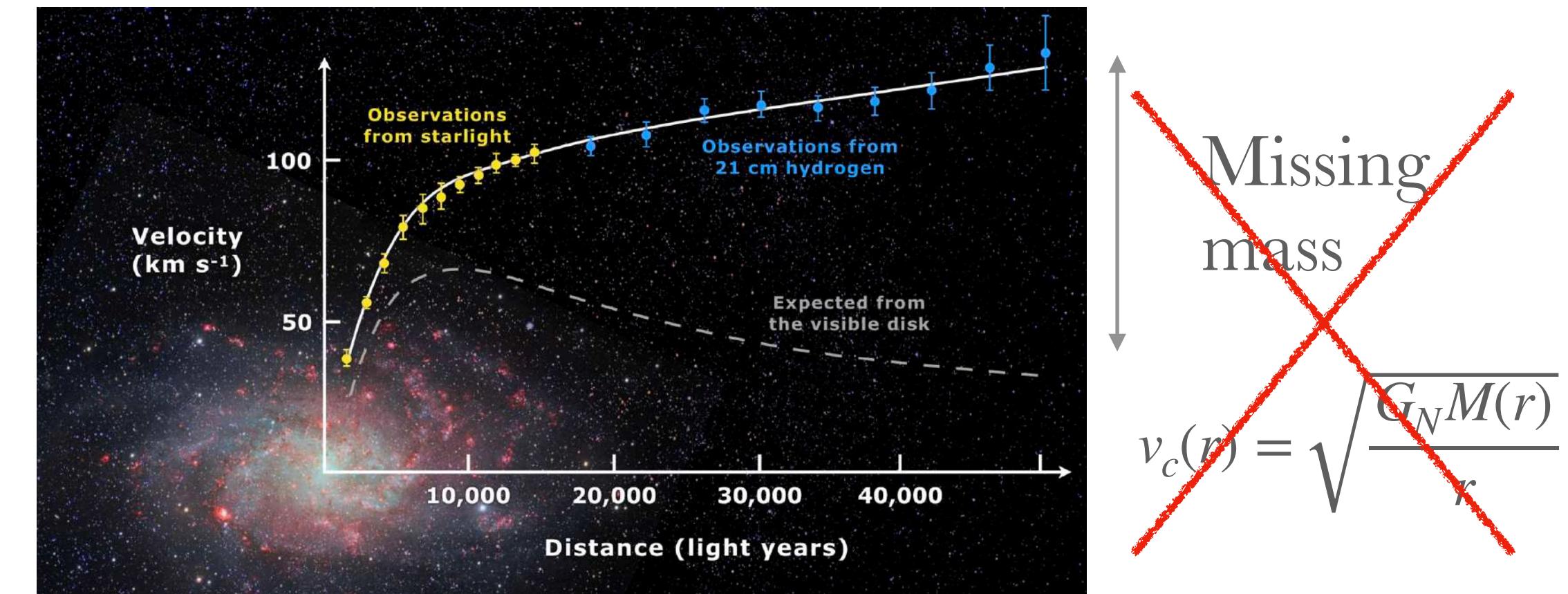
Modified Newtonian Dynamics

Empirical relation

$$a = \begin{cases} a_N^b, & a_N^b \gg a_0, \\ \sqrt{a_N^b a_0}, & a_N^b \ll a_0. \end{cases}$$

$$a_N^b = \frac{G_N M_b}{r^2}$$

$$a_0 \simeq 1.2 \times 10^{-8} \text{ cm/s}^2$$



Curiosity: Baryons lead the dynamics!

Works really well to: (1) Fit galaxy rotation curves; (2) Explain the scaling relations

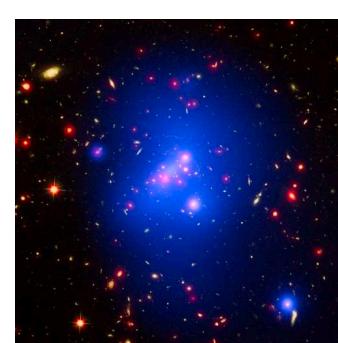
BUT: Modified theory of gravity

Milgrom, 1983.

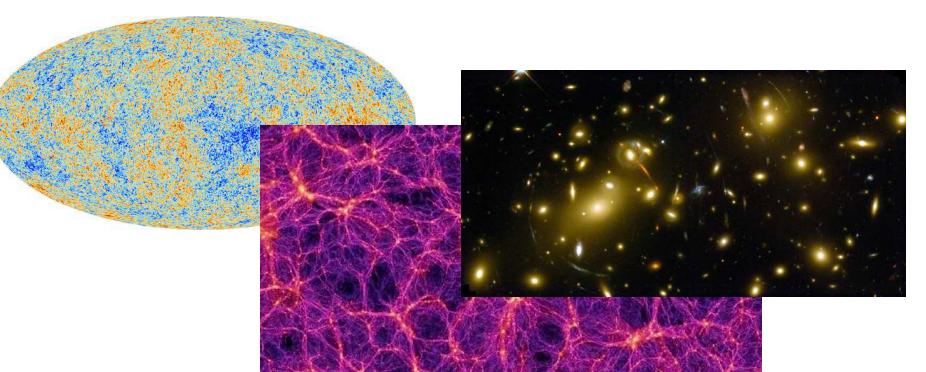
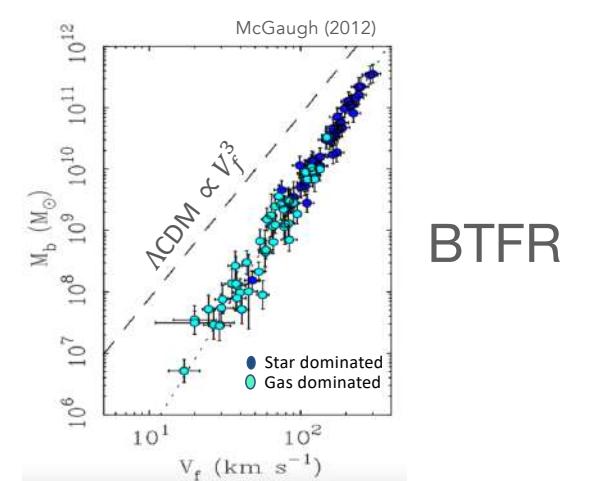
Relativistic extension: TeVeS, (BIMOND)

2020: “A new relativistic theory for Modified Newtonian Dynamics”, C. Skordis, T. Zlosnik → Agreement with CMB

~~MOND without DM~~



Clusters



Large scales

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Lecture 2: DM models

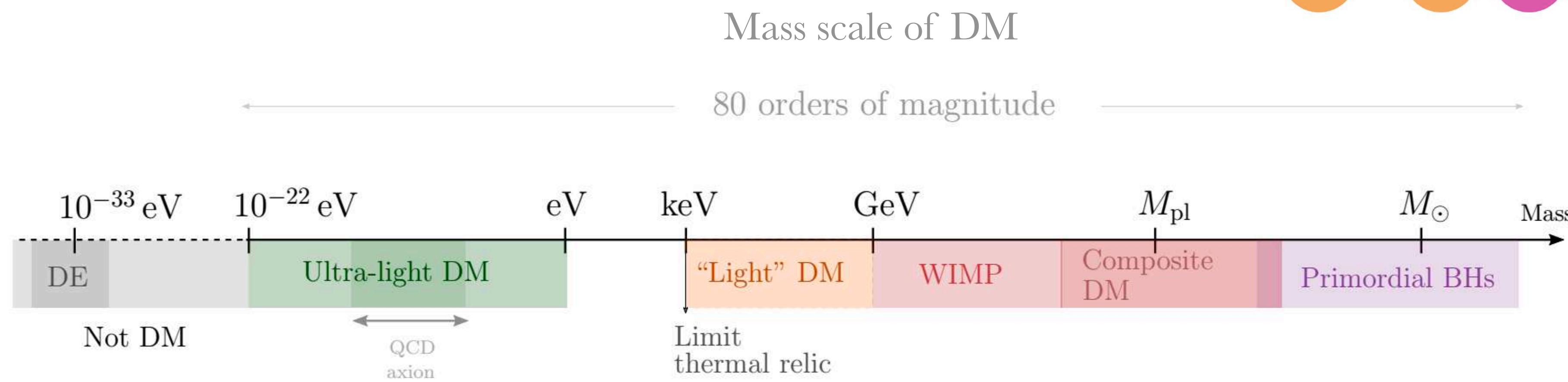
- DM models
 - Particle DM: WIMPS, WIMPzillas
 - Macroscopic DM: MACHOS, Primordial BHs
 - Wave DM

* Biased review of the DM models

Landscape of dark matter models

- What is DM? What is the nature of DM?

State of the “art”



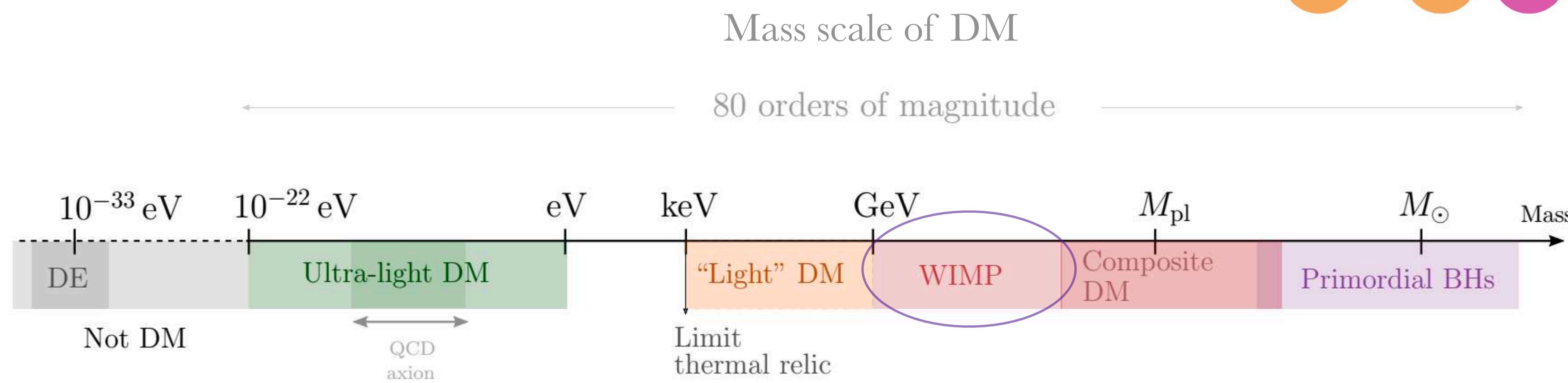
Particle DM



Landscape of dark matter models

- What is DM? What is the nature of DM?

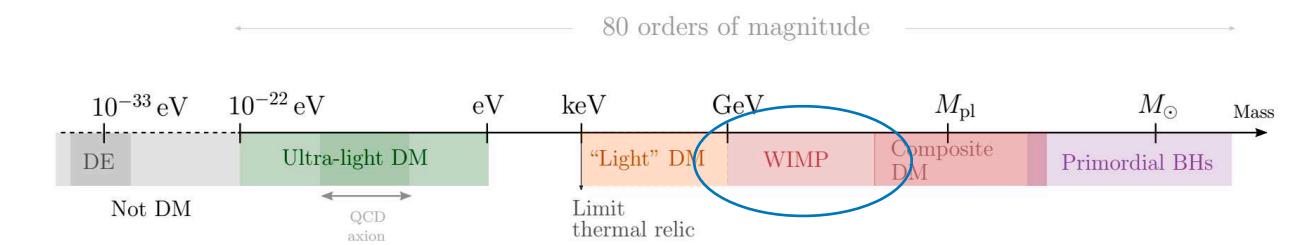
State of the “art”



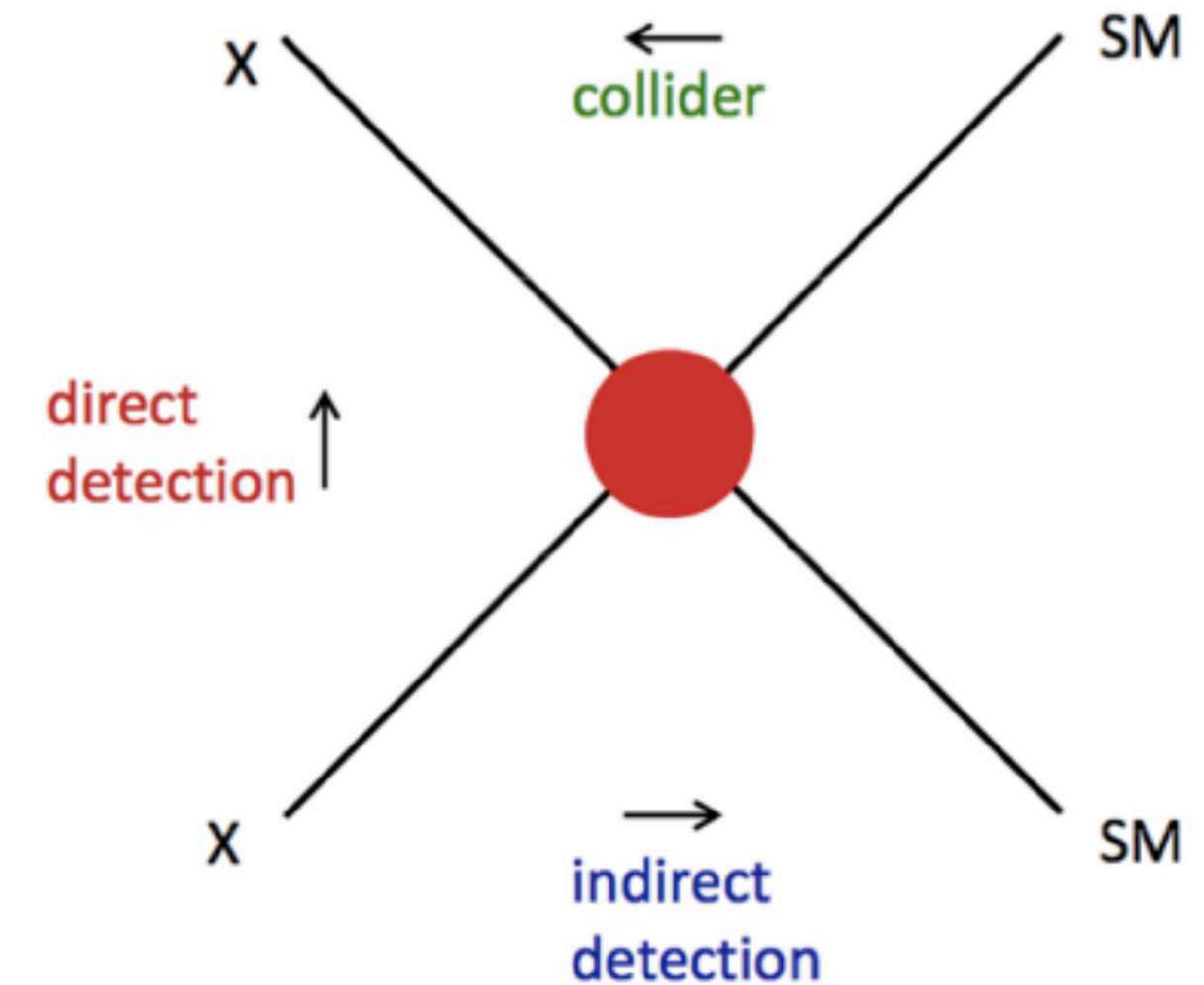
WIMPS

weakly interacting massive particles

WIMP - weakly interacting massive particle

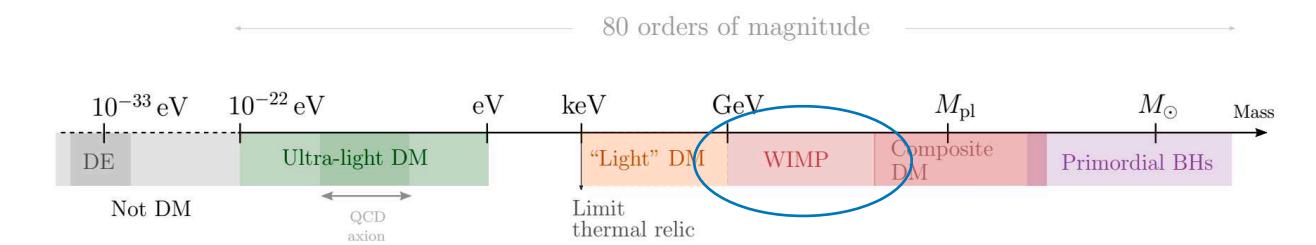


- Most accepted candidate
- (Beyond standard model) massive particle
- "WIMP miracle"
 - Thermally produced
 - $m \sim$ weak scale \rightarrow abundance of DM



Credito: F. Iocco

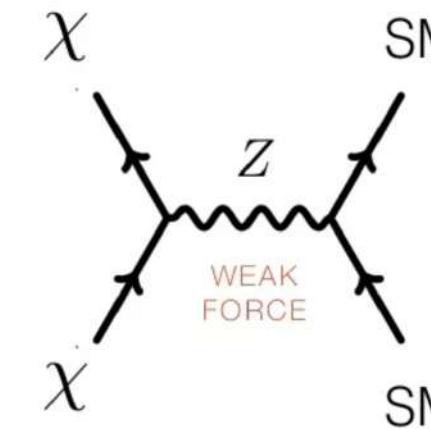
WIMP miracle



A thermal relic with cross-section \sim weak interaction would freeze out with the \sim density of the obs. DM today

$$\Omega_\chi h^2 \simeq 0.1 \left(\frac{3 \times 10^{-9} \text{ GeV}^{-2}}{\langle \sigma v \rangle} \right)$$

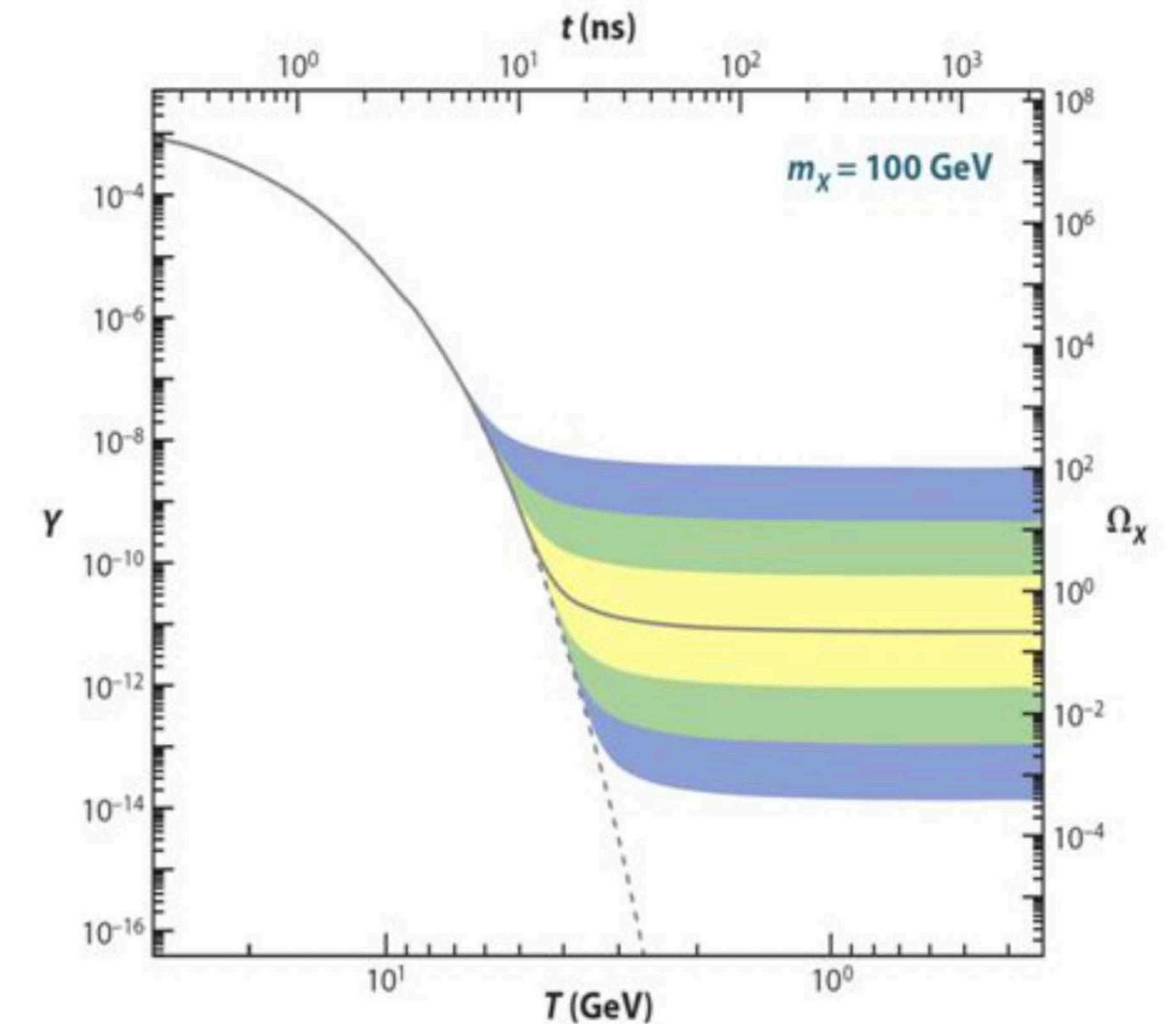
Annihilation cross-section



So we can have the correct abundance today:

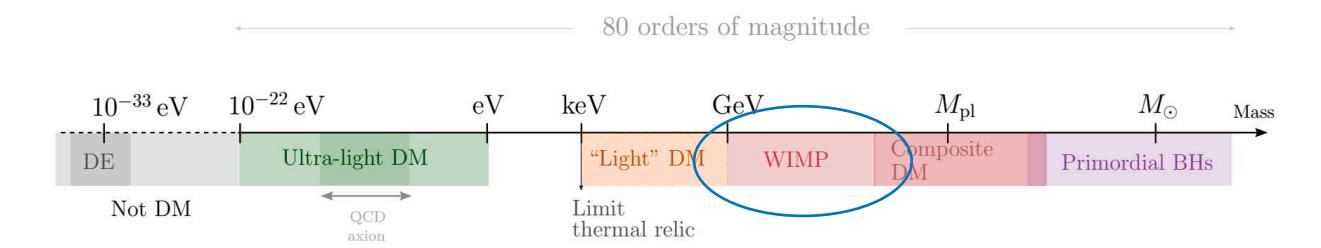
$$\langle \sigma v \rangle \simeq 3 \times 10^{-9} \text{ GeV}^{-2} \simeq G_F \times \frac{v_{wimp}}{c}$$

Expected cross-section for the weak interactions!

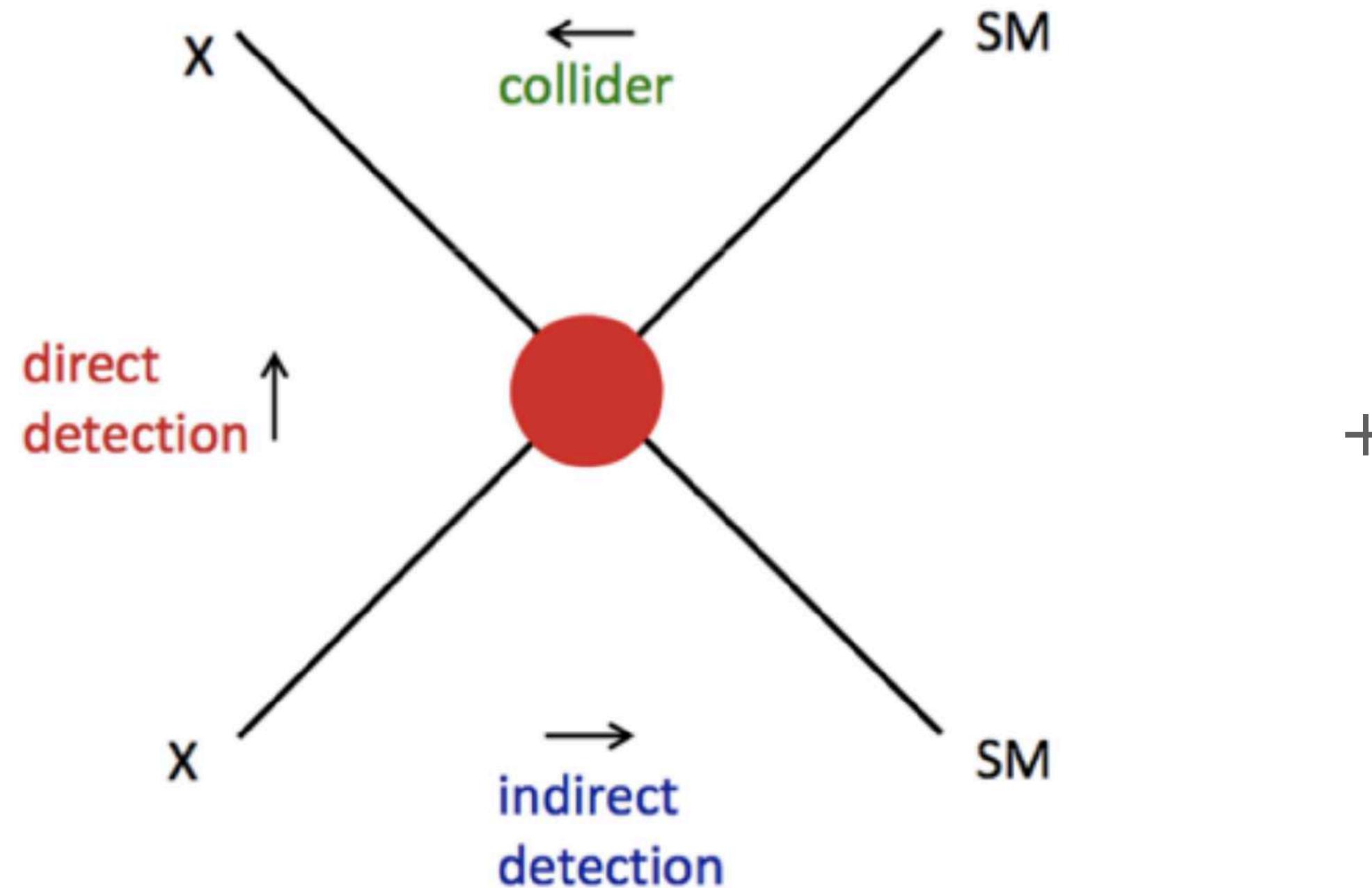


Therefore, if DM interacts through the weak force, we have the correct abundance of DM! Miracle!?

WIMP - weakly interacting massive particle



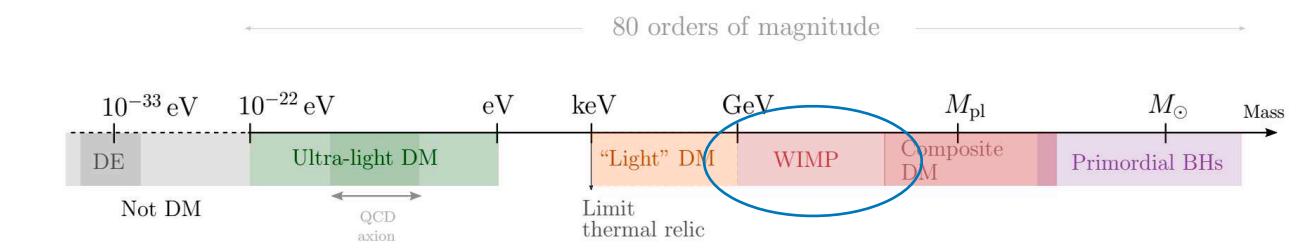
How can we measure it?



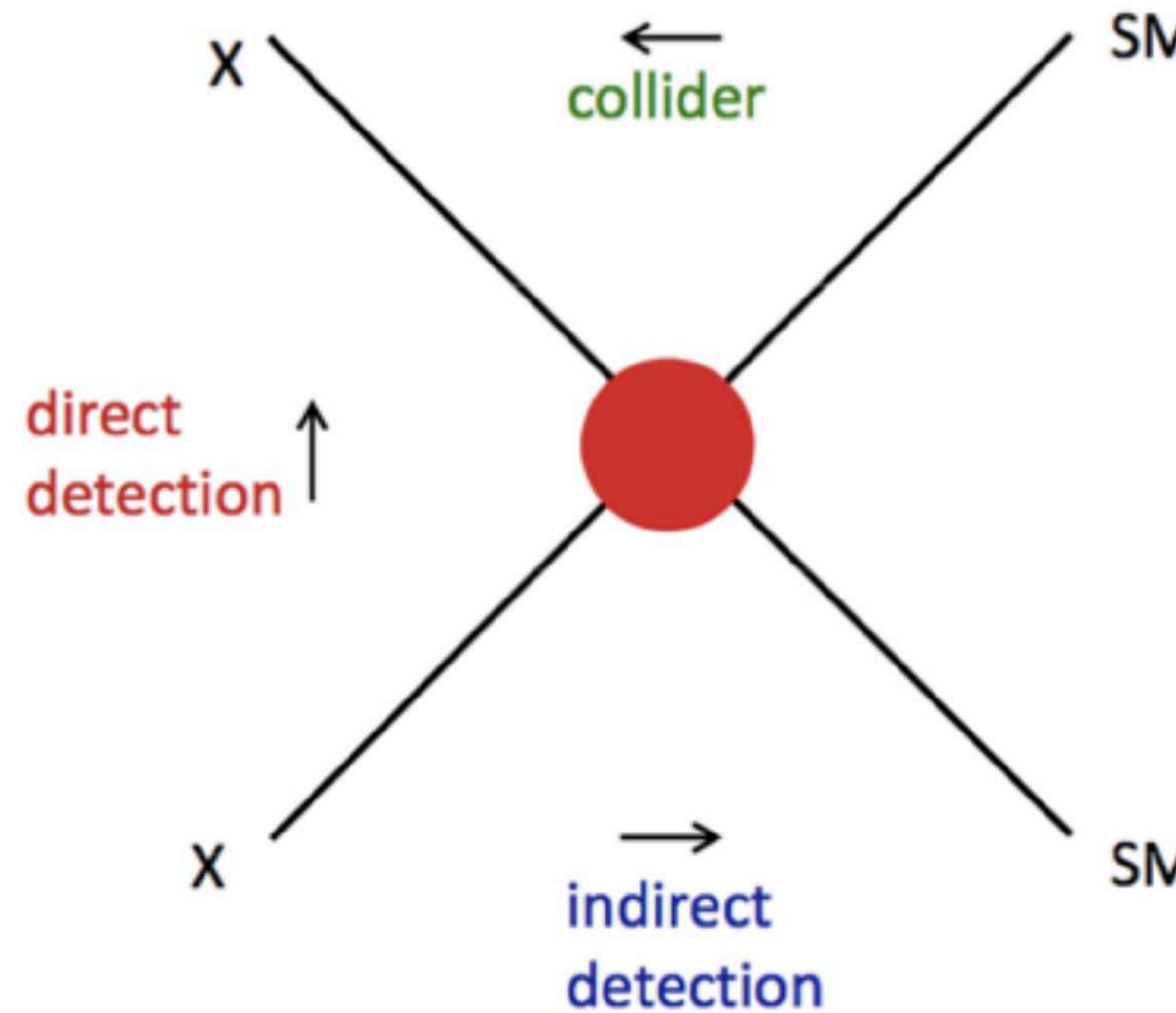
Gravitationally
Cosmological and astrophysical searches

Credito: F. Iocco

WIMP - weakly interacting massive particle



How can we measure it?



Credito: F. Iocco

+ Gravitationally
Cosmological and astrophysical searches

Direct detection:

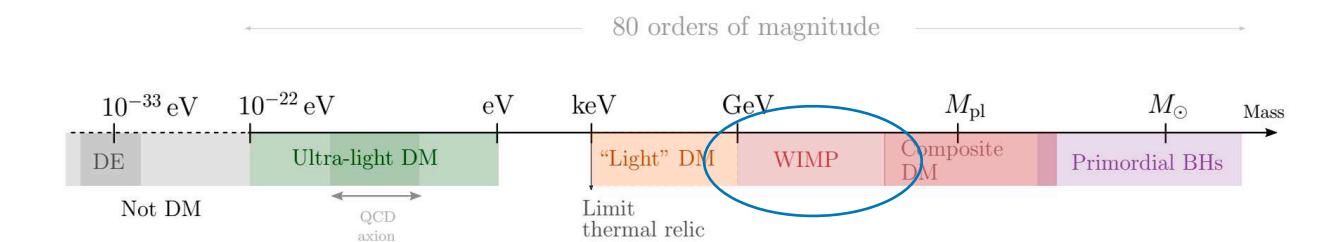
- DM scattering against nuclei, recoil

Indirect detection:

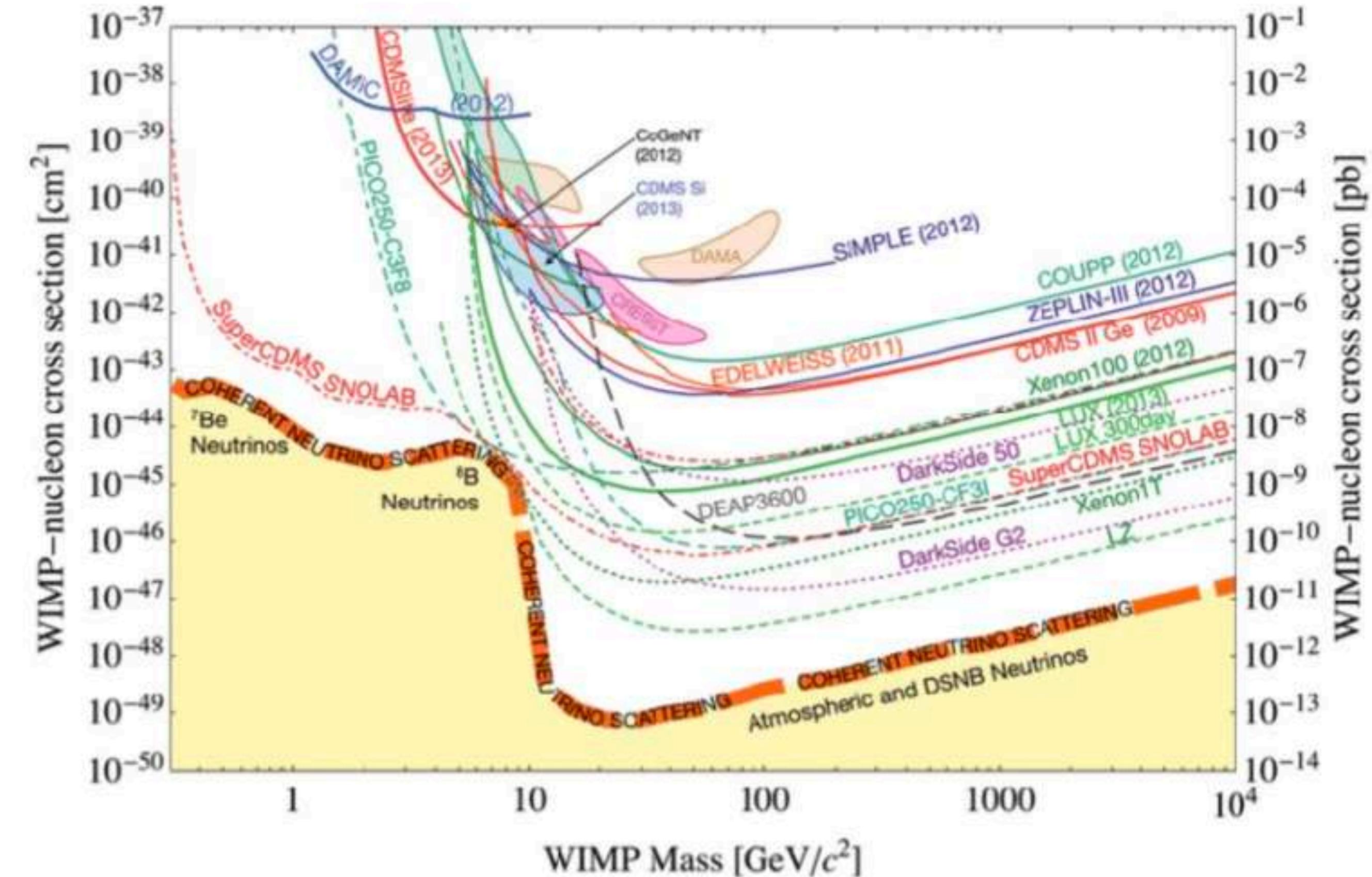
- Annihilation in astrophysical environment
- Observation of SM products of annihilation

Production at collider (LHC)

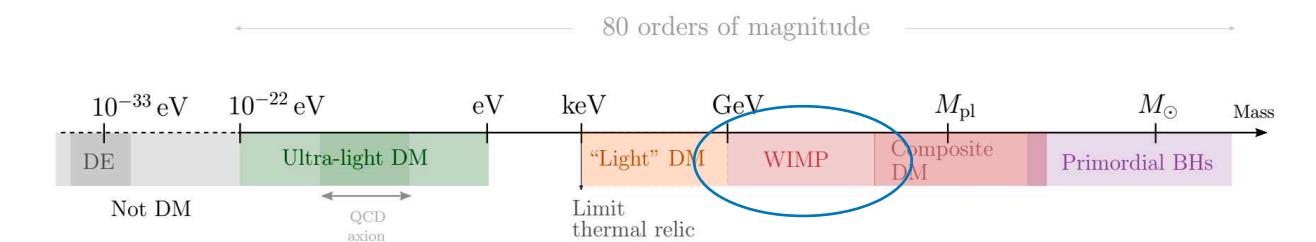
WIMP - weakly interacting massive particle



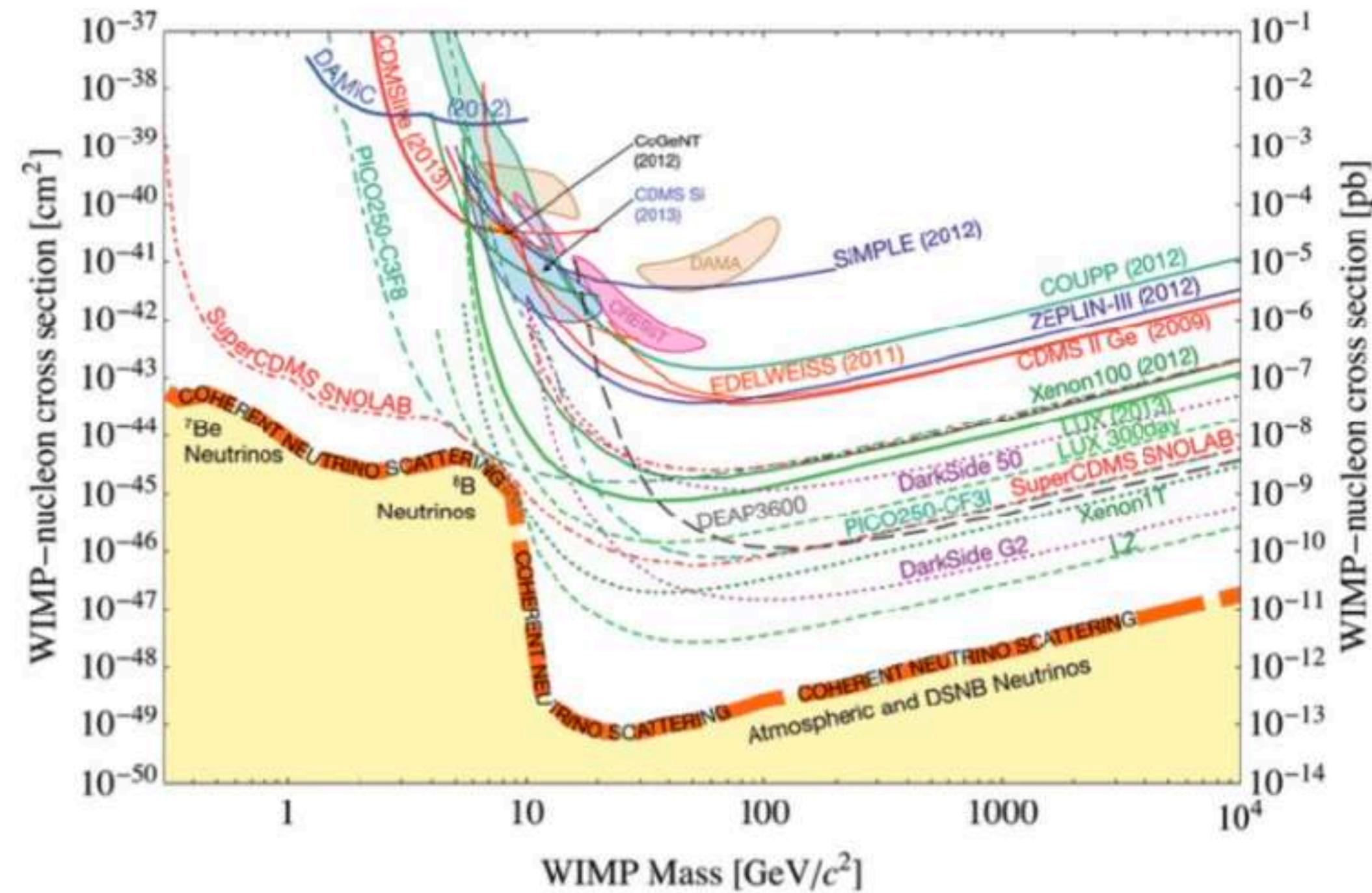
Bounds



WIMP - weakly interacting massive particle



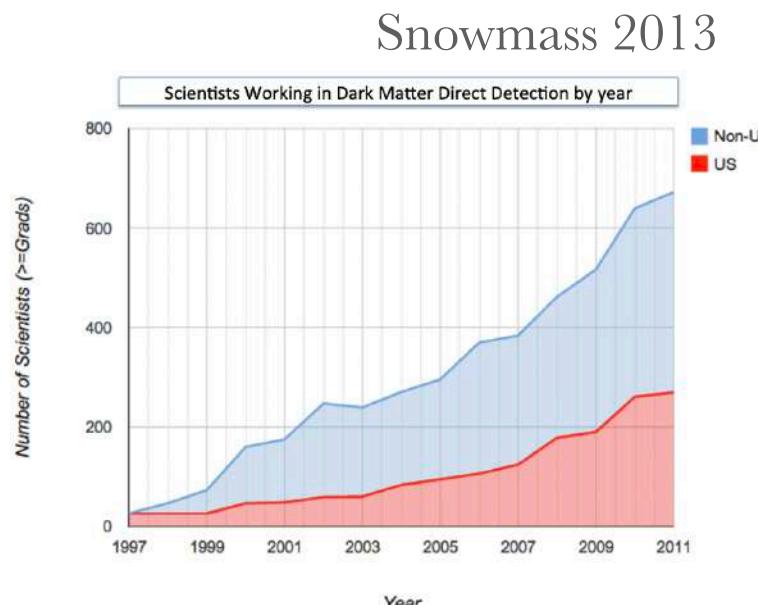
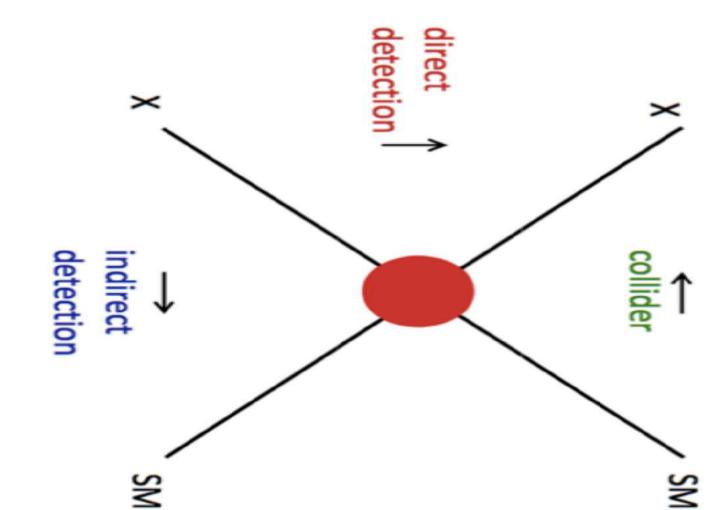
Bounds



HUGE experimental effort for discovery/bound

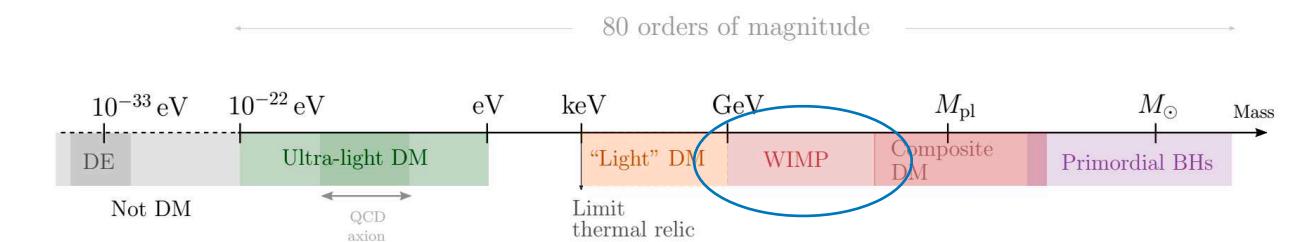
Still not detected!

Exclusion limits/bounds are of difficult interpretation for WIMPs



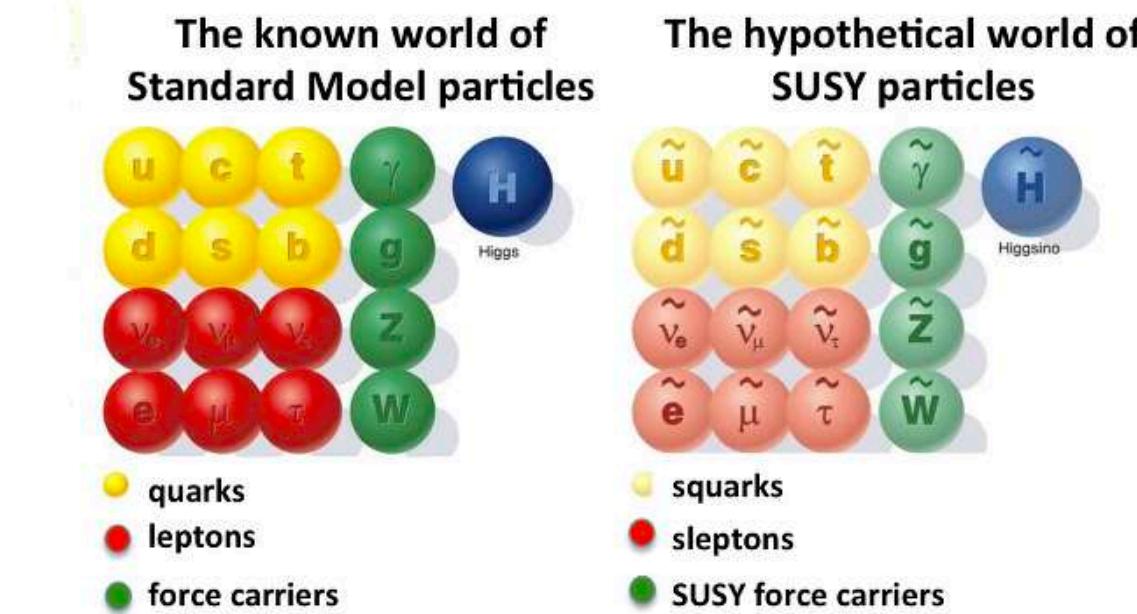
Supersymmetry

DM from supersymmetry

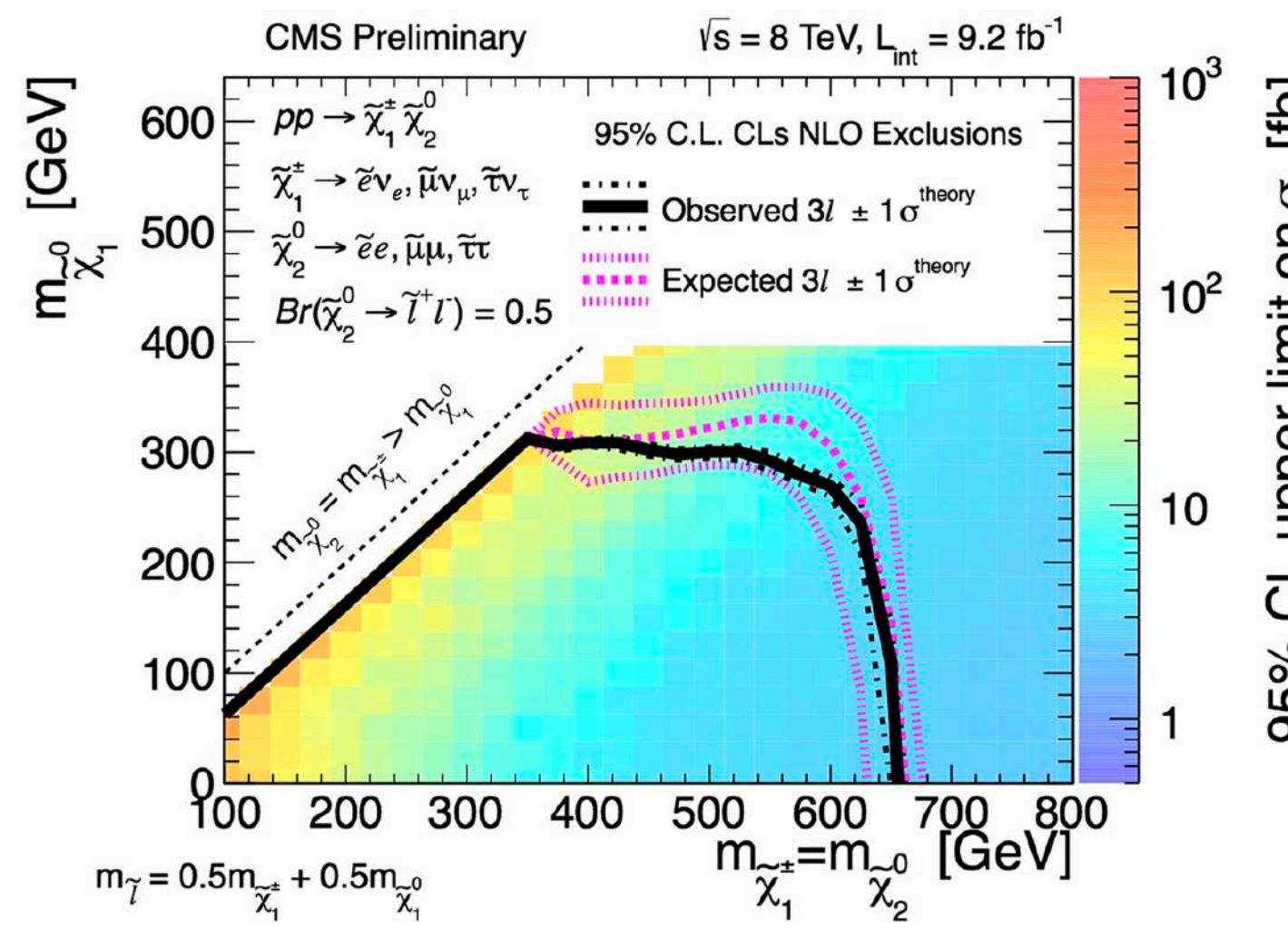


Lightest supersymmetric partner is stable

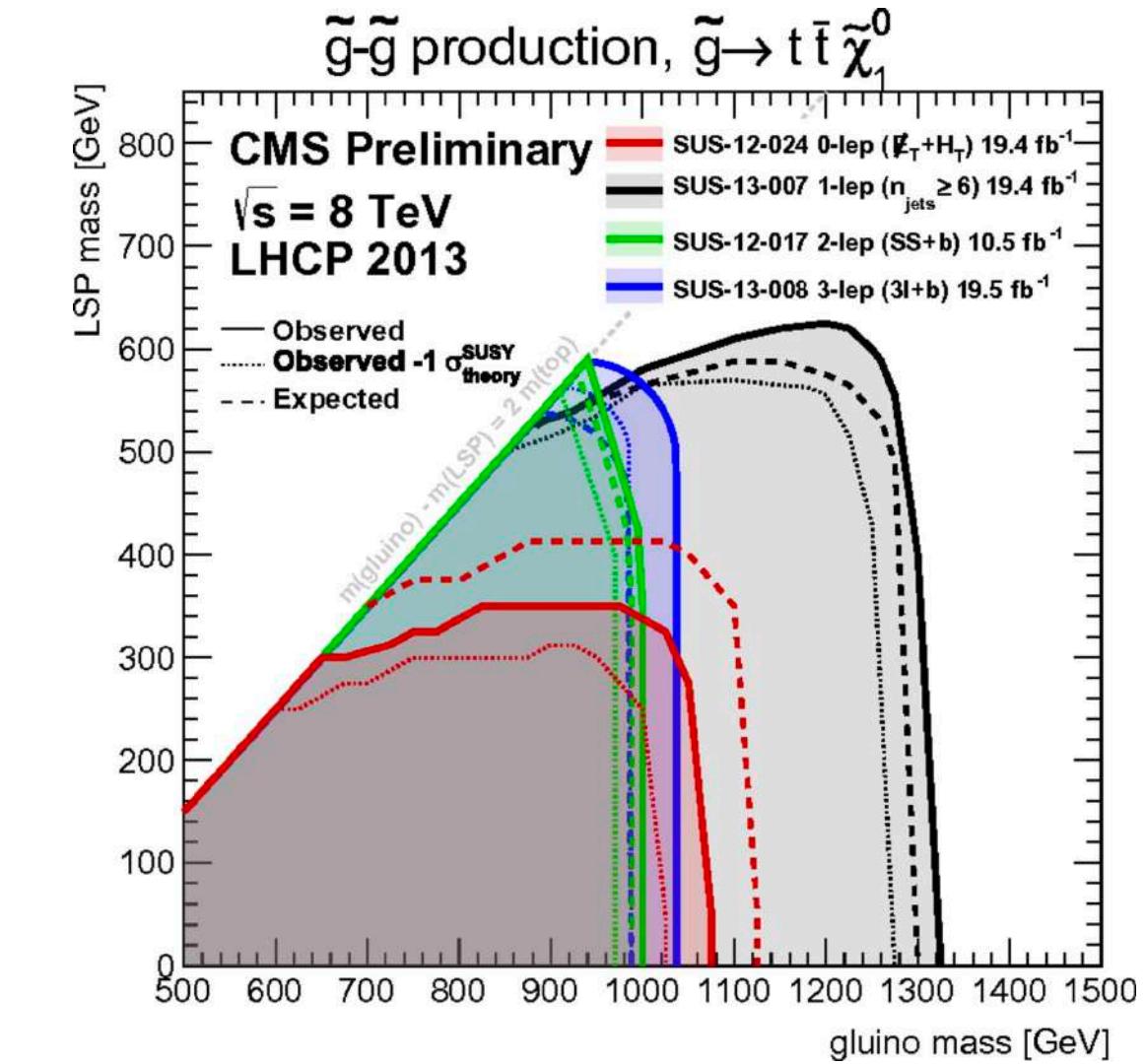
- Neutralino
- Gravitino
- Chargino
- Bino
- ...



Search at colliders!

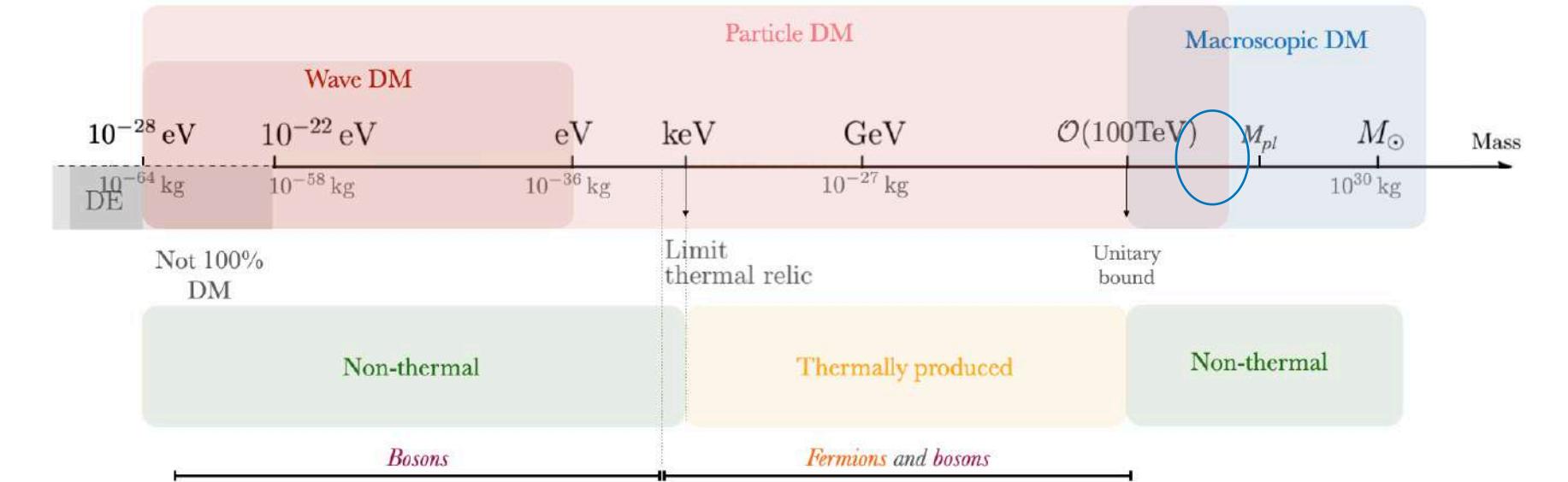


Neutralino and chargino



Gluino

WIMPzillas



WIMPzillas

Not a lot of superheavy candidates between $\mathcal{O}(100) \text{ TeV} - \mathcal{O}(10^{40}) \text{ eV}$

WIPMzillas: superheavy **particle**

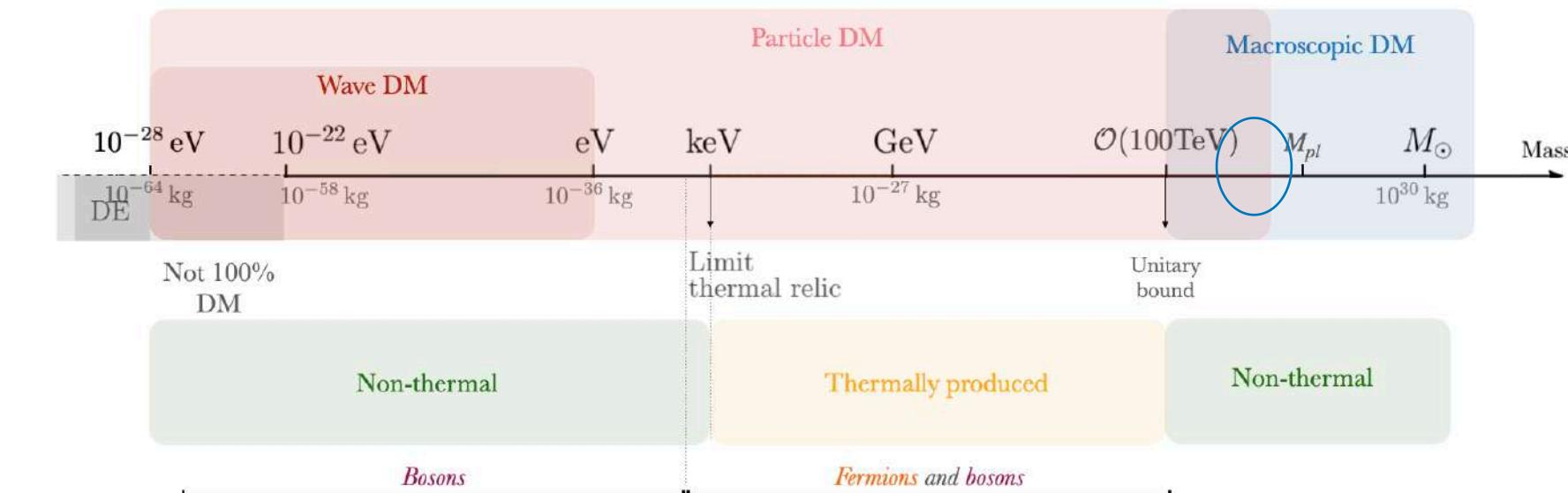
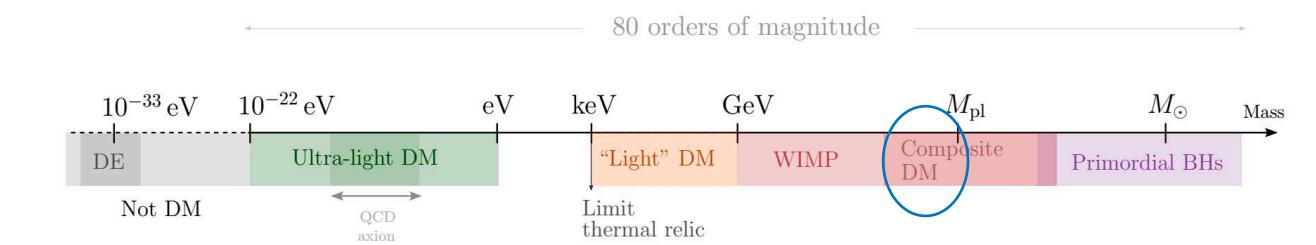
Ref.: Kolb et al 1998

2 necessary conditions:

- Must be stable (Condition for being particle DM)
- Must not have been in equilibrium when it froze out (i.e., it is **not** a thermal relic), otherwise $\Omega_X h^2$ would be much larger than one
- A sufficient condition for nonequilibrium is that the annihilation rate (per particle) must be smaller than the expansion rate: $n_X \sigma v < H$ (Condition for being non-thermal relic)

\Rightarrow Produced during inflation - $10^9 \text{ GeV} - 10^{16} \text{ GeV}$ (GUT scale)

There are no experiments looking for WIMPzillas



Produced non-thermally!!

(Not subjected to the unitary bound)

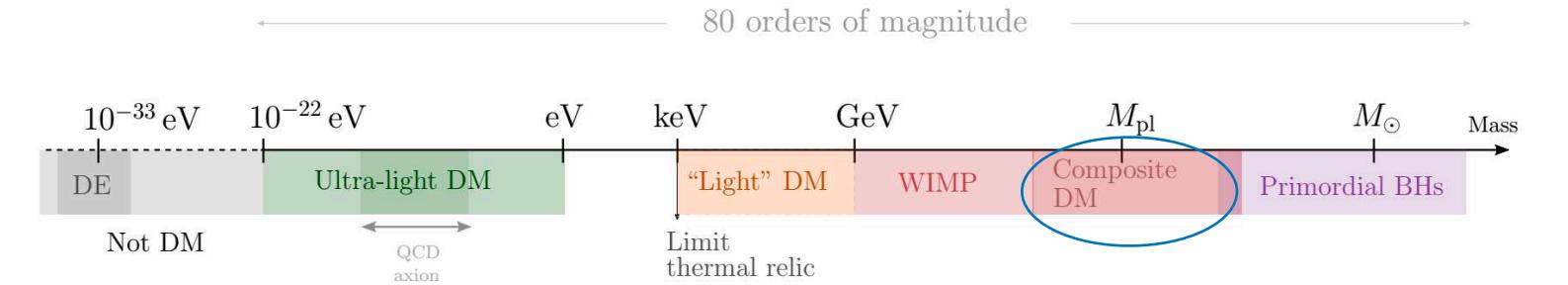


(Size does matter)

$$M_{pl} \sim 10^{19} \text{ GeV}$$

Macroscopic/composite DM



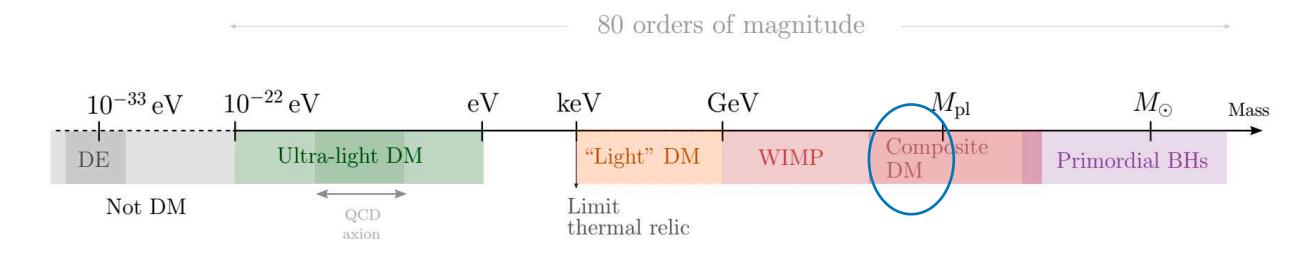


Composite DM

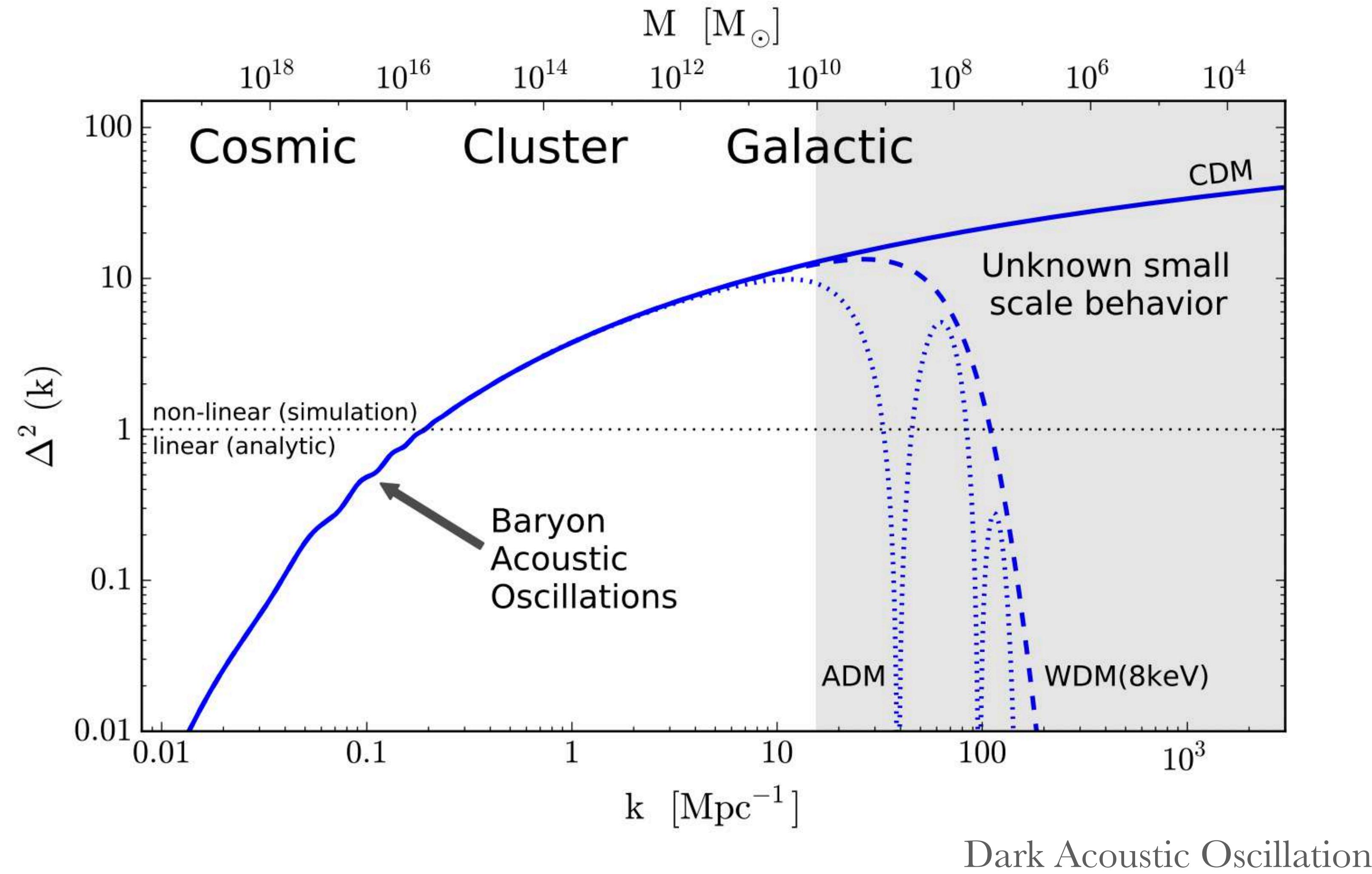
Dark atoms, (dark) glueballs, nuggets of baryons or other fermions ...

(an entire "SM" dark sector)

Atomic dark matter

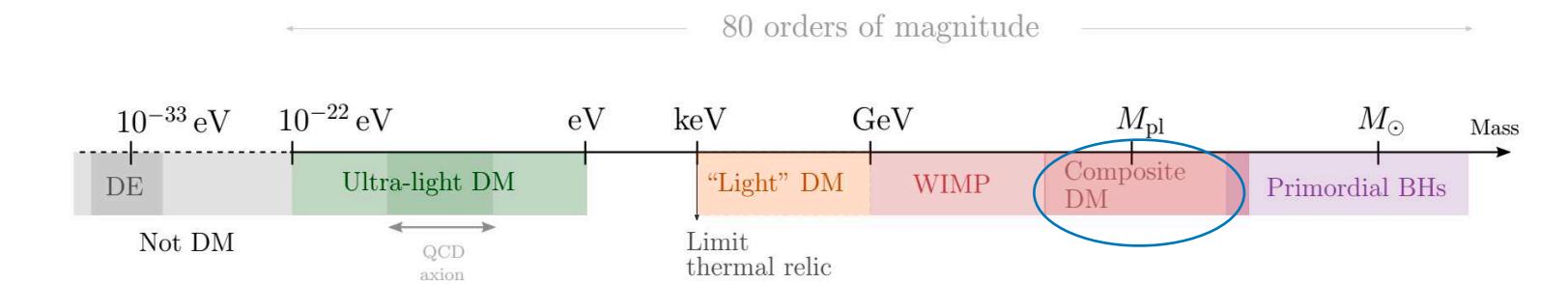


Ref.: Kaplan et al 2009
Cyr-Racine et al 2012



MACHOS

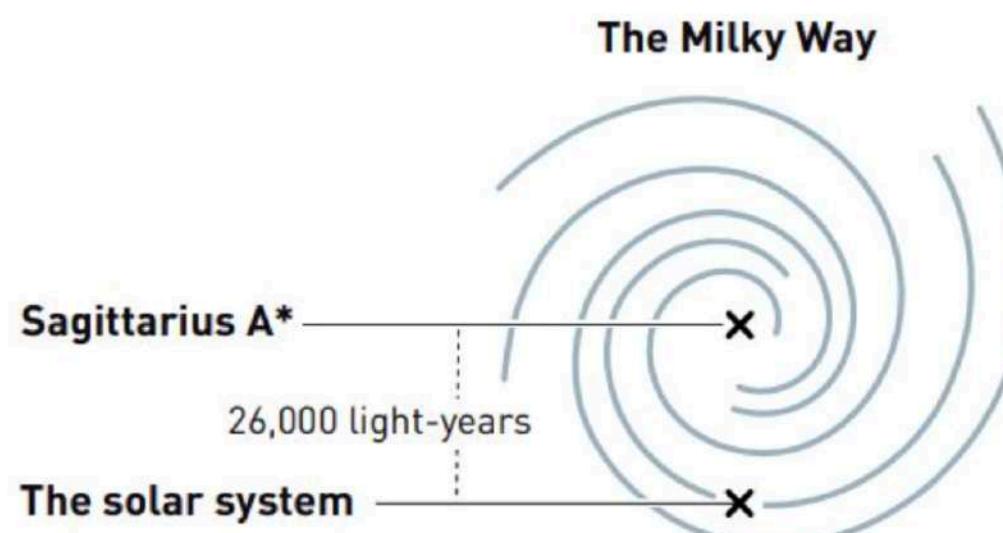
massive compact halo object



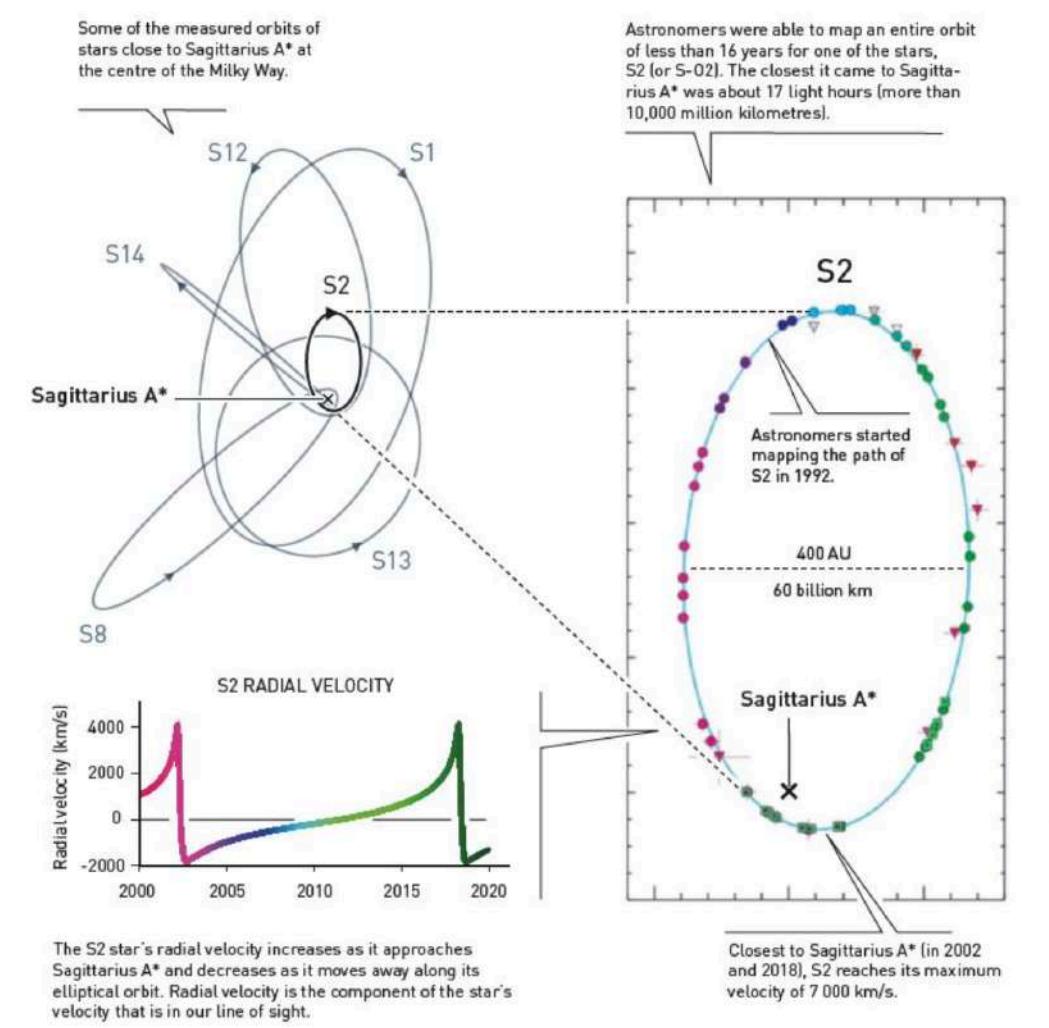
Primordial Black Holes

We know BHs exist!

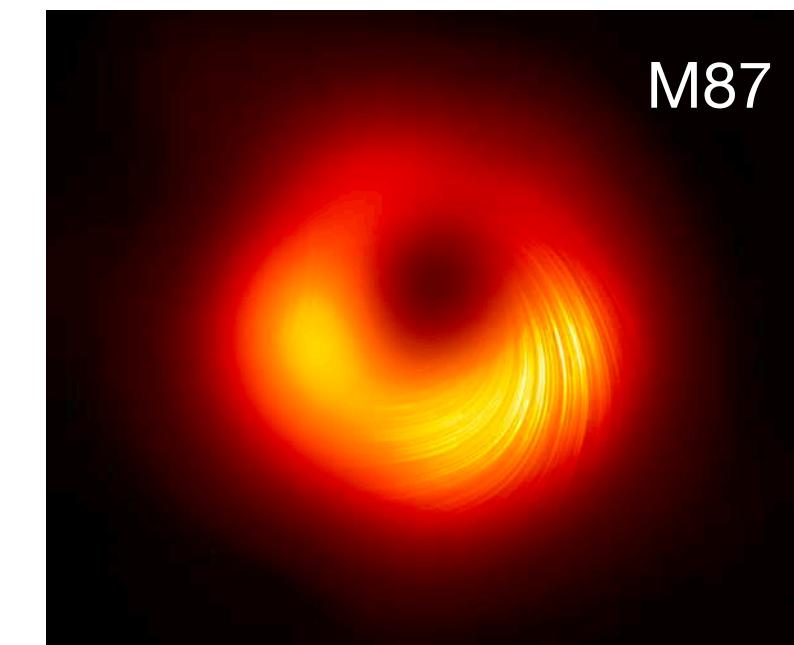
Star motion



Nobel prize (2021)

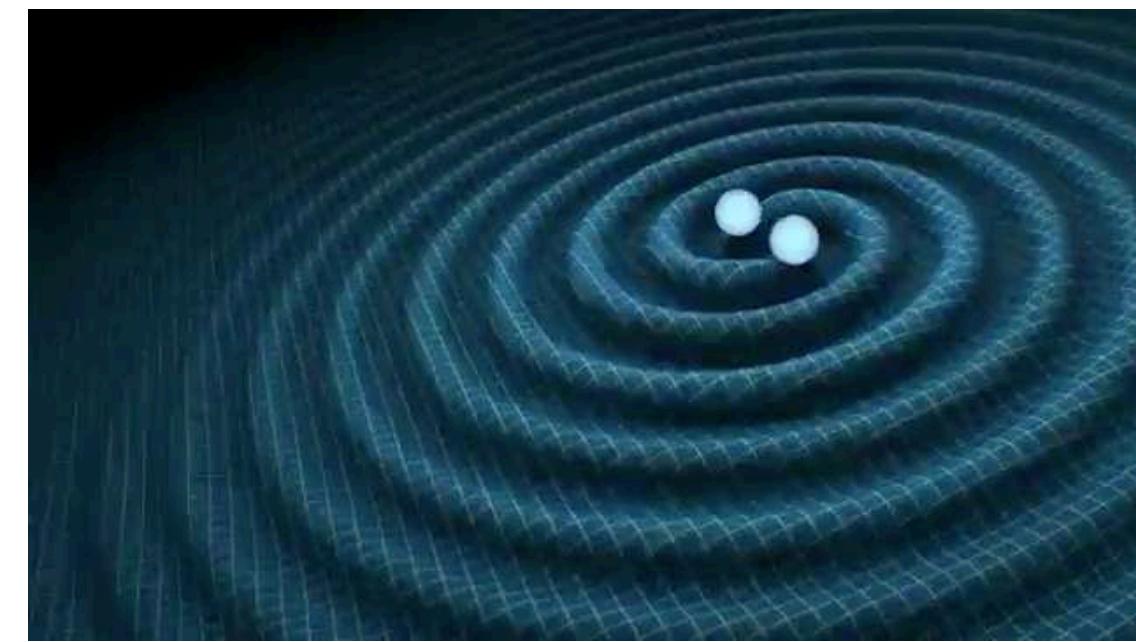


Event Horizon
Telescope



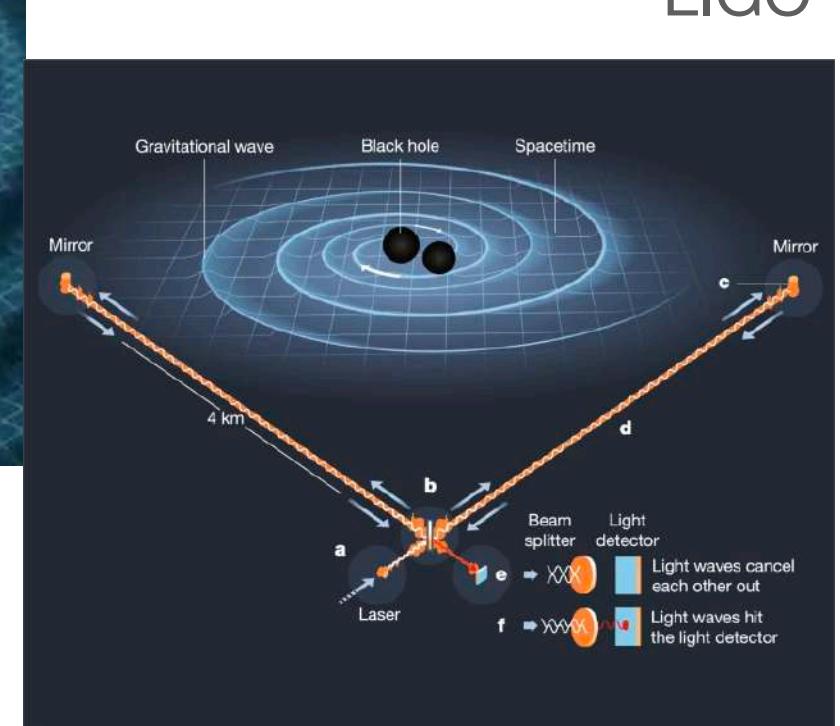
Crédito: ESO

Gravitational waves



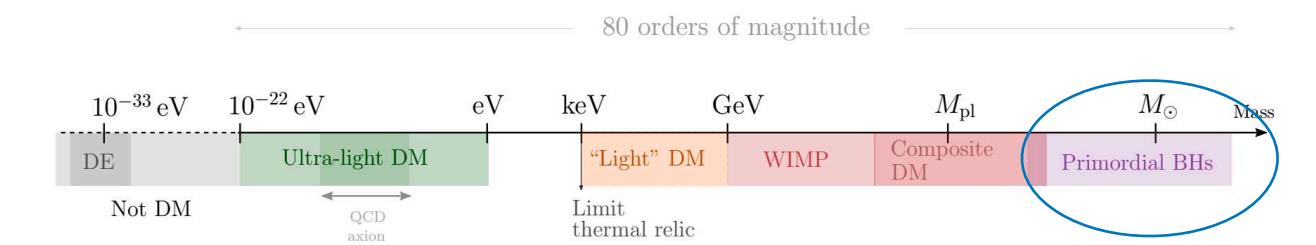
Crédito: ABC Science

+ VIRGO + KAGRA



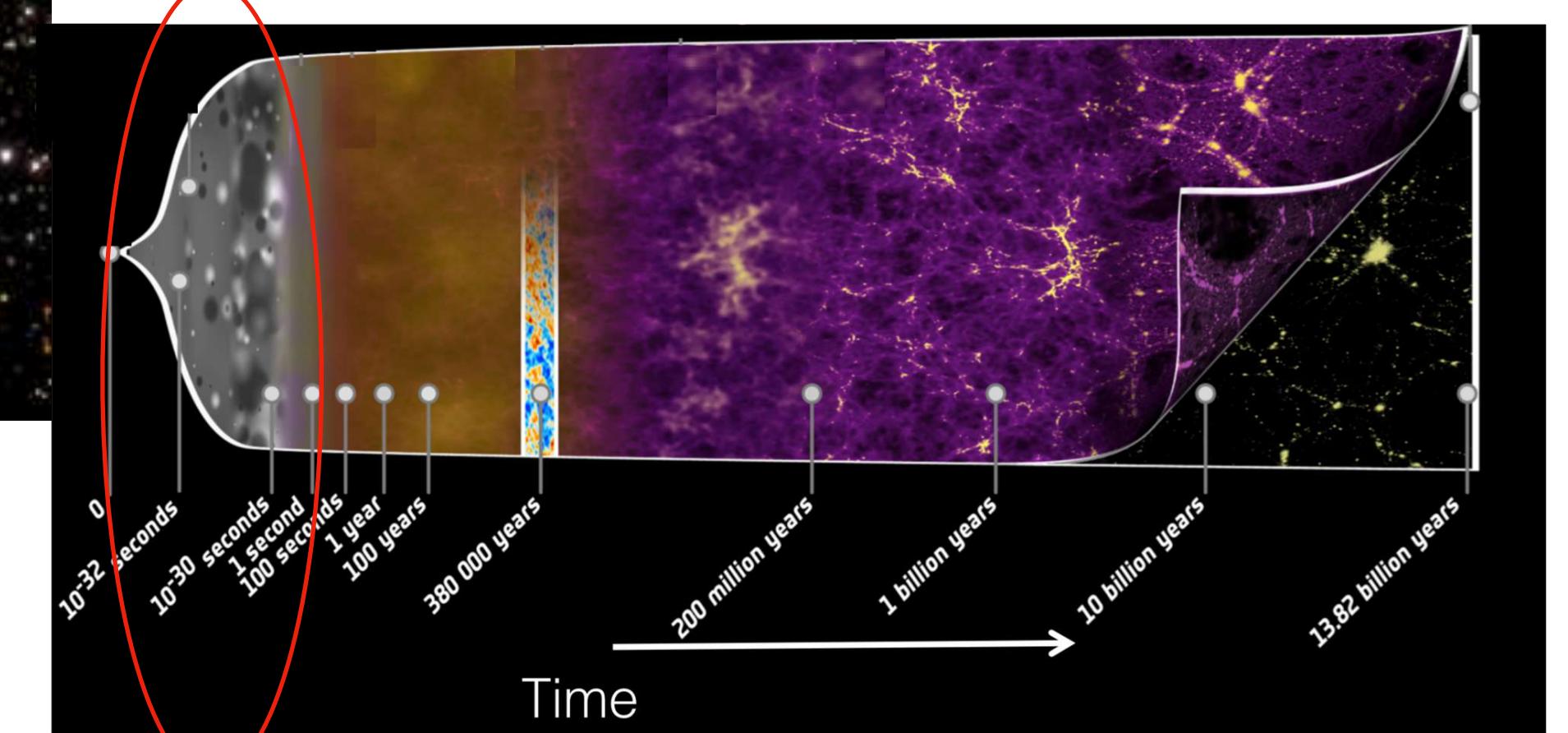
Crédito: ©Johan Jarnestad/The Royal Swedish Academy of Sciences.

Primordial Black Holes

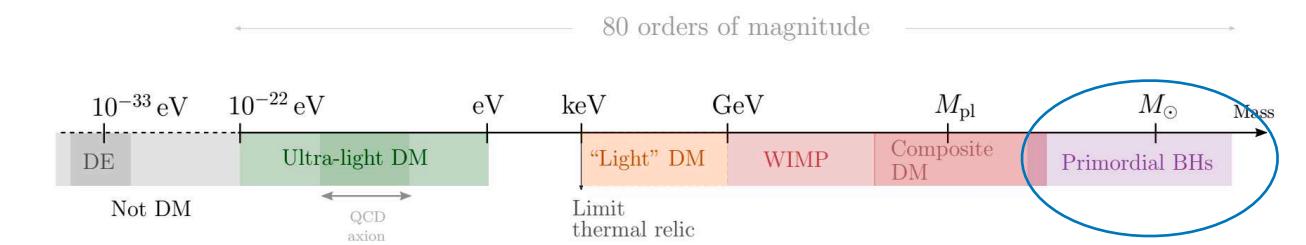


$$M_{\odot} \sim 2 \times 10^{30} \text{ kg}$$

- BHS formed at early times
- BHs with mass $10^{-15} M_{\odot} \lesssim M_{\text{PBH}} \lesssim \mathcal{O}(1) M_{\odot}$
- Can explain part or all of the DM



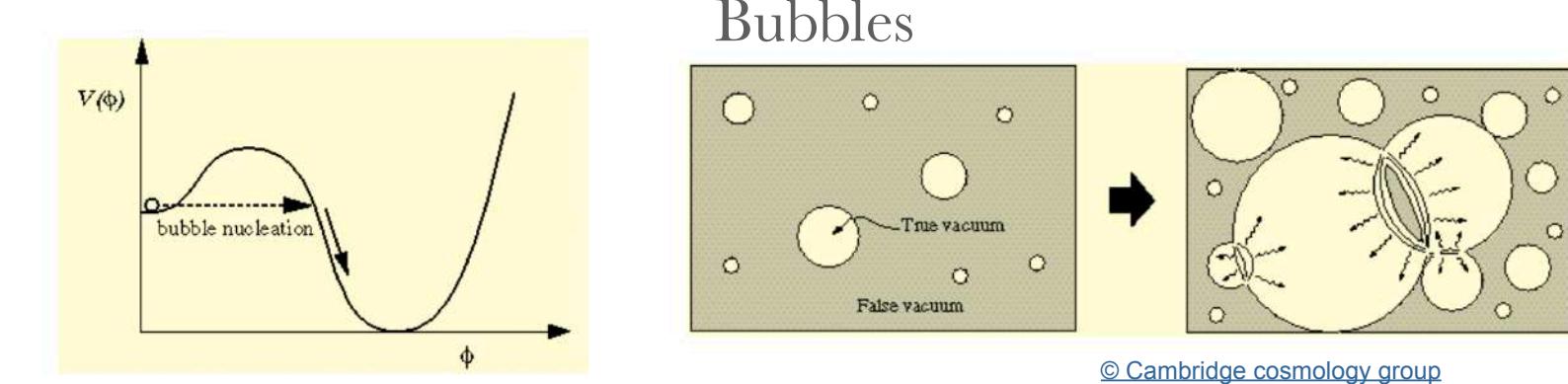
Primordial Black Holes



Formation mechanisms

- Bubble collision (Hawking et al, 1982)

1st order phase transitions occur via the nucleation of bubbles

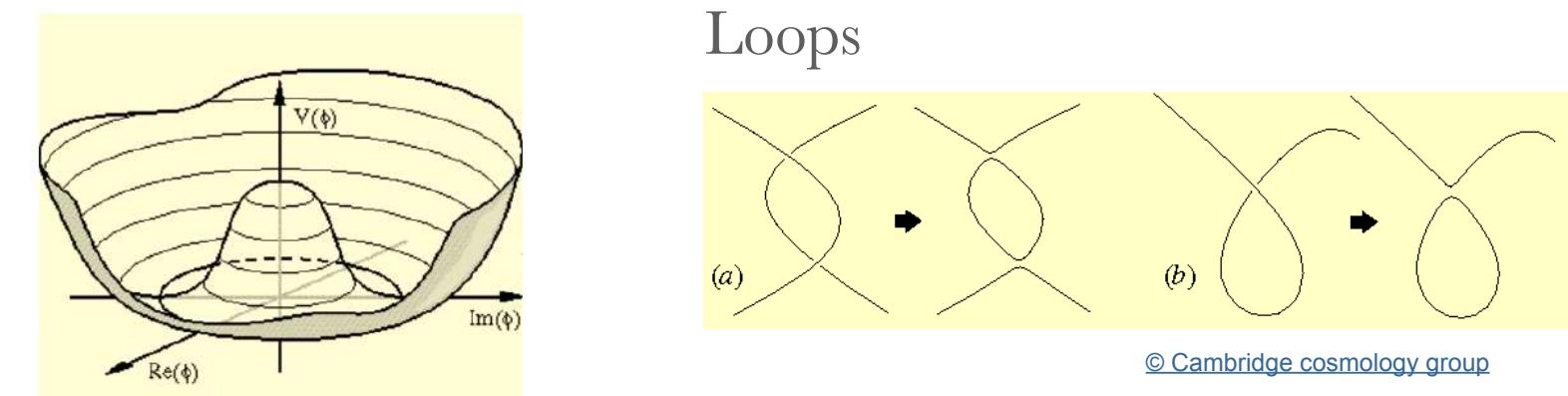


PBHs can form when bubbles collide (but bubble formation rate must be fine tuned)

⇒ PBH mass \sim order horizon mass at phase transition.

- Cosmic string loops (Hawking 1987)

Cosmic strings: 1d topological defects formed during symmetry breaking phase transition

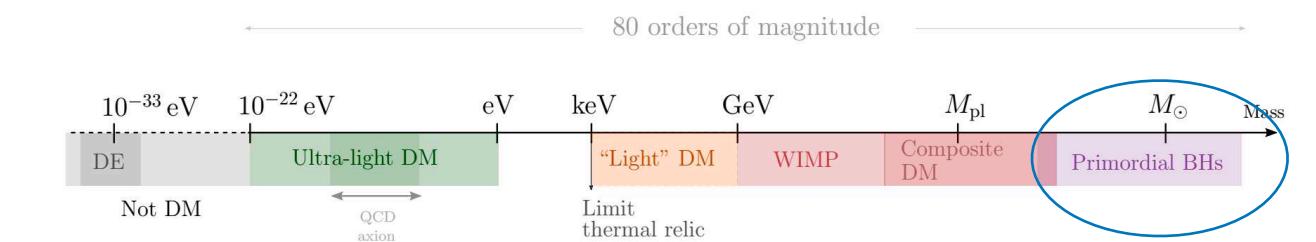


Small probability that loop will get into configuration of size \sim Schwarzschild radius

⇒ hence collapse to from a PBH with mass of order the horizon mass at that time

- Collapse of density perturbations (Carr and Hawking 1974)

Primordial Black Holes



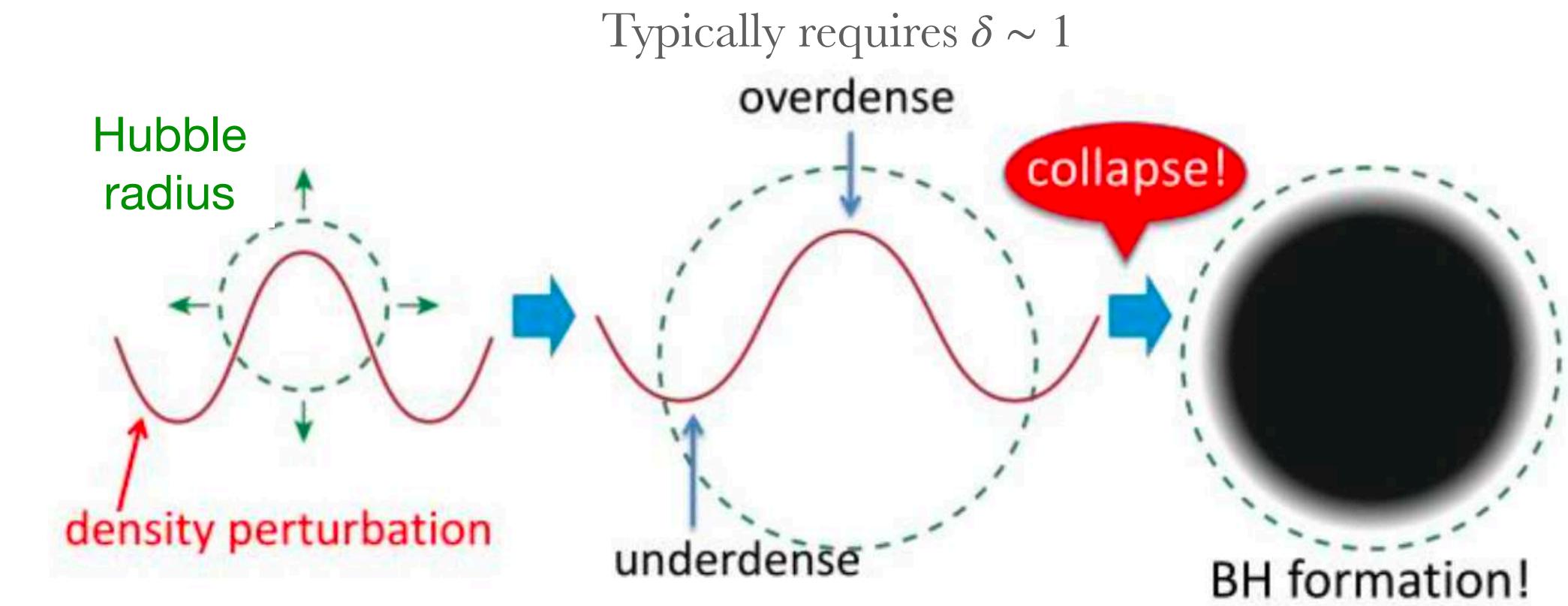
Formation mechanisms: collapse of large density perturbations (during radiation domination)

(0th order argument)

If a density perturbation is sufficiently large (at Hubble radius entry) it can collapse to form a **PBH**

Threshold for formation:

$$\delta_{hc} \geq \delta_c \sim w = \frac{P}{\rho} = \frac{1}{3}$$



⇒ Form a **PBH** with $M_{\text{PBH}} \sim M_{R_H}$

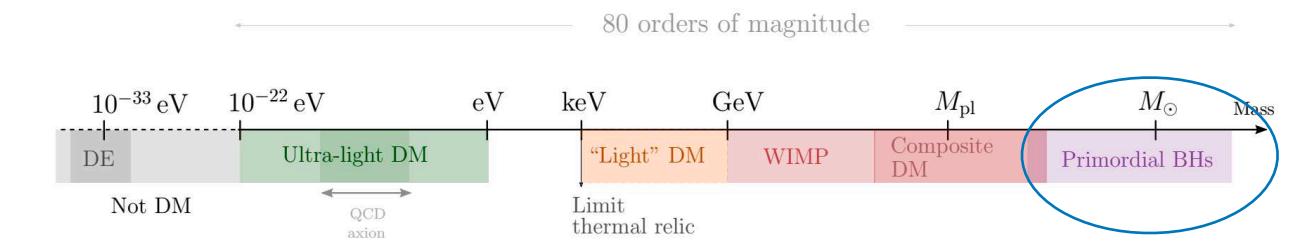
$$M_{R_H} = \frac{4\pi}{3} \rho(cH^{-1}) = \frac{c^3}{2GH} = \frac{tc^3}{G} \sim 10^{15} \text{ g} \left(\frac{t}{10^{-23} \text{ s}} \right) \sim M_\odot \left(\frac{t}{10^{-6} \text{ s}} \right) \sim M_{\text{PBH}}$$

Mass contained in the Hubble radius

* Here not in natural units!!

Kawasaki et al 2012

Primordial Black Holes



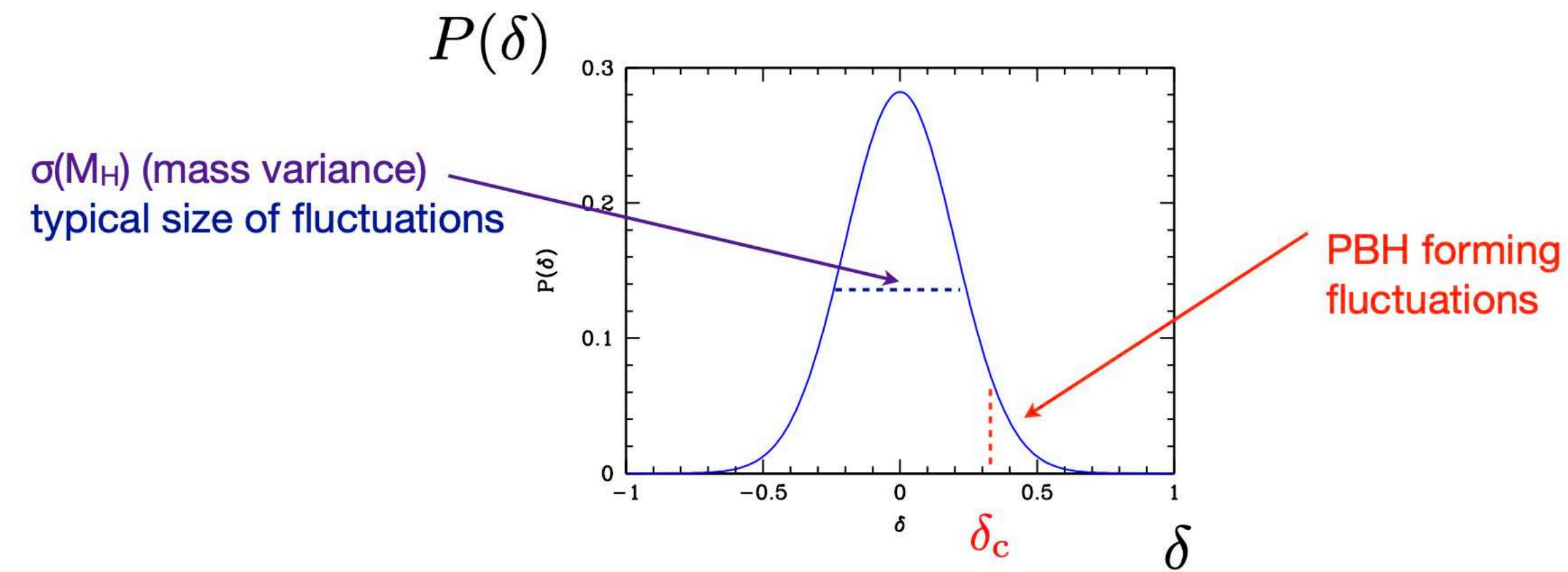
Initial PBH mass fraction: fraction of universe in regions dense enough to form PBHs

(0th order argument)

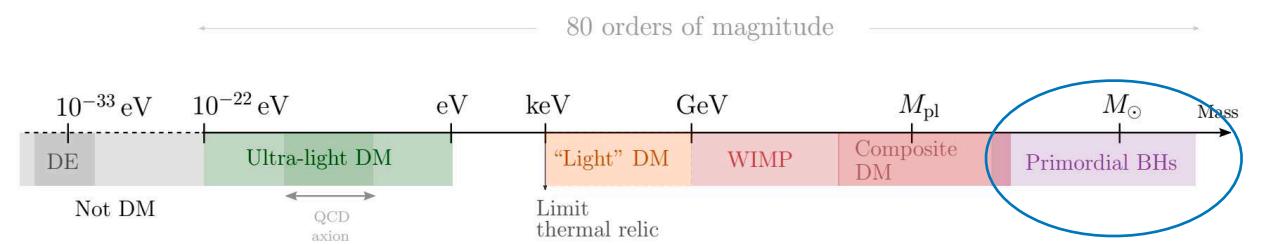
$$\beta(M) = \left(\frac{\rho_{\text{PBH}}}{\rho_{\text{tot}}} \right)_i \sim \int_{\delta_c}^{\infty} P(\delta(M_{R_H})) d\delta(M_{R_H}) \sim \sigma(M_{R_H}) \exp \left(-\frac{\delta_c^2}{2\sigma^2(M_{R_H})} \right)$$

density contrast,
smoothed on a
scale R_H

Assuming Gaussian
prob. distribution



Primordial Black Holes



PBH abundance

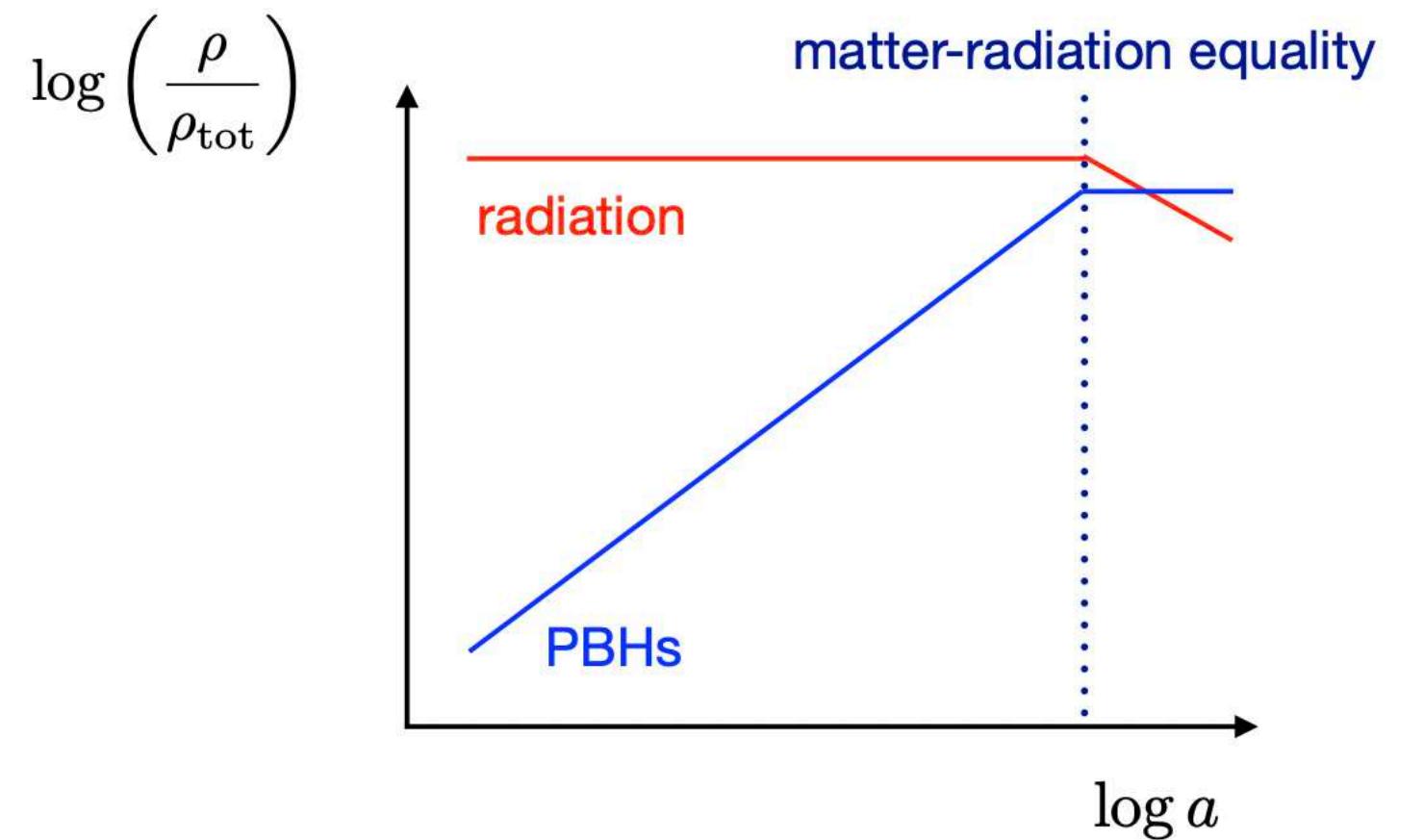
Since PBHs are matter, during radiation domination the fraction of energy in PBHs grows

$$\frac{\rho_{\text{PBH}}}{\rho_{\text{rad}}} \propto \frac{a^{-3}}{a^{-4}} \propto a$$

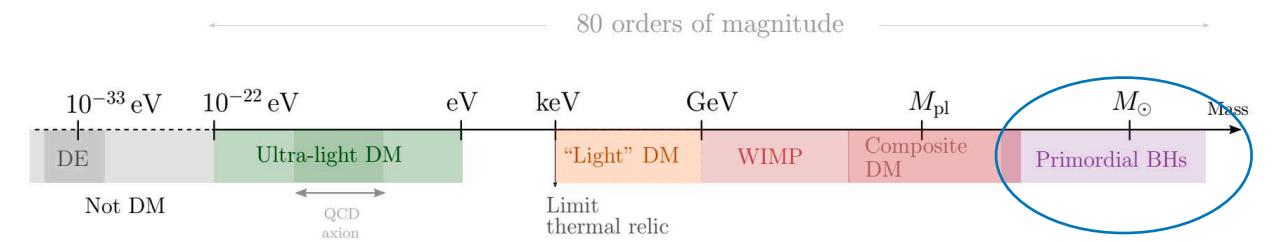
The **PBH** initial mass fraction, β , and **fraction of DM in form of PBH** are related by:

$$\beta(M) \sim 10^{-9} f_{\text{PBH}} \left(\frac{M}{M_\odot} \right)^{1/2}$$

⇒ initial mass fraction must be small, but non-negligible.



Primordial Black Holes



PBH abundance

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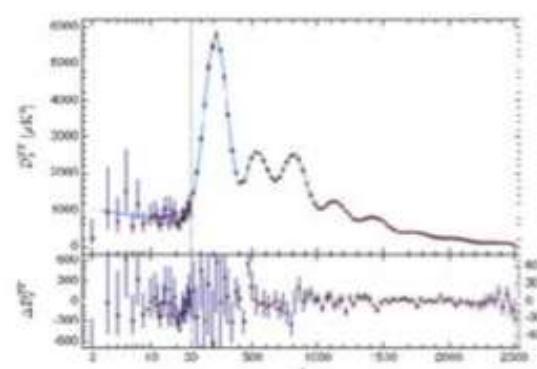
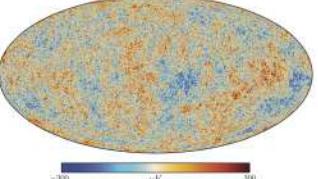
$$\beta(M) \sim 10^{-9} f_{\text{PBH}} \left(\frac{M}{M_\odot} \right)^{1/2}$$

\Rightarrow initial mass fraction must be small, but non-negligible.

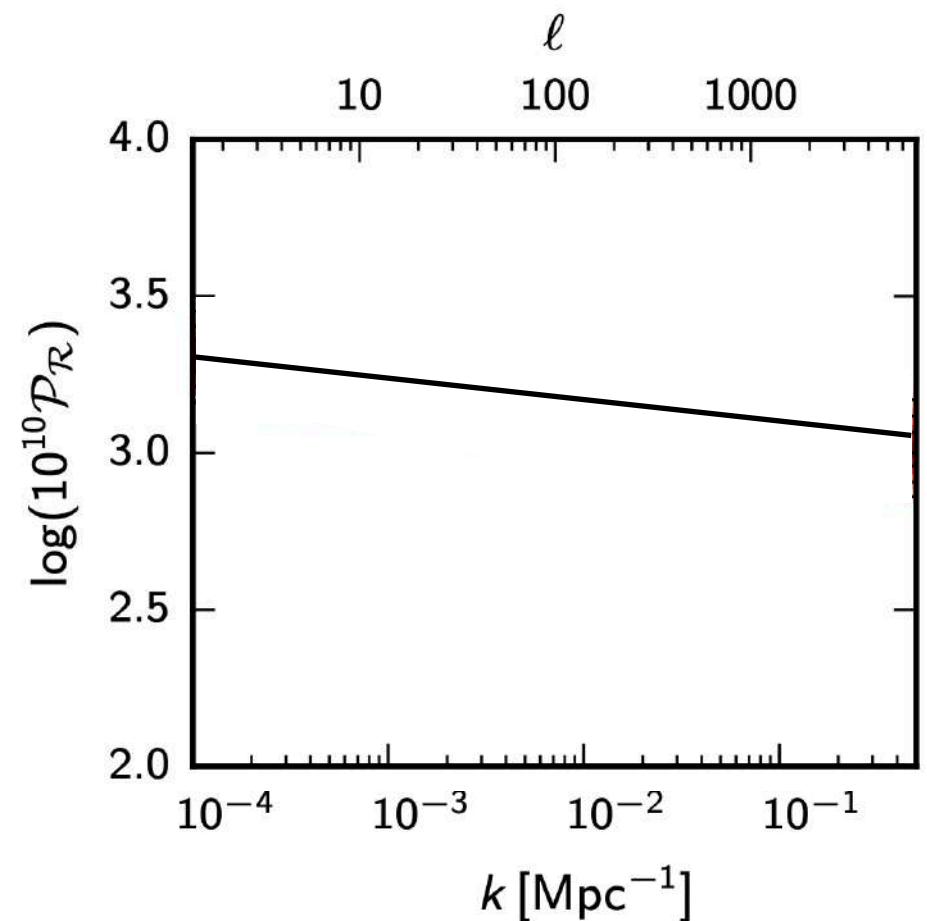
Initial perturbation:

Can the (nearly) scale-invariant primordial pert from early times (same that is the seed to LSS) source **PBH** and give a sizeable initial fraction? NO!

From CMB: $\sigma(M_{R_H}) \sim 10^{-5}$



Power spectrum of primordial curvature perturbation



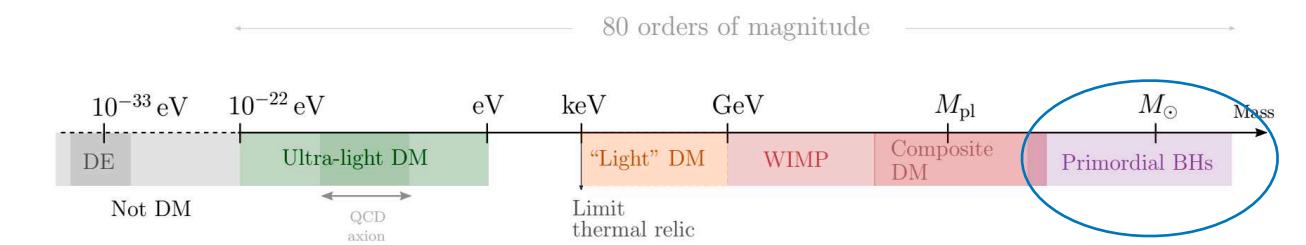
$$P_R = A_S \left(\frac{k}{k_p} \right)^{n_s - 1}$$



$$\beta(M) = \sigma(M_{R_H}) \exp \left(-\frac{\delta_c^2}{2\sigma^2 M_{R_H}} \right)$$

$$\sim \exp(10^{-10}) \ll 1 \quad \text{Negligible!}$$

Primordial Black Holes



PBH abundance

The **PBH** initial mass fraction, β , and **fraction of DM in form of PBH** are related by:

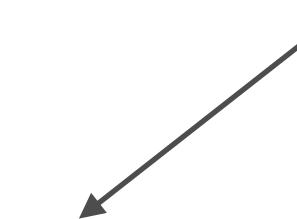
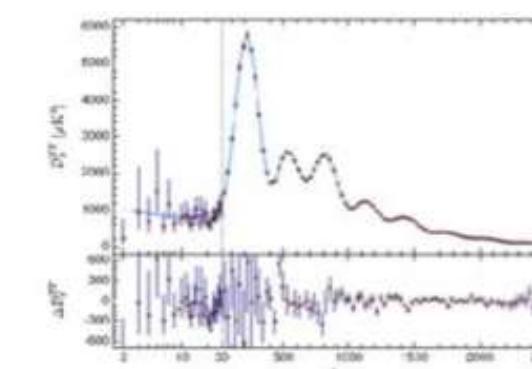
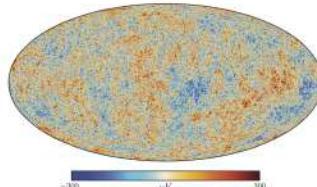
$$\beta(M) \sim 10^{-9} f_{\text{PBH}} \left(\frac{M}{M_\odot} \right)^{1/2}$$

\Rightarrow initial mass fraction must be small, but non-negligible.

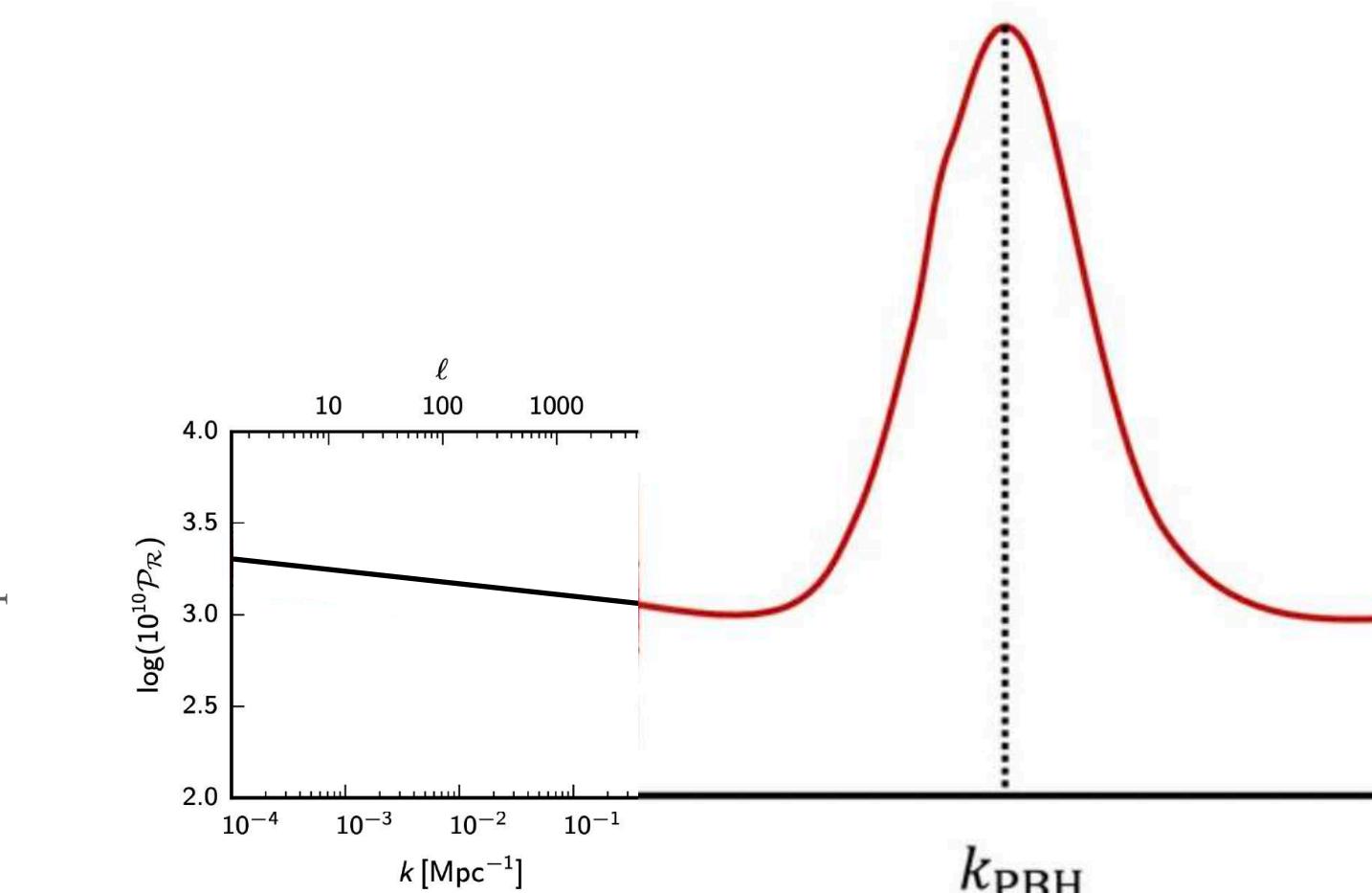
Initial perturbation:

Can the (nearly) scale-invariant primordial pert from early times (same that is the seed to LSS) source **PBH** and give a sizeable initial fraction? NO!

From CMB: $\sigma(M_{R_H}) \sim 10^{-5}$



Power spectrum of primordial curvature perturbation

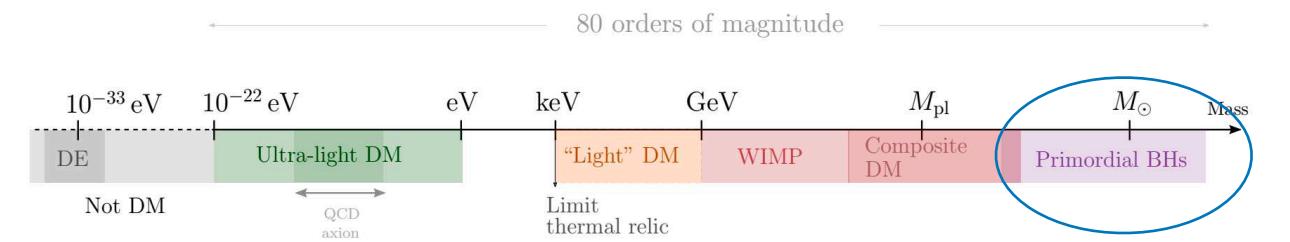


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$$\sim \exp(10^{-10}) \ll 1 \quad \text{Negligible!}$$

To form an interesting number of PBHs amplitude of primordial perturbations must be 2-3 orders of larger on small scales than on cosmological scales and fine-tuned.

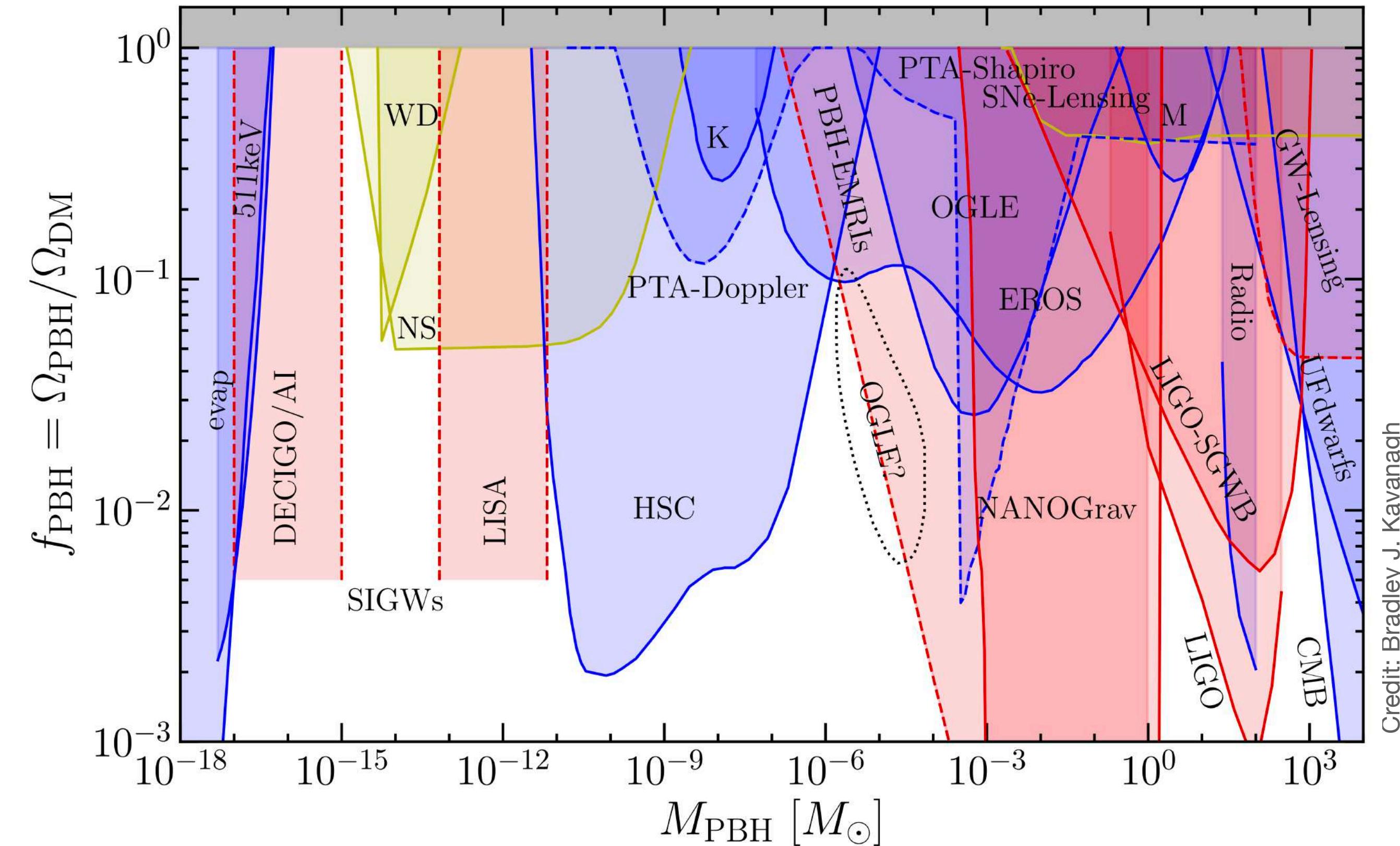
Primordial Black Holes



On cosmological scales PBH DM would behave like particle DM, however on galactic and smaller scales its granularity can have *observable consequences*.

Primordial Black Holes

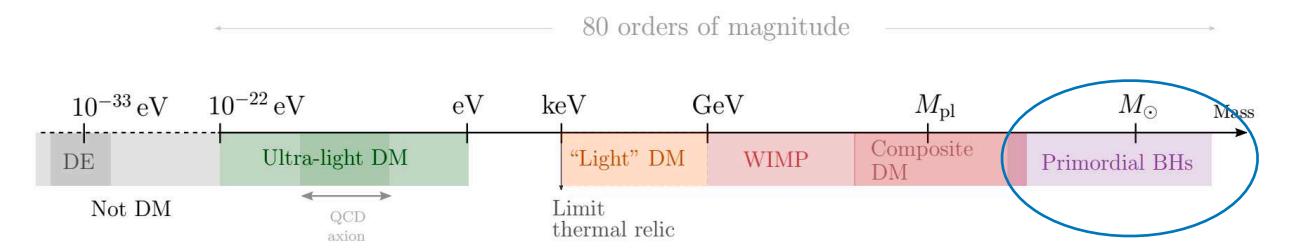
Bounds



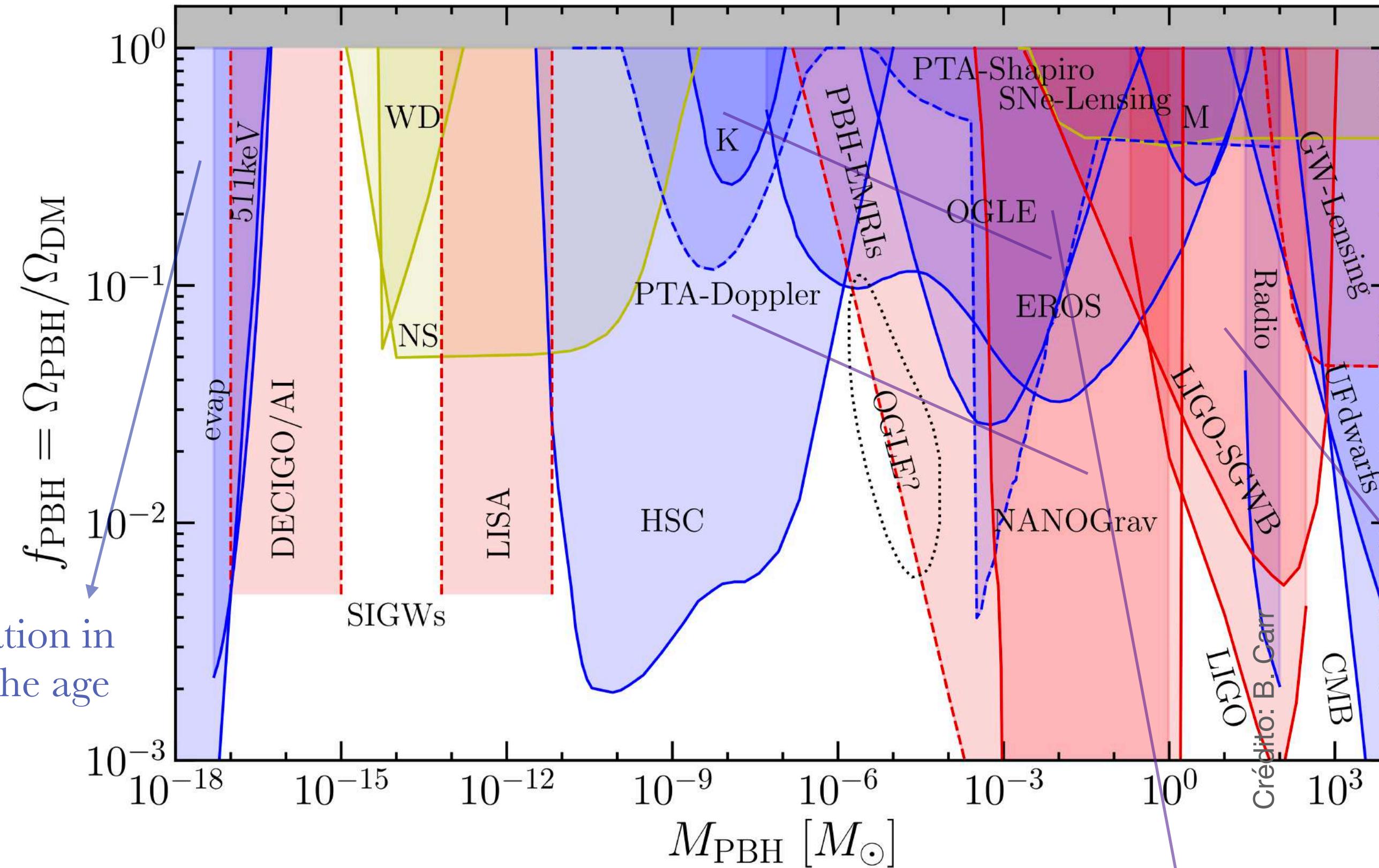
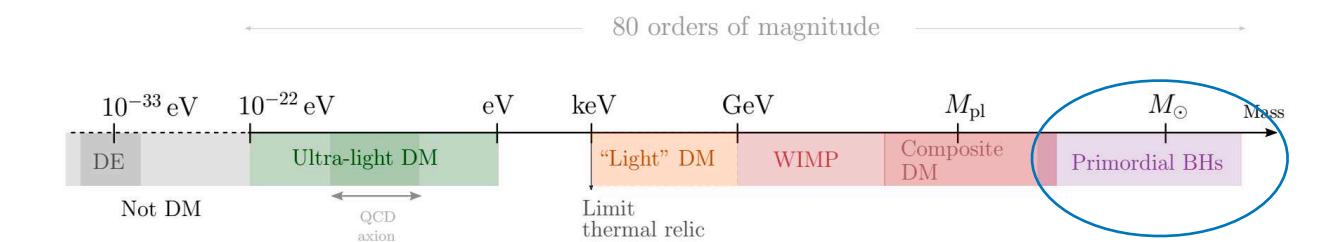
Credit: Bradley J. Kavanagh

$$M_{\odot} \sim 2 \times 10^{30} \text{ kg}$$

Notebook to plot the PBH bounds: <https://github.com/bradkav/PBHbounds>



Primordial Black Holes



PBH evaporate, emits gamma radiation in scales or the order of smaller than the age of the universe

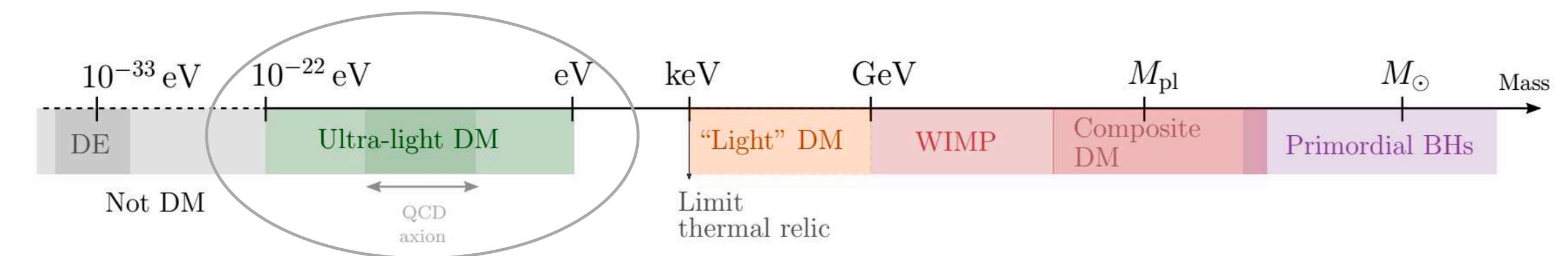
$$M_{\odot} \sim 2 \times 10^{30} \text{ kg}$$

Direct searches via microlensing in our galaxy and M31,... (does not require that it is a BH - scalar bound system)

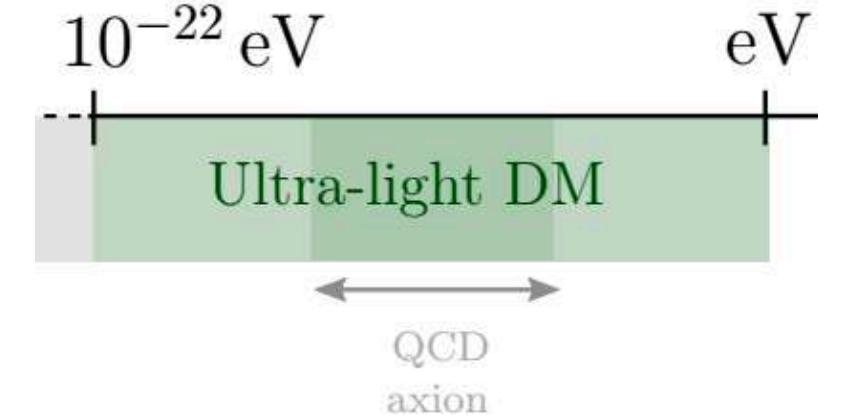
Wave DM



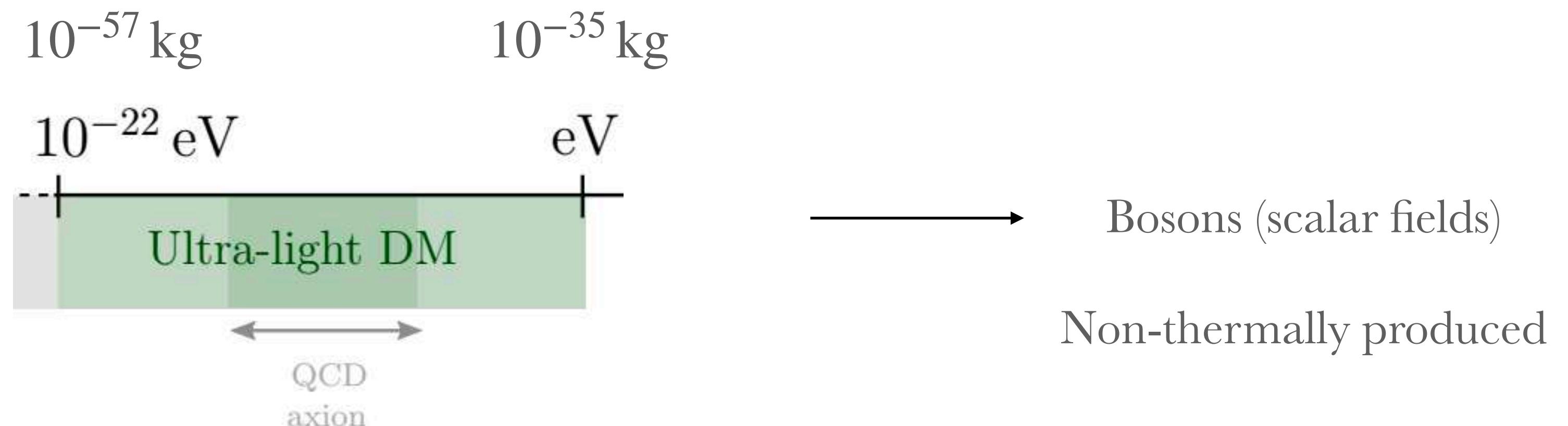
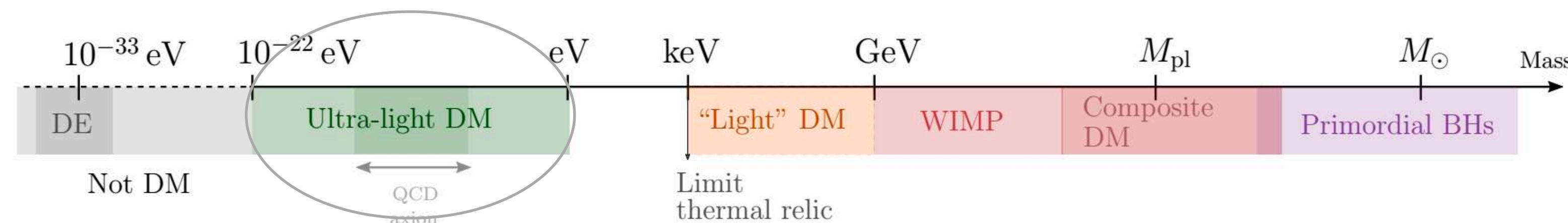
Ultra-light dark matter



Ultra-light dark matter



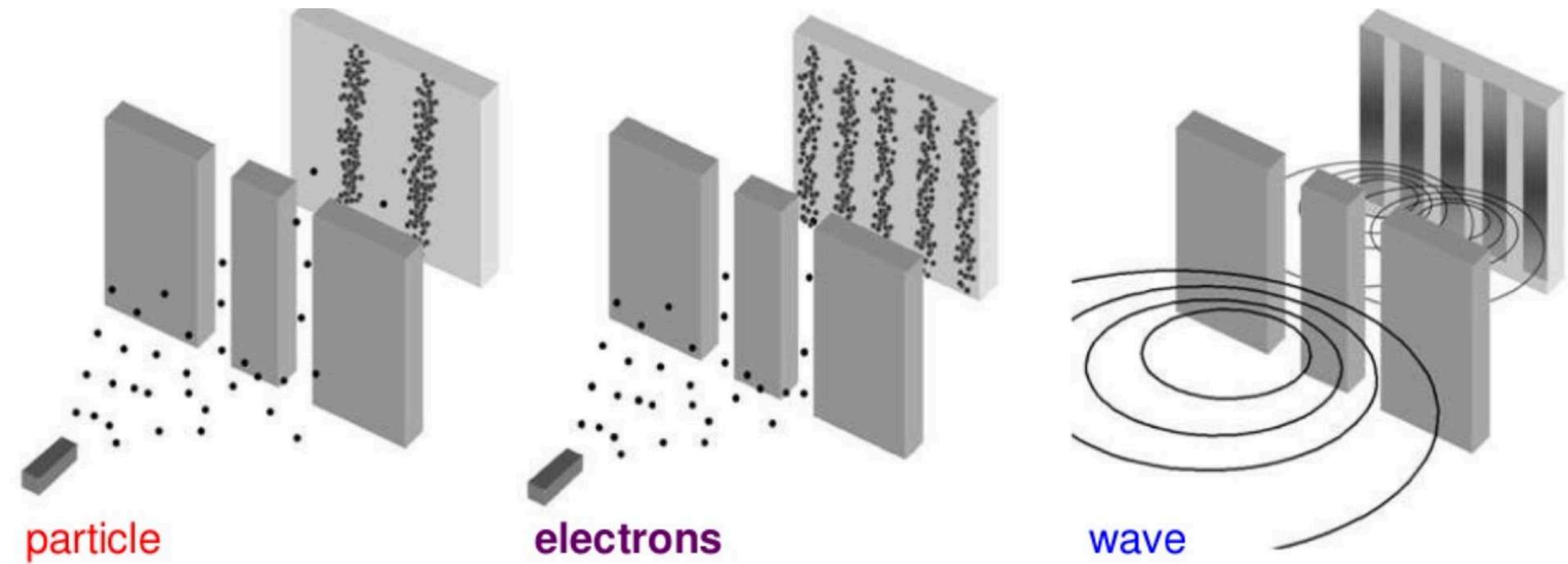
Ultra-light candidate, cold \longrightarrow Large $\lambda_{\text{dB}} \sim 1/mv$
 Lightest possible candidate for DM



Wave-Particle duality

All matter exhibits a wave behaviour

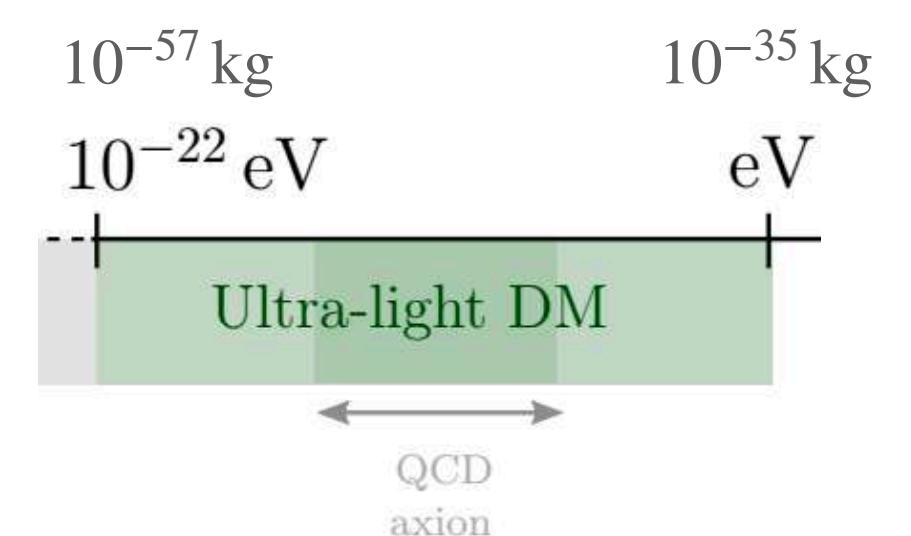
De Broglie 1924



$$\lambda_{dB} \sim \frac{1}{mv}$$

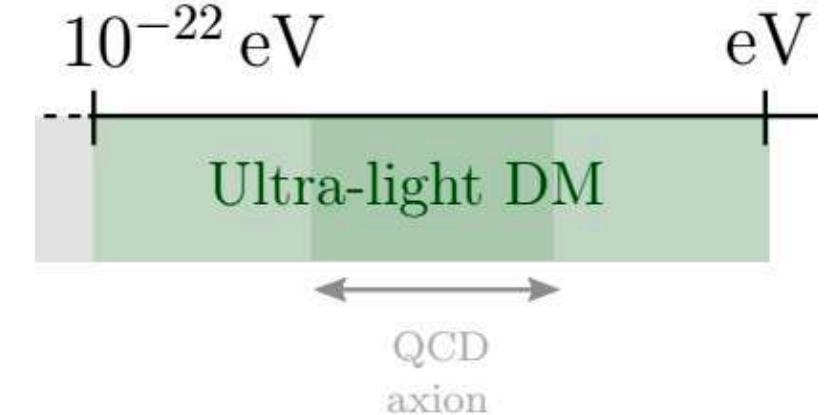
$$\lambda_{dB} \sim 1/\sqrt{2\pi mk_B T}$$

	Mass (kg)	Speed (m/s)	λ_{dB} (m)
Accelerated e-	9.1×10^{-31}	5.9×10^6	1.2×10^{-10}
Golf ball	0.045	220	4.8×10^{-30}



$$\lambda_{dB}^{ULDM} \sim \text{pc} - \text{kpc}$$

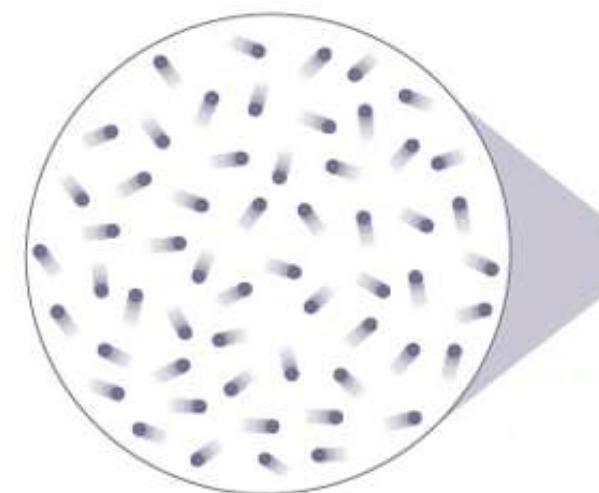
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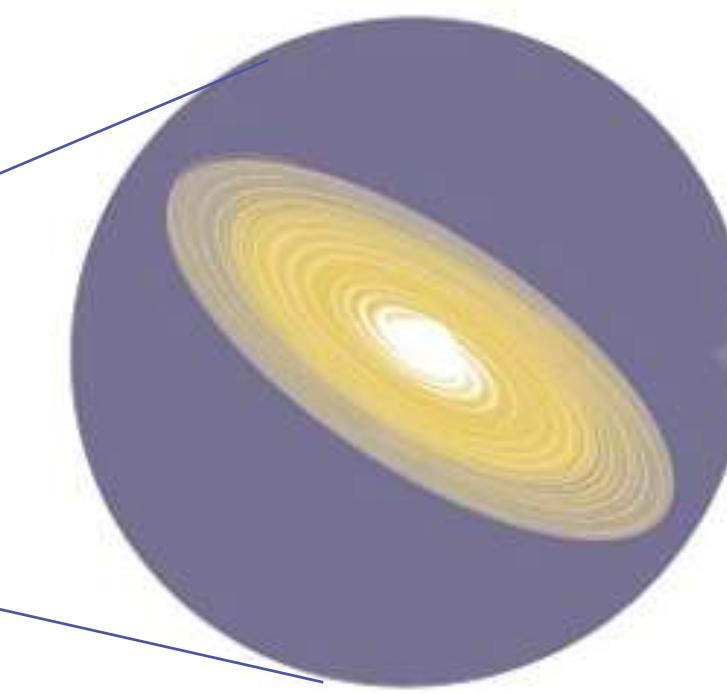
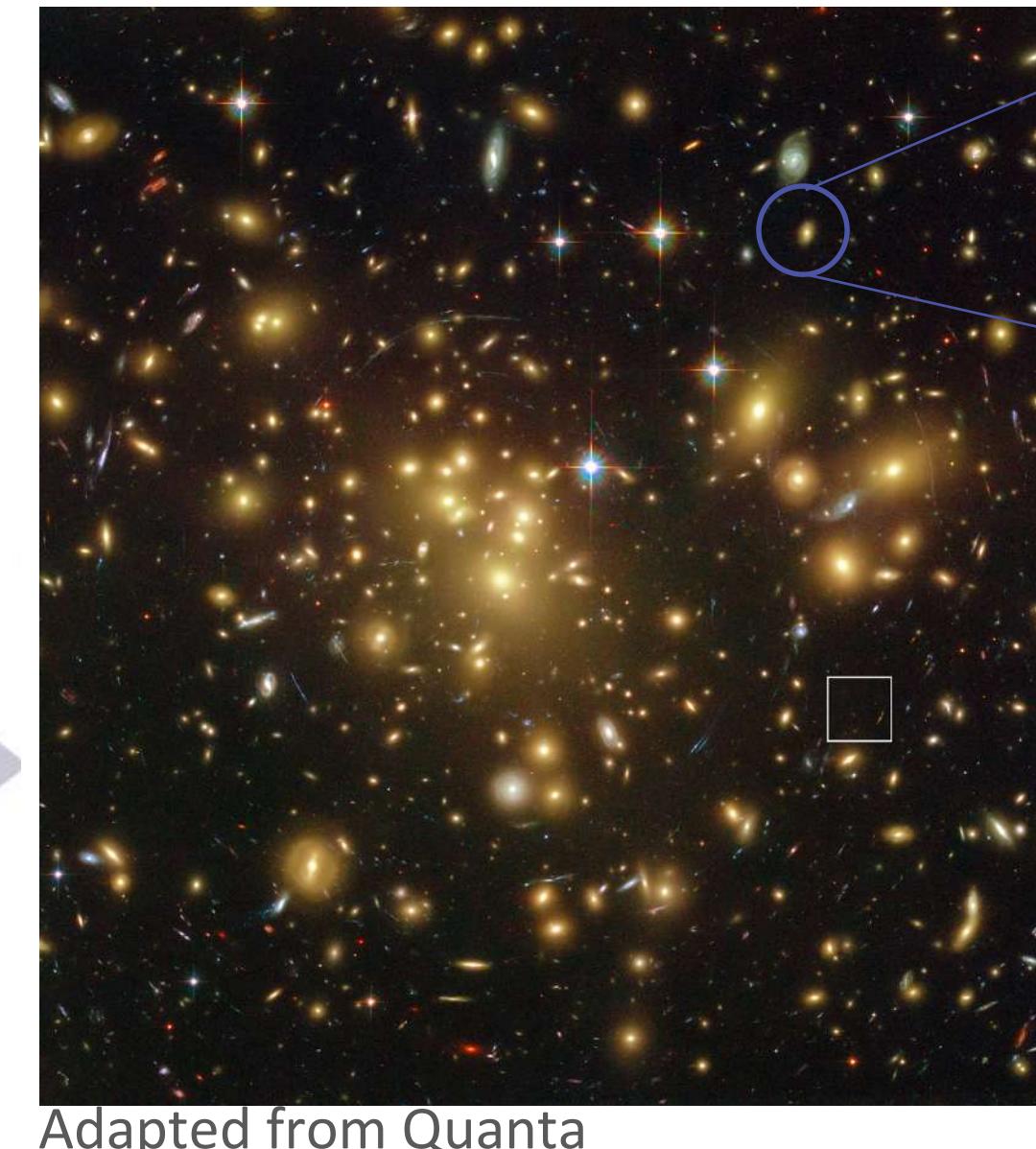
Ultra-light candidate

Large $\lambda_{dB} \sim 1/mv$

Large scales:
DM behaves like standard
particle DM (**CDM**).

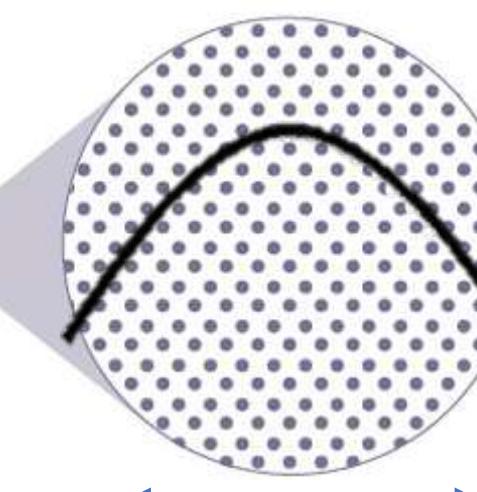


DM: particles
 $d \gg \lambda_{dB}$



Galaxy halo

DM: wave behaviour



λ_{dB}
 $d \ll \lambda_{dB}$

Small scales:
DM behaves like a **wave**

$$10^{-60} \text{ kg} \quad 10^{-35} \text{ kg}$$
$$10^{-25} \text{ eV} \lesssim m \lesssim \text{eV}$$

$$\lambda_{dB}^{ULDM} \sim \text{pc} - \text{kpc}$$

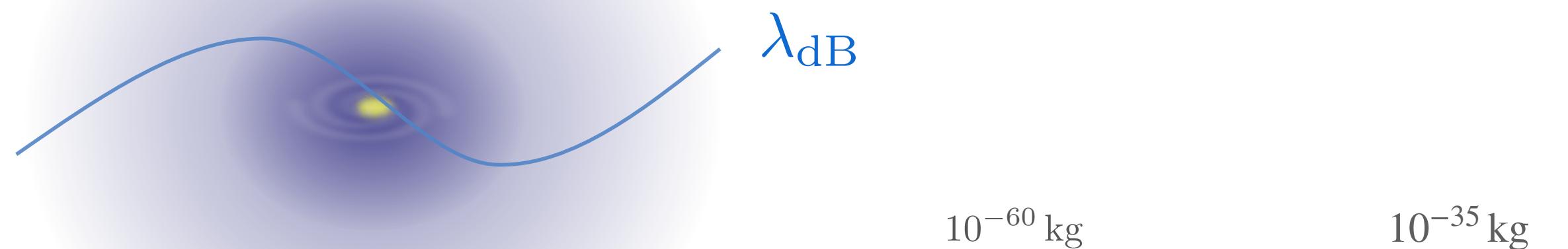
How light is ultra-light?

“Ultra-light dark matter”, EF, 2020.

Behave as wave on galactic scales:

- λ_{dB} must be **smaller** than the halo

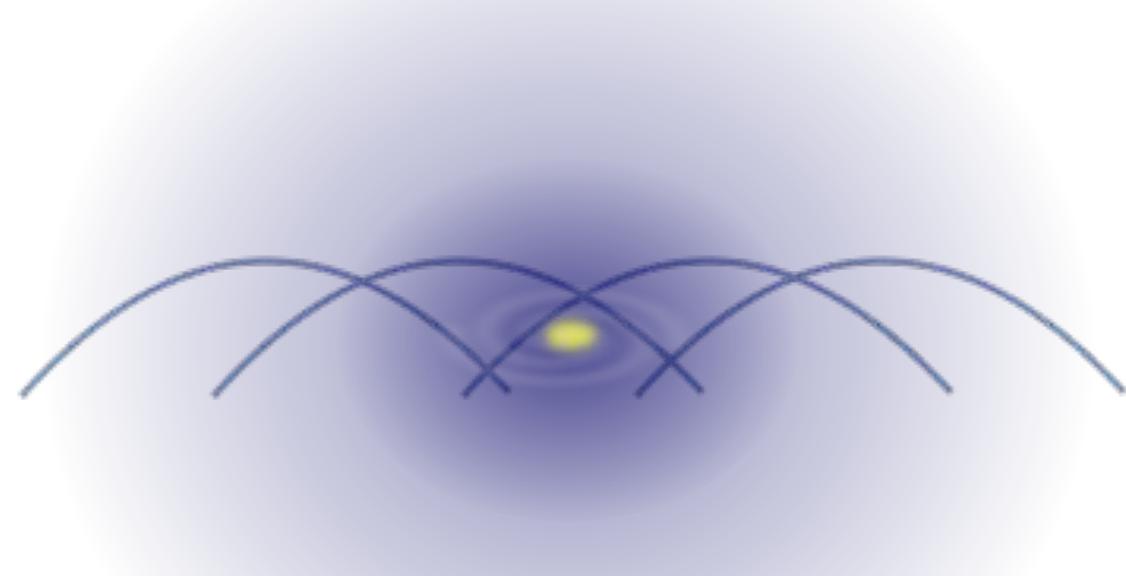
$$\lambda_{dB} < R_{\text{halo}}$$
$$\Rightarrow m \gtrsim 10^{-25} \text{ eV}$$



$$10^{-25} \text{ eV} \lesssim m \lesssim \text{eV}$$

- λ_{dB} **overlap** to be of halo size

$$\lambda_b \sim \frac{1}{mv} \geq d \sim \left(\frac{m}{\rho_{vir}} \right)^{\frac{1}{3}}$$
$$\Rightarrow m \leq 2 \text{ eV}$$

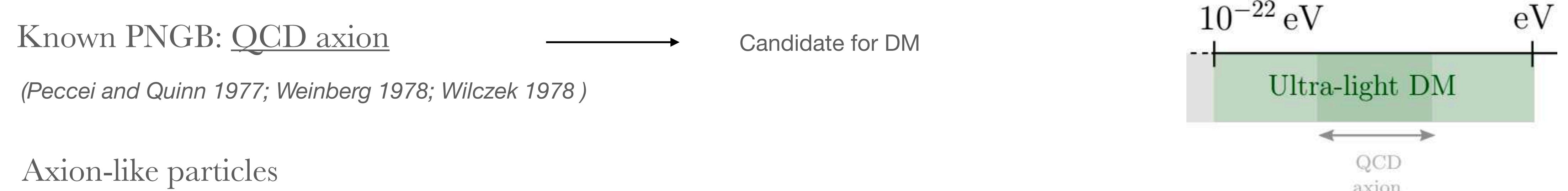


$$\lambda_{dB}^{ULDM} \sim \text{pc} - \text{kpc}$$

Motivation: particle physics

ULDM candidates

- Natural candidate for a light scalar field is a pseudo-Nambu Goldstone boson (breaking of an approximate symmetry)



Axions or Axion like particles (ALP)

Axions and ALPs are pseudo Nambu Goldstone bosons from the spontaneous symmetry breaking of a $U_{\text{PQ}}(1)$ ($U(1)$) symmetry, and are described by the complex field: $\Psi = v e^{i\phi/f_a}$

$$v_{0,ssb} = f_a/\sqrt{2} \quad \longrightarrow \quad \phi \rightarrow \phi + c$$

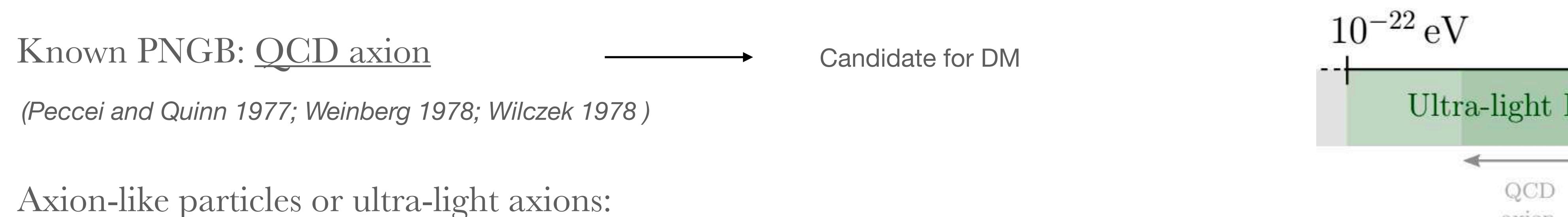
Non-perturbative effects (from string theory or instantons) induce a potential:

$$V(\phi) = \Lambda_a^4 [1 - \cos(\phi/f_a)] \xrightarrow{\phi \ll f_a} \frac{1}{2} m^2 \phi^2 + \frac{g}{4} \phi^4 + \dots$$

Motivation: particle physics

ULDM candidates

- Natural candidate for a light scalar field is a pseudo-Nambu Goldstone boson



Axion-like particles or ultra-light axions:

- ALPs expected in string theory (*Arvanitaki et al., Svrcek, Witten*)
- Can generate PNGB that are ultra-light
- Formation mechanism: needs to have a relic abundance that gives the correct DM abundance

Non-thermal mechanism (e.g. mis-alignement)

$$\Omega_{axion} \sim 0.15 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \theta_1^2$$

$$\Omega_{ALP} \sim 0.1 \left(\frac{f_a}{10^{17} \text{ GeV}} \right)^2 \left(\frac{m}{10^{-22} \text{ eV}} \right)$$

Motivation: particle physics

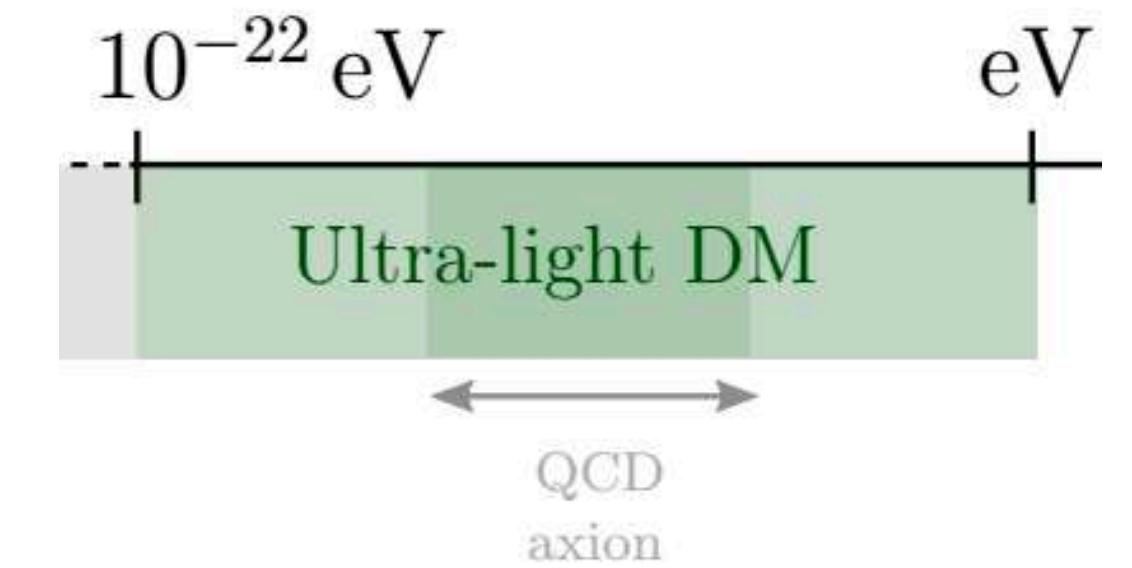
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Known PNGB: QCD axion Candidate for DM
(*Peccei and Quinn 1977; Weinberg 1978; Wilczek 1978*)

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Vector FDM: challenging in the ultra-light regime

(e.g. from misalignment requires non-minimal couplings to Ricci scalar -> viol. of unitarity long. graviton-photon scattering; oscillating Higgs or oscillating misaligned axion - resonant production - choices for couplings for right abundance)

Spin 2 FDM: (e.g bigravity)

Motivation: particle physics

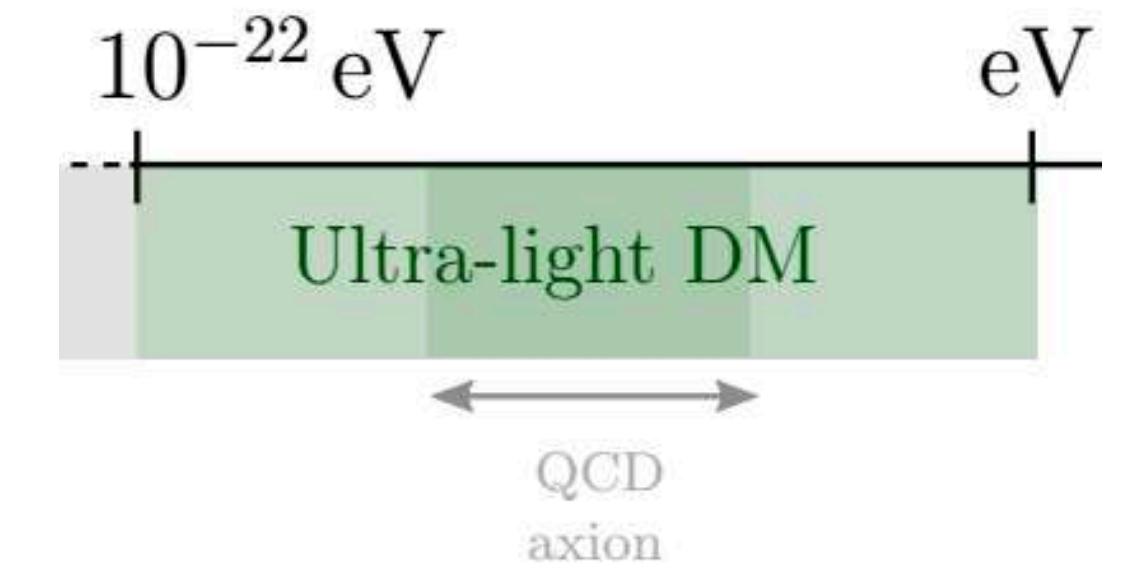
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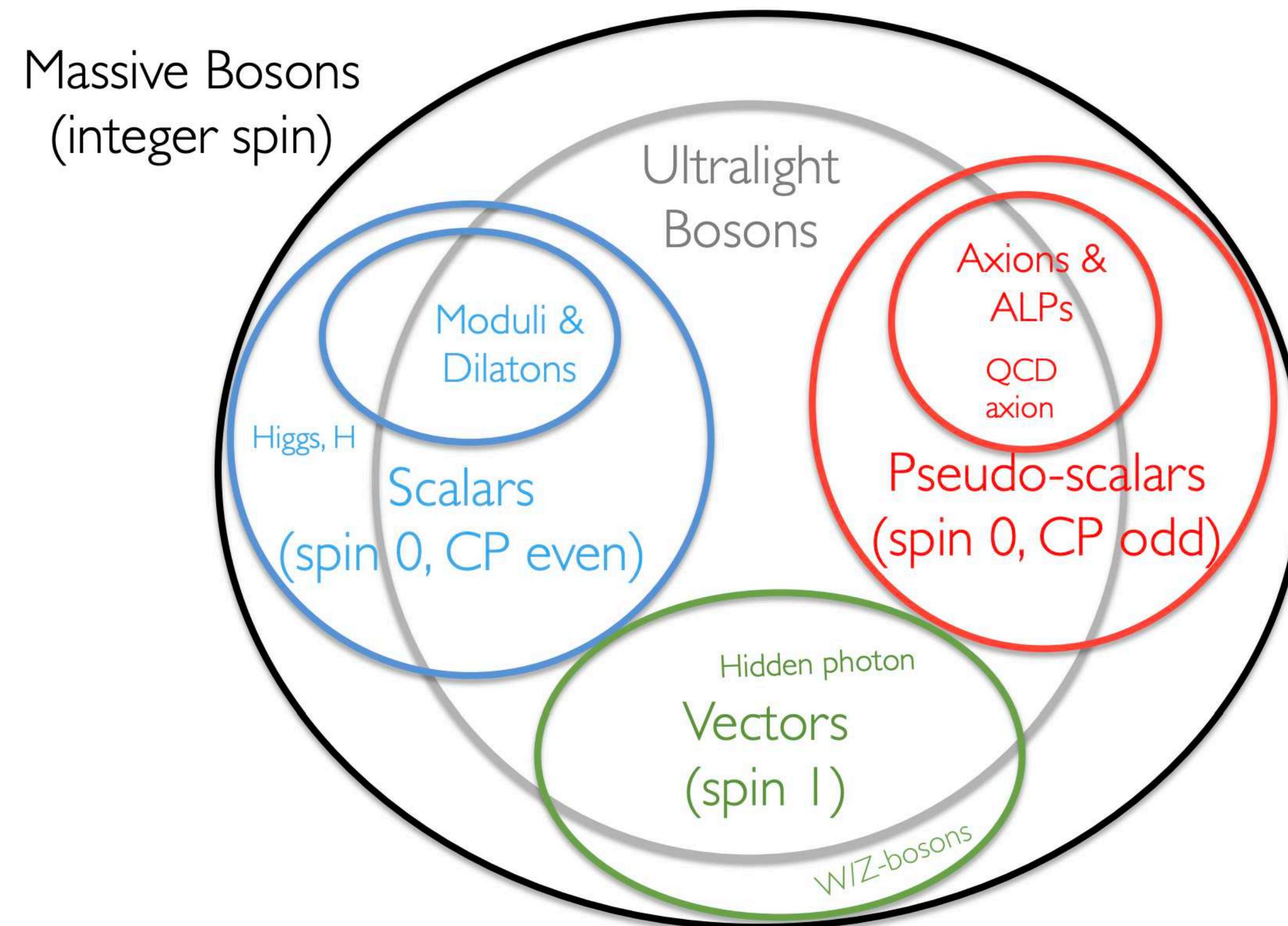
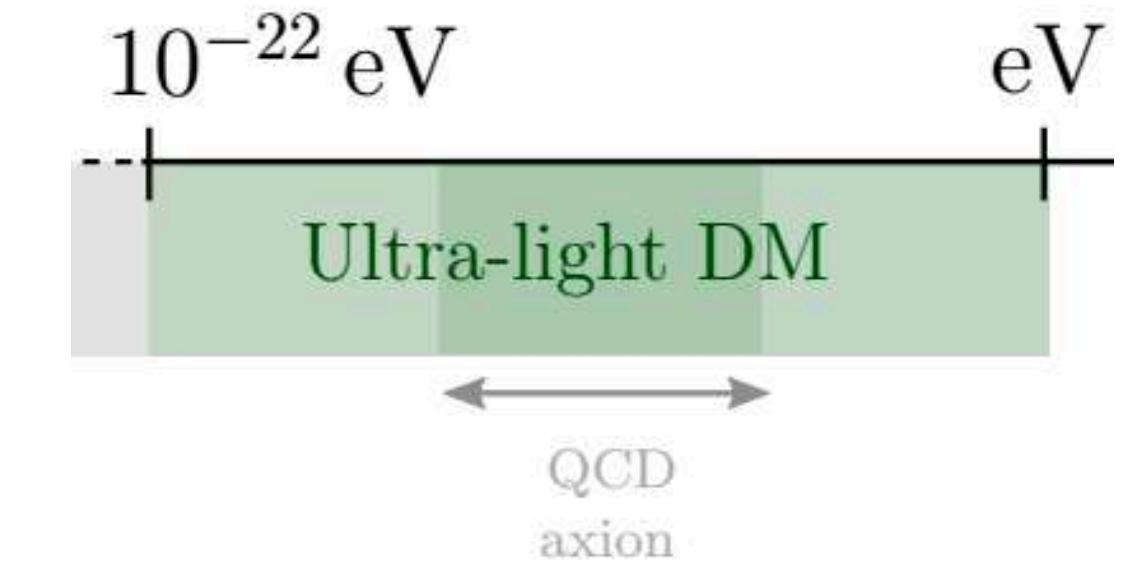
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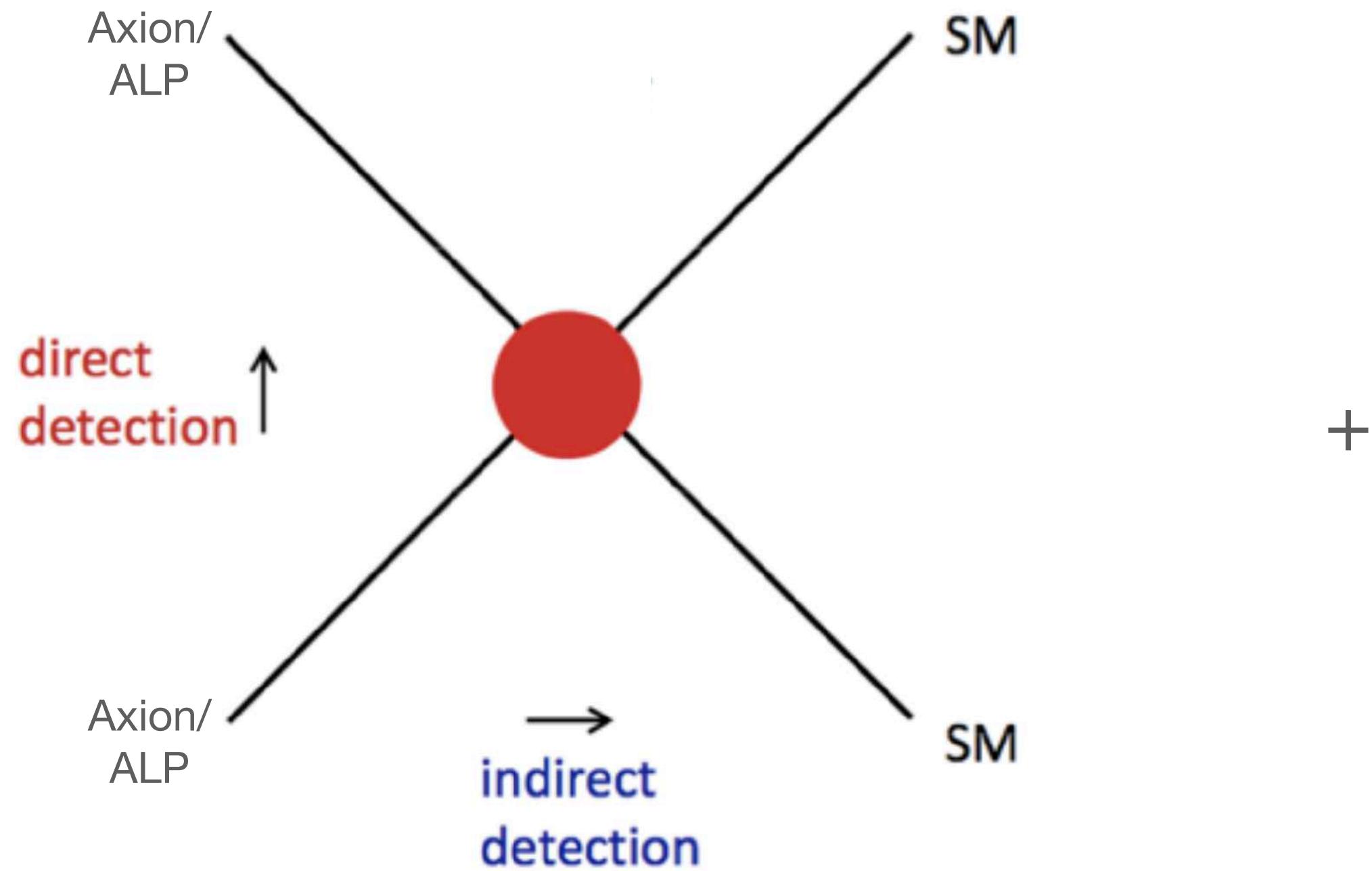
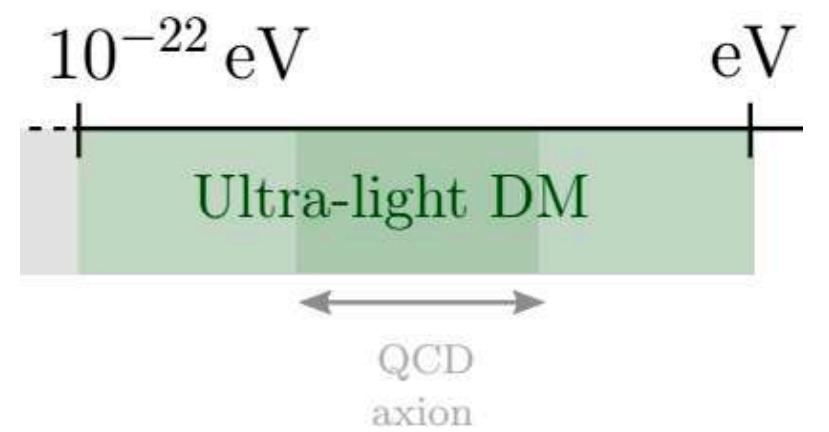
Motivation: particle physics

ULDM candidates

Many extensions of the Standard Model predict additional massive bosons

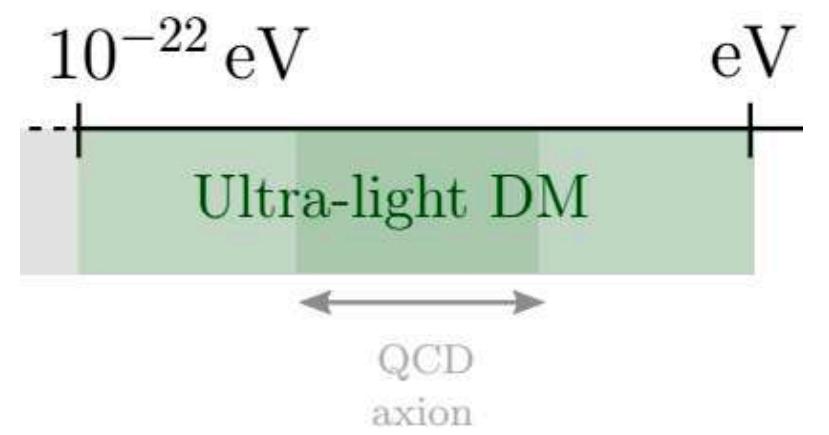


How to search for axions/ALPs?

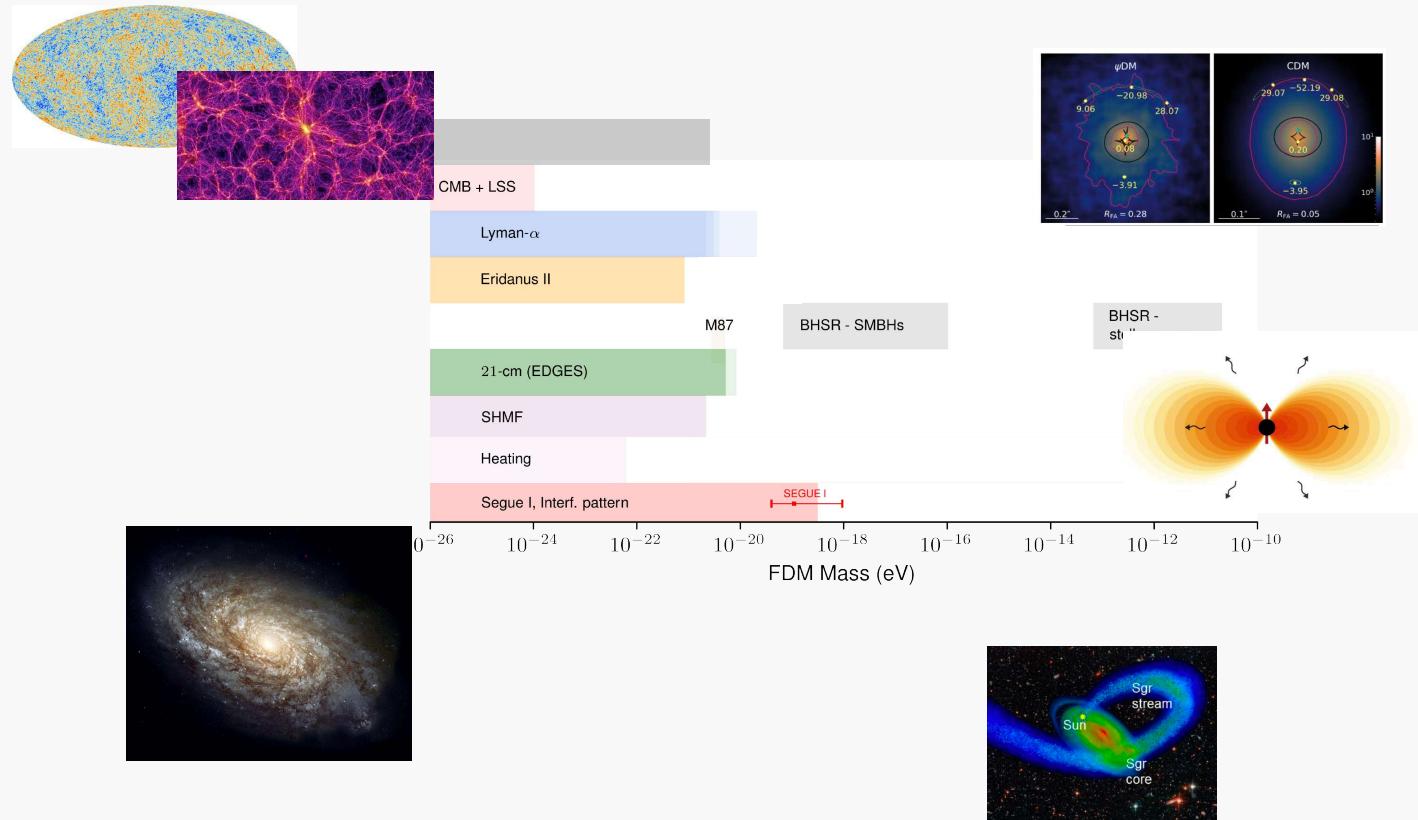


Gravitationally
Cosmological and astrophysical searches

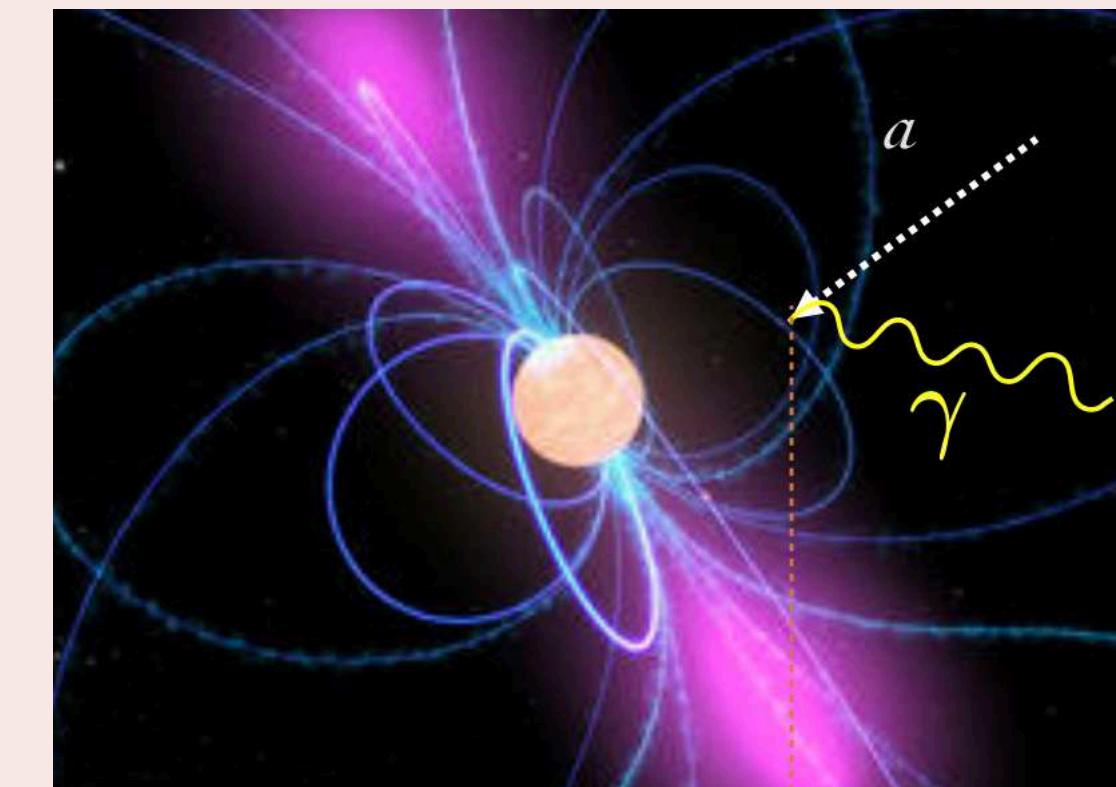
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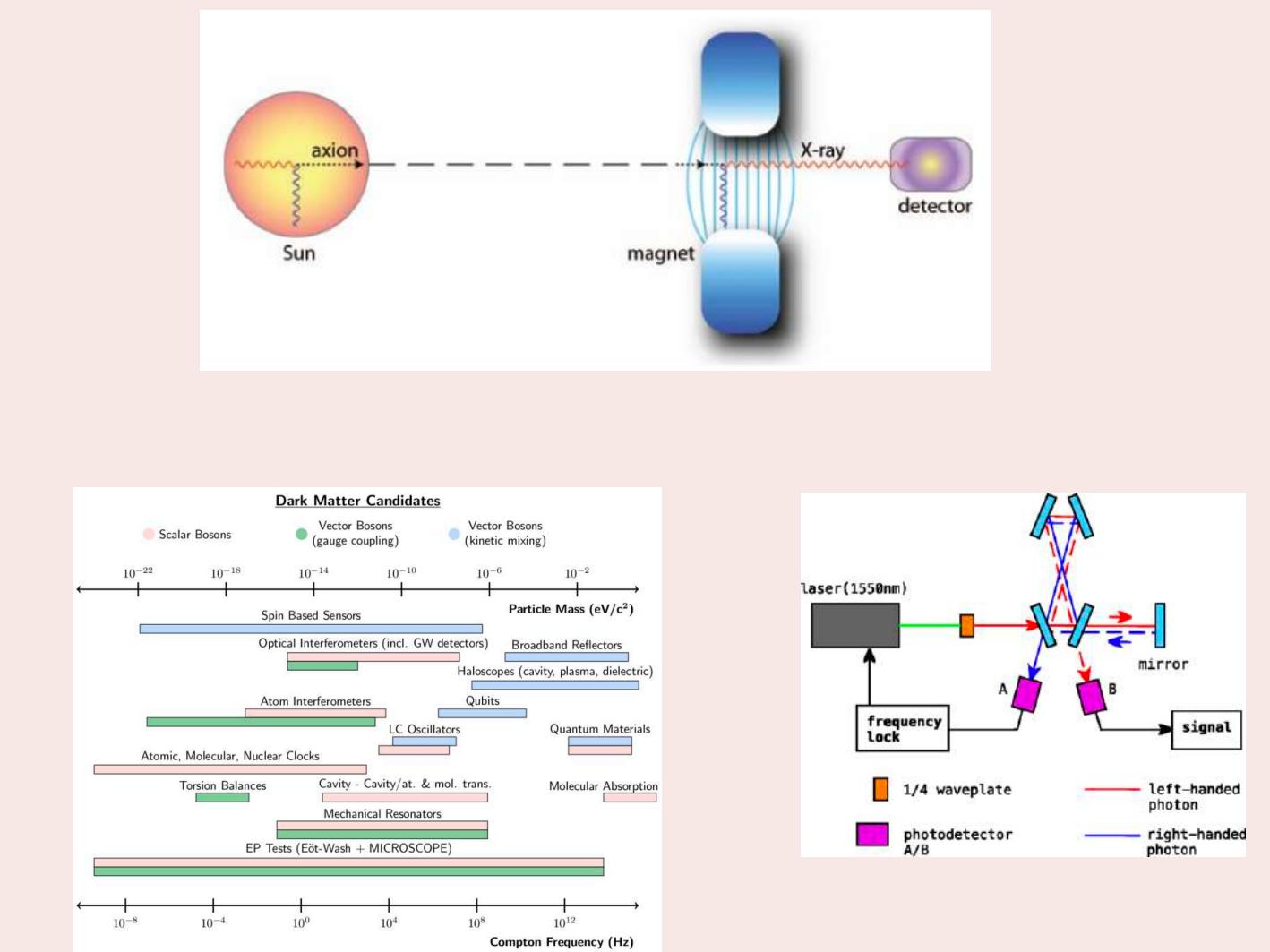
Cosmological and astrophysical searches



Indirect detection

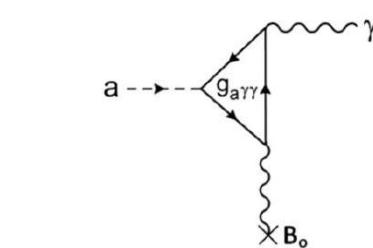
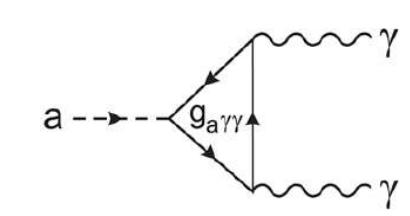


"Direct detection"
Axion/ALPs experiments

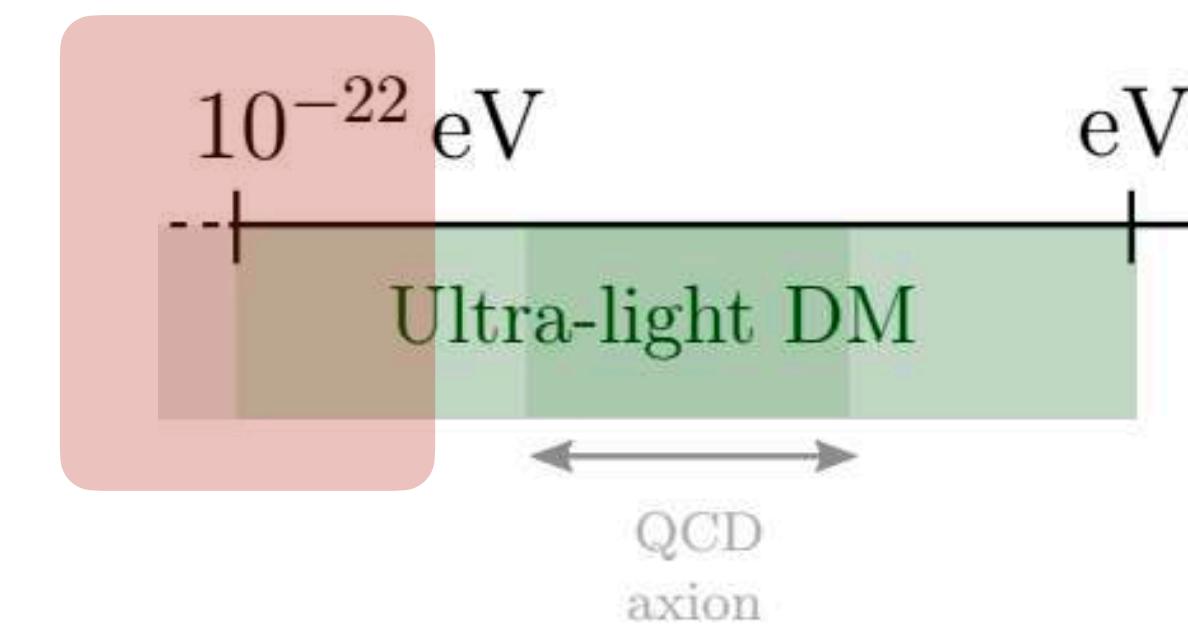


Gravitational

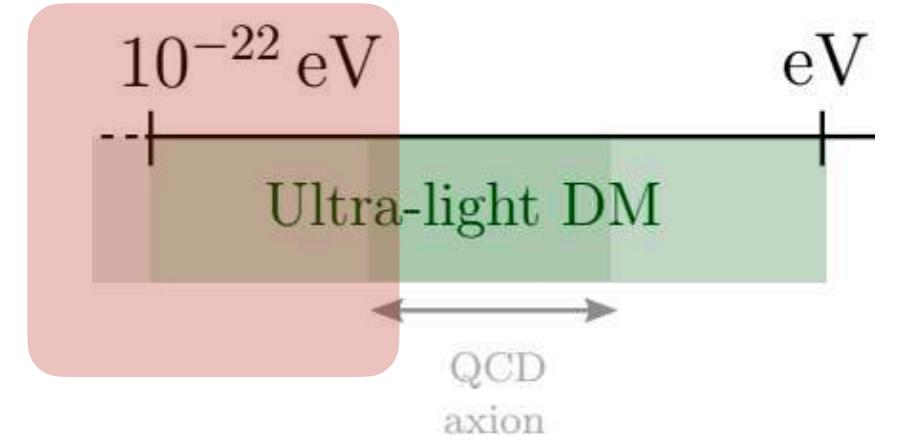
Interactions with the SM



Cosmological signatures



Ultra-light Dark Matter -classes



3 classes:

Fuzzy DM (FDM)

- Gravitationally bounded ultra-light scalar field model
- Condensation under gravity (BEC)

m

DOFs

Self Interacting FDM (SIFDM)

- Presence of (weakly) self-interaction
- Condensation under gravity + SI (superfluid)

$m \quad g$

DM Superfluid

- Forms a superfluid in galaxies
- MOND behaviour interior of galaxies

Axion and ALP (axion like particles)

$$i\dot{\psi} = \left(-\frac{1}{2m} \nabla^2 + \frac{g}{8m^2} |\psi|^2 - m\Phi \right) \psi$$

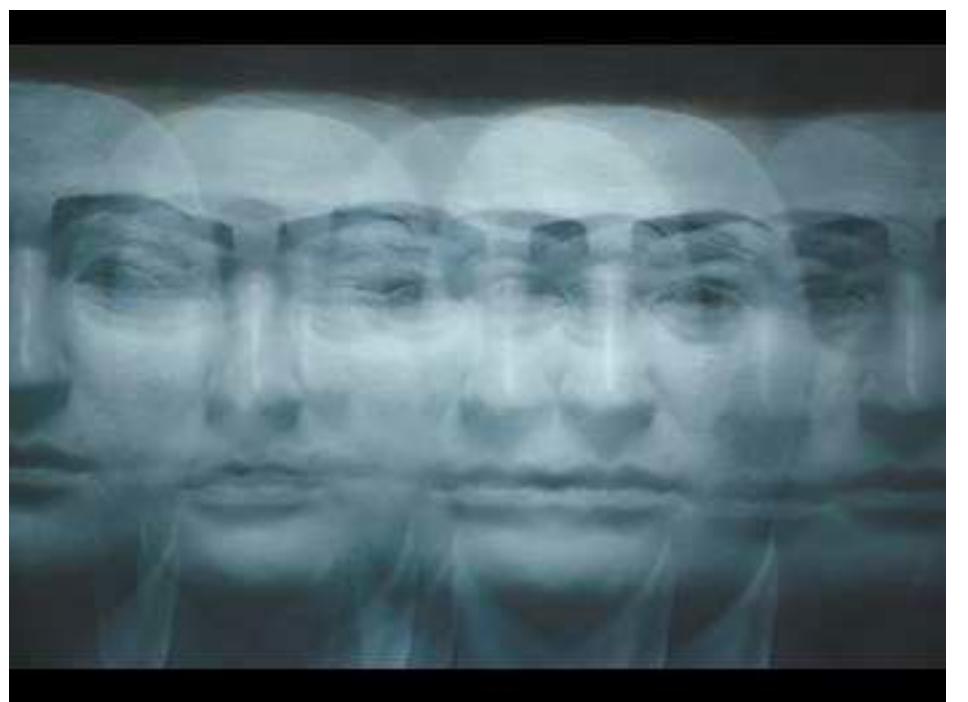
$$\mathcal{L} = P(X)$$

→ Connection with condensed matter and particle physics!

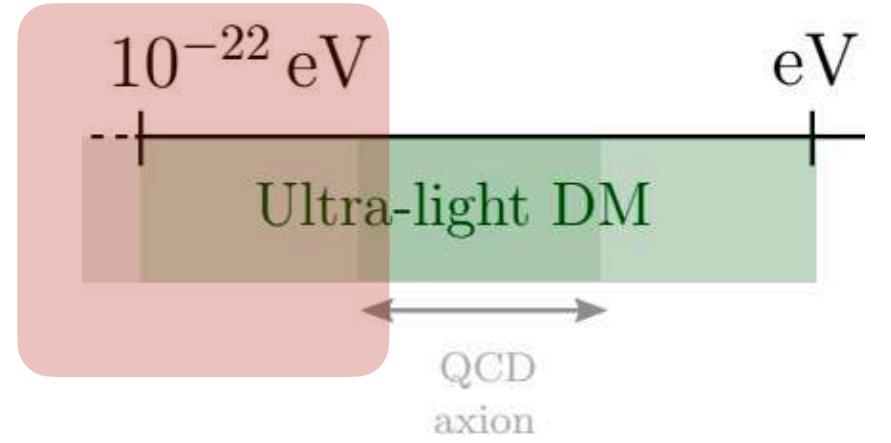
“*Ultra-light dark matter*”, **E.Ferreira**, 2020. The Astronomy and Astrophysics Review.

Fuzzy dark matter

Self interacting fuzzy dark matter



Fuzzy dark matter



Fuzzy DM (FDM)

- Gravitationally bounded ultra-light scalar field model
- Condensation under gravity (BEC)

m

Wave DM Ultra-light axions

Self Interacting FDM (SIFDM)

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Hu W, Barkana R, Gruzinov A (2000 a,b)

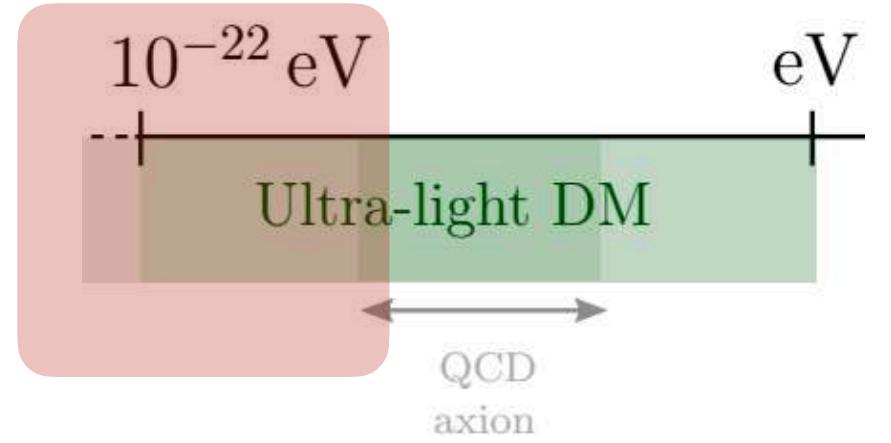
(Reviews: EF (2021), J. Niemeyer (2019), L. Hui (2021))

Idea:

$$m_{\text{fdm}} \sim 10^{-22} \text{ eV}$$

address the small scale problems+ rich phenom.

Fuzzy dark matter



Fuzzy DM (FDM)

- Gravitationally bounded ultra-light scalar field model
- Condensation under gravity (BEC)

m

Wave DM Ultra-light axions

Focus in spin 0 particles here!

(Some of the grav. phenom. is carried for vectors, for example)

- Spin 0 - FDM
- Spin 1 - Vector FDM
- Higher spin FDM

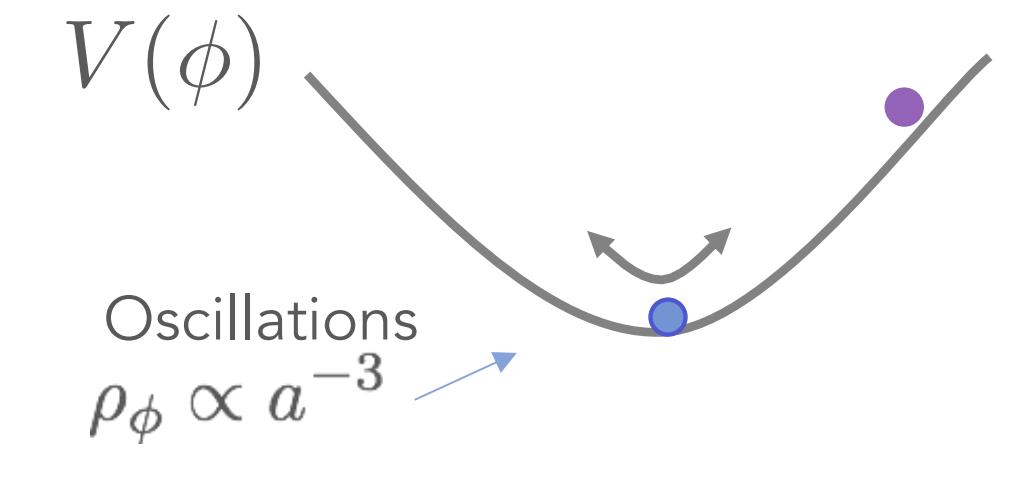
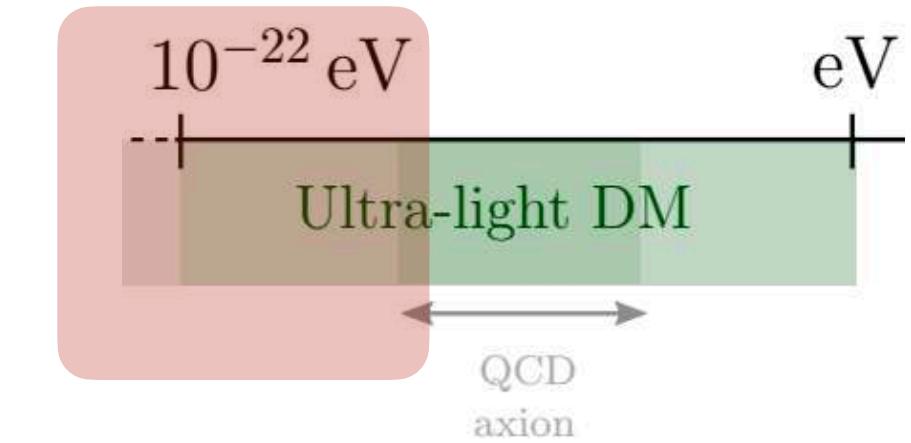
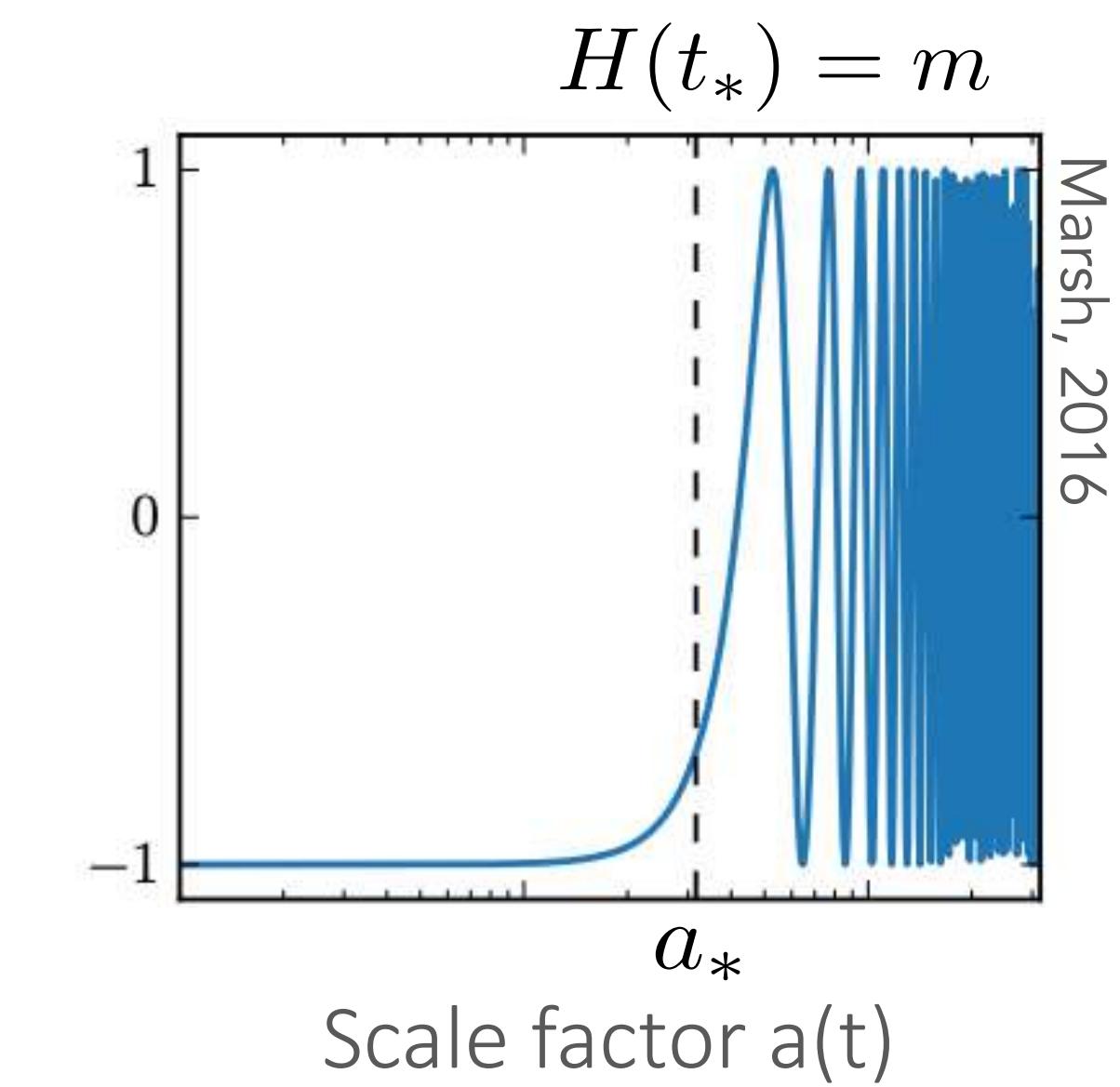
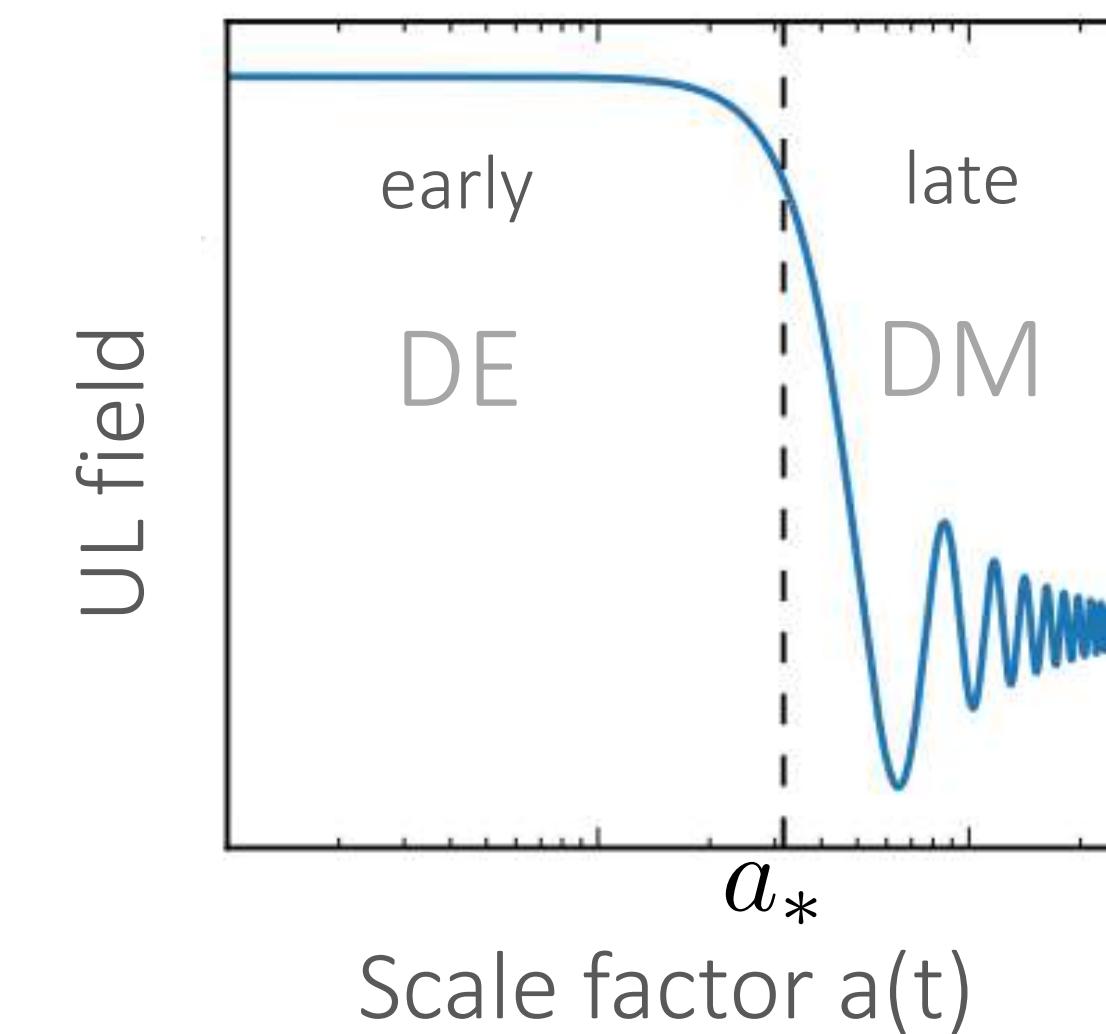
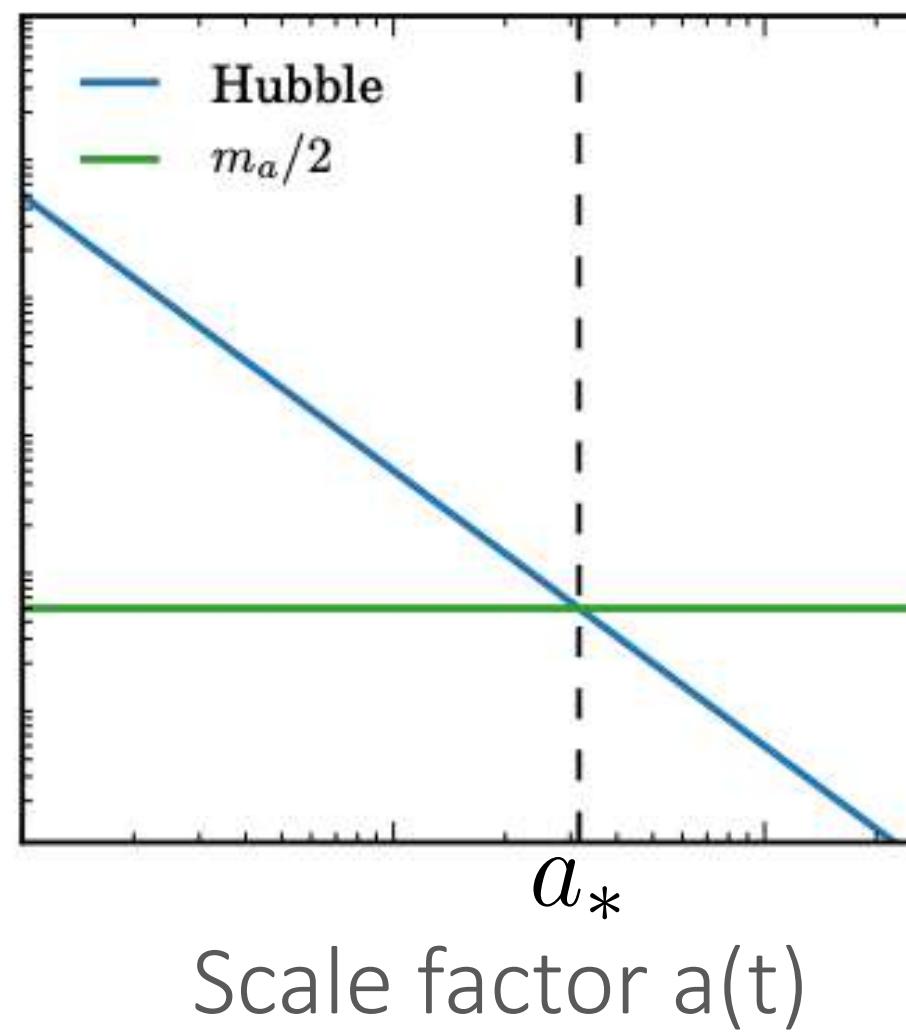
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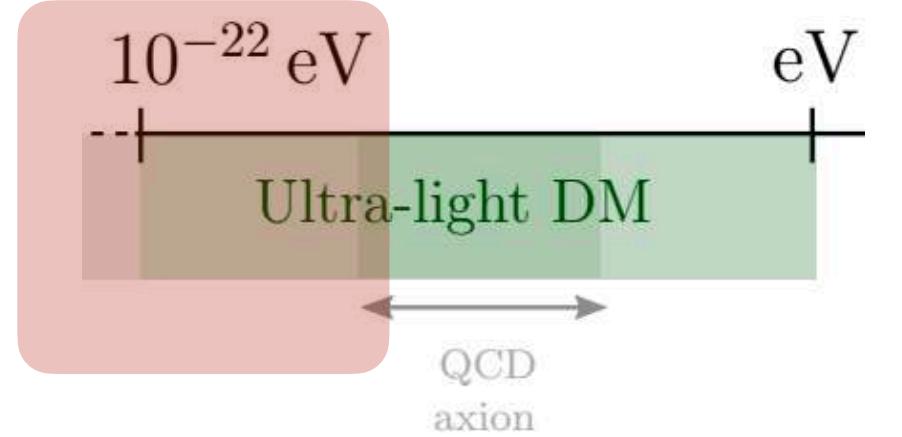
Cosmological evolution

$$\ddot{\phi} + 3H\dot{\phi} + m^2\phi = 0$$

$$\left[\begin{array}{ccc} H \gg m & \xrightarrow{\quad} & \phi_{\text{early}} = \phi(t_i) & \xrightarrow{\quad} & \omega = -1 & \text{DE} \\ H \ll m & \xrightarrow{\quad} & \phi_{\text{late}} \propto e^{imt} & \xrightarrow{\quad} & \langle \omega \rangle = 0 & \text{DM} \end{array} \right]$$



Structure formation - non-relativistic regime



Evolution on small scales: take non-relativistic regime of the theory, relevant for structure formation.

Schrödinger-Poisson system : describe the FDM and the SIFDM

$$\left\{ \begin{array}{l} i\dot{\psi} = \left(-\frac{1}{2m}\nabla^2 + \frac{g}{8m^2}|\psi|^2 - m\Phi \right) \psi \\ \nabla^2\Phi = 4\pi G(m|\psi|^2 - \bar{\rho}) \end{array} \right.$$

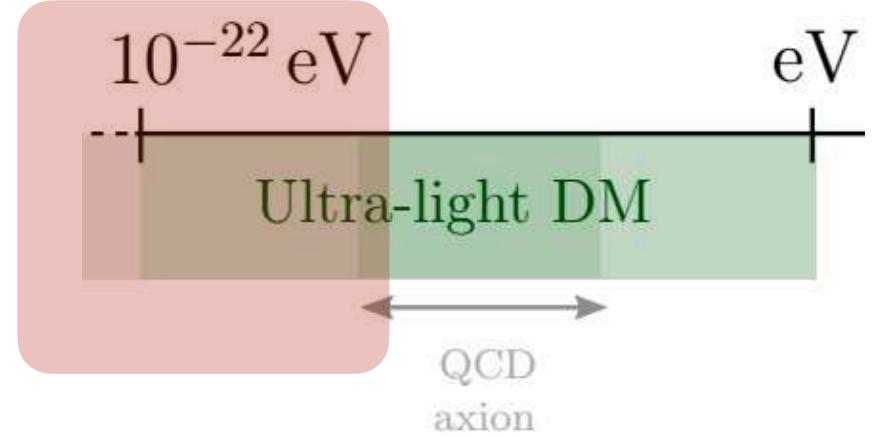
Schrödinger equation
(Gross-Pitaevskii)

Poisson equation

$g = 0 \longrightarrow$ FDM
 $g \neq 0 \longrightarrow$ SIFDM

Fundamentally different than
CDM/WDM/SIDM!

Structure formation - non-relativistic regime



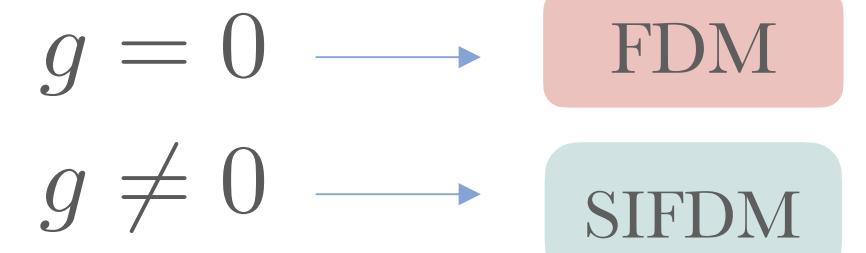
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Schrödinger equation
(Gross-Pitaevskii)

Poisson equation



Fundamentally different than
CDM/WDM/SIDM!

Madelung equations $(\psi \equiv \sqrt{\rho/m} e^{i\theta} \text{ and } \mathbf{v} \equiv \nabla\theta/m)$

$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\dot{\mathbf{v}} + (\mathbf{v} \cdot \nabla)\mathbf{v} = -\frac{1}{m} \left(V_{grav} - P_{int} - \frac{1}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right)$$

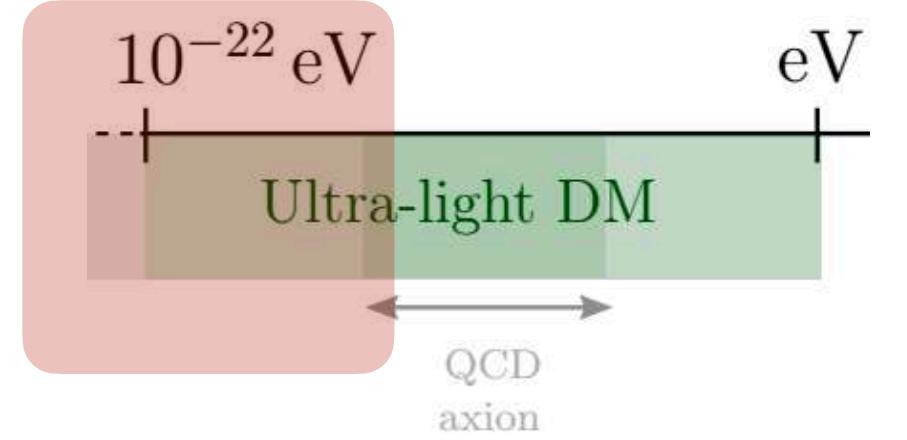
$$P_{int} = K\rho^{(j+1)/j} = \frac{g}{2m^2}\rho^2$$

$$\frac{1}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}}$$

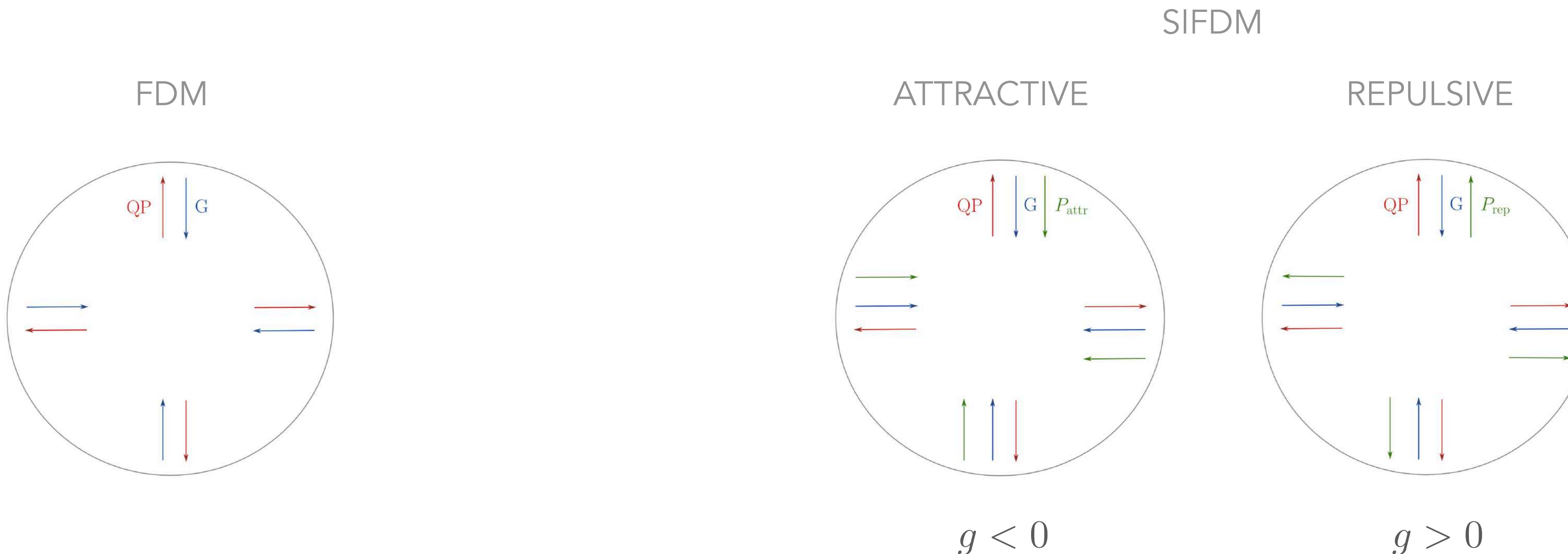
Quantum pressure

FLUID
DESCRIPTION

Structure formation - perturbation and stability



Competition between gravity and pressure (quantum pressure and interaction)



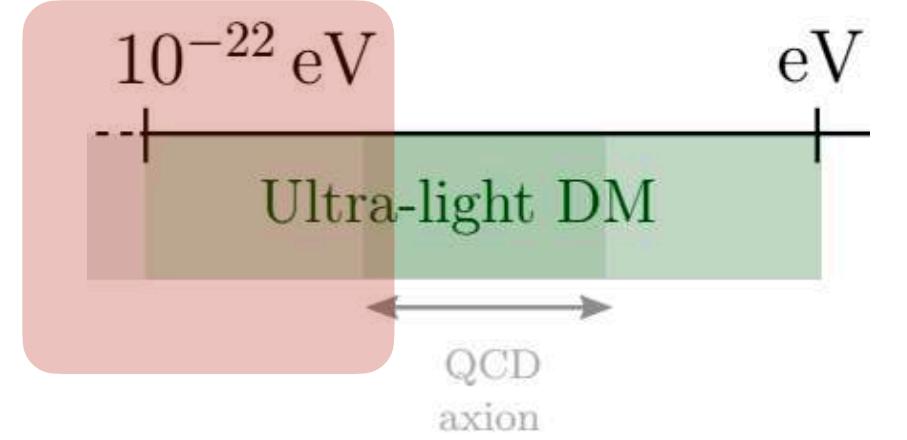
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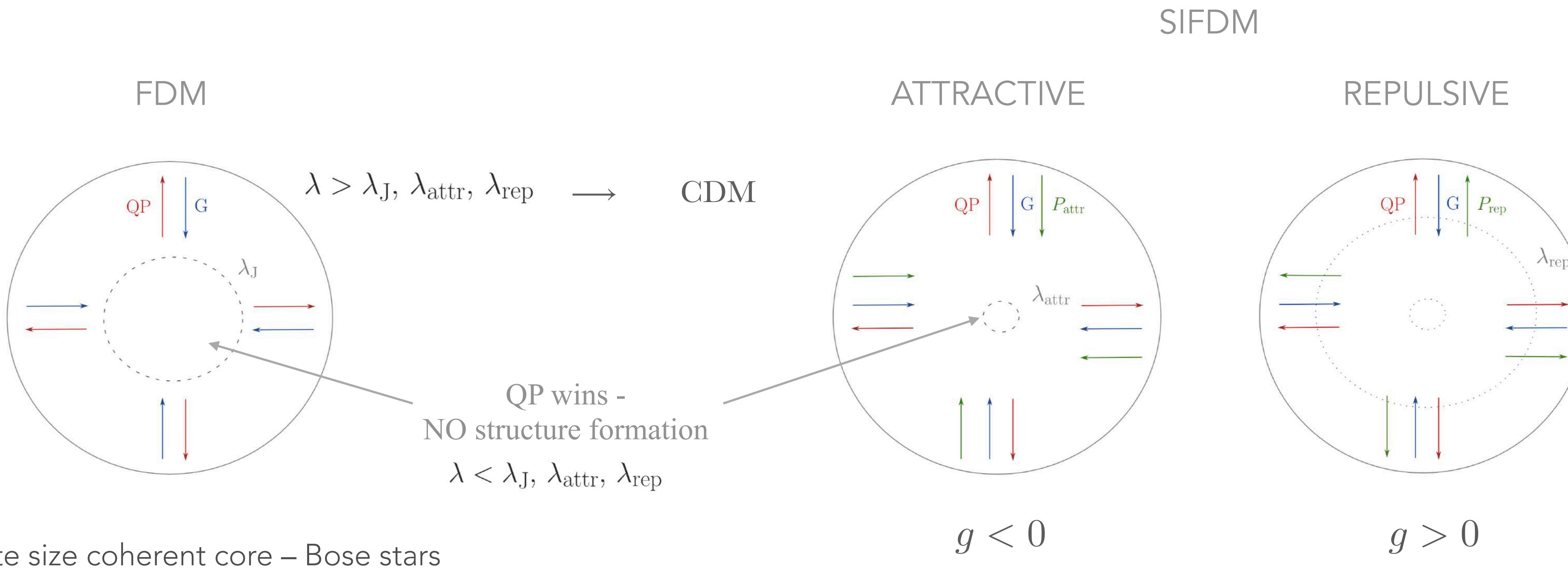
$$P_{int} = \frac{g}{2m^2} \rho^2$$

Quantum pressure

Structure formation - perturbation and stability



Finite clustering scale - no structure formation on small scales



$$\lambda_J = 55 \left(\frac{m}{10^{-22} \text{ eV}} \right)^{-1/2} \left(\frac{\rho}{\bar{\rho}} \right)^{-1/4} (\Omega_m h)^{-1/4} \text{ kpc}$$

$$m \leq 10^{-20} \text{ eV} \Rightarrow \lambda_{dB} > \mathcal{O}(\text{kpc})$$

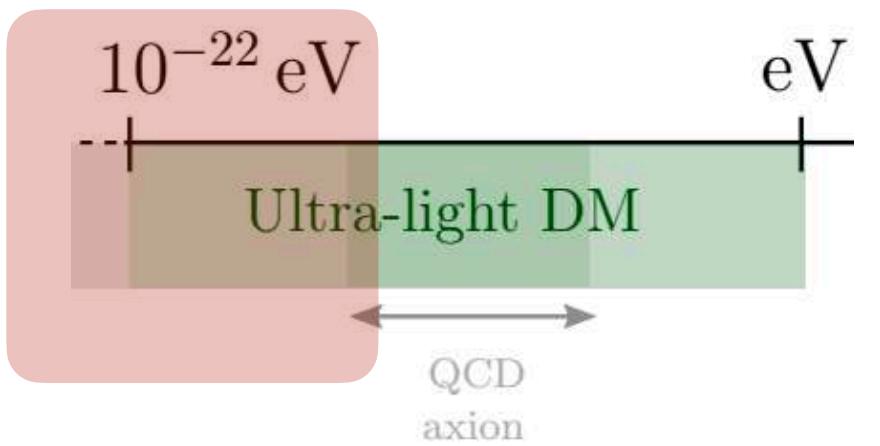
Galactic scales

For **attractive** interactions can only form **localized clumps** (solitons)

QCD axion: $m \sim 10^{-5} \text{ eV}$
 $\lambda_a \sim -10^{-48}$ $\rightarrow l_{\text{soliton}} \sim 10^{-5} \text{ kpc}$

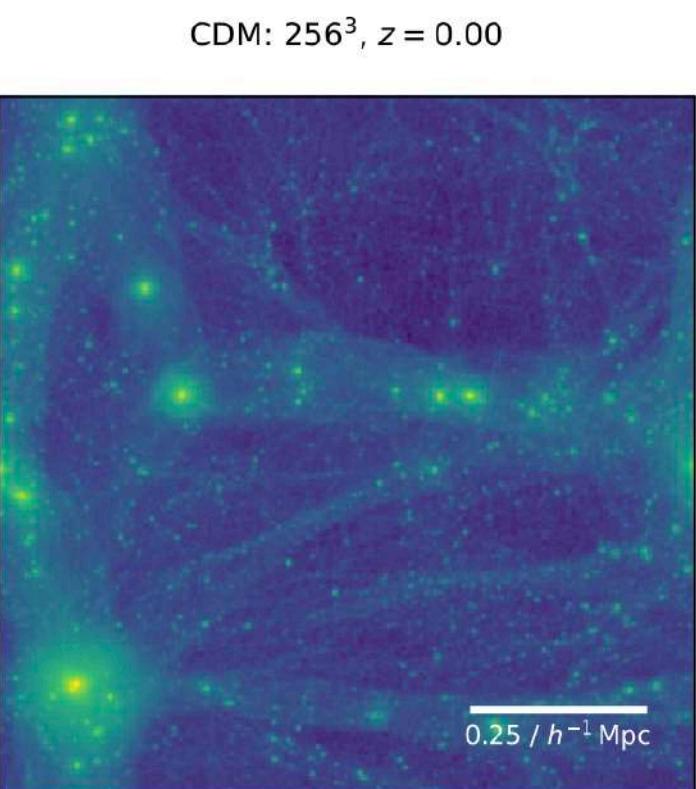
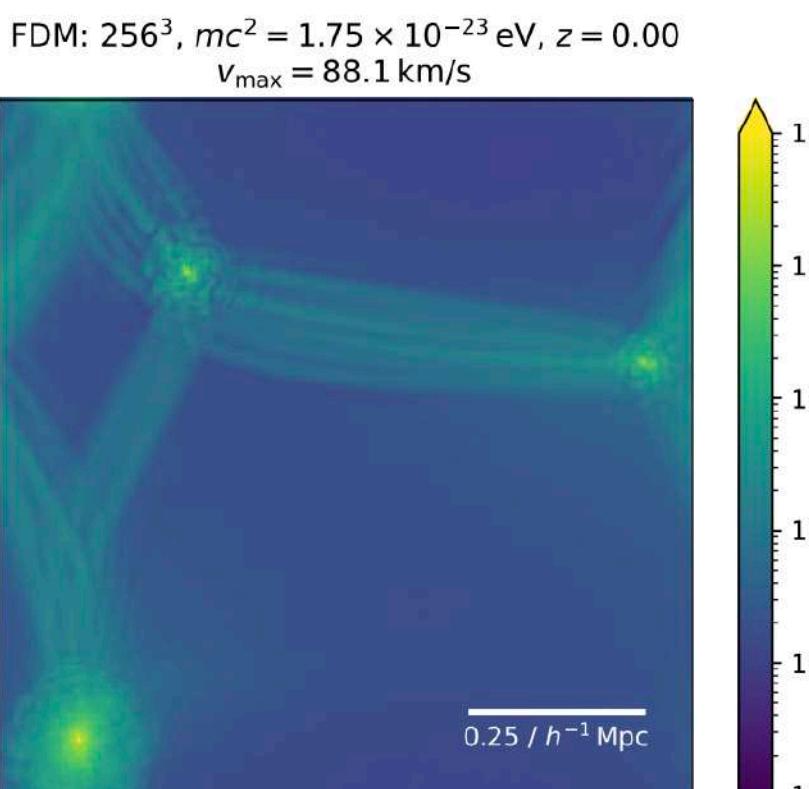
Phenomenology

RICH PHENOMENOLOGY ON SMALL SCALES



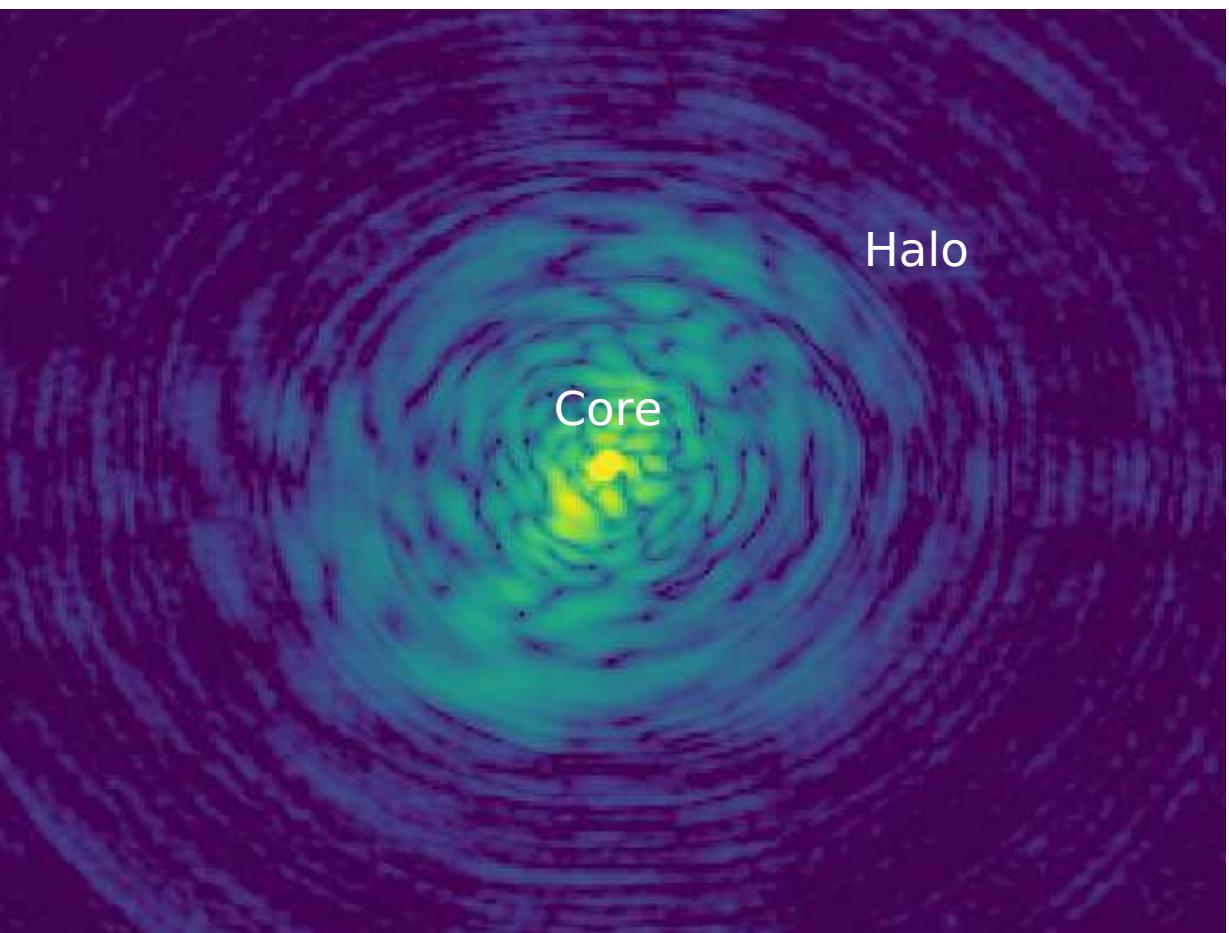
* Focus only in gravitational signatures

Suppression of small structures

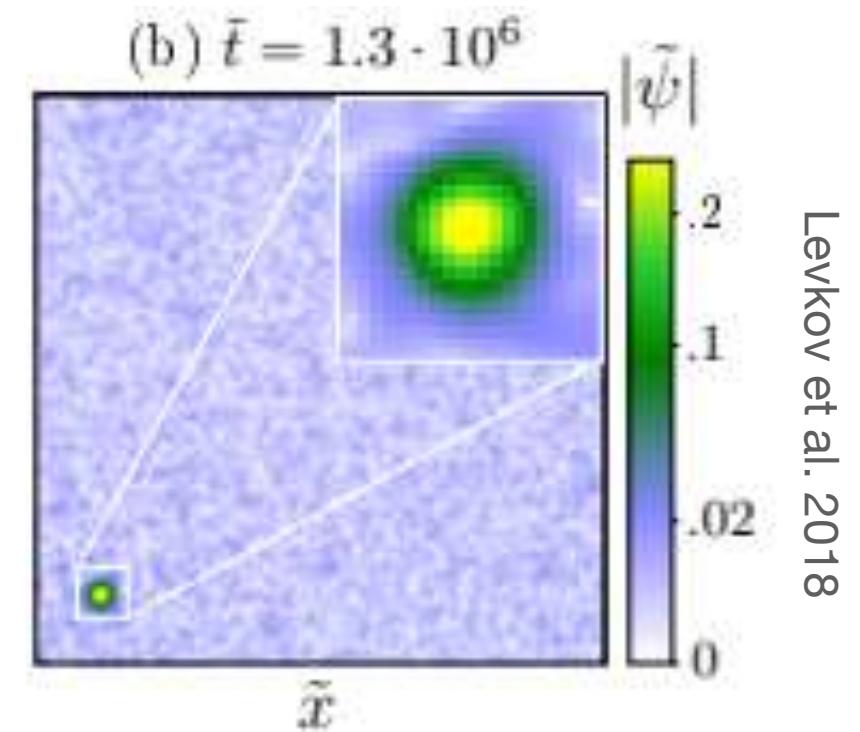


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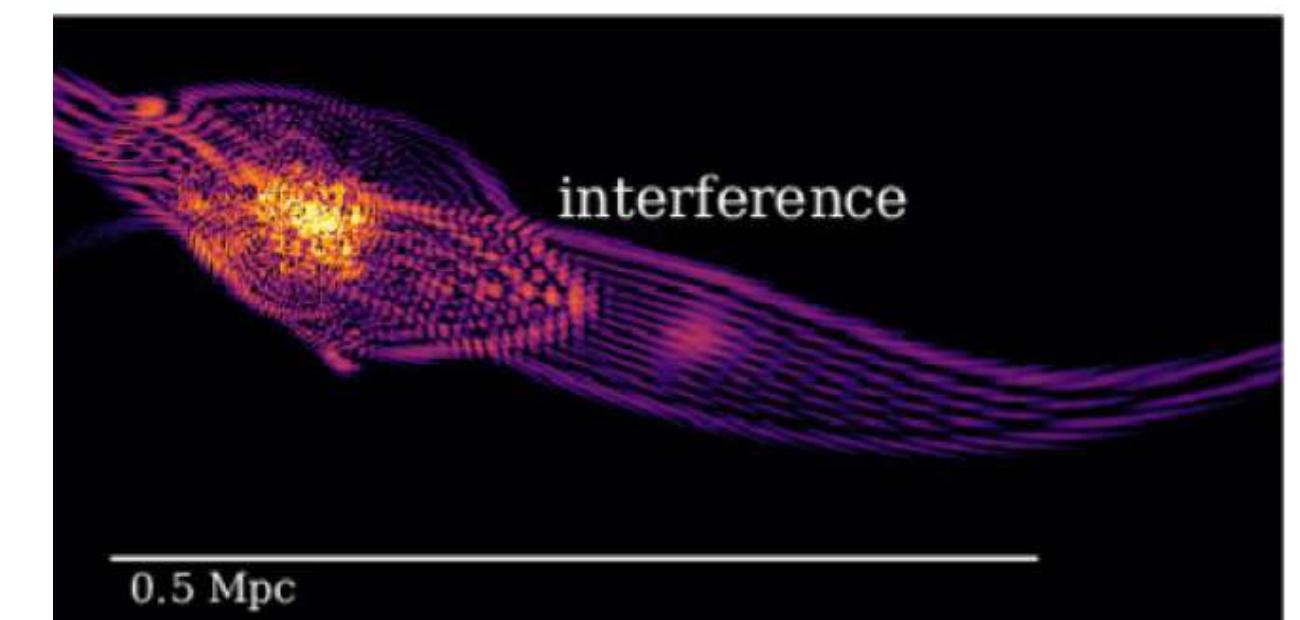
Formation of a solitonic core



Dynamical effects

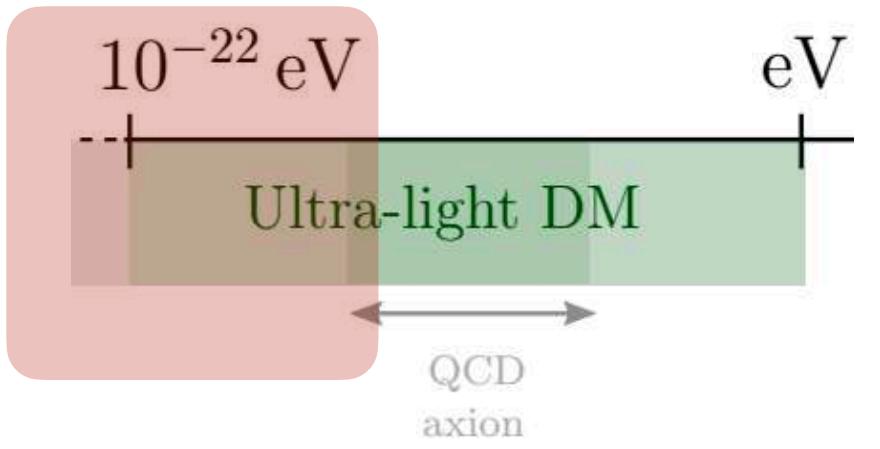


Wave interference

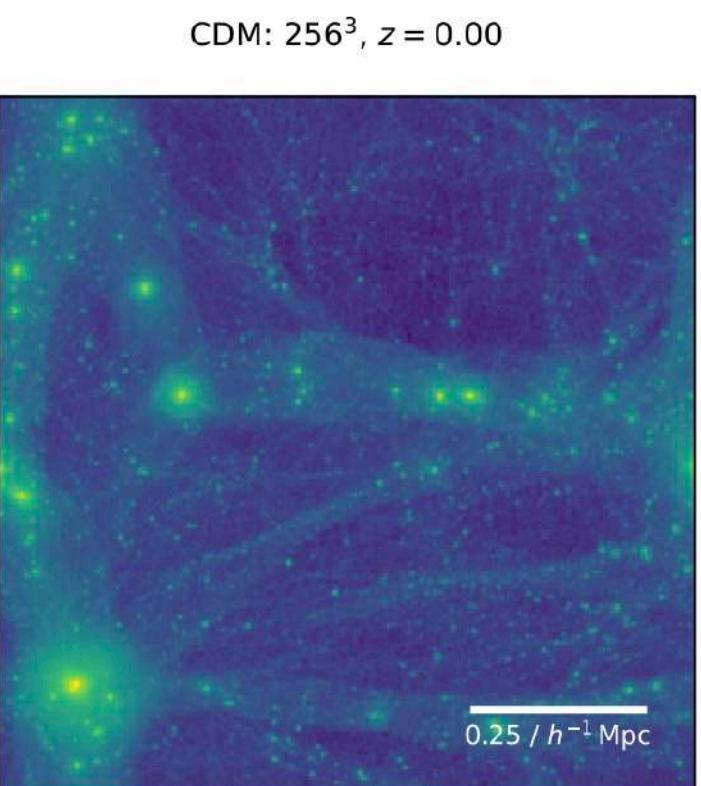
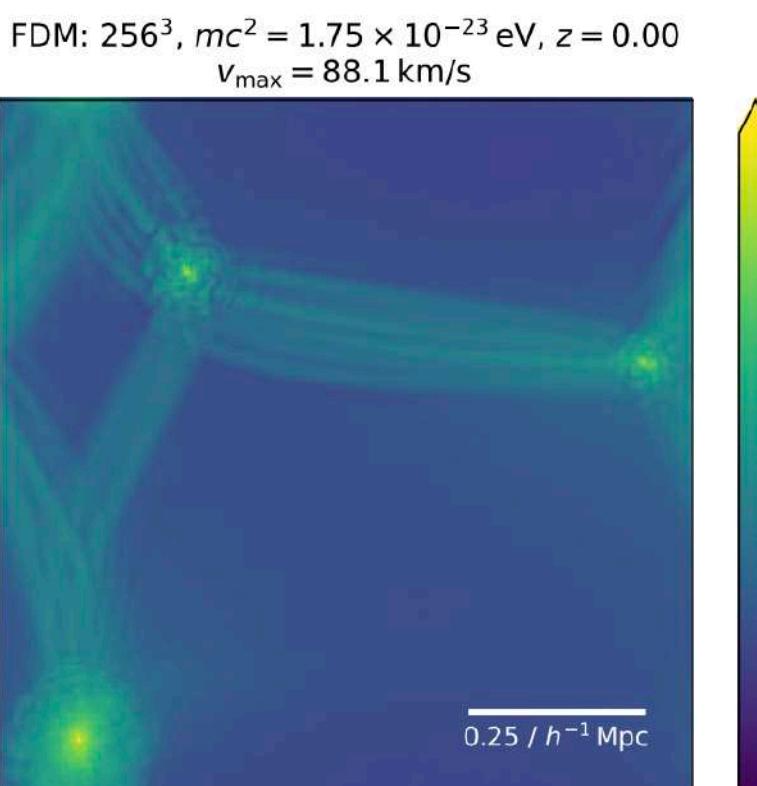


Phenomenology

RICH PHENOMENOLOGY ON SMALL SCALES

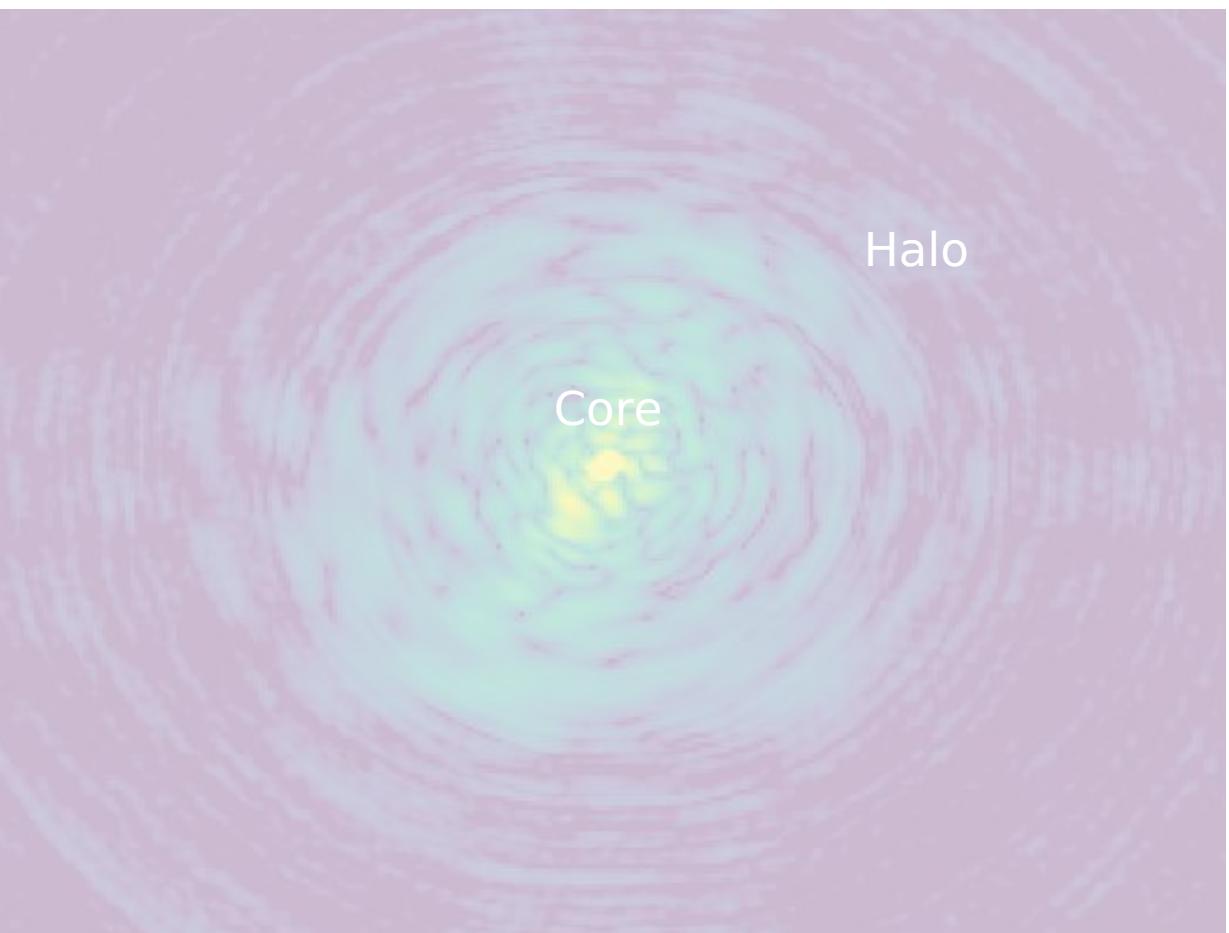


Suppression of small structures

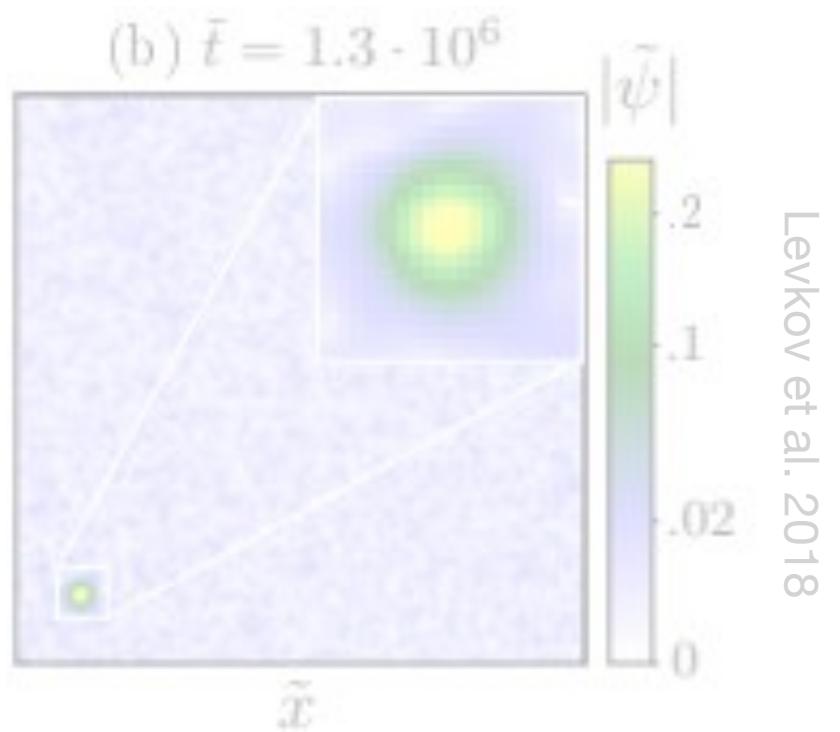


S. May et al. 2021

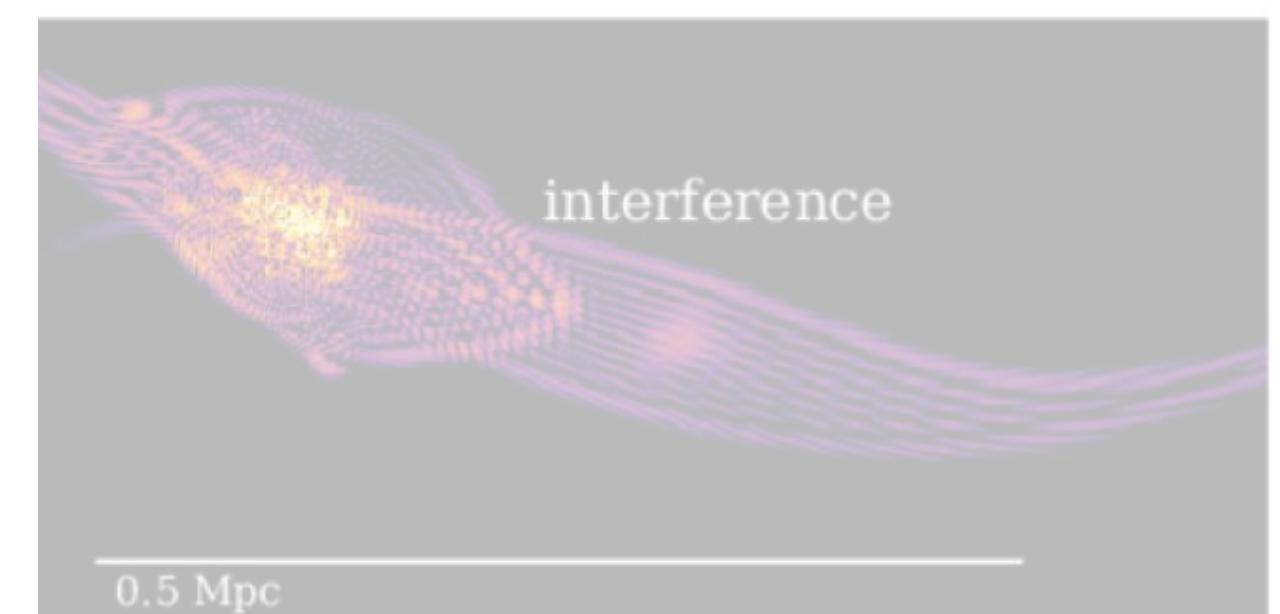
Formation of a solitonic core



Dynamical effects



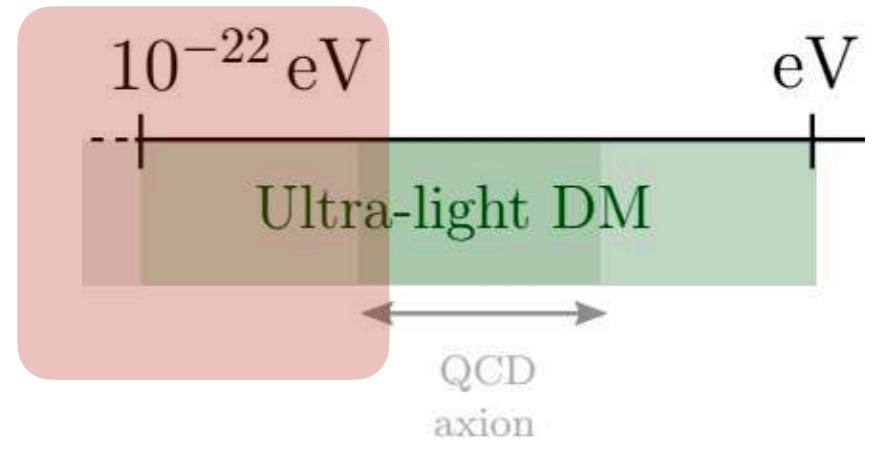
Wave interference



Mocz et al. 2017

Phenomenology

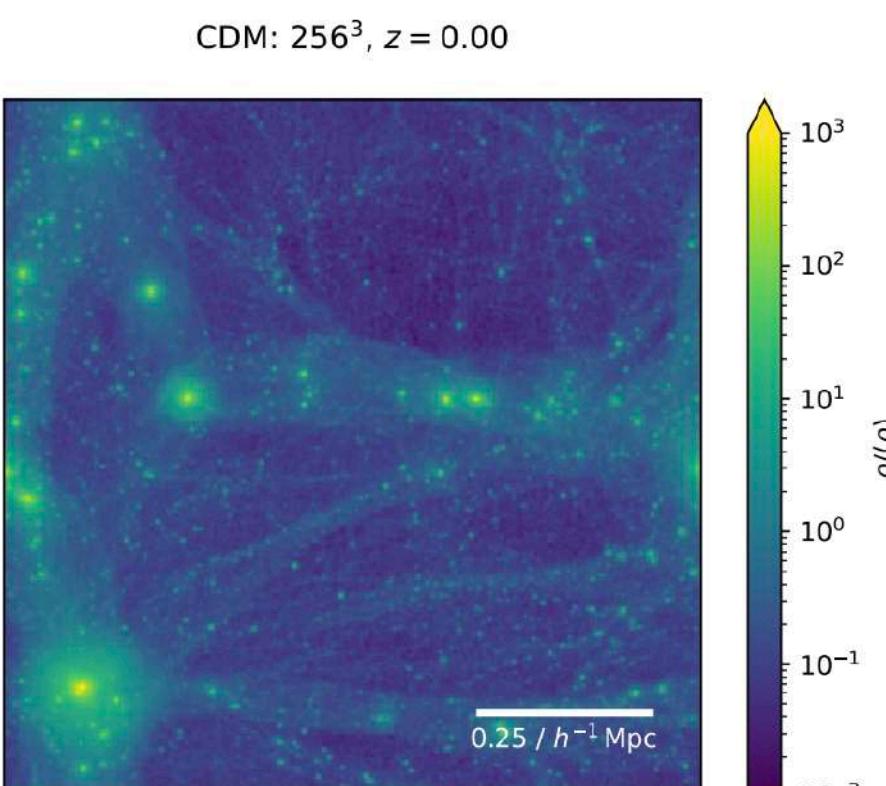
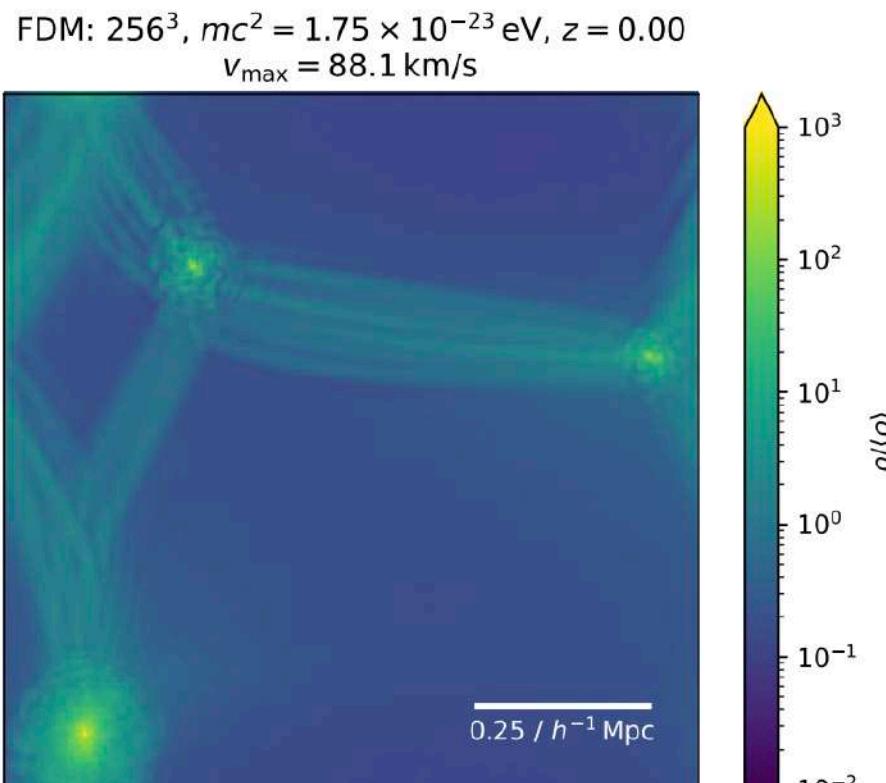
Suppression of small structures



Finite Jeans length λ_J or $\lambda_{\text{attr}}, \lambda_{\text{rep}}$

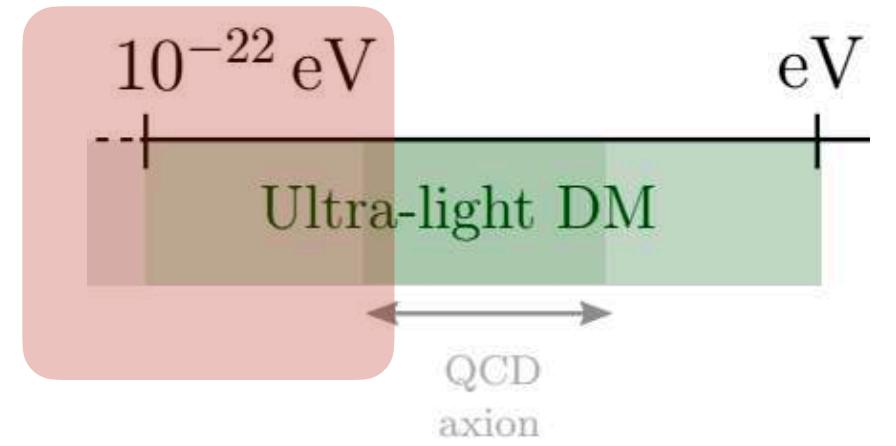


No small scale structure



Phenomenology

Suppression of small structures

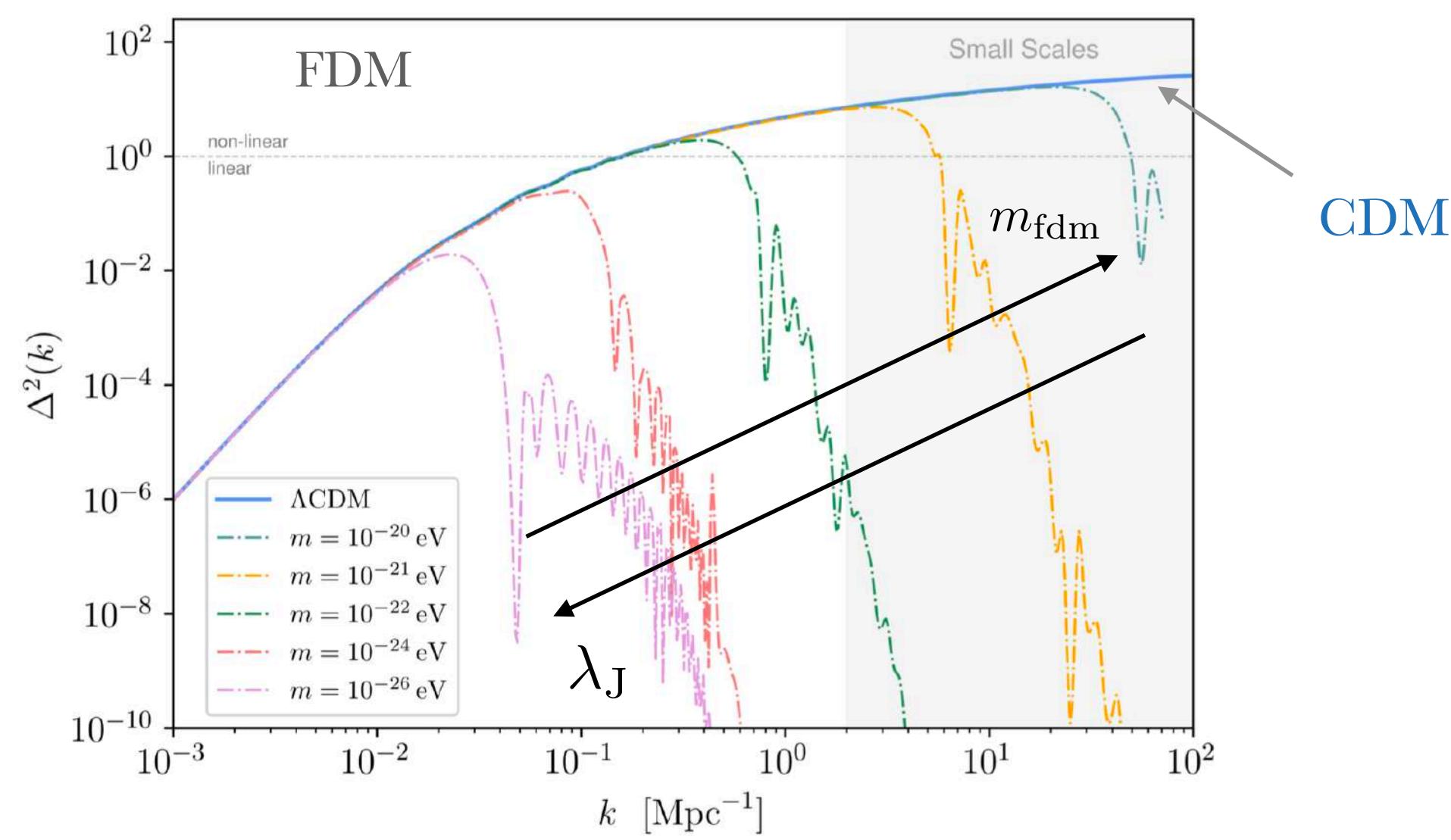


Finite Jeans length λ_J or $\lambda_{\text{attr}}, \lambda_{\text{rep}}$

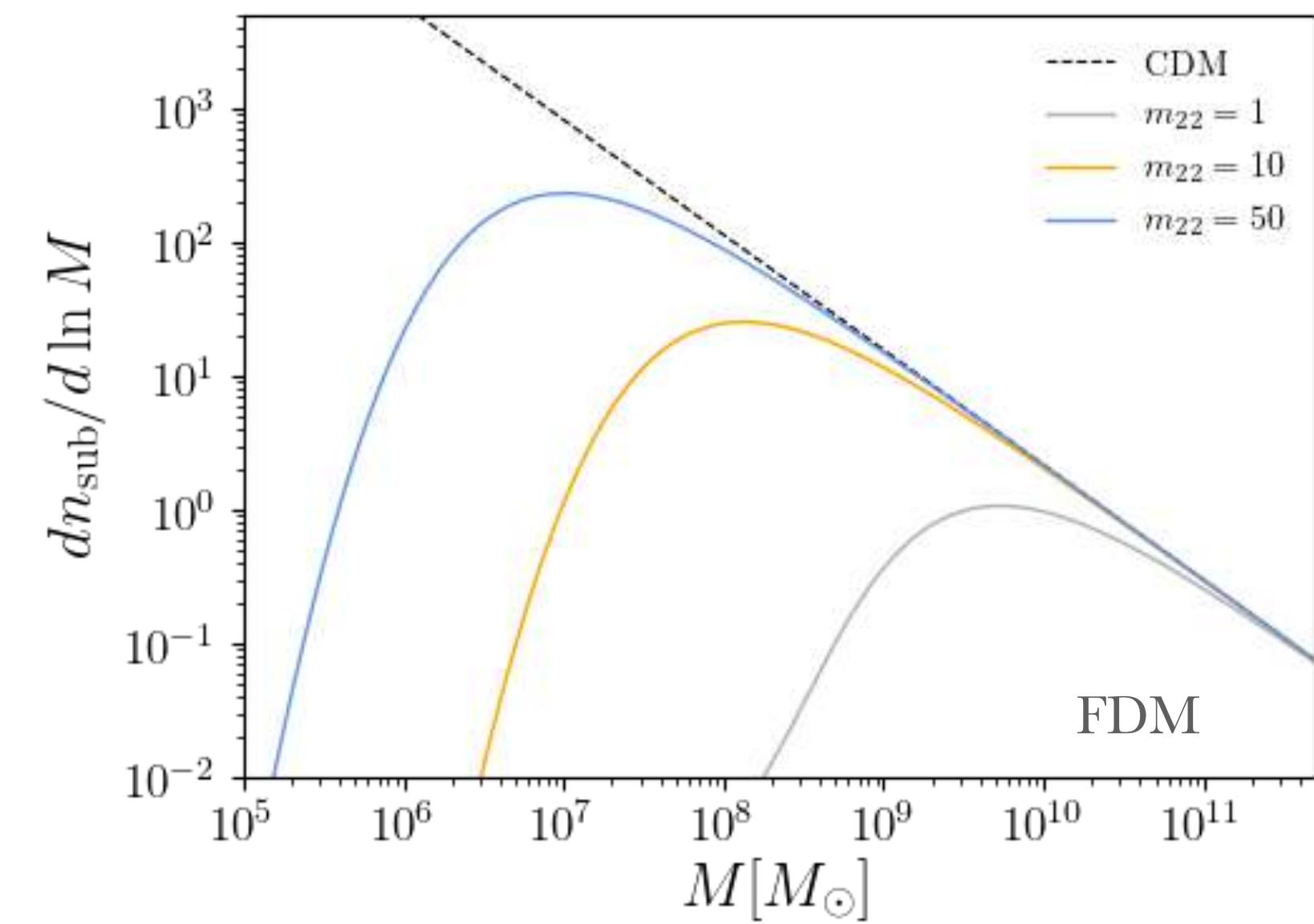


Suppresses small scale structure

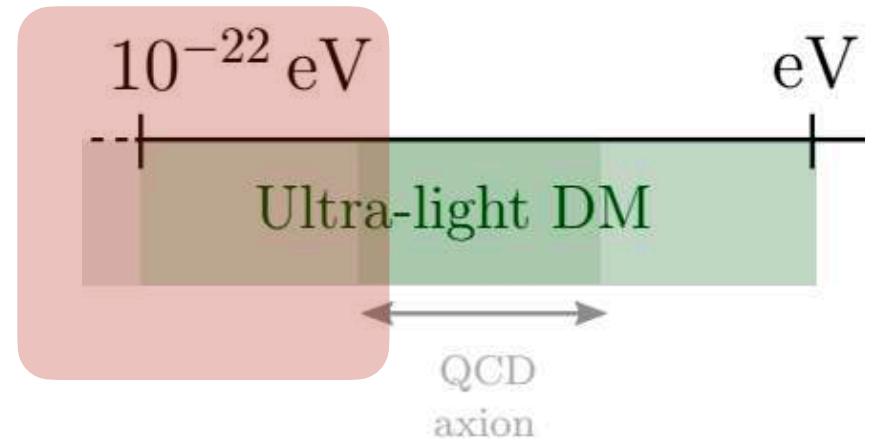
POWER SPECTRUM



(sub) HALO MASS FUNCTION

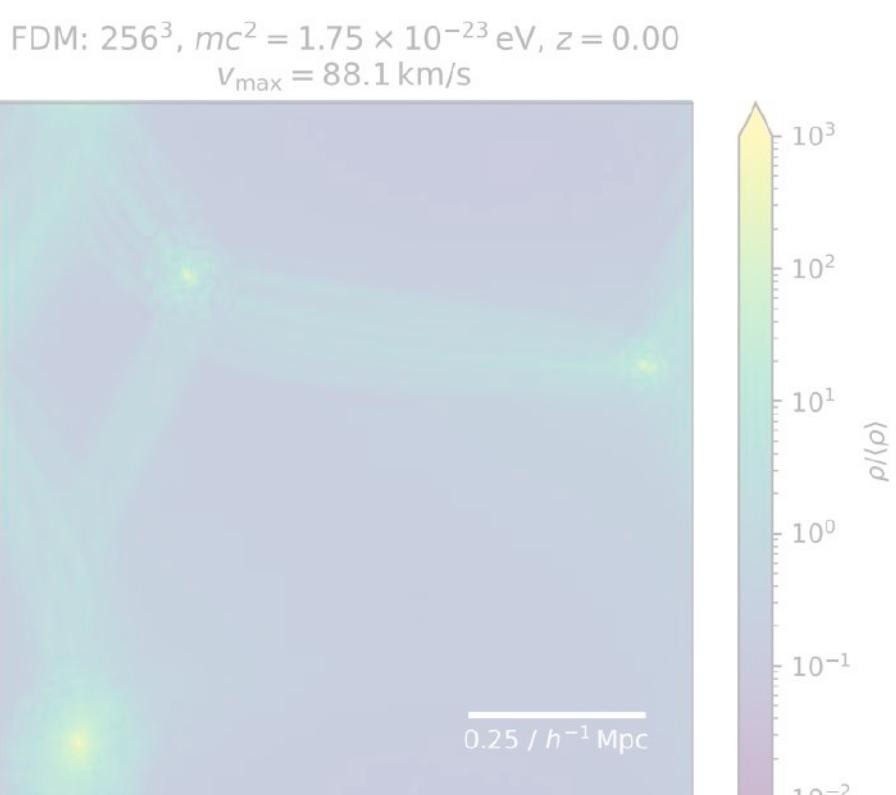


Phenomenology



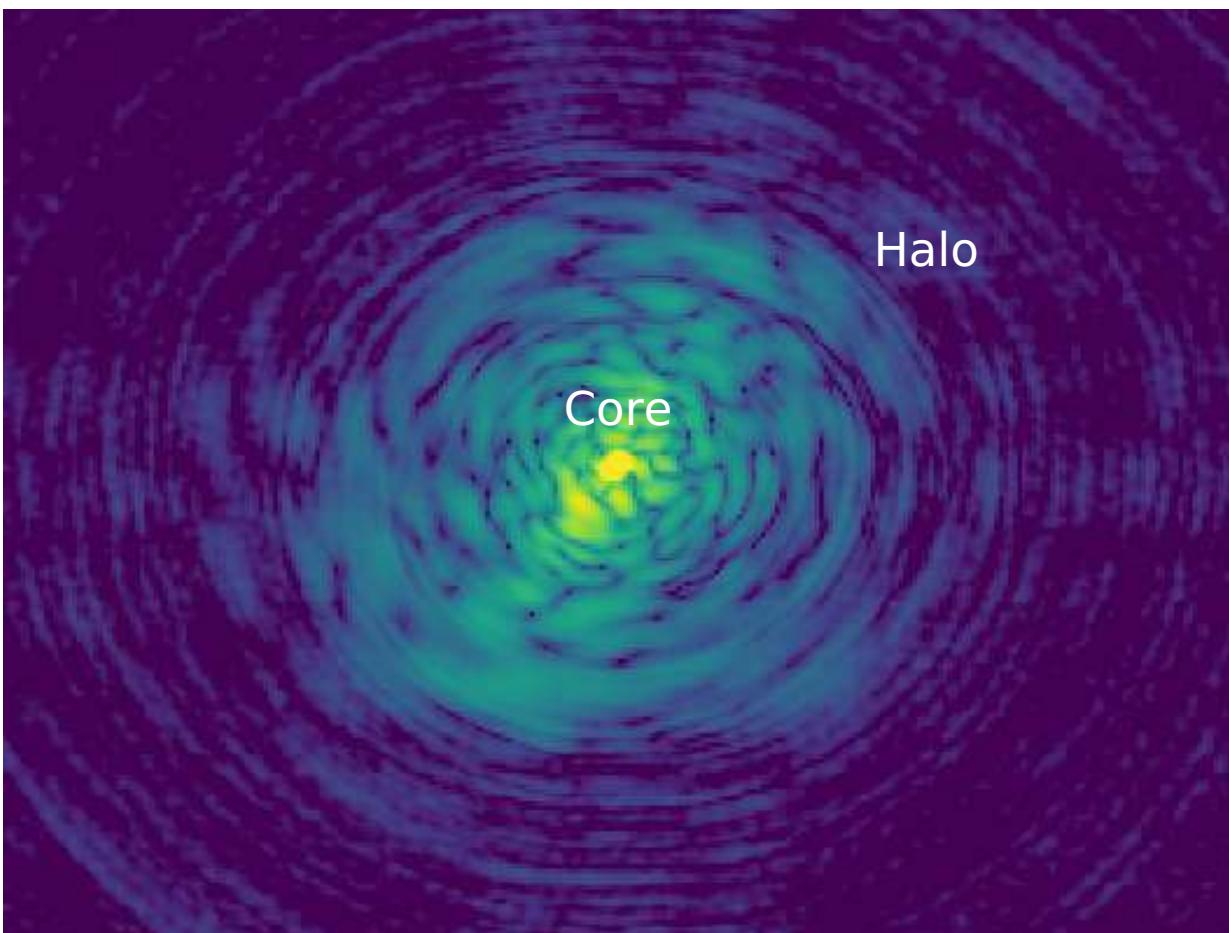
RICH PHENOMENOLOGY ON SMALL SCALES

Suppression of small structures

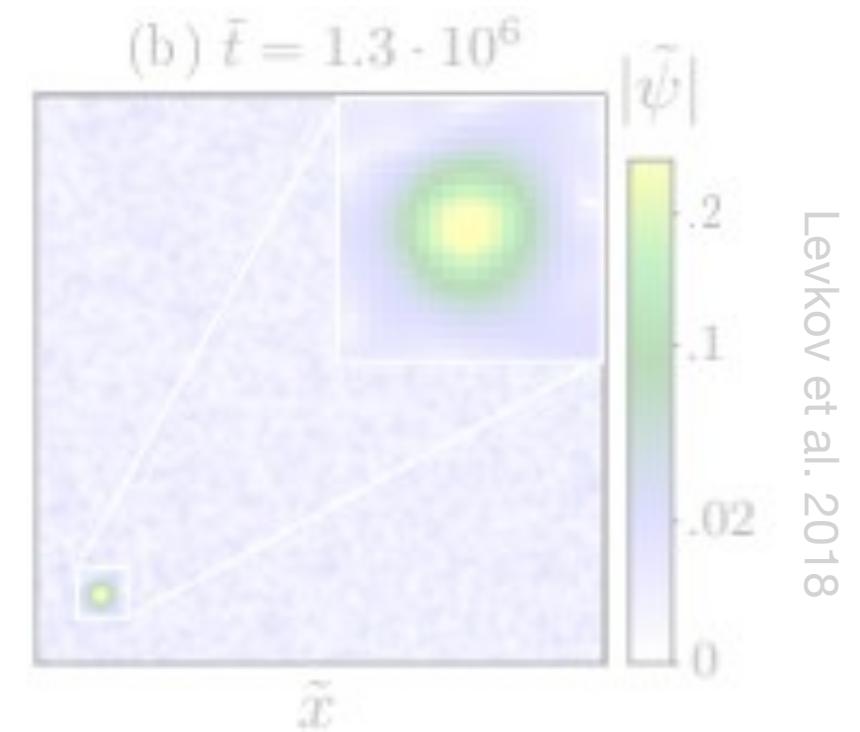


S. May et al. 2021

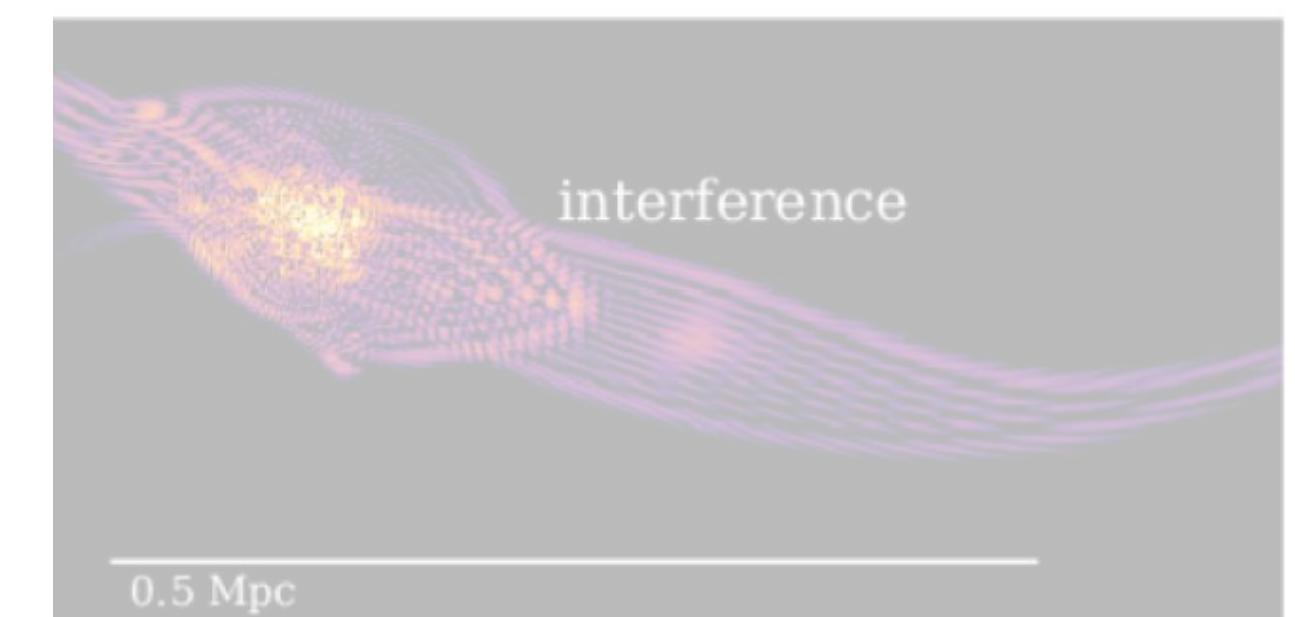
Formation of a solitonic core



Dynamical effects



Wave interference

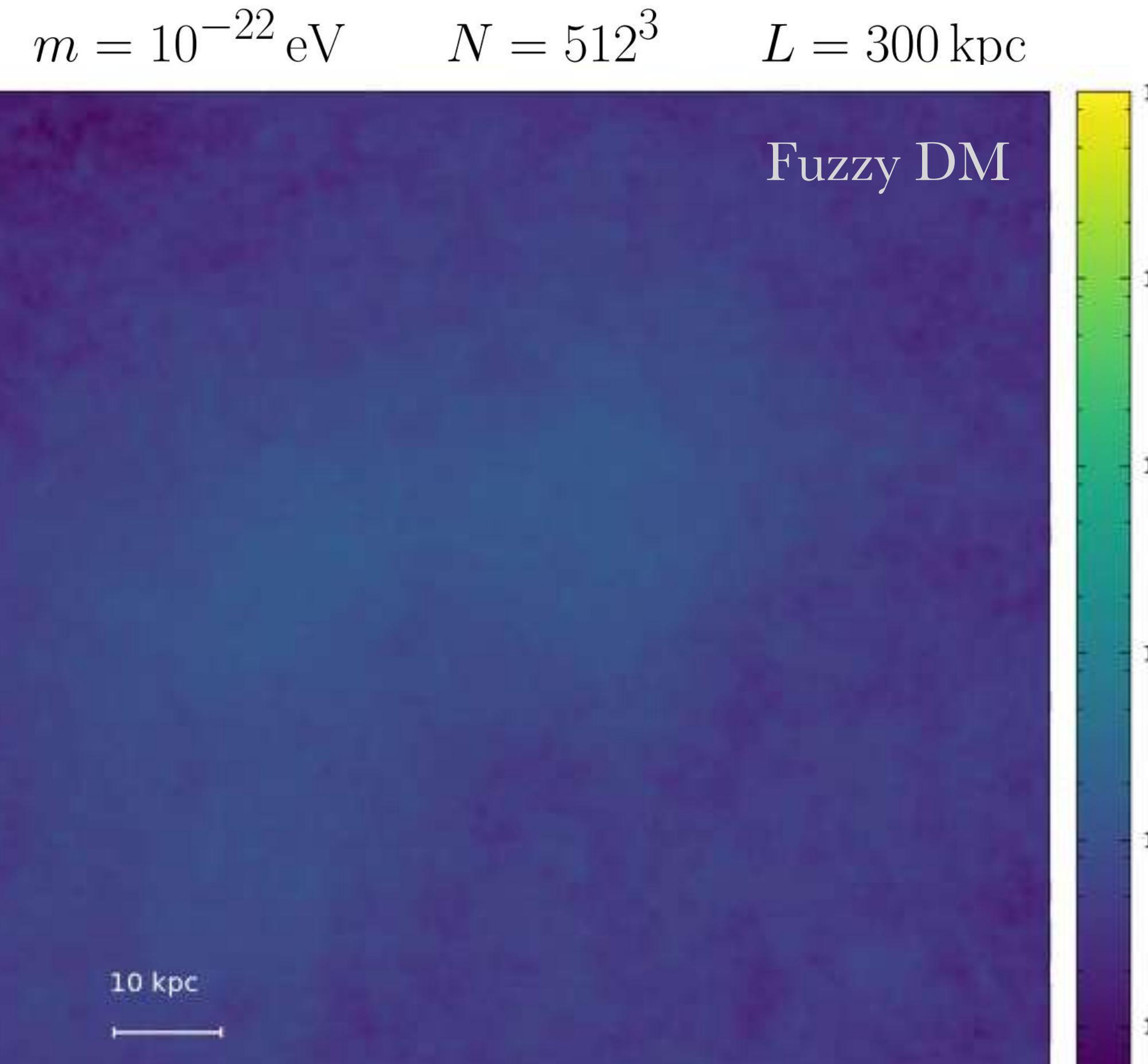


Mocz et al. 2017

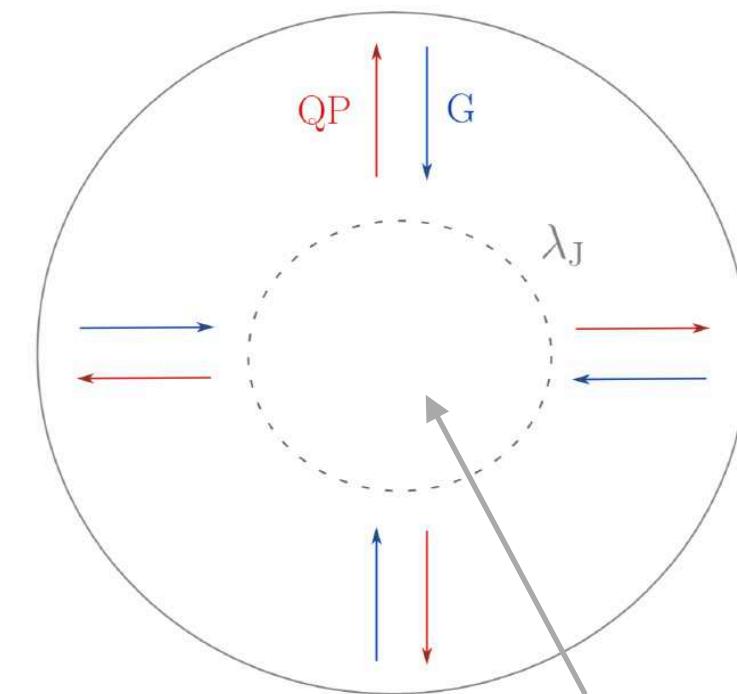
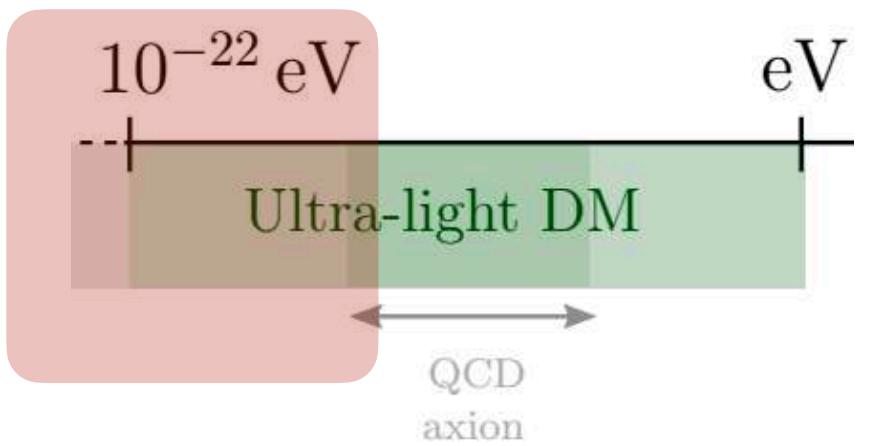
Phenomenology

Formation of cores

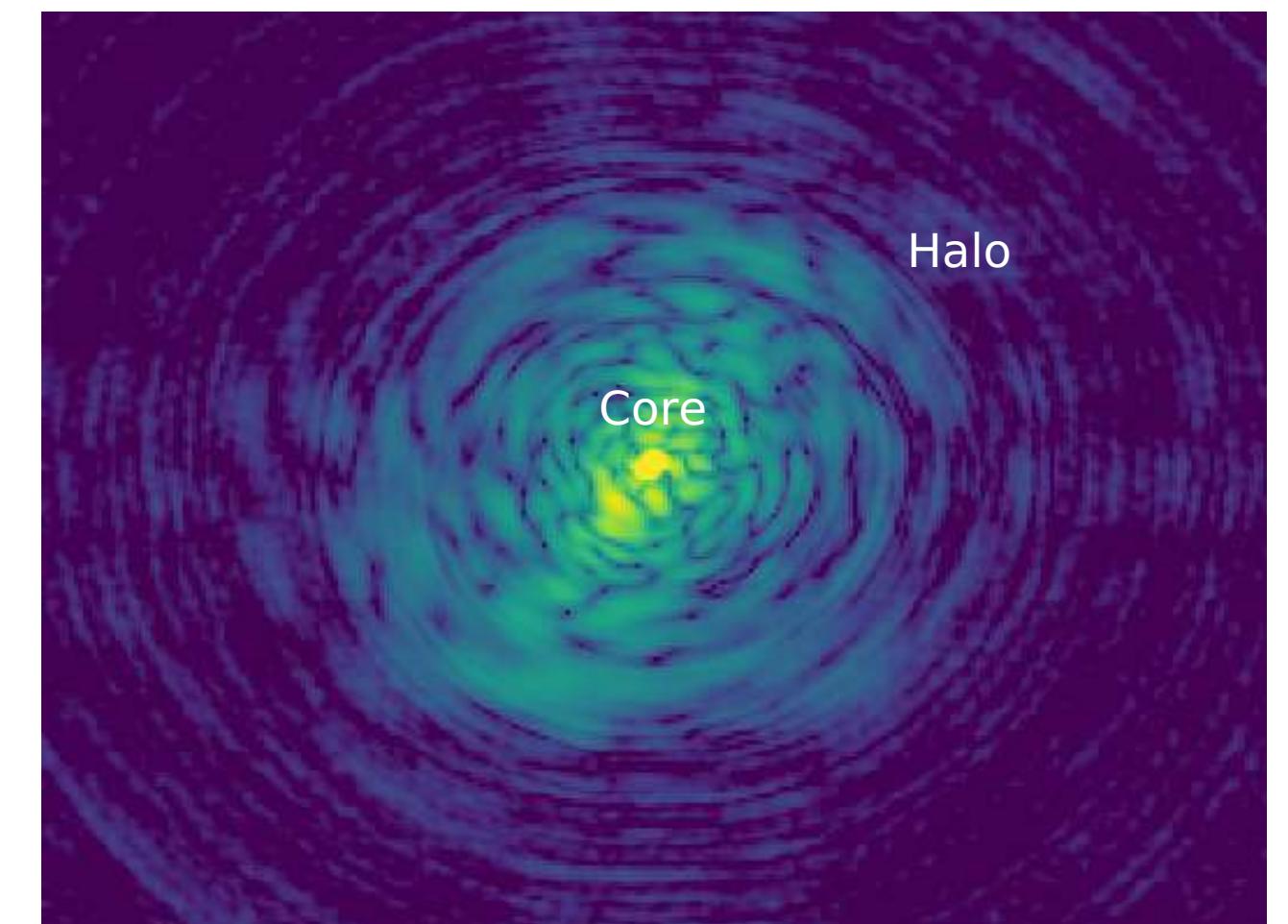
NON-LINEAR
evolution: need
simulations



Simulation by Jowett Chan

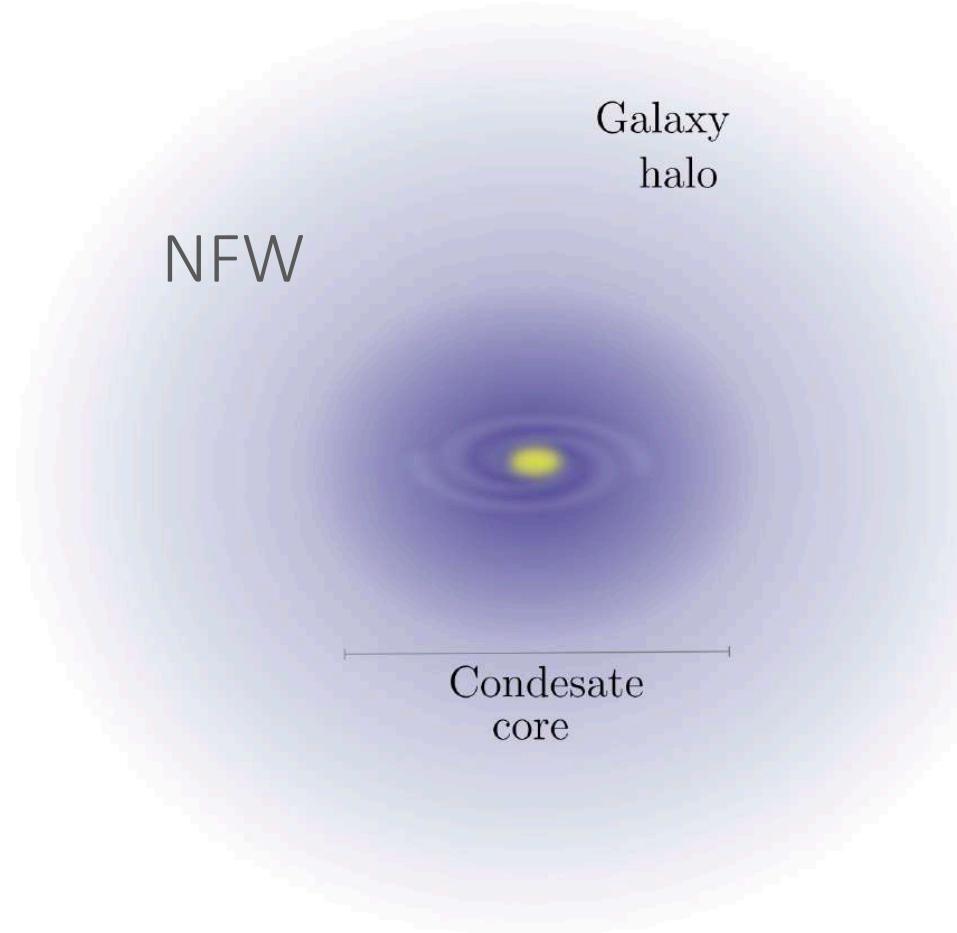


NO structure formation
Stable, oscillating solution

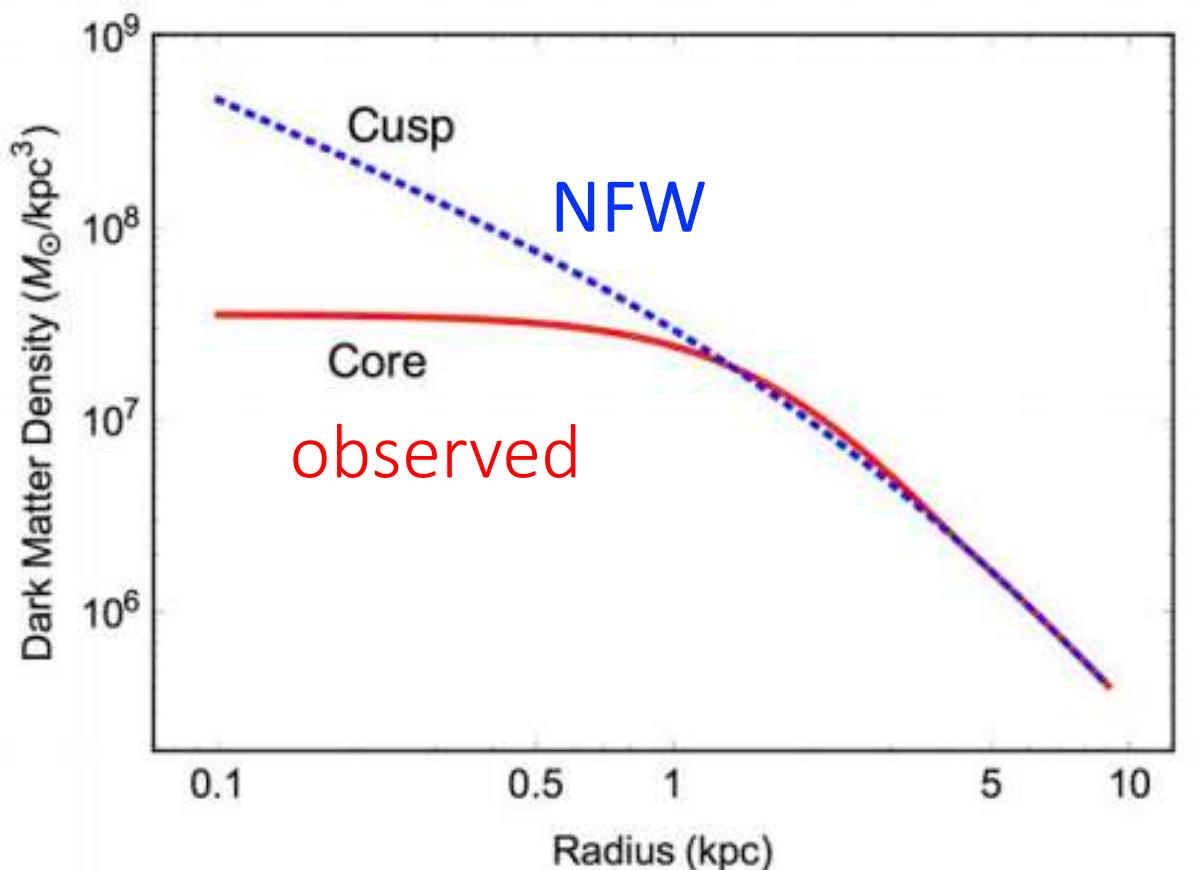
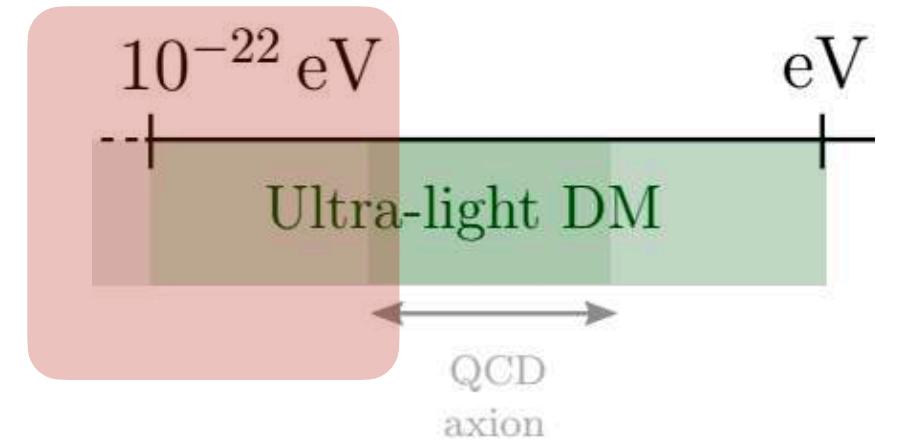


Phenomenology

Formation of cores



$$\rho(r) \simeq \begin{cases} \rho_c & \text{for } r \leq r_c \\ \rho_{\text{NFW}} & \text{for } r \geq r_c \end{cases}$$



FDM

From simulations Schive et al. 2014, fitting function:

Stable core solution

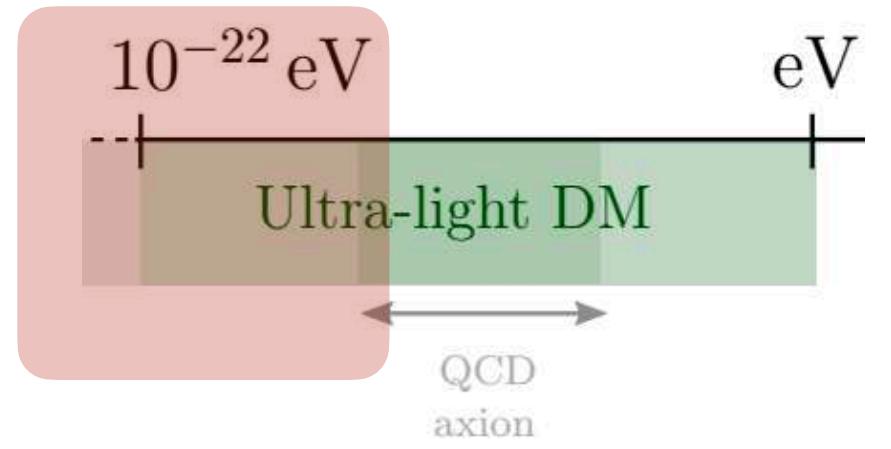
$$\rho_c \simeq \frac{1.9 \times 10^{-2}}{[1 + 0.091(r/R_{1/2,c})^2]^8} \left(\frac{m}{10^{-22} \text{ eV}}\right)^{-2} \left(\frac{r_c}{\text{kpc}}\right)^{-4} M_\odot \text{ pc}^{-3},$$

$$r_c \simeq 0.16 \left(\frac{m}{10^{-22} \text{ eV}}\right)^{-1} \left(\frac{M}{10^{12} M_\odot}\right)^{-1/3} \text{ kpc}.$$

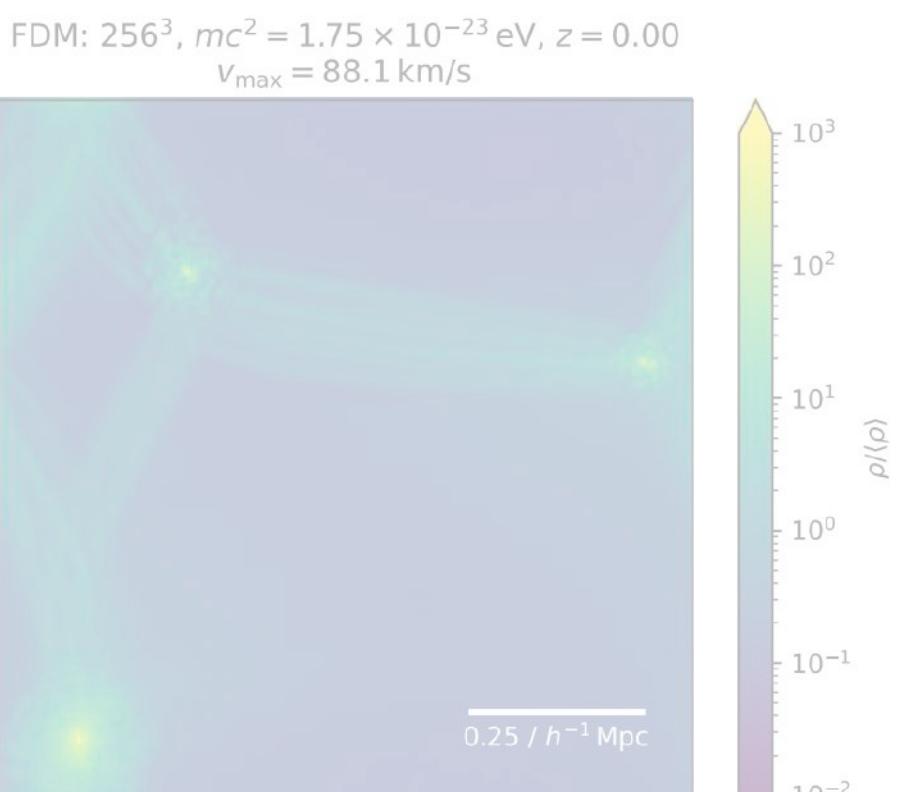
Relations used to compare
with observations

Phenomenology

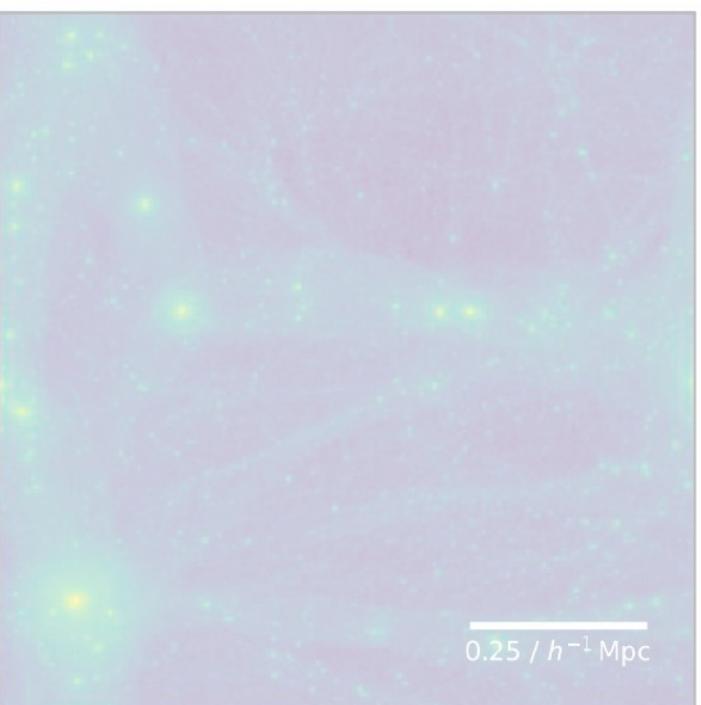
RICH PHENOMENOLOGY ON SMALL SCALES



Suppression of small structures

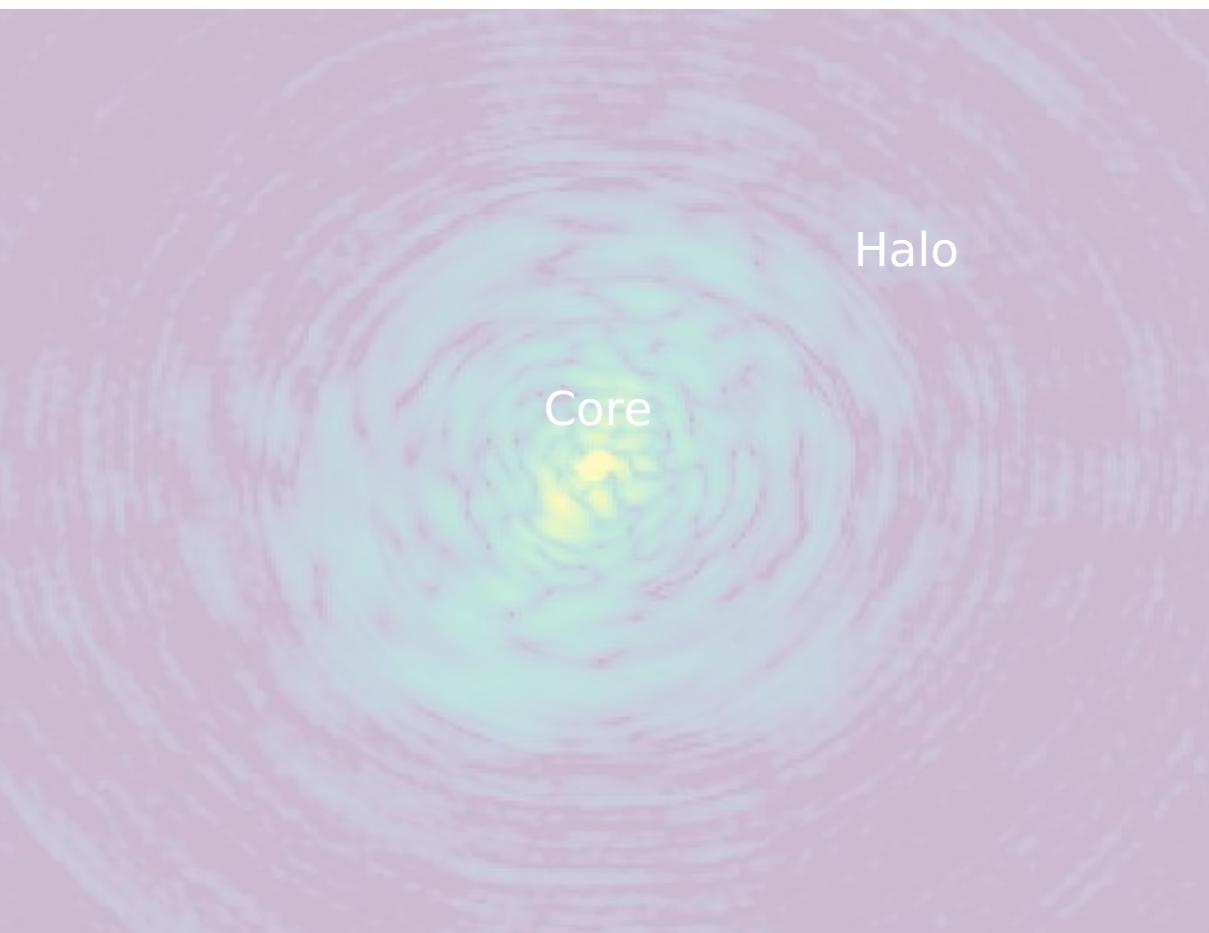


CDM: 256^3 , $z = 0.00$

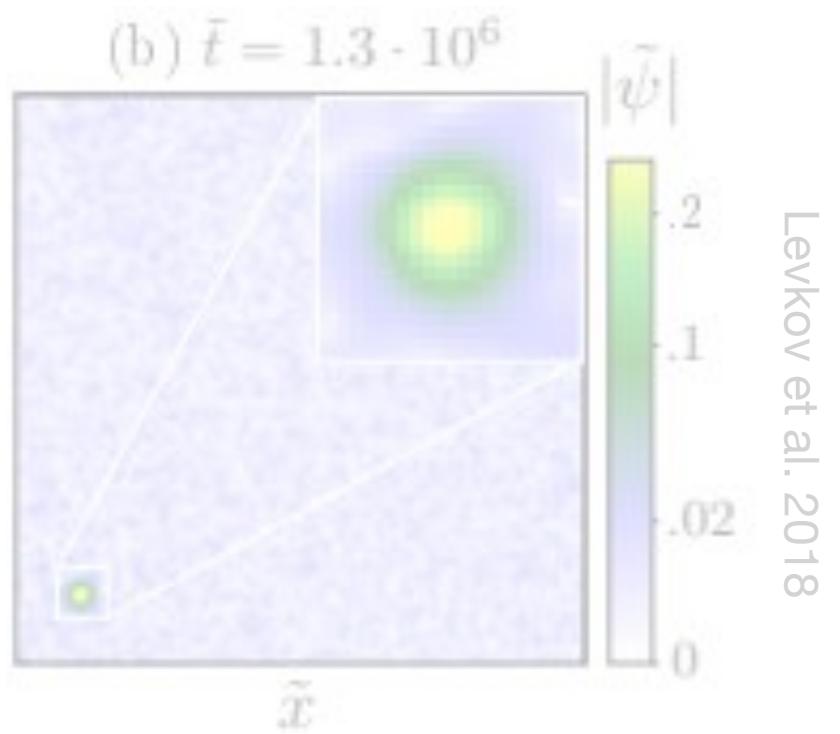


S. May et al. 2021

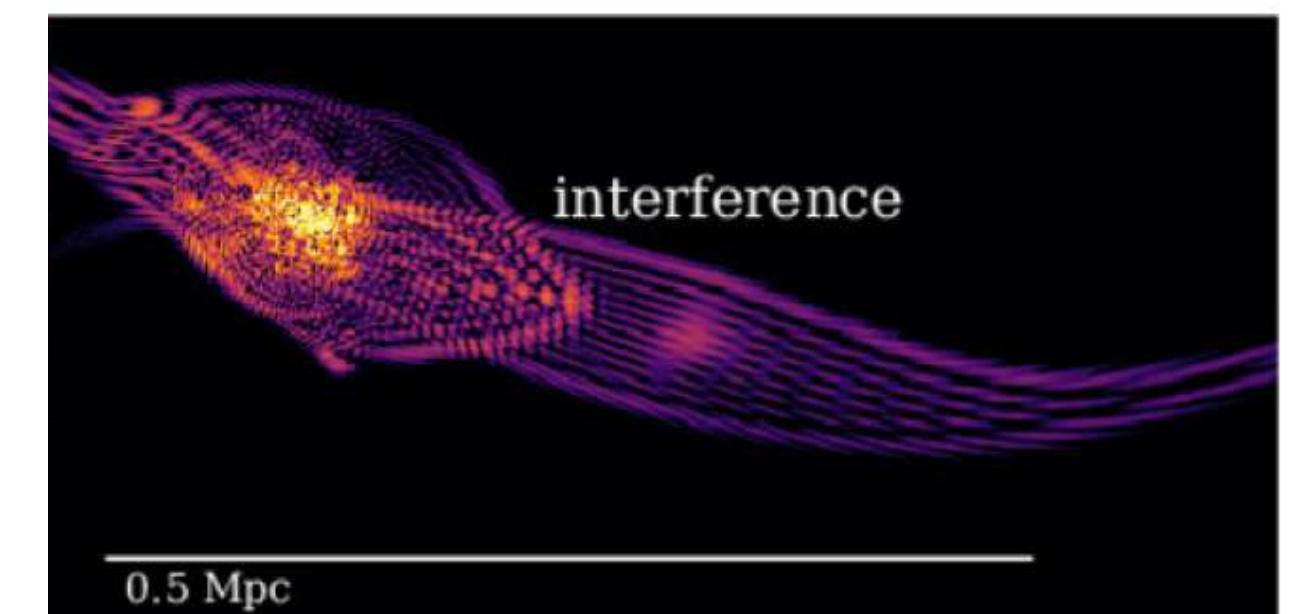
Formation of a solitonic core



Dynamical effects



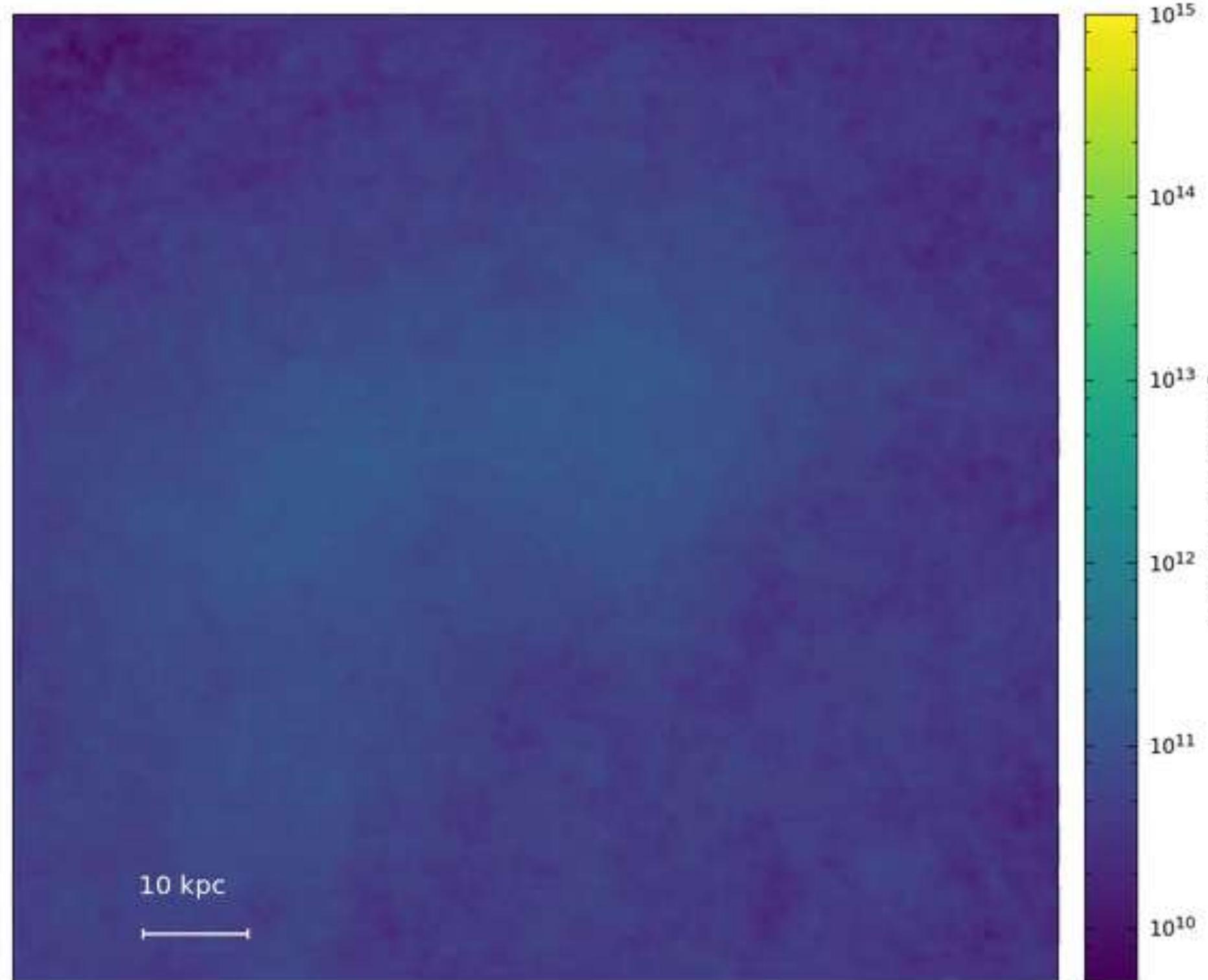
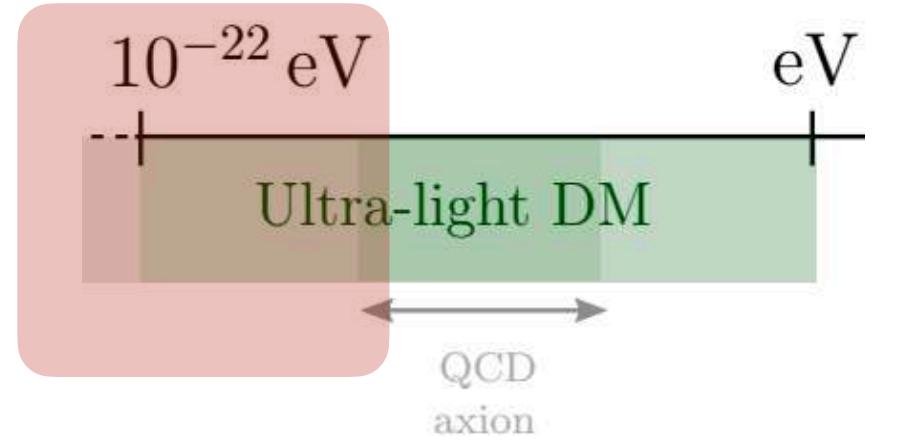
Wave interference



Mocz et al. 2017

Phenomenology

Wave interference: granules and vortices

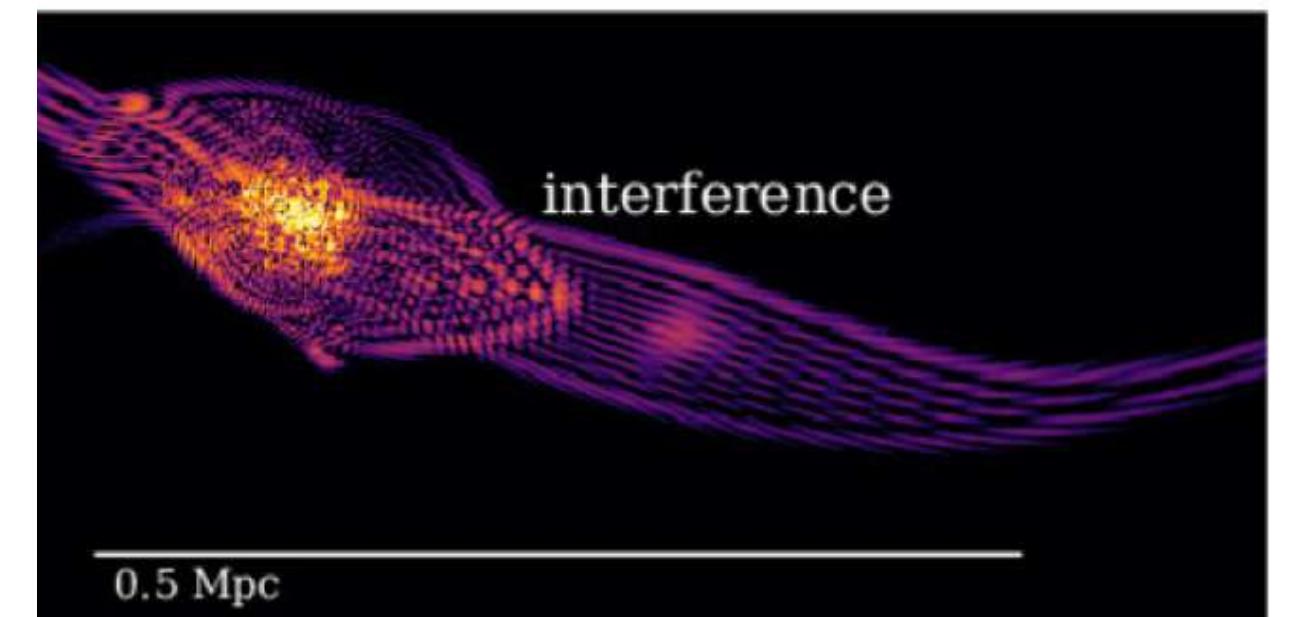
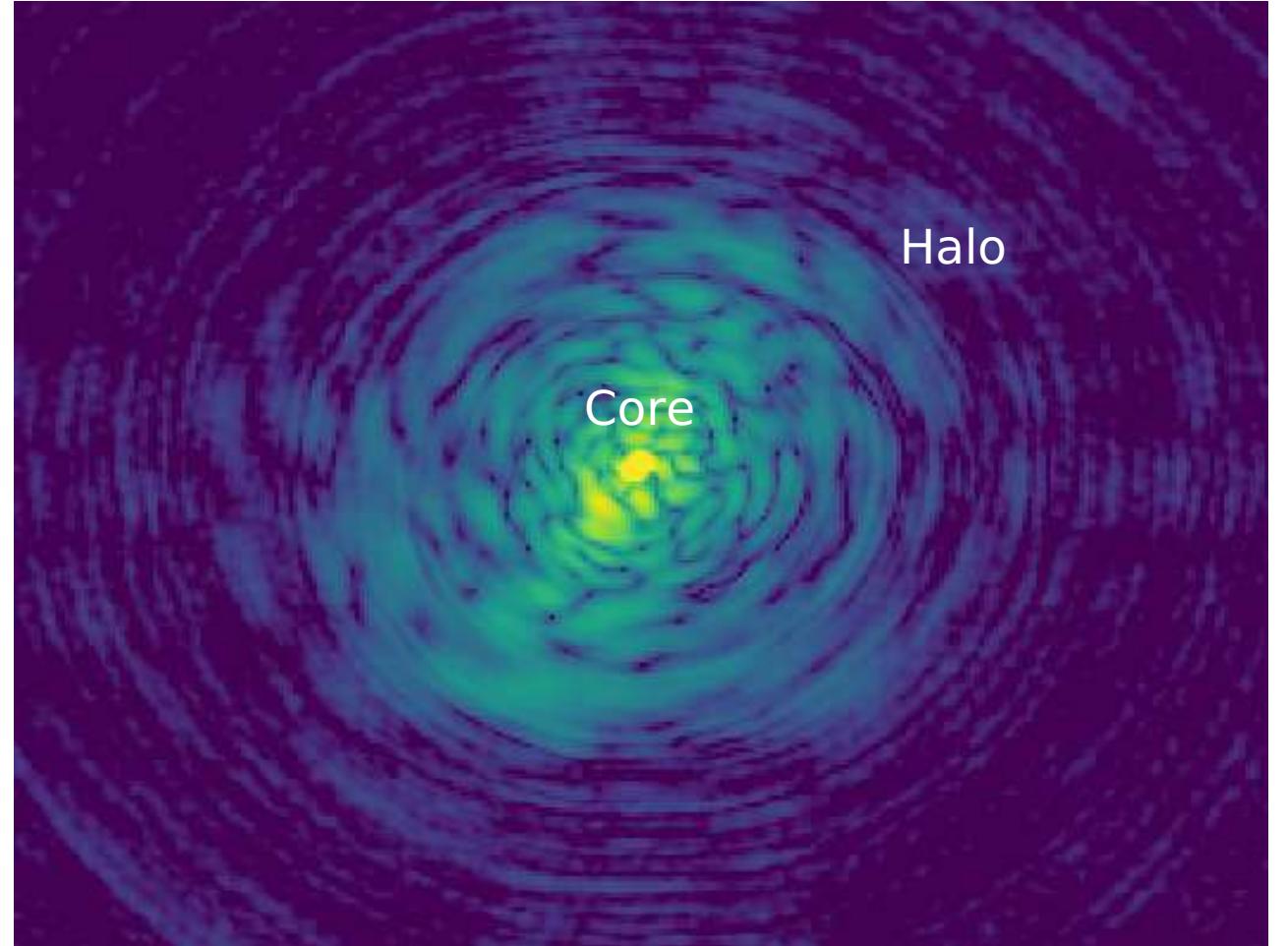


Simulation by Jowett Chan

Order one fluctuations in density \longrightarrow

Constructive interference: **granules**
Destructive interference

$$\sim \lambda_{\text{dB}}$$



Mocz et al. 2017

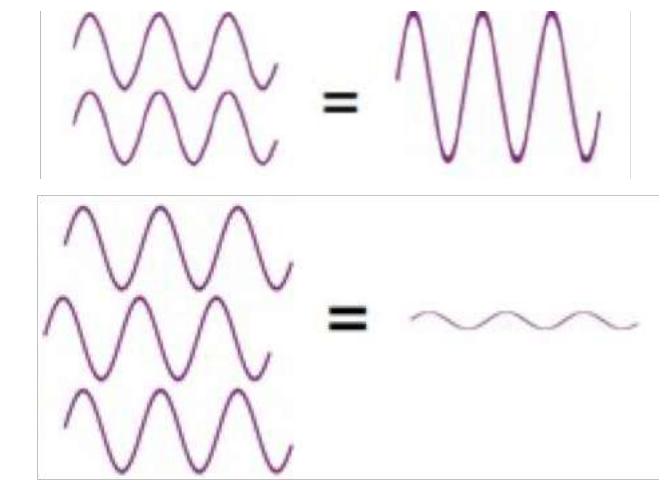
Hard to observe!

Vector, higher spin or multicomponent FDM

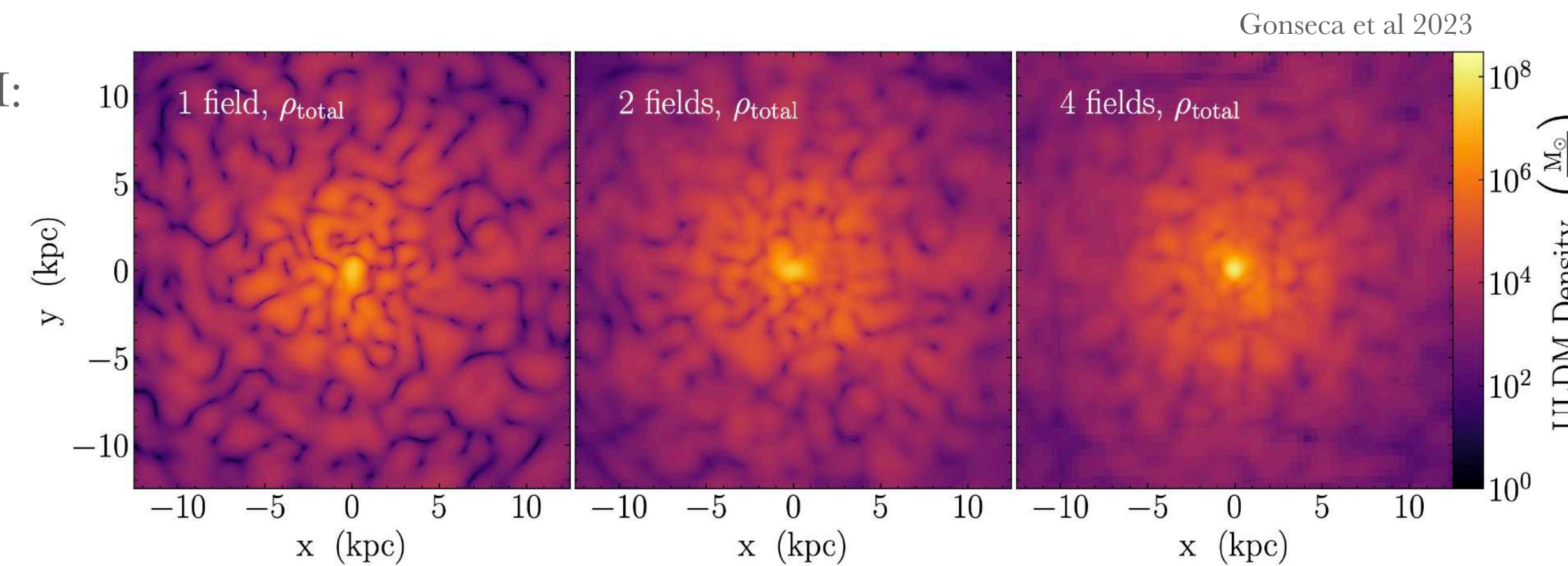
ULDM or ULA are a coherent wave - same frequency and constant phase difference

Multiple coherent waves

Interference patterns



For ULDM:



Gonseca et al 2023

Multiple FDM or VFDM (or higher spin s FDM)
attenuates the granule amplitude by

$$\frac{[\delta\rho/\rho]_{\text{nfdm},s}}{[\delta\rho/\rho]_{\text{fdm}}} \propto \frac{1}{\sqrt{(2s+1)}} = \frac{1}{\sqrt{N}}$$

(Amin et al 2022)

Vector (and higher-spin) FDM Amin et al 2022

(Vector FDM = 3 x same mass FDM (spin 0))

Multicomponent FDM Gonseca et al 2023

Phenomenology

Vortices

Vortices are sites where the fluid velocity has a non-vanishing curl

Two ways:

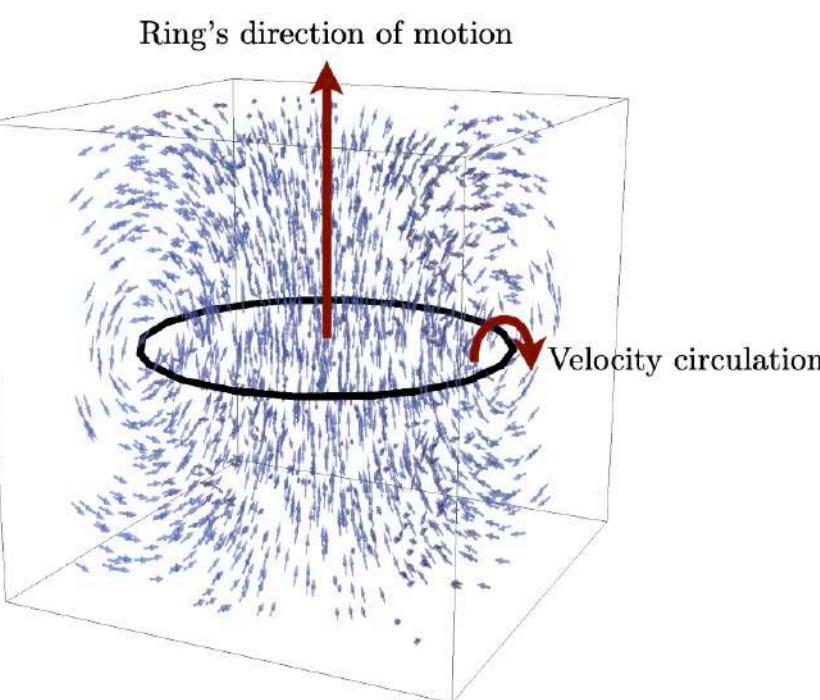
- regions where the density vanishes
- transfer of angular momentum (superfluids only)

Fuzzy DM

Interference of waves leads to **vortices** - where there is **destructive interference**

General defet in 3D

$$\mathcal{C} = \frac{1}{m} \oint_{\partial A} d\theta = \frac{2\pi n}{m}$$



$$(\psi \equiv \sqrt{\rho/m} e^{i\theta} \text{ and } \mathbf{v} \equiv \nabla\theta/m)$$

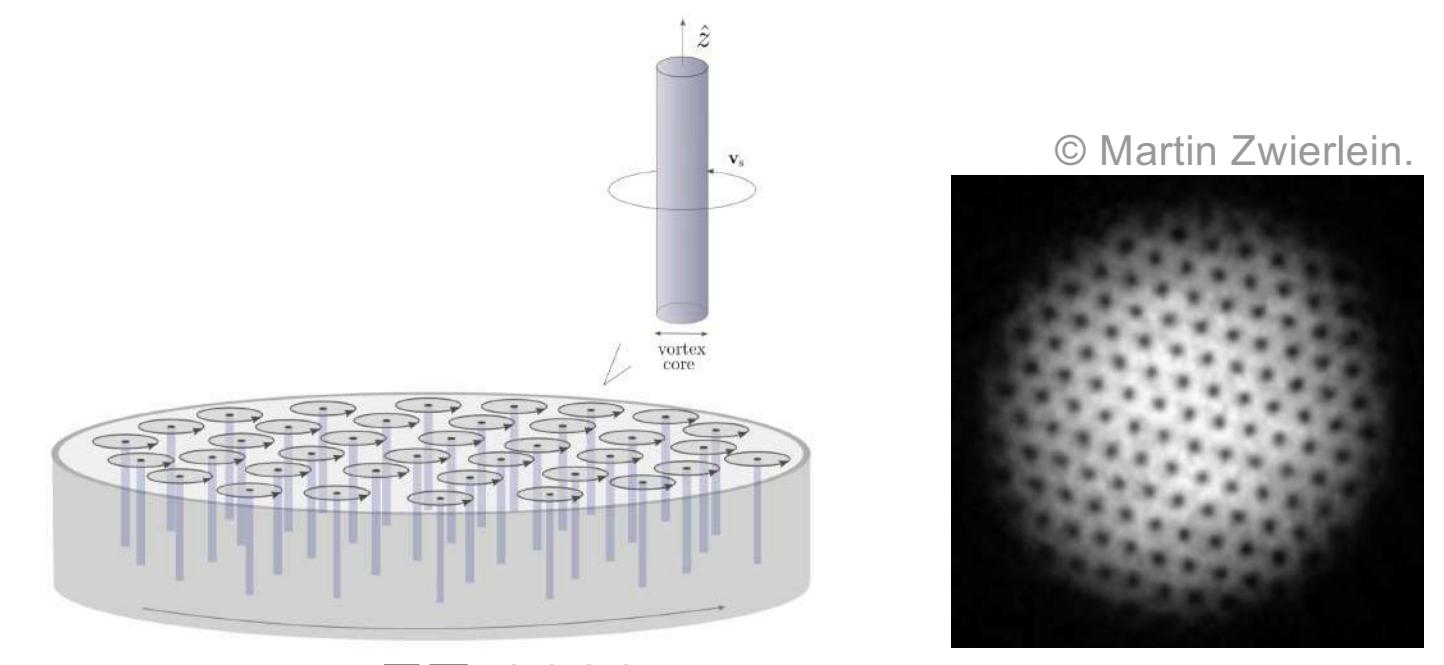
$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\dot{\mathbf{v}} + (\mathbf{v} \cdot \nabla)\mathbf{v} = -\frac{1}{m} \left(V_{grav} - P_{int} - \frac{1}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right)$$

Vel. field is a gradient flow \longrightarrow irrotational fluid, no vorticity

Self-interacting Fuzzy DM

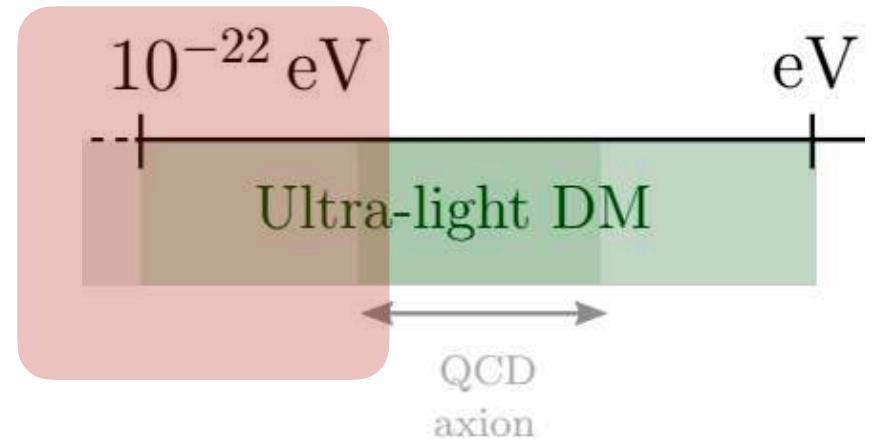
Superfluid cannot rotate uniformly. If the superfluid rotates faster than the critical vel., network of vortices are formed.



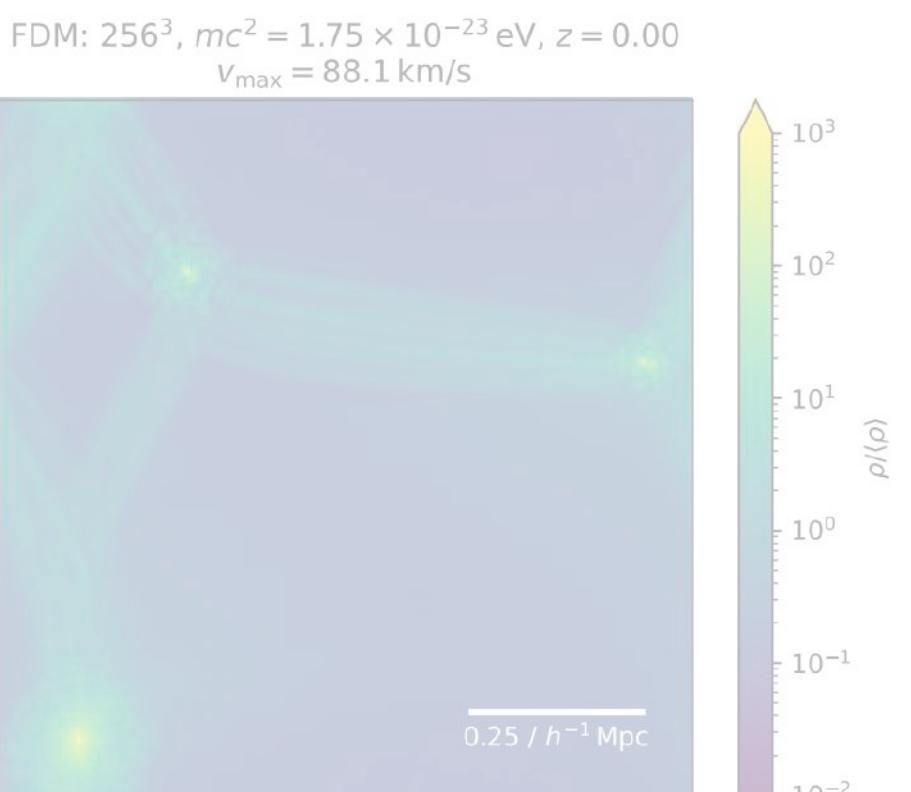
© Martin Zwierlein.

Phenomenology

RICH PHENOMENOLOGY ON SMALL SCALES

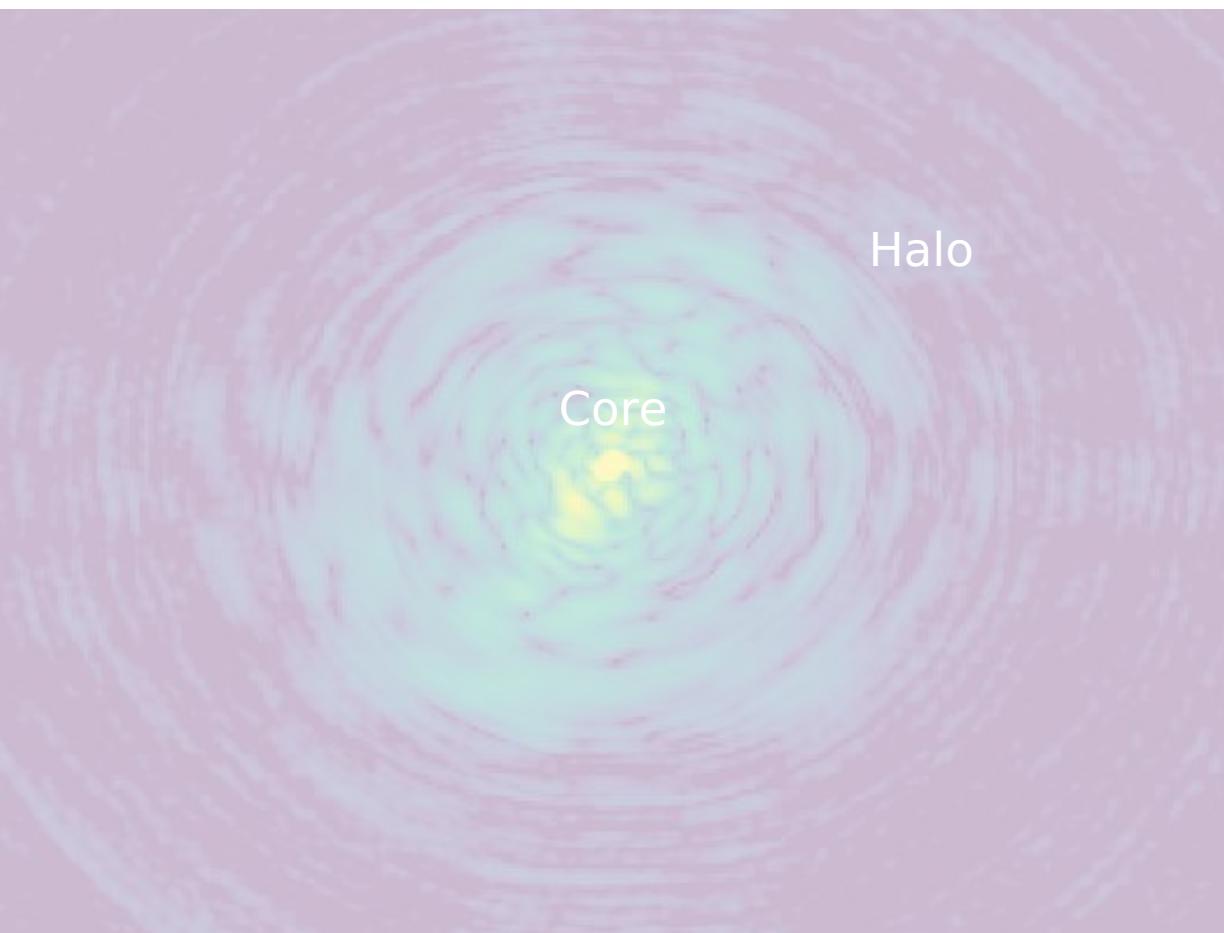


Suppression of small structures

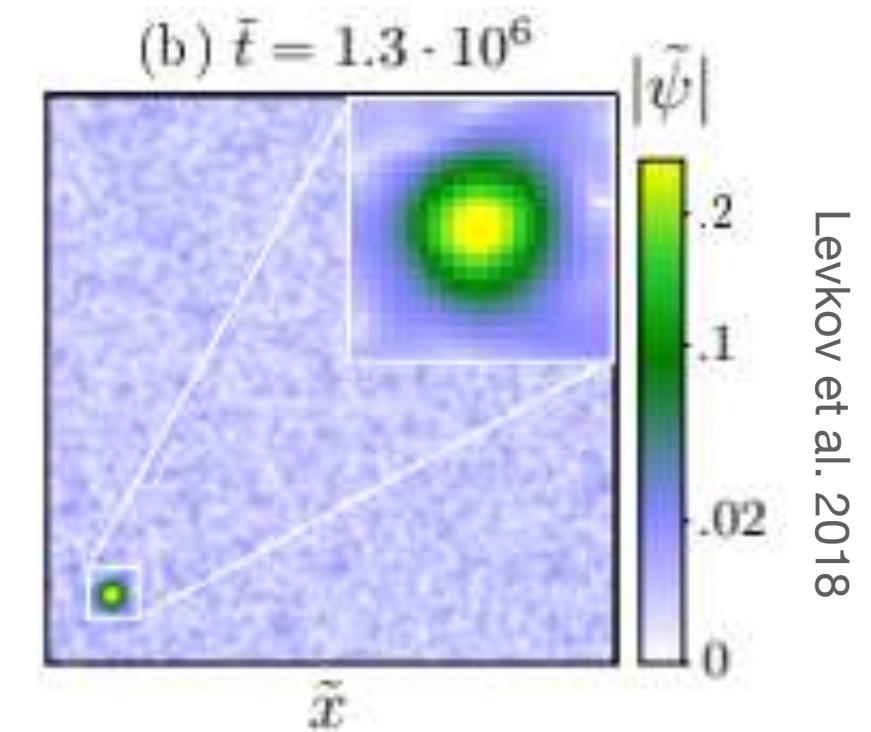


S. May et al. 2021

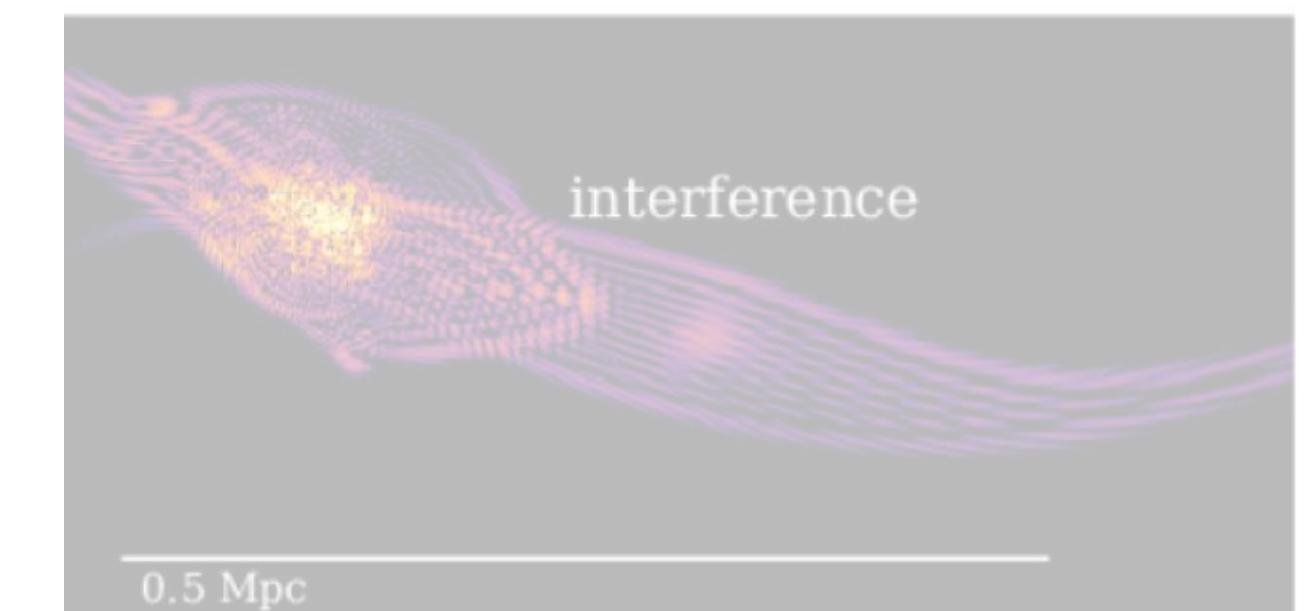
Formation of a solitonic core



Dynamical effects



Wave interference



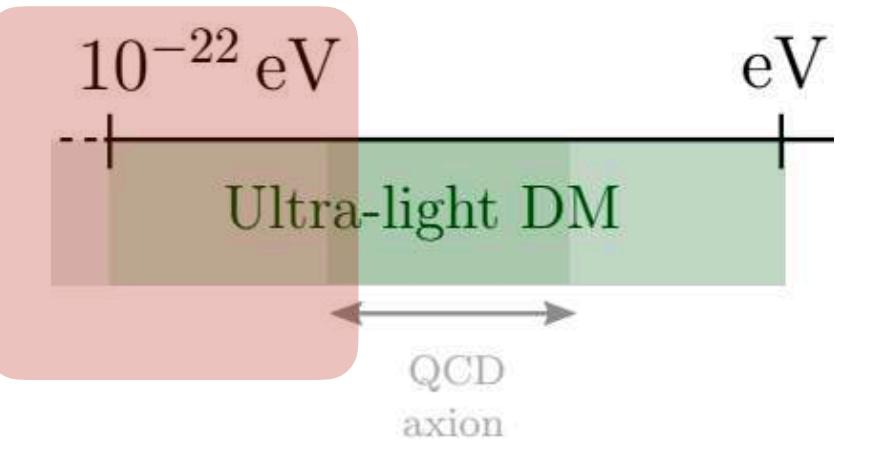
Mocz et al. 2017

S. May et al. 2021

Phenomenology

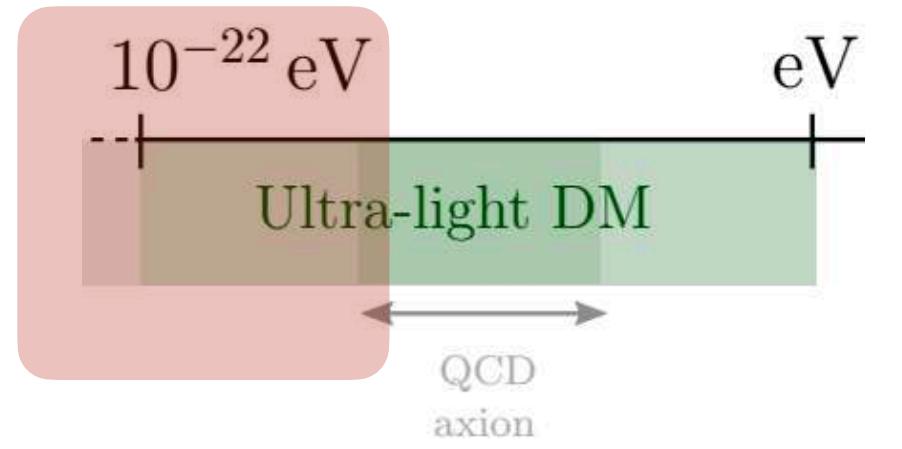
Dynamical effects

Relaxation, oscillation, friction, and heating

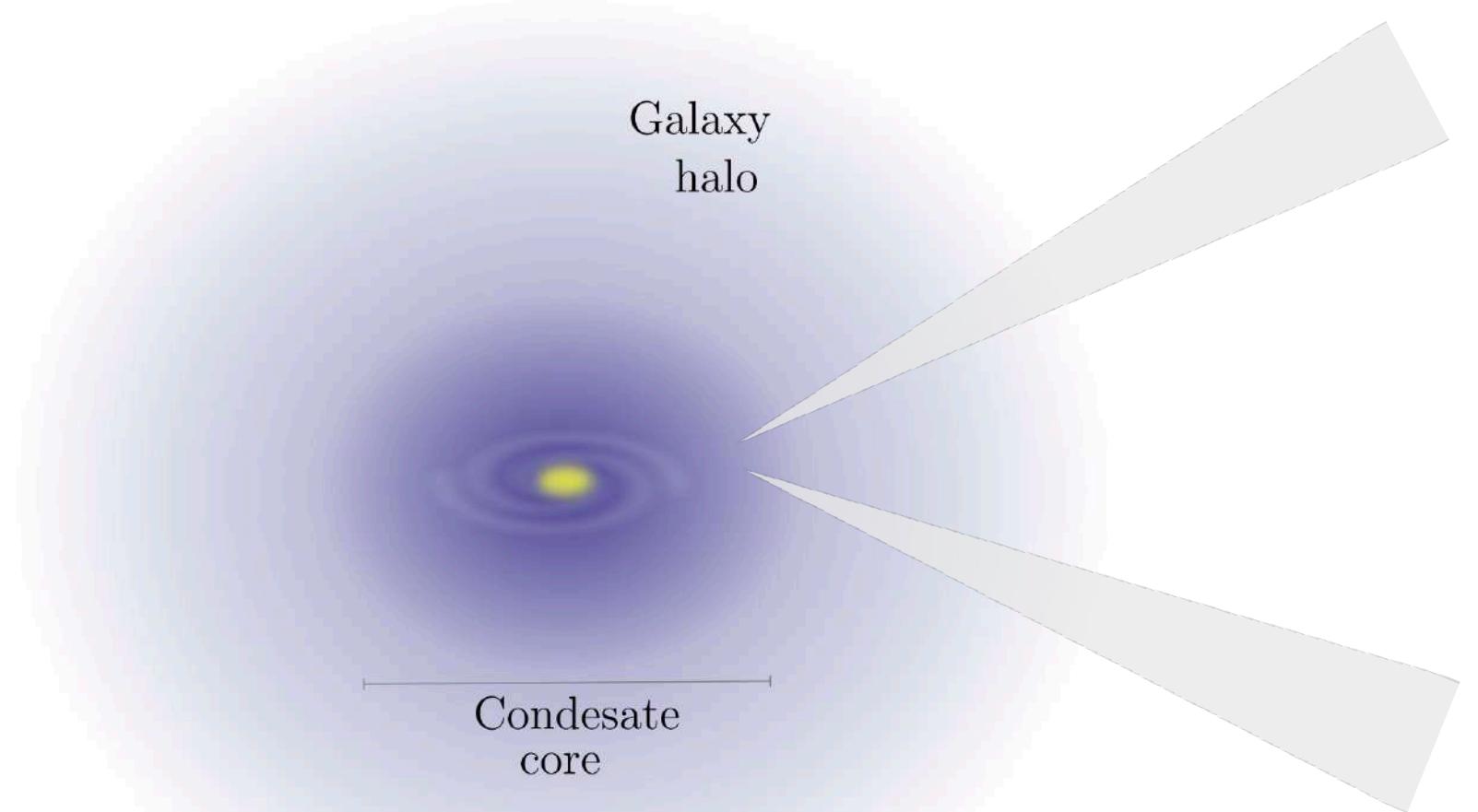


Phenomenology

Dynamical effects



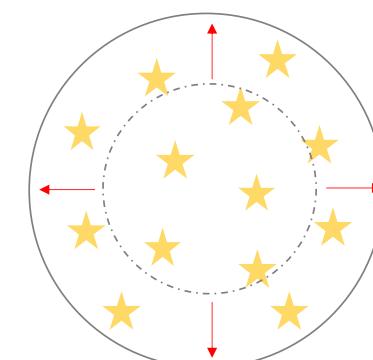
Relaxation, oscillation, friction, and heating



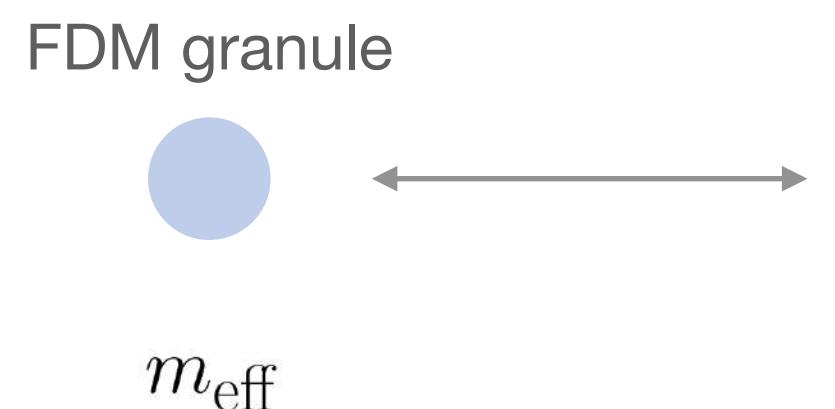
Heating



System (star)
gains energy



Friction

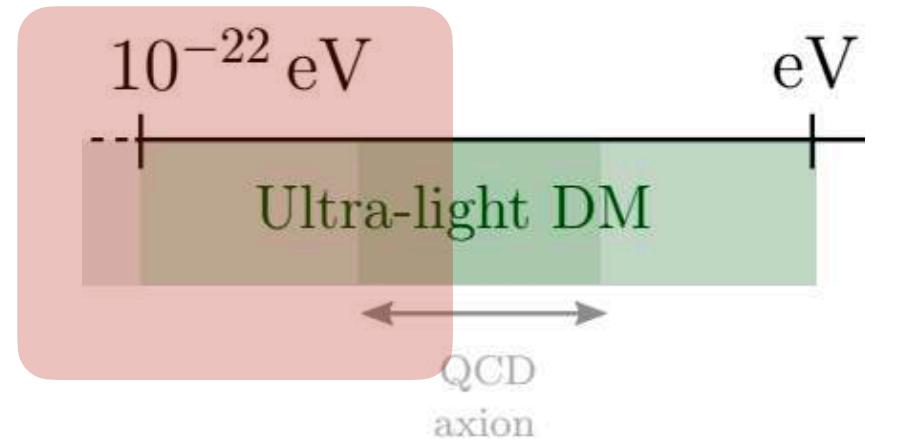


System (GC or BH)
loses energy



Globular cluster

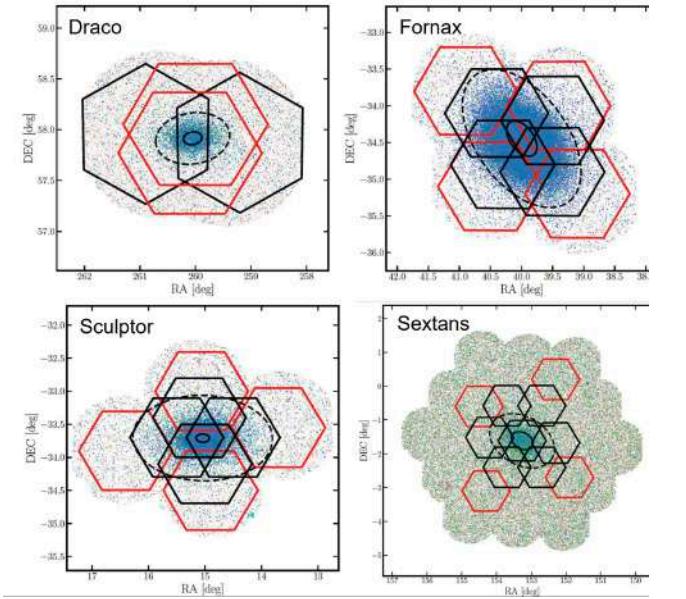
Observational implications and constraints



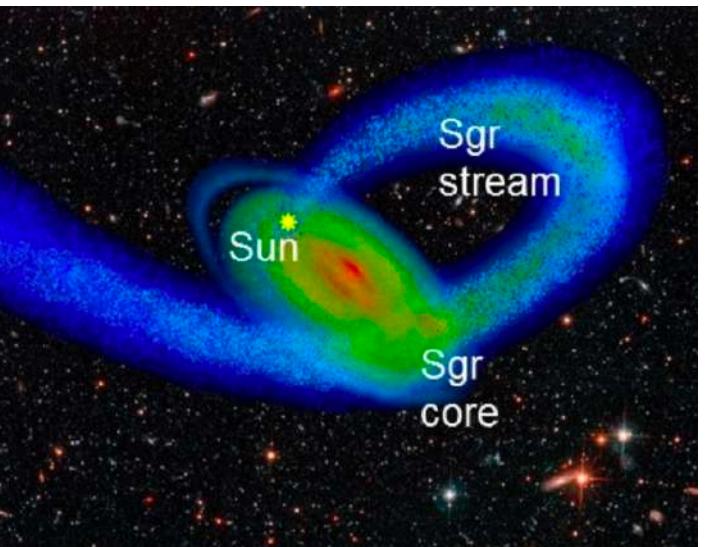
Galaxies



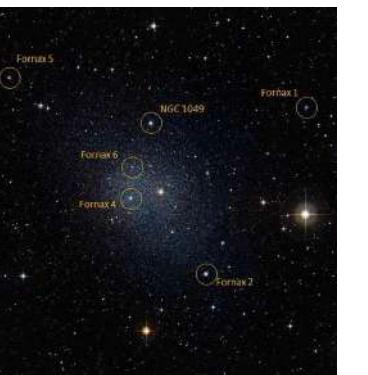
Dwarfs



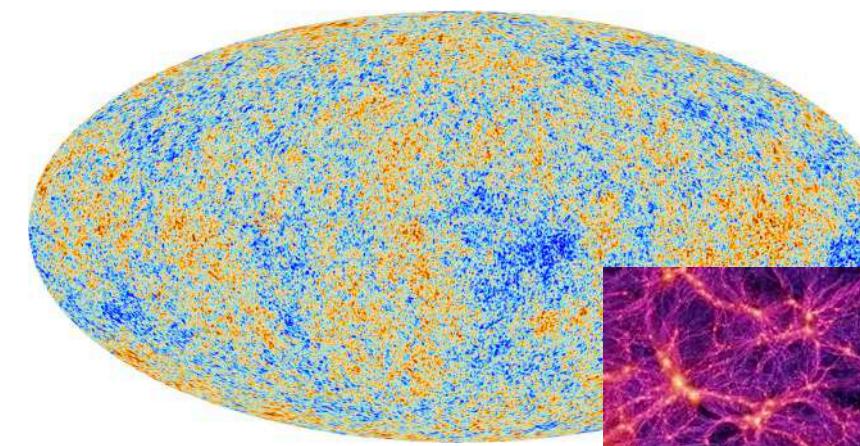
Stellar stream



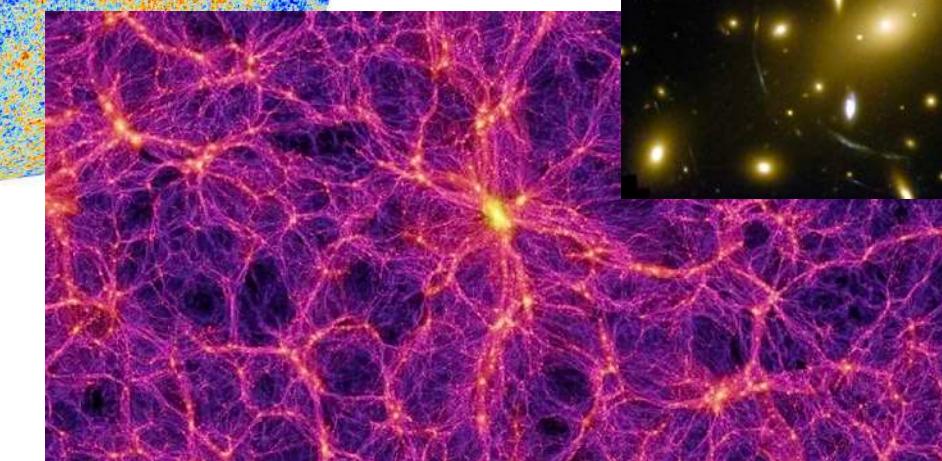
Globular clusters



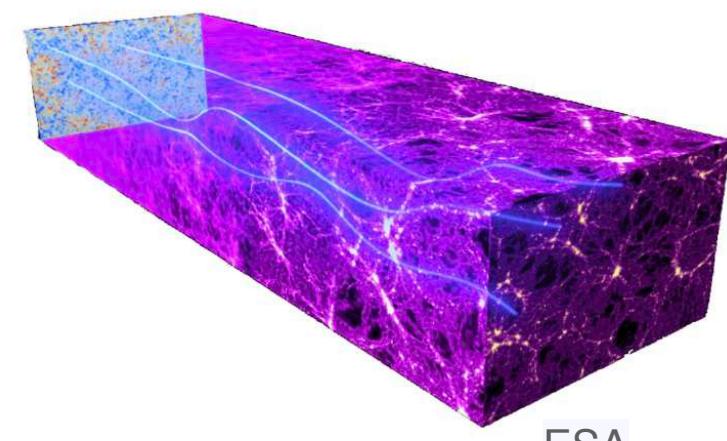
CMB+LSS



ESA and the Planck Collaboration

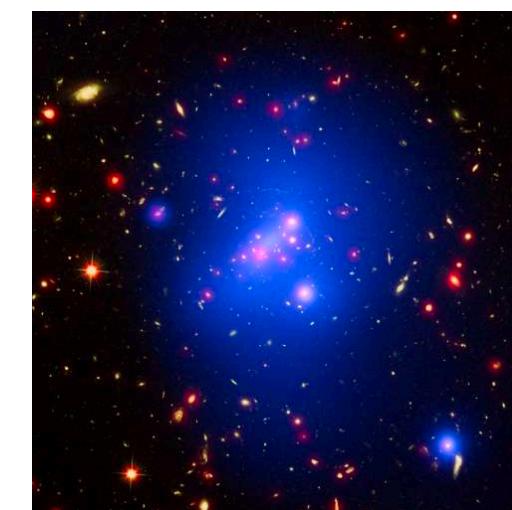


Springel & others / Virgo Consortium



NASA and ESA

Clusters

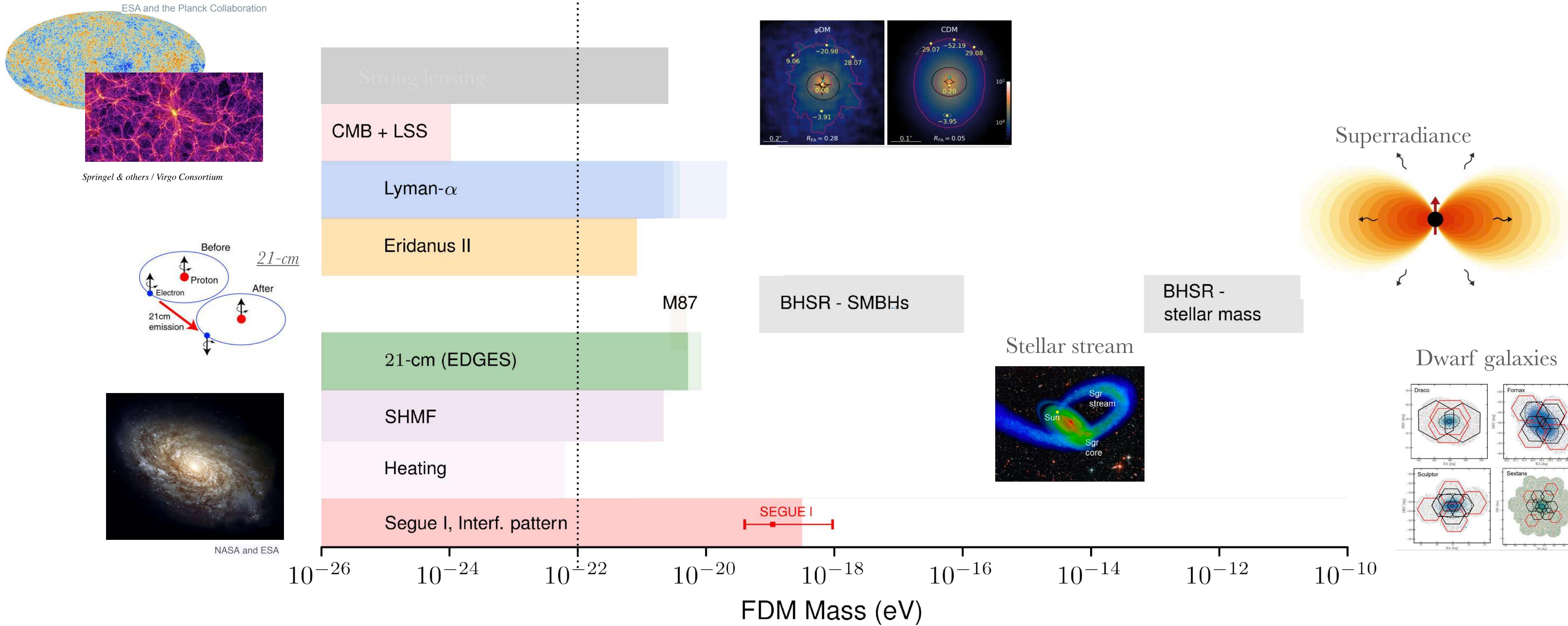
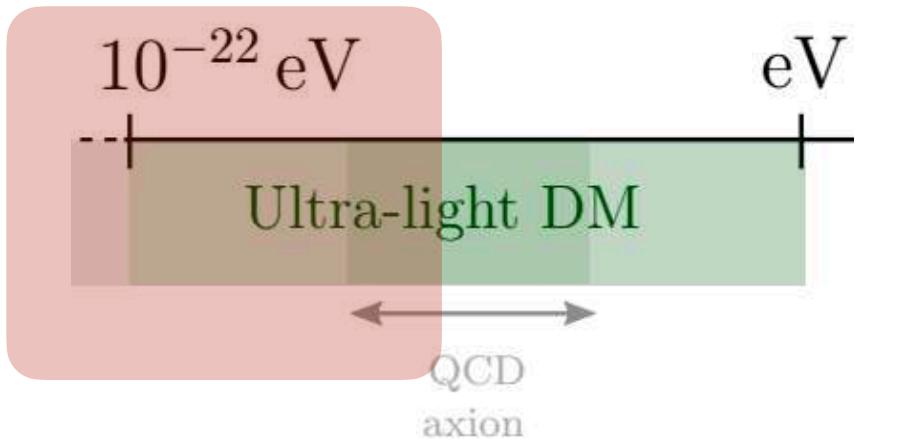


CC BY 4.0

ESA

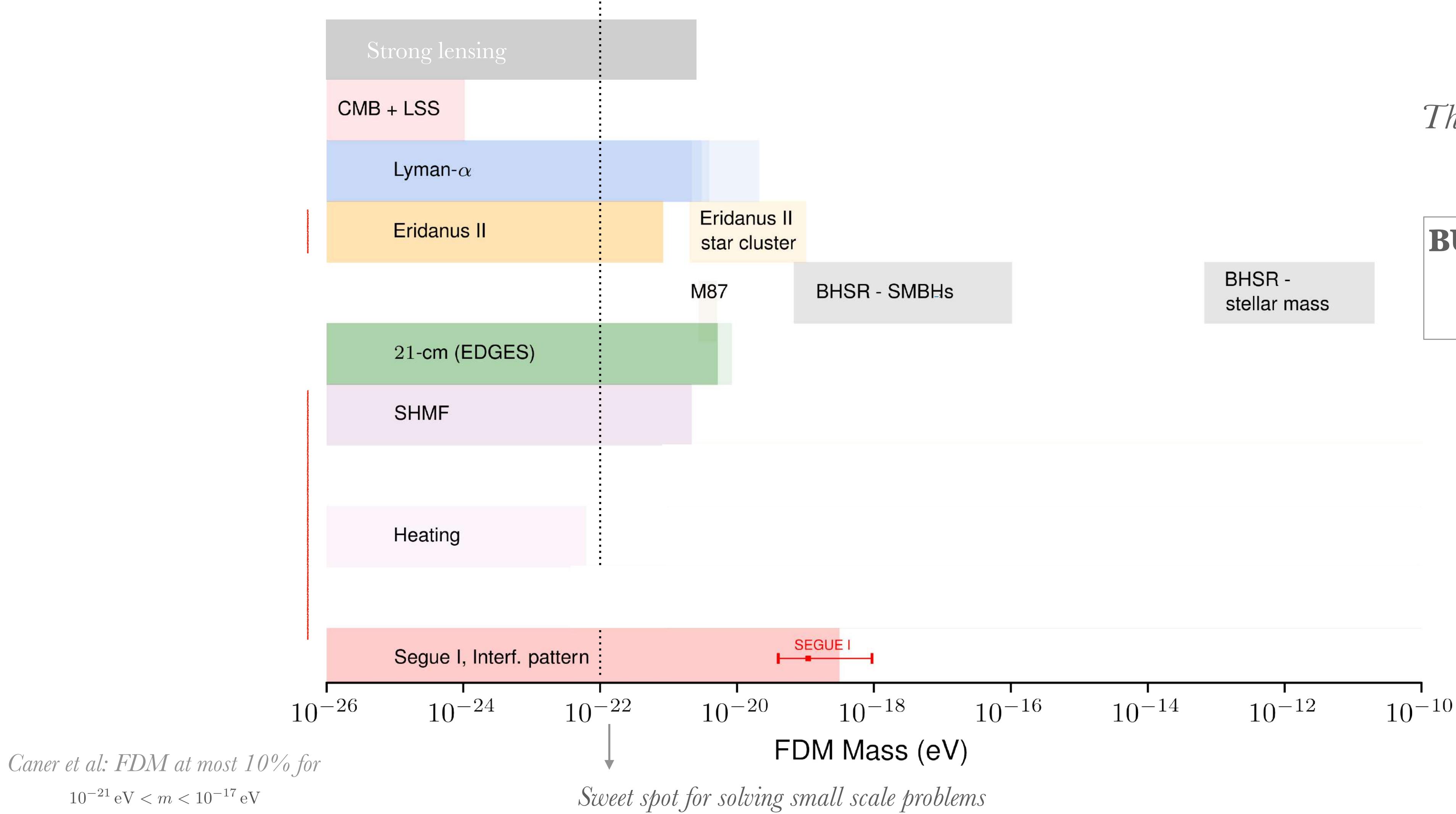
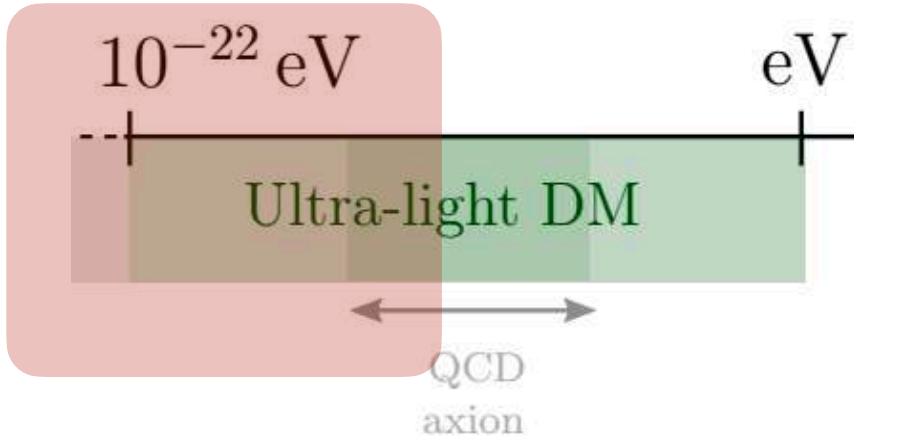
Current status

Fuzzy Dark Matter - bounds on the mass



Current status

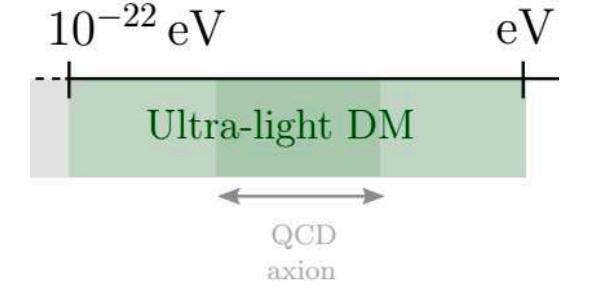
Fuzzy Dark Matter - bounds on the mass



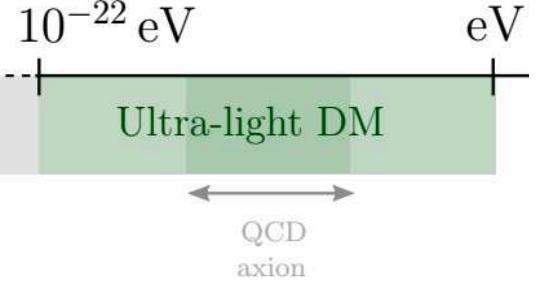
These models can be constrained

BUT: - systematic effects!!
 - dynamics of FDM not fully understood.

- Need:
- Observations
 - Improve sims
 - New observables
 - New probes



Axion and ALPs interaction with the SM



Axion and ALPs interaction with the SM

Axions and ALPs interact with the standard model particles

$$\boxed{\begin{aligned} F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu \\ \tilde{F}^{\mu\nu} &= \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta} \\ \mathbf{E} &= -\nabla A_0 - \dot{\mathbf{A}} \\ \mathbf{B} &= \nabla \times \mathbf{A} \end{aligned}}$$

Minimal definition: New light pseudoscalar, with coupling to photons and/or derivative couplings to fermions

$$\mathcal{L} = \frac{1}{2}(\partial_\mu a)(\partial^\mu a) - \frac{1}{2}m_a^2 a^2 - \frac{g_{a\gamma}}{4}aF_{\mu\nu}\tilde{F}^{\mu\nu} + \partial_\mu a \sum_\psi \frac{g_{a\psi}}{2m_\psi} (\bar{\psi}\gamma^\mu\gamma^5\psi)$$

Not considering here

+ a few model-dependent assumptions

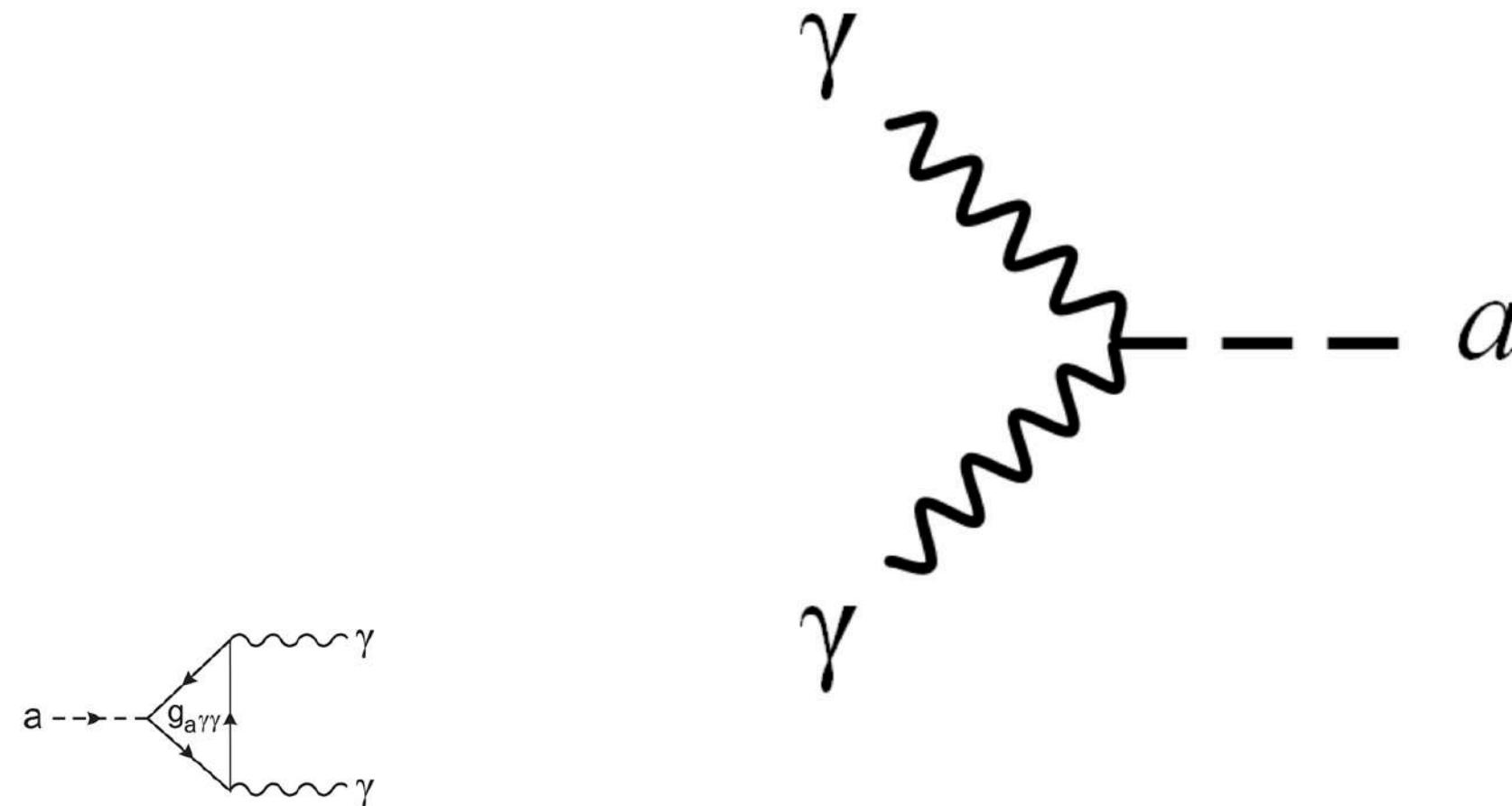
Axion and ALPs interaction with the SM

Axions and ALPs interact with the standard model particles

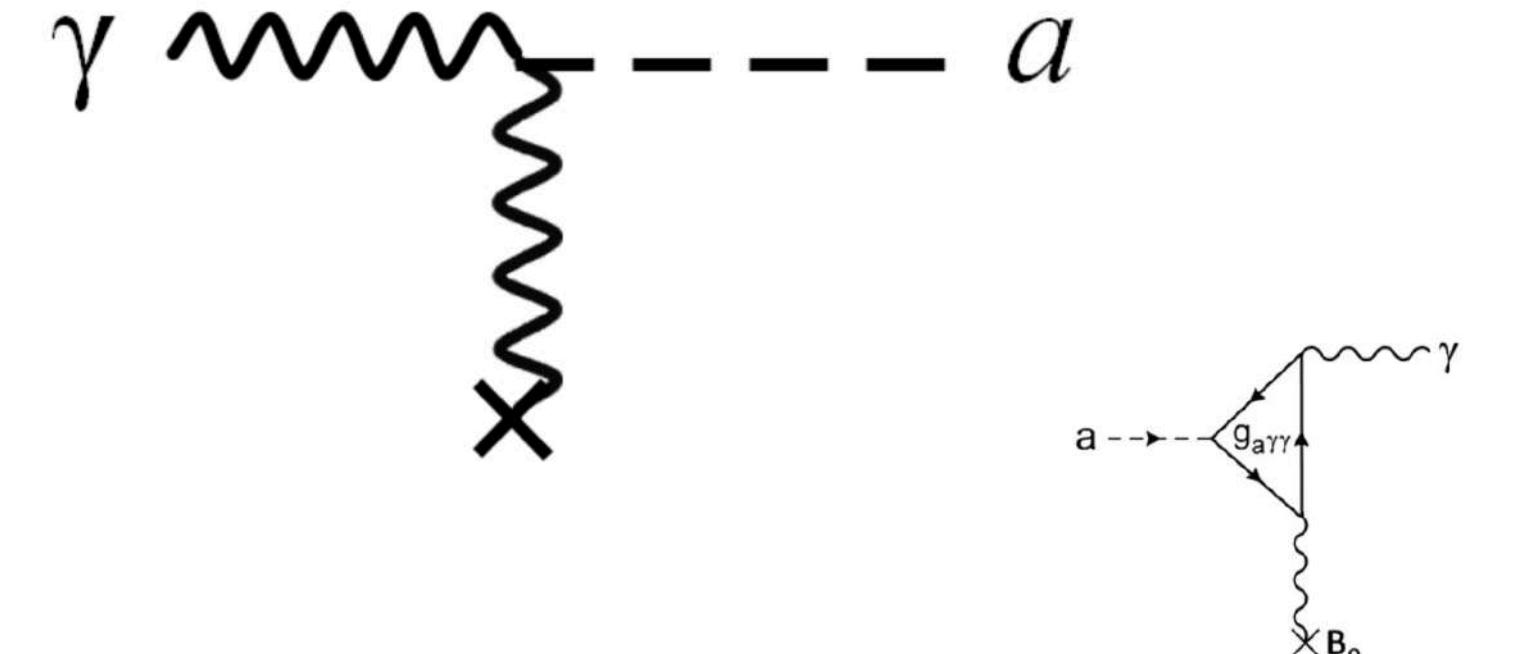
Photon - Axion electrodynamics

$$\boxed{\begin{aligned} F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu \\ \tilde{F}^{\mu\nu} &= \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta} \\ \mathbf{E} &= -\nabla A_0 - \dot{\mathbf{A}} \\ \mathbf{B} &= \nabla \times \mathbf{A} \end{aligned}}$$

$$\mathcal{L}_{ALP} = \frac{1}{2}\partial^\mu a\partial_\mu a - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}g_{a\gamma\gamma}F_{\mu\nu}\tilde{F}^{\mu\nu} = \frac{1}{2}\partial^\mu a\partial_\mu a - \frac{1}{2}m_a^2 a^2 + g_{a\gamma\gamma}\mathbf{E} \cdot \mathbf{B} a$$



Photon-photon-ALP vertex with coupling constant $g_{a\gamma\gamma}$



$\gamma \rightarrow a$ conversion in the external magnetic field \mathbf{B}
(Primakoff effect)

Other diagrams...

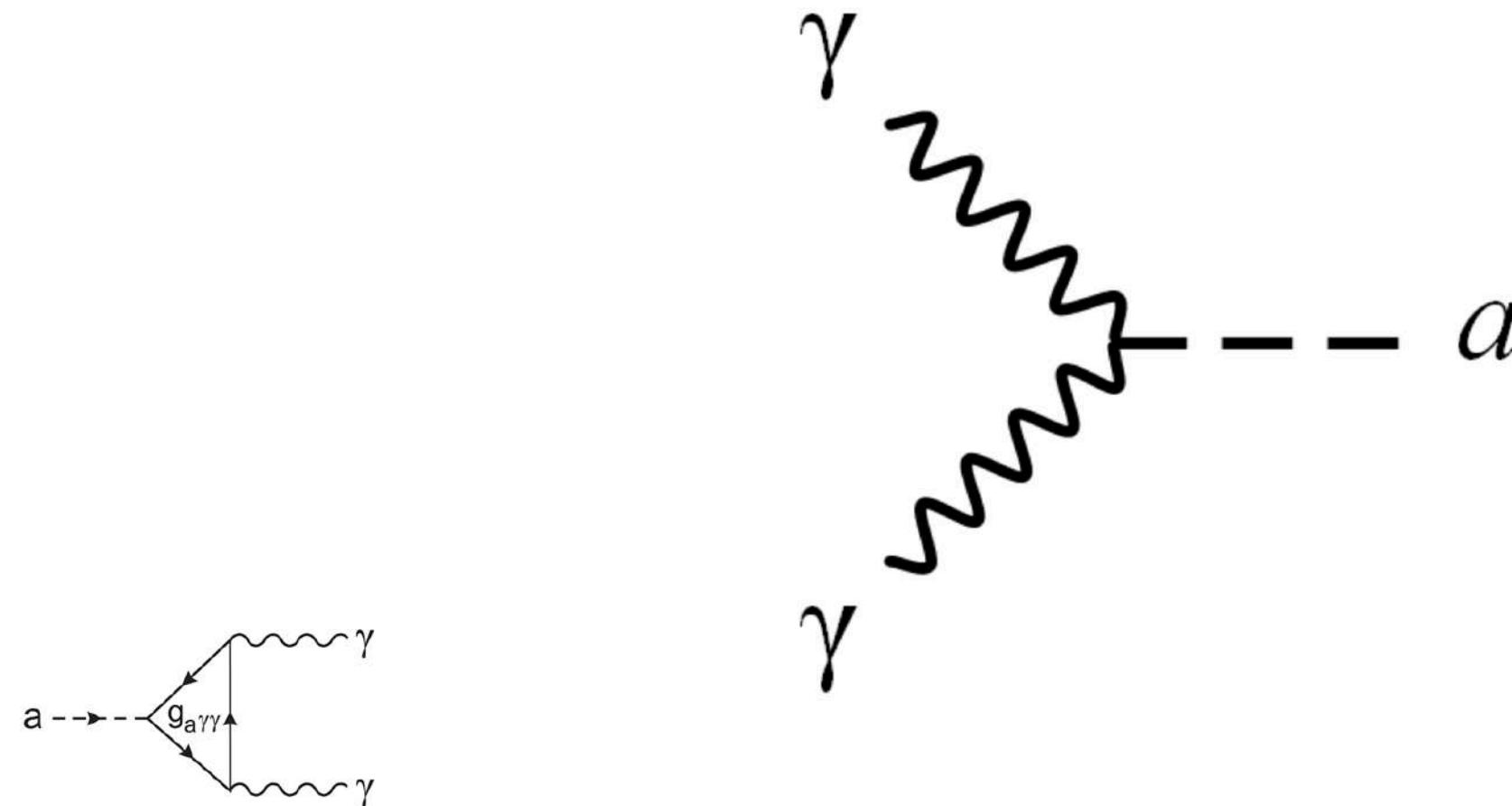
Axion and ALPs interaction with the SM

Axions and ALPs interact with the standard model particles

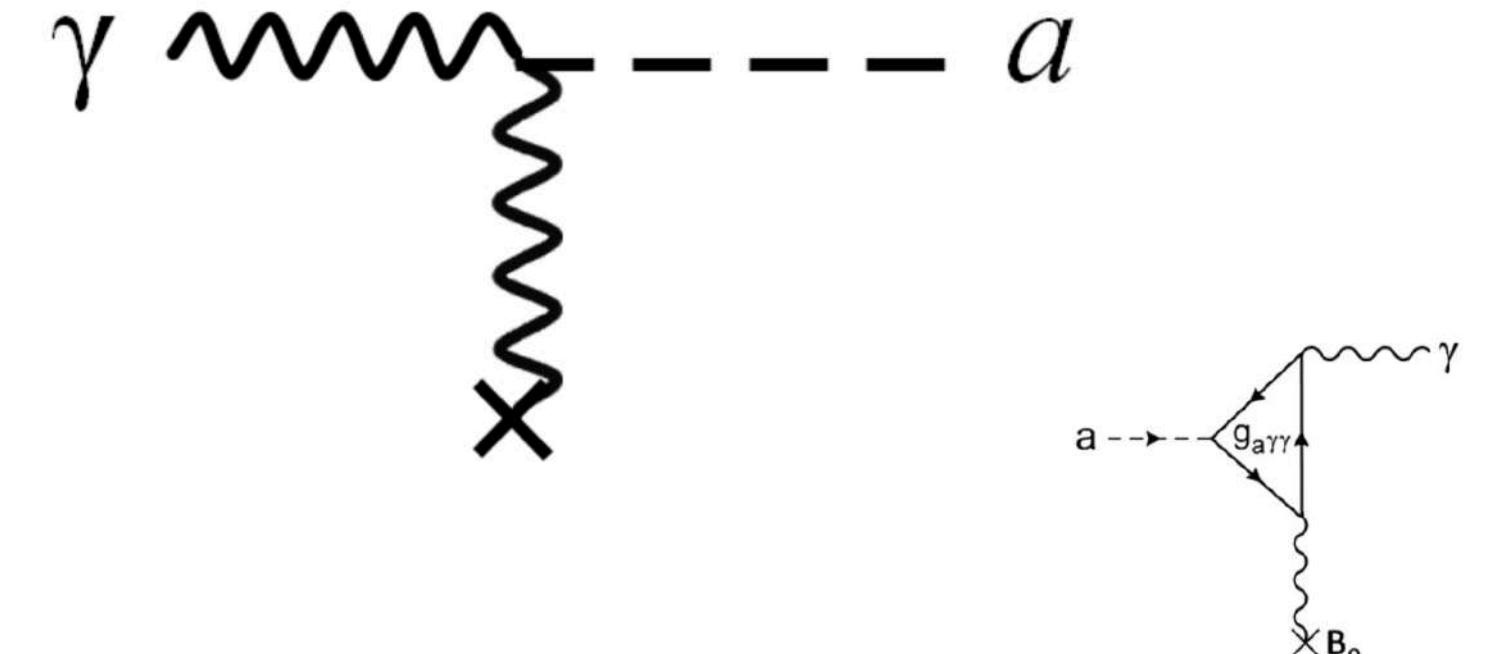
Photon - Axion electrodynamics

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(Primakoff effect)

Other diagrams...

Axion and ALPs interaction with the SM

Axions and ALPs interact with the standard model particles

Axion electrodynamics

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - J^\mu A_\mu - \frac{g_{a\gamma}}{4}F_{\mu\nu}\tilde{F}^{\mu\nu}a$$

$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$
$\tilde{F}^{\mu\nu} = \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta}$
$\mathbf{E} = -\nabla A_0 - \dot{\mathbf{A}}$
$\mathbf{B} = \nabla \times \mathbf{A}$

- We can interpret axion as the source of an effective current:

$$\partial_\mu F^{\mu\nu} = J^\nu - g_{a\gamma} \underbrace{\tilde{F}_{\mu\nu}\partial_\mu a}_{} \downarrow \\ J_a^\mu = g_{a\gamma}(-\mathbf{B} \cdot \nabla a, -\mathbf{E} \times \nabla a + \partial_t a \mathbf{B})$$

Maxwell's equations:

$$\nabla \cdot \mathbf{E} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} + \mathbf{J}$$

Extended Maxwell's equations:

$$\nabla \cdot \mathbf{E} = \rho - g_{a\gamma} \mathbf{B} \cdot \nabla a$$

$$\nabla \cdot \mathbf{B} = 0$$

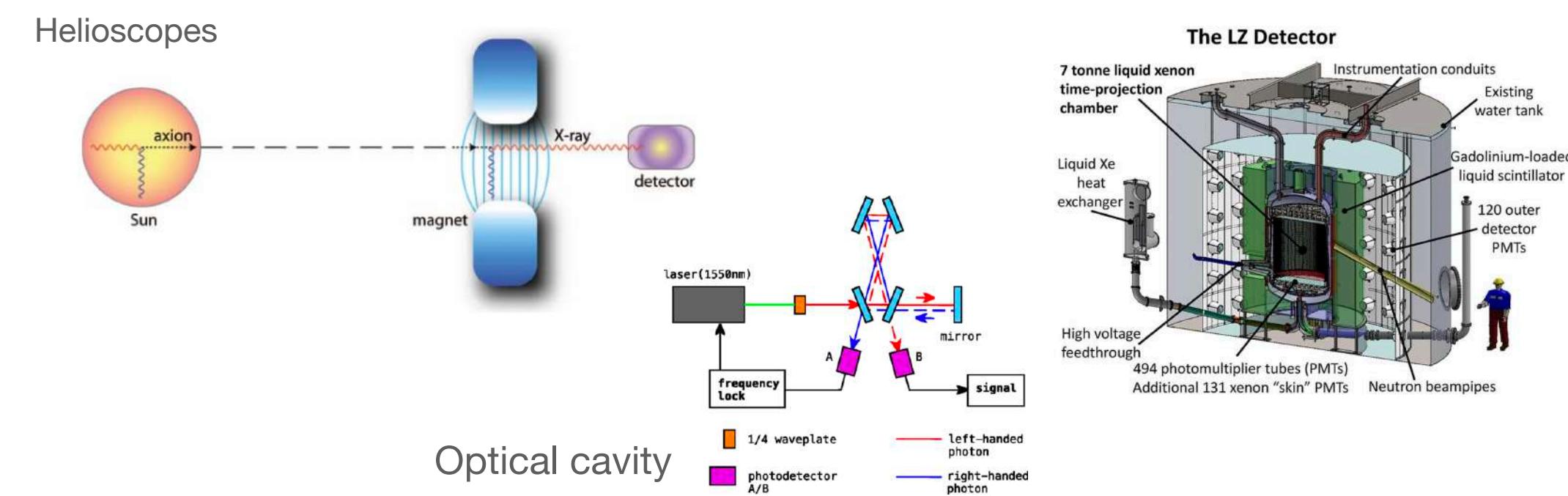
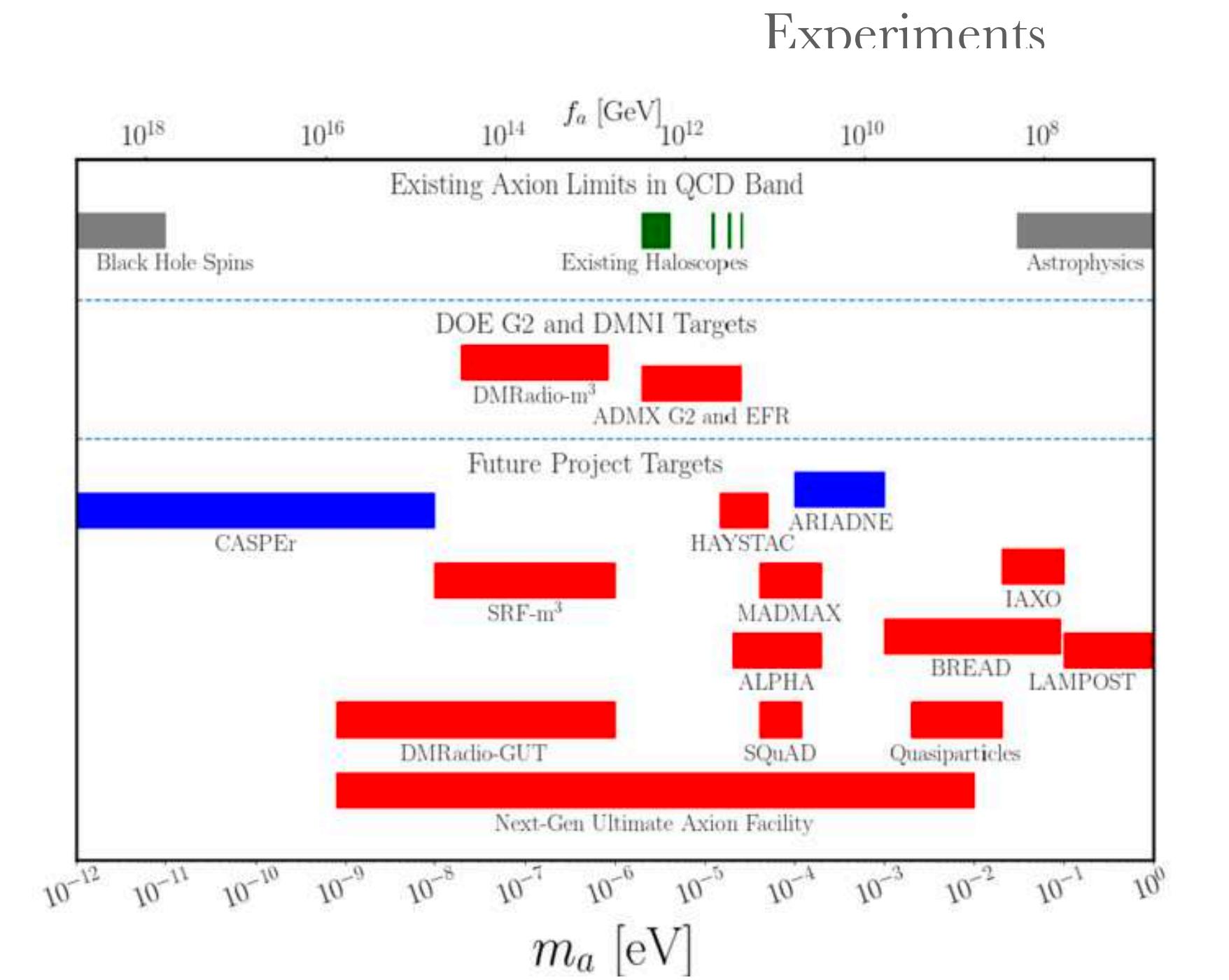
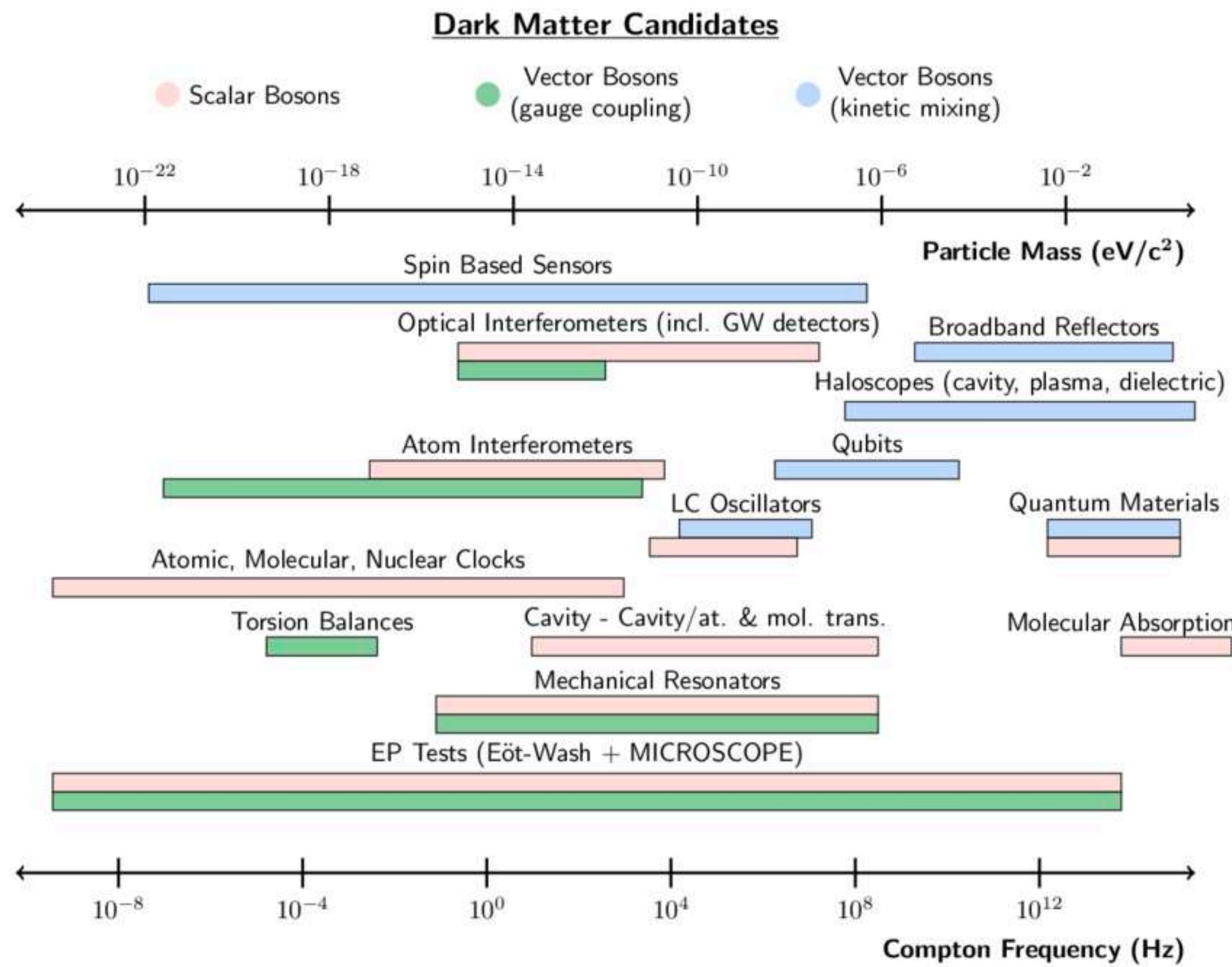
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} + \mathbf{J} - g_{a\gamma} \left(\mathbf{E} \times \nabla a - \frac{\partial a}{\partial t} \mathbf{B} \right)$$

Axion and ALPs interaction with the SM

"Direct Detection": axion/ALPs experiments

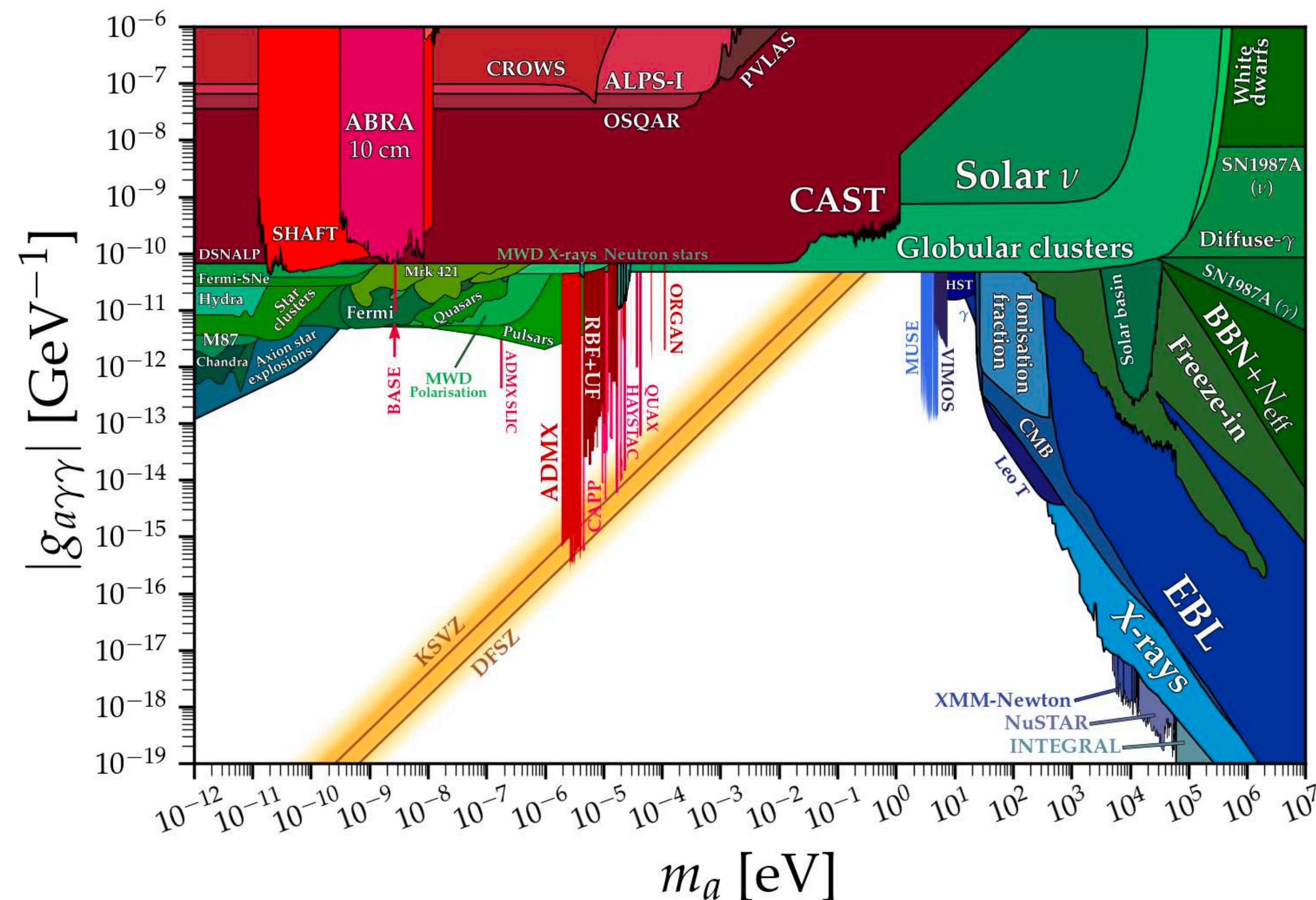
Overview of experimental techniques and the mass ranges they target



Axion and ALPs interaction with the SM

Bounds on Axion-photon coupling

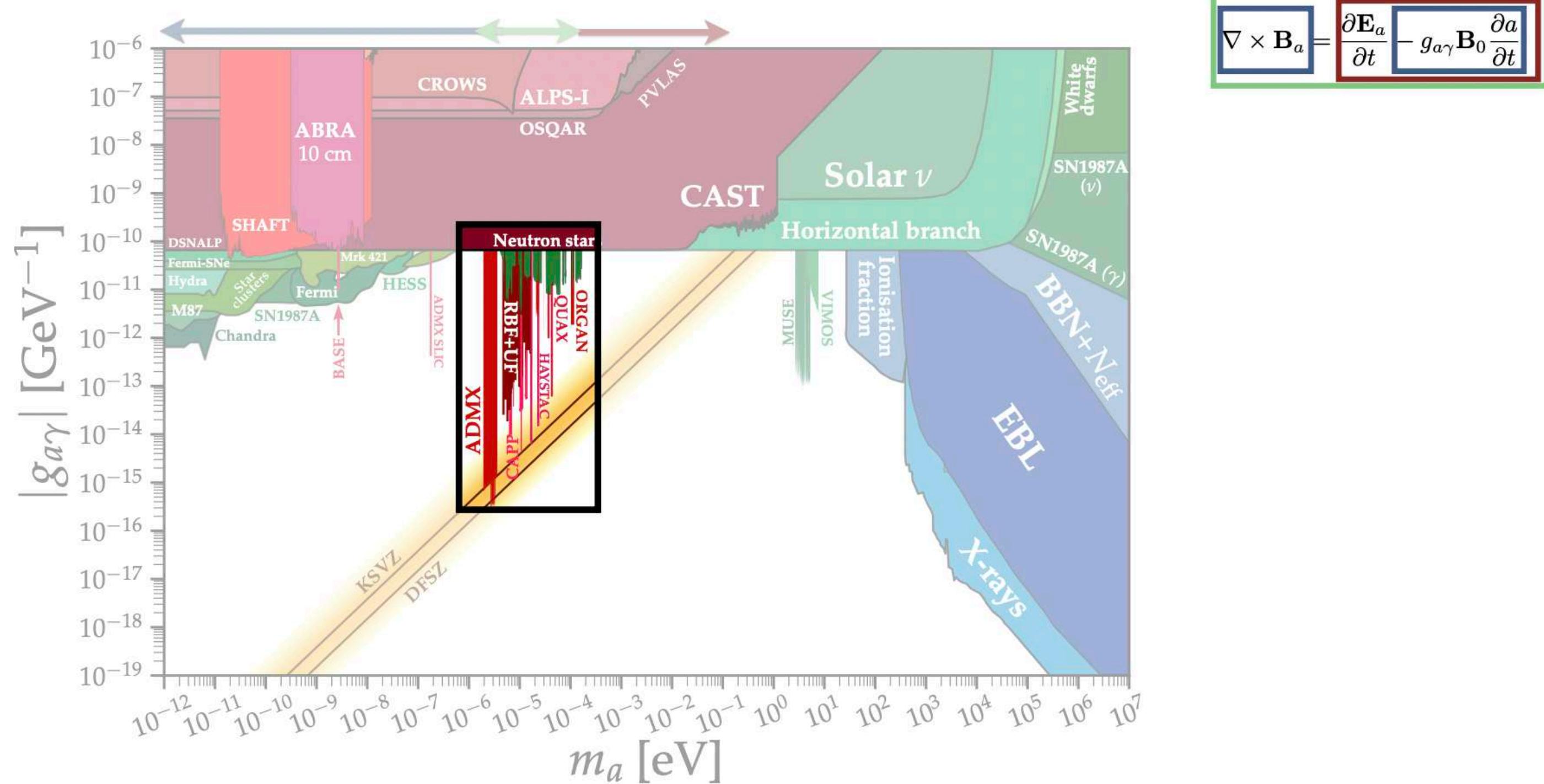
Includes direct and indirect detection



Website with up-to-date with axion/ALP bounds: <https://cajohare.github.io/AxionLimits>

Axion and ALPs interaction with the SM

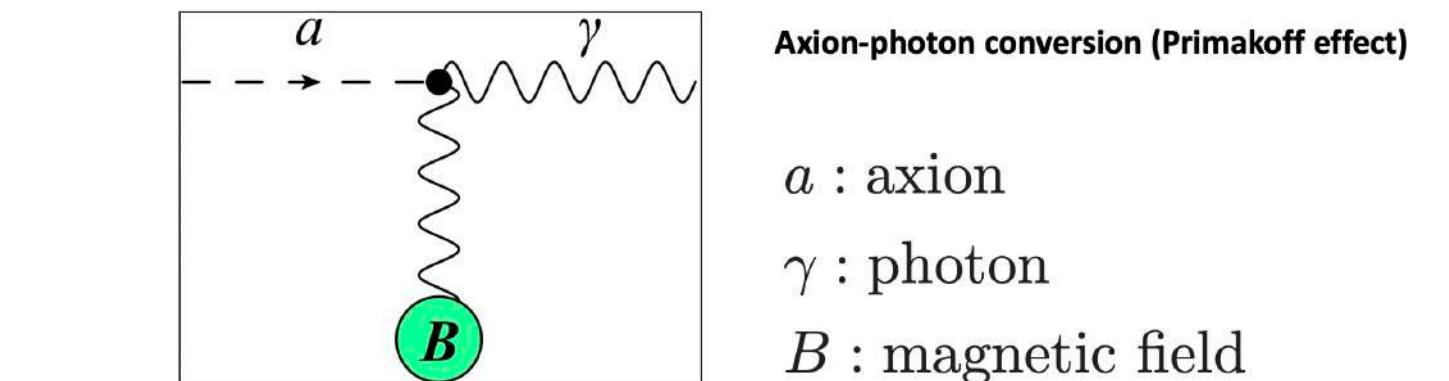
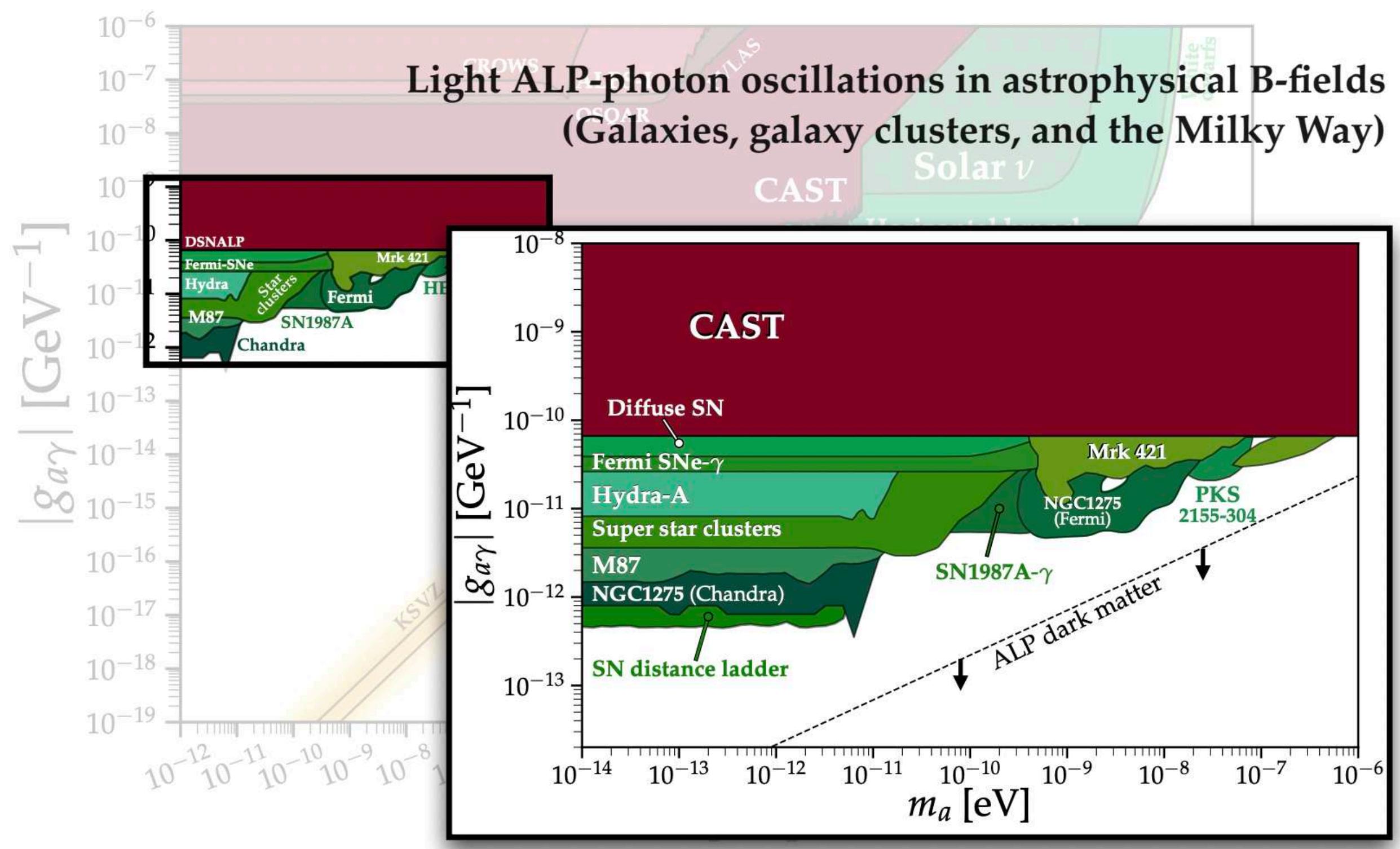
Bounds on Axion-photon coupling



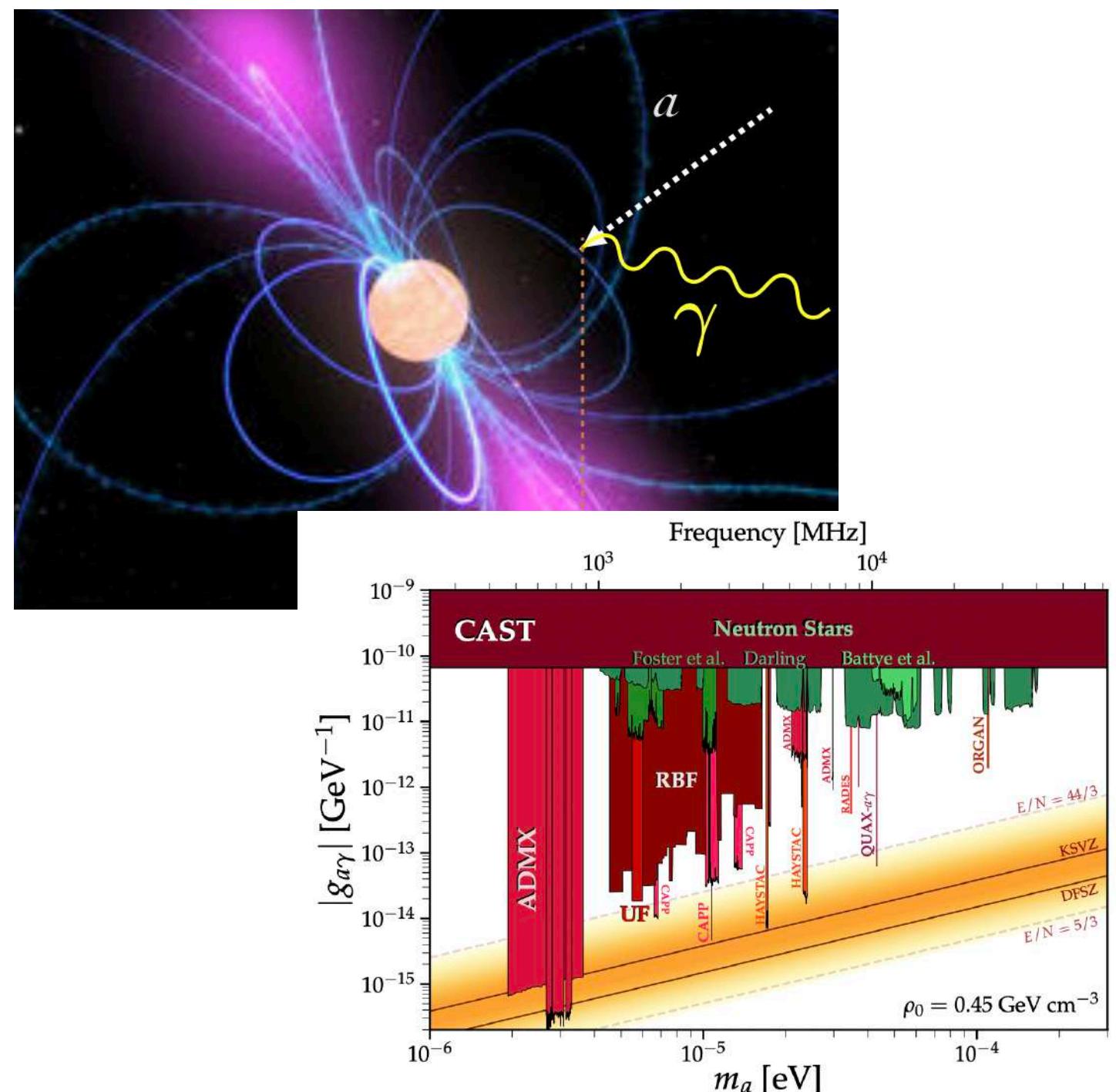
Axion and ALPs interaction with the SM

Indirect Detection

In astrophysical systems



DM axions in neutron star magnetospheres

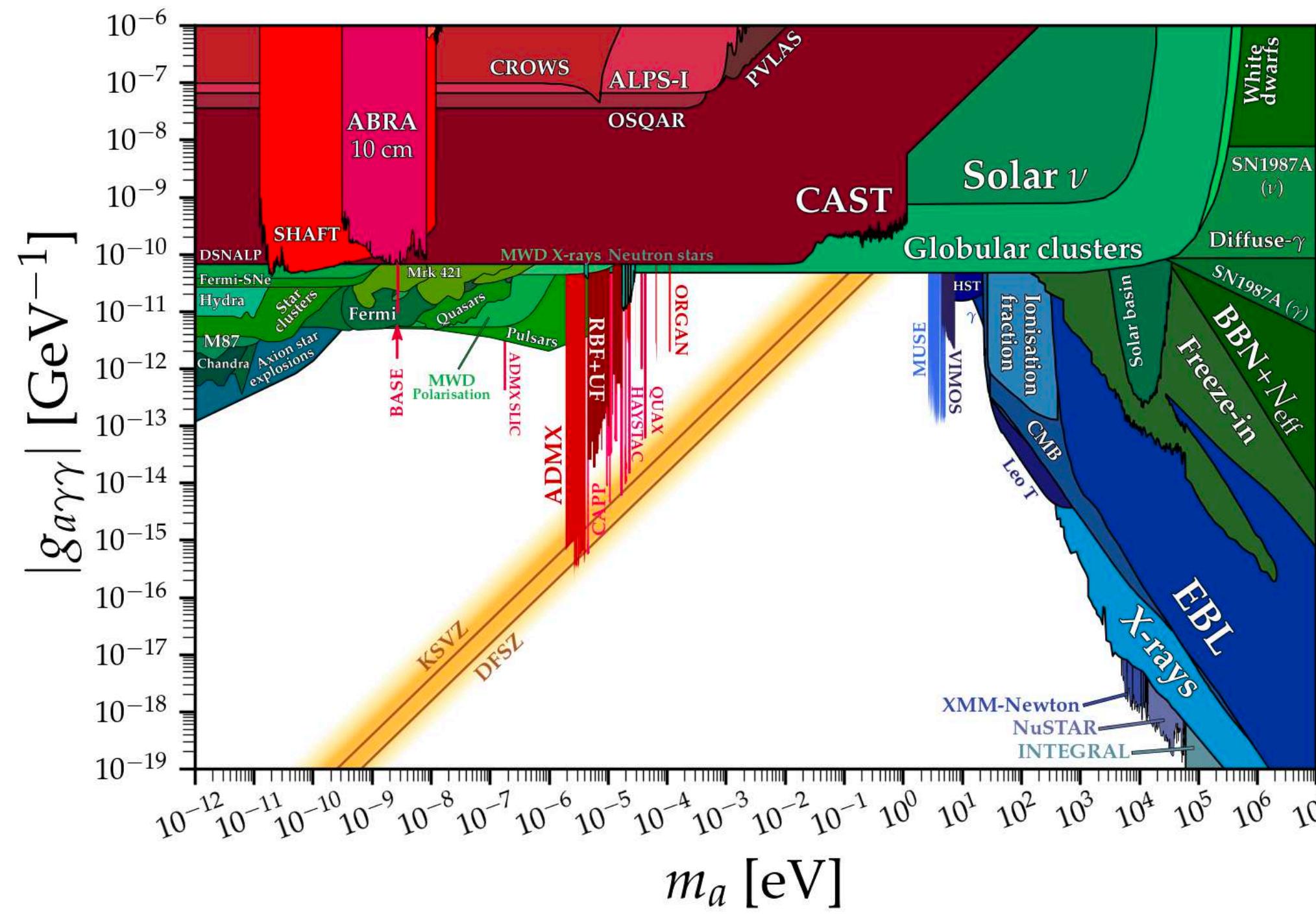


Axion and ALPs interaction with the SM

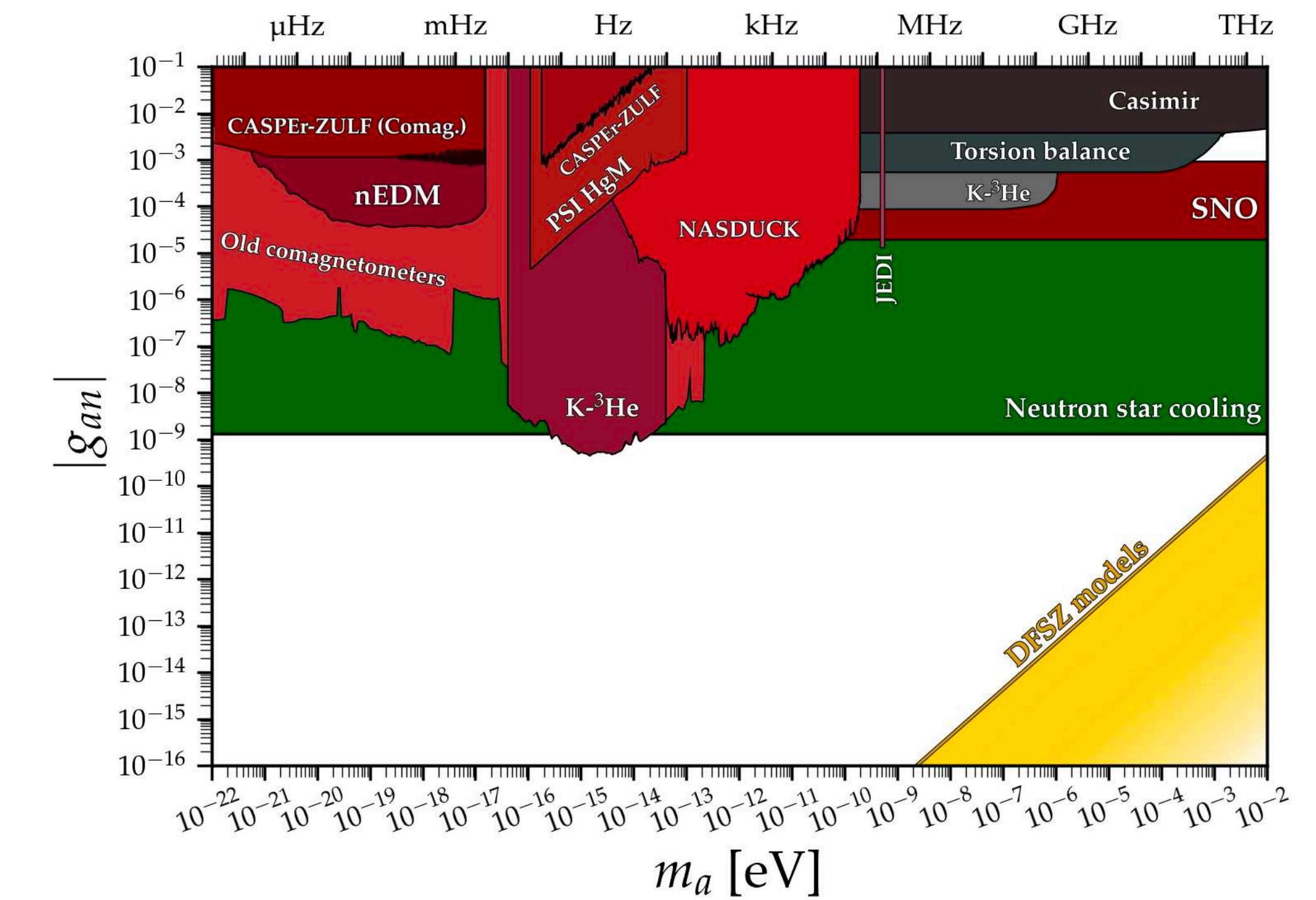
Bounds

Includes direct and indirect detection

Axion-electron coupling



Axion-neutron coupling



+ many more: axion-proton, dark photon, ...

Website with up-to-date with axion/ALP bounds: <https://cajohare.github.io/AxionLimits>
(Includes notebooks)

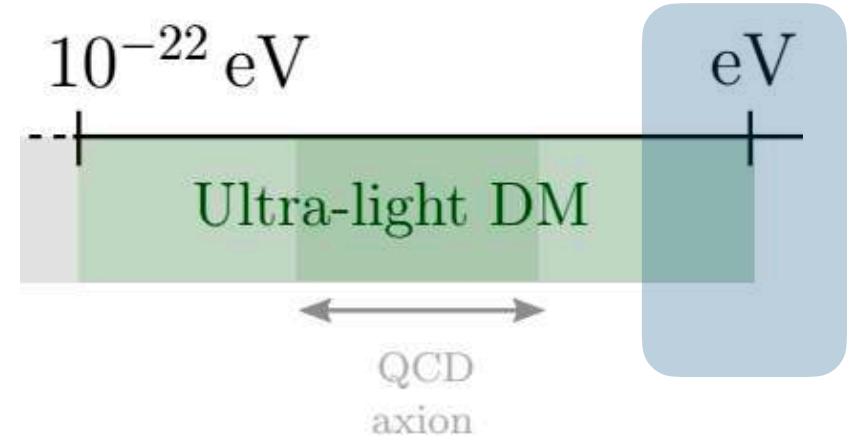
Superfluid Dark Matter



MakeAGIF.com

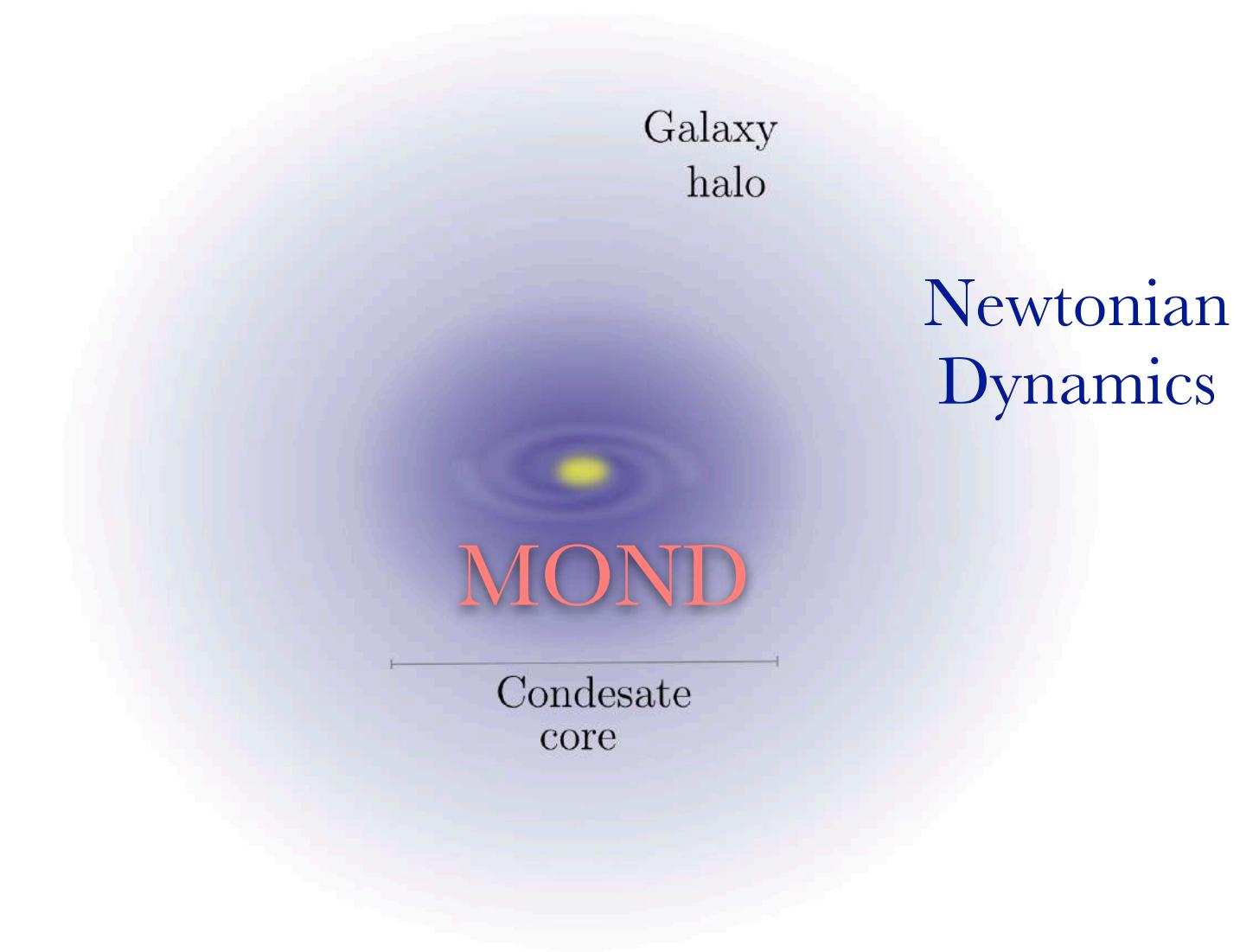


Superfluid Dark Matter



Lasha Berezhiani and Justin Khoury (2016)

Large scales:
DM behaves like standard
particle DM (**CDM**).



Galactic scales:
DM forms a **superfluid**
→ emergent **MOND** dynamics
in galaxies

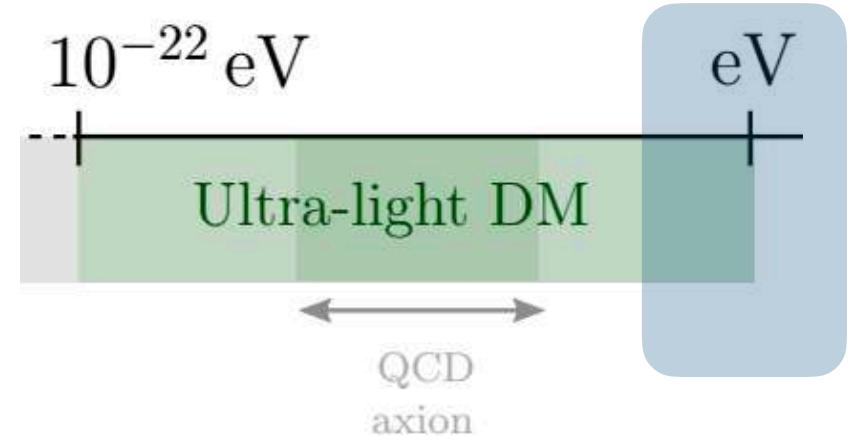


$$a = \begin{cases} a_N^b, & a_N^b \gg a_0, \\ \sqrt{a_N^b a_0}, & a_N^b \ll a_0. \end{cases}$$

Similar phenomenology than the FDM & SIFDM + explains the **rotations curves and scaling relations**

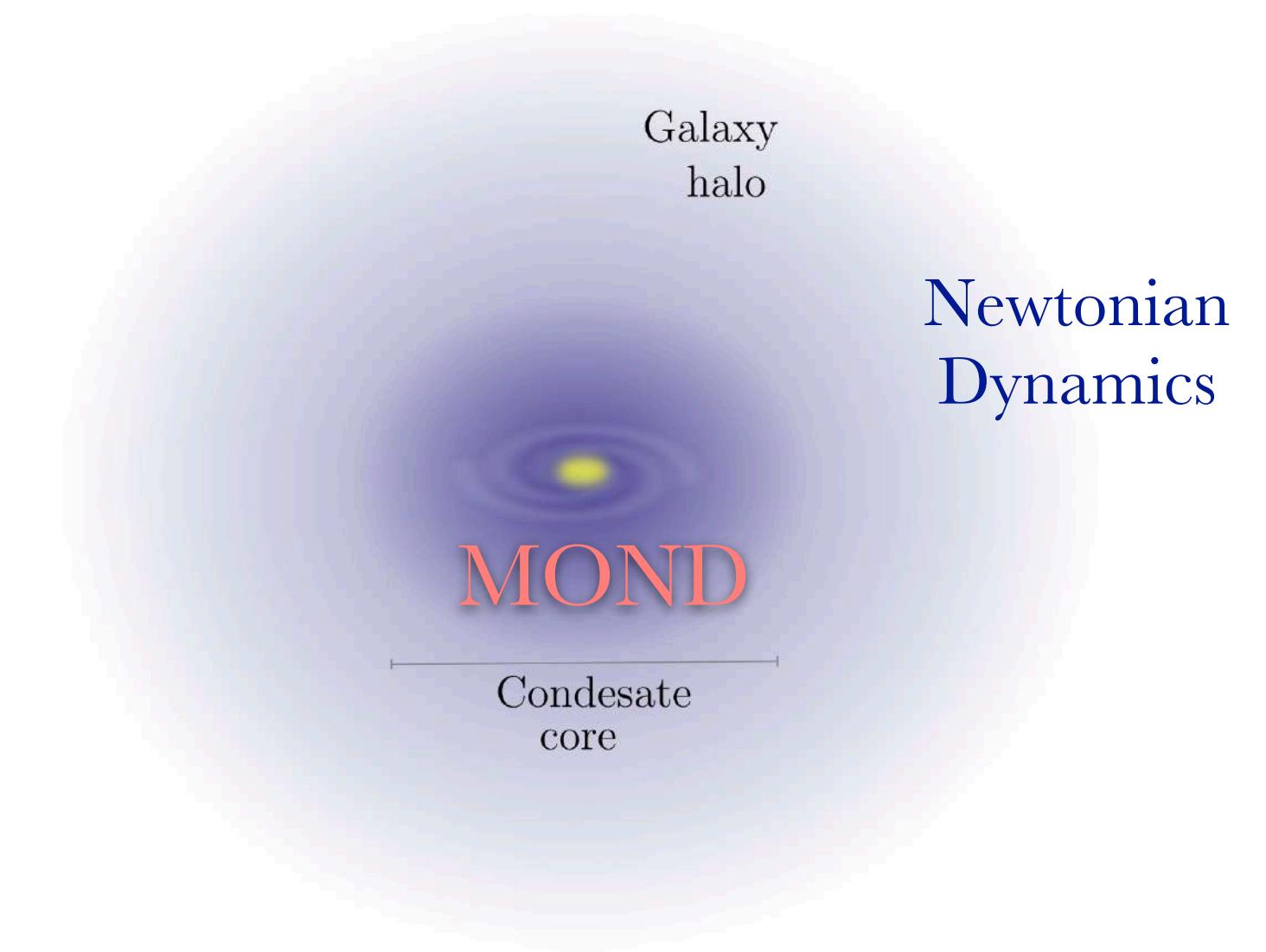
Suppresses small structures, dyn. effects, formation of cores

Superfluid Dark Matter



Lasha Berezhiani and Justin Khoury (2016)

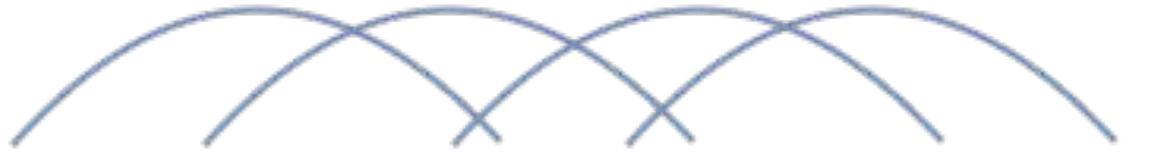
Large scales:
DM behaves like standard
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To describe non-relativistic MOND, it is imposed that:

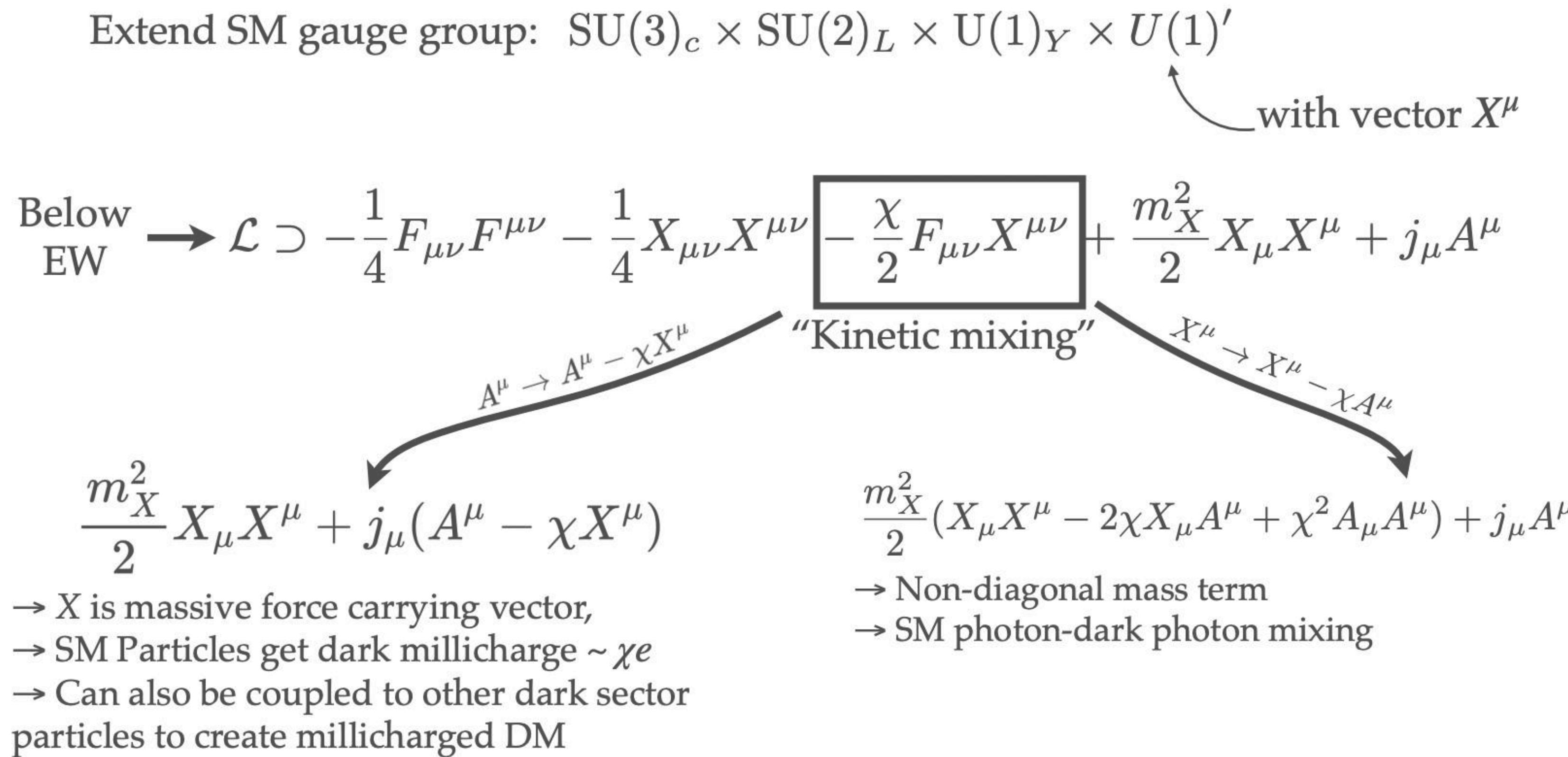
$$\mathcal{L} = P(X), \quad P(X) = \frac{2\Lambda}{3} (2m)^{3/2} X \sqrt{|X|}$$

Galactic scales:
DM forms a **superfluid**
→ emergent **MOND** dynamics
in galaxies



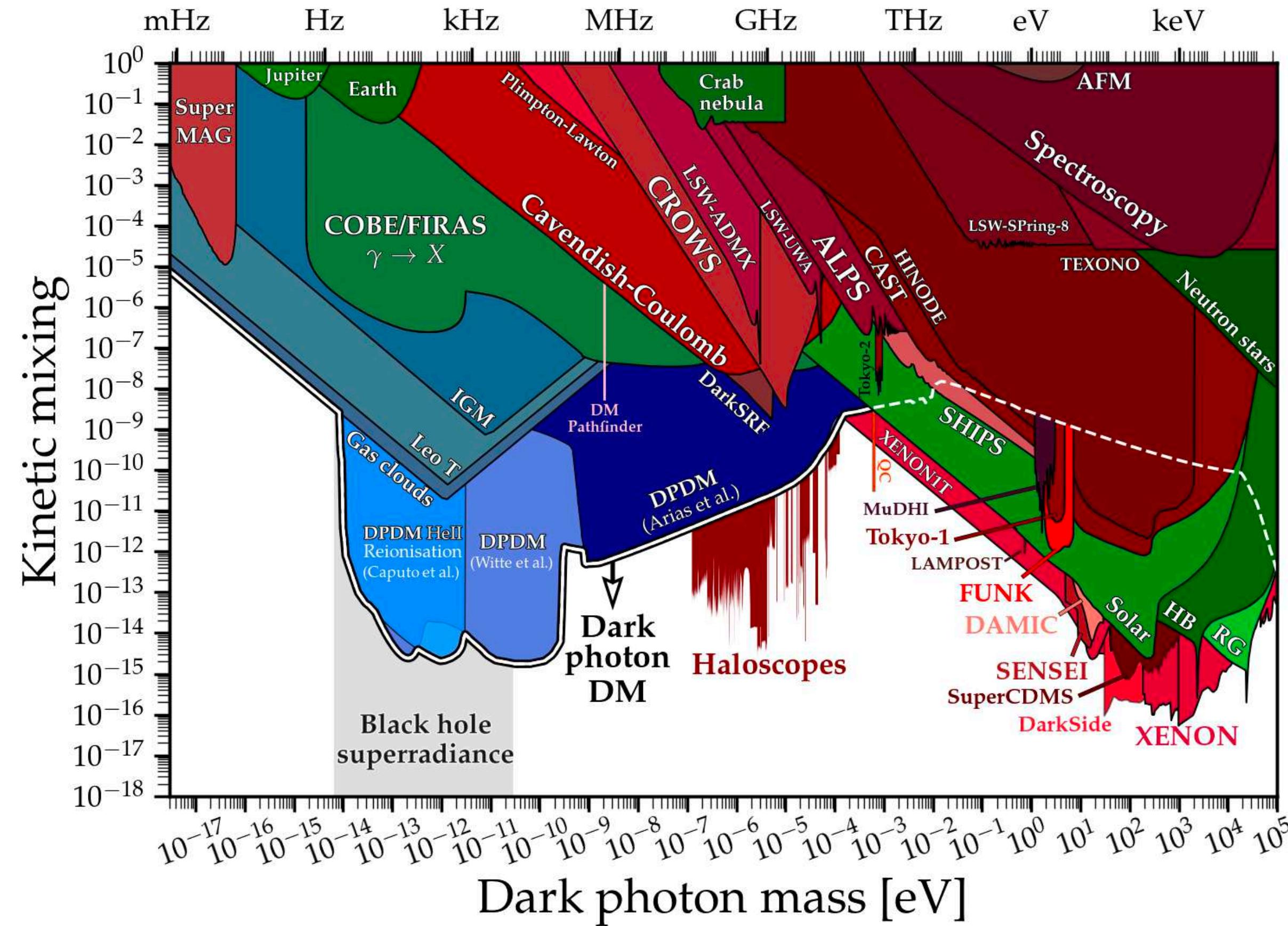
Dark photon

New gauge dark sector - spin 1



Dark photon

New gauge dark sector - spin 1



Website with up-to-date with axion/ALP bounds: <https://cajohare.github.io/AxionLimits>
(Includes notebooks)

S8 tension

Ref.: K. Rogers et al 2023

Changes in the small scale paradigm can change the behaviour of DM in many scales, including cosmology

Ex.: Fuzzy DM

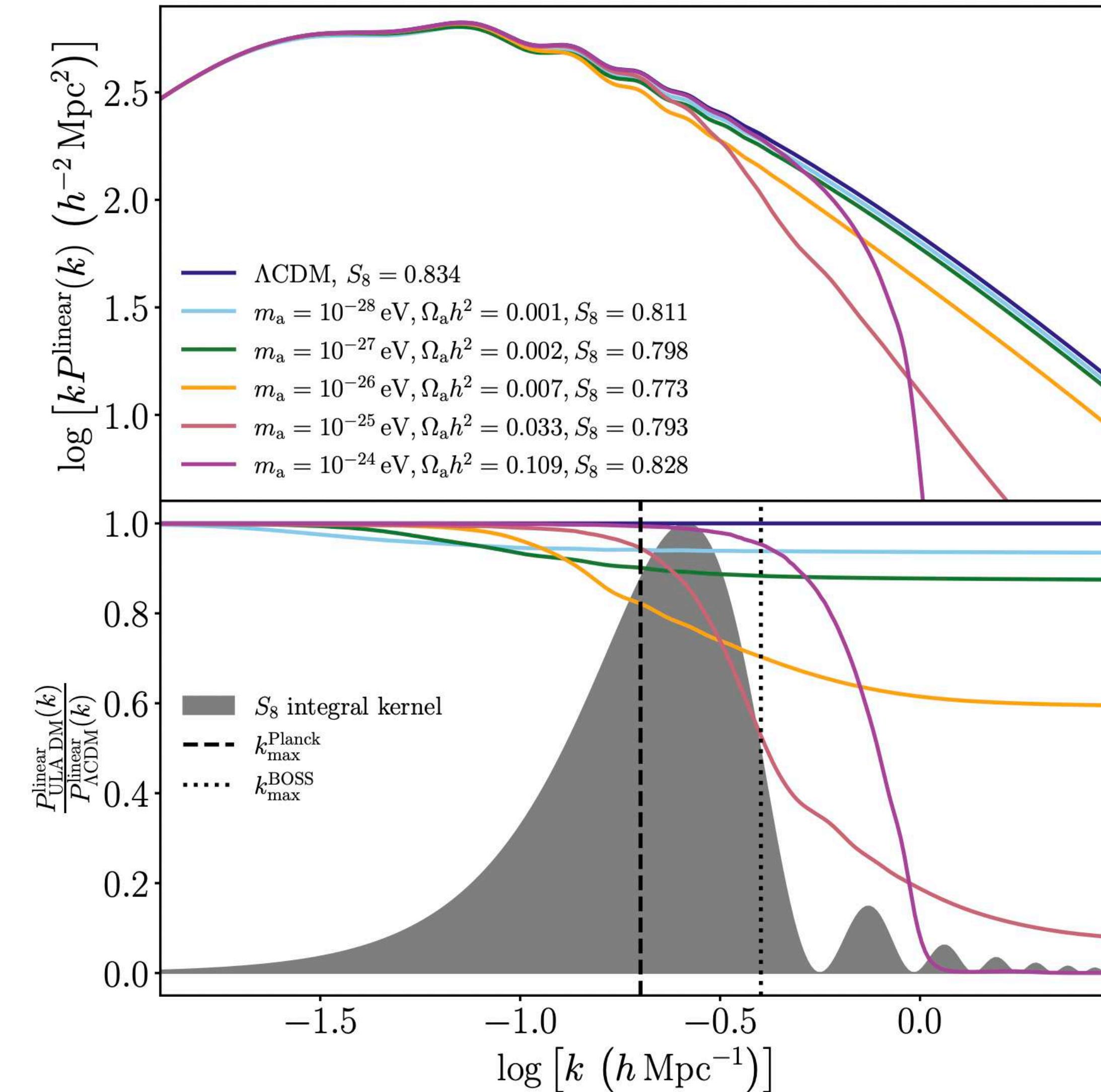
$$\sigma_8 = \int d\ln k \frac{k^3}{2\pi} W^2(k) P^{\text{linear}}(k)$$

$$S_8 = \sqrt{\frac{\Omega_m}{0.3}} \sigma_8$$

The presence of ULAs can significantly lowers S8 for:

$$m_a \in [10^{-27}, 10^{-25}] \text{ eV}$$

S8 is lowered because the Jeans scale today for $m_a = 10^{-25} - 10^{-26}$ eV is about $\lambda_J = 4 - 12 h^{-1} \text{ Mpc}$



S₈ tension

Ref.: K. Rogers et al 2023

Ex.: Fuzzy DM

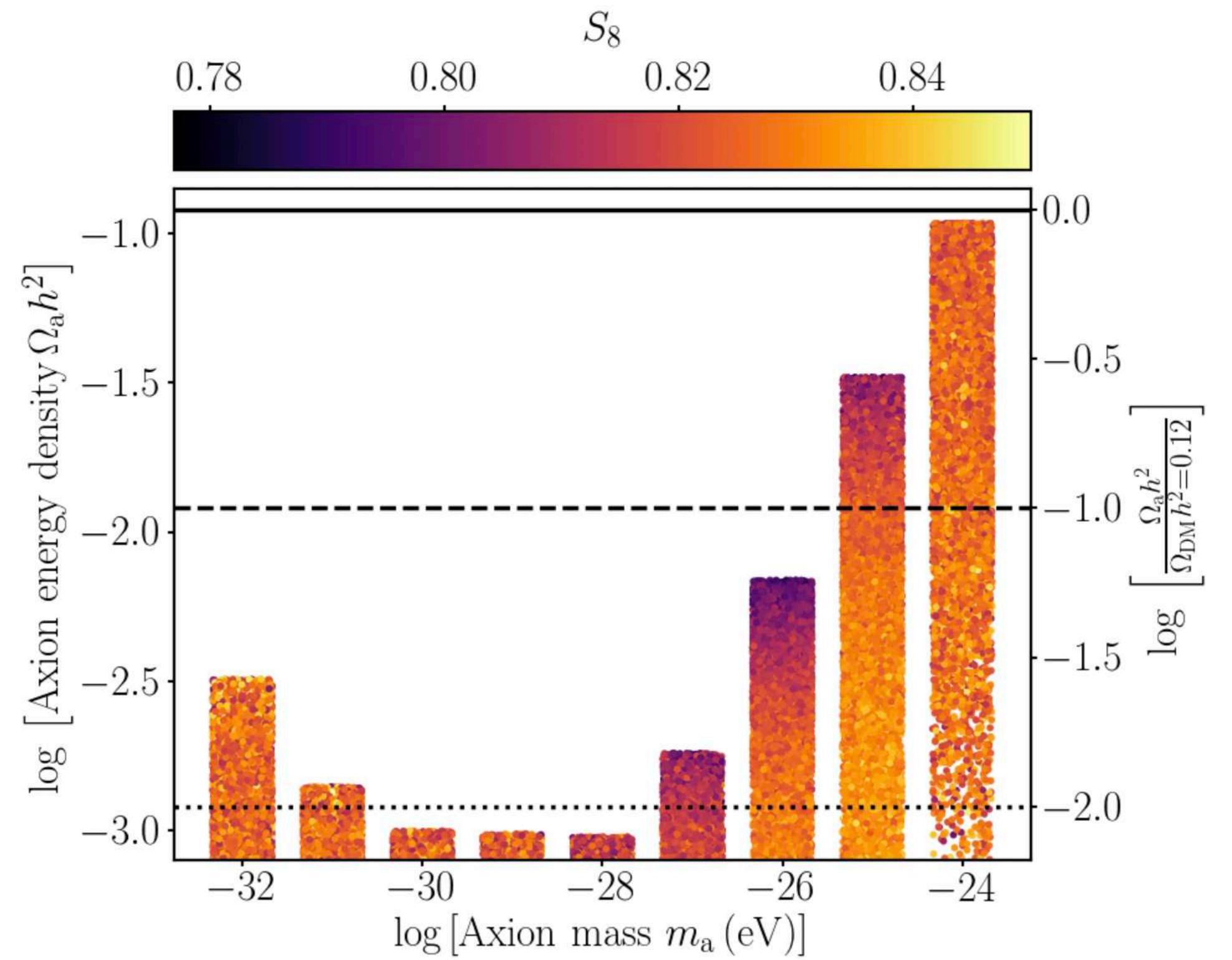
Planck + BOSS

The presence of ULAs with mass

$$10^{-28} \text{ eV} \leq m_a \leq 10^{-25} \text{ eV}$$

can improve consistency between CMB and galaxy clustering
(reduce the S₈ discrepancy)

from 2.6σ to 1.7σ



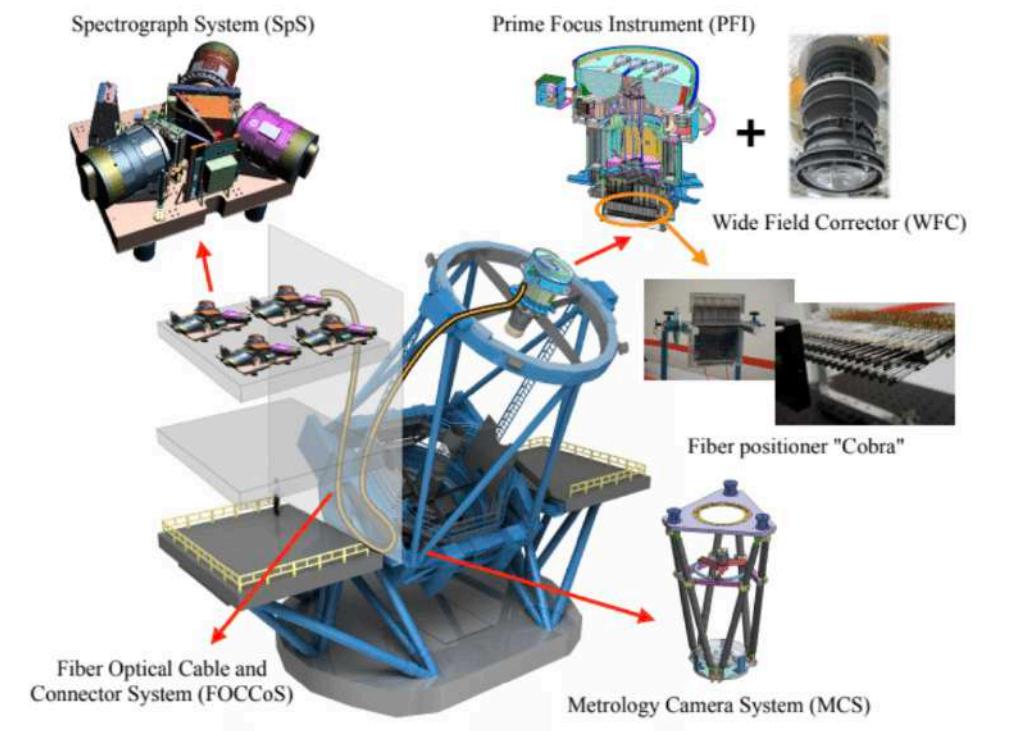
Ex.:

- H0 tension: Early dark energy - axion-like particle
 - Model address H0 and S8 tensions
- "Chameleon EDE", Karwal et al 2021

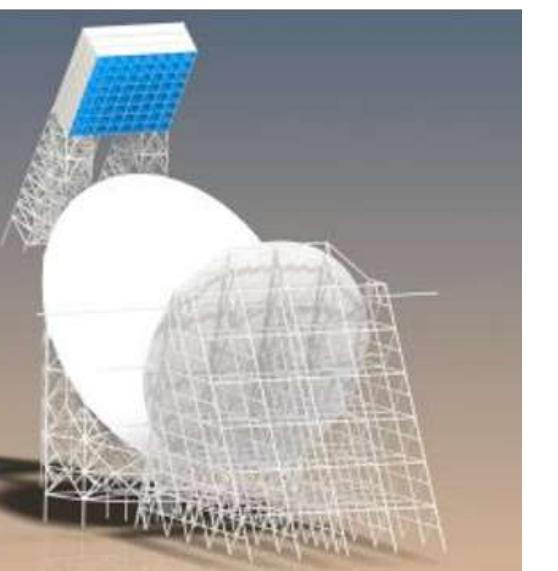
Future

Observations

Prime Focus Spectrograph (PFS)



BINGO telescope

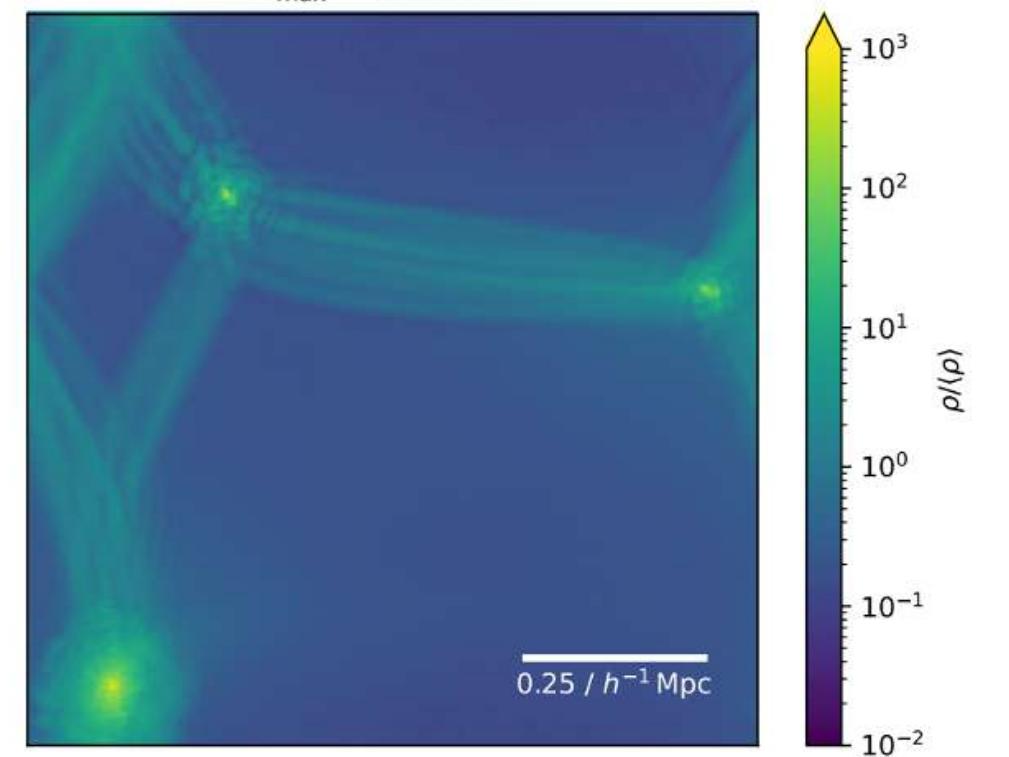


Vera Rubin observatory (LSST)



Simulations

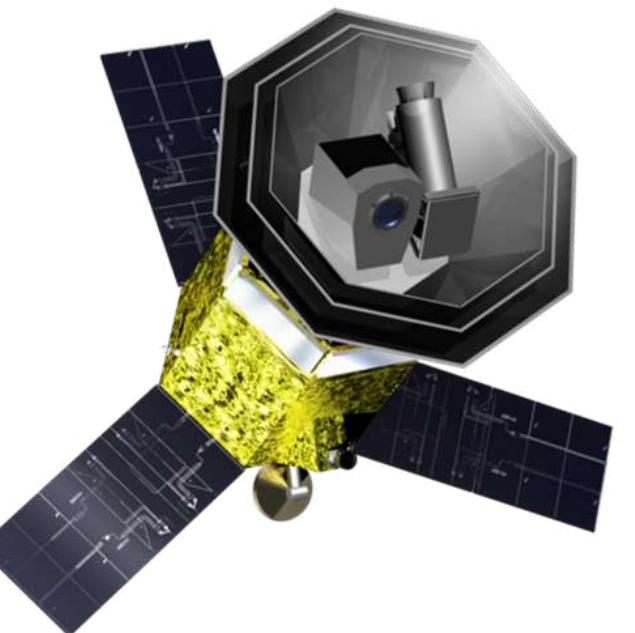
FDM: 256^3 , $mc^2 = 1.75 \times 10^{-23}$ eV, $z = 0.00$
 $v_{\max} = 88.1$ km/s



CMB-S4



LiteBIRD

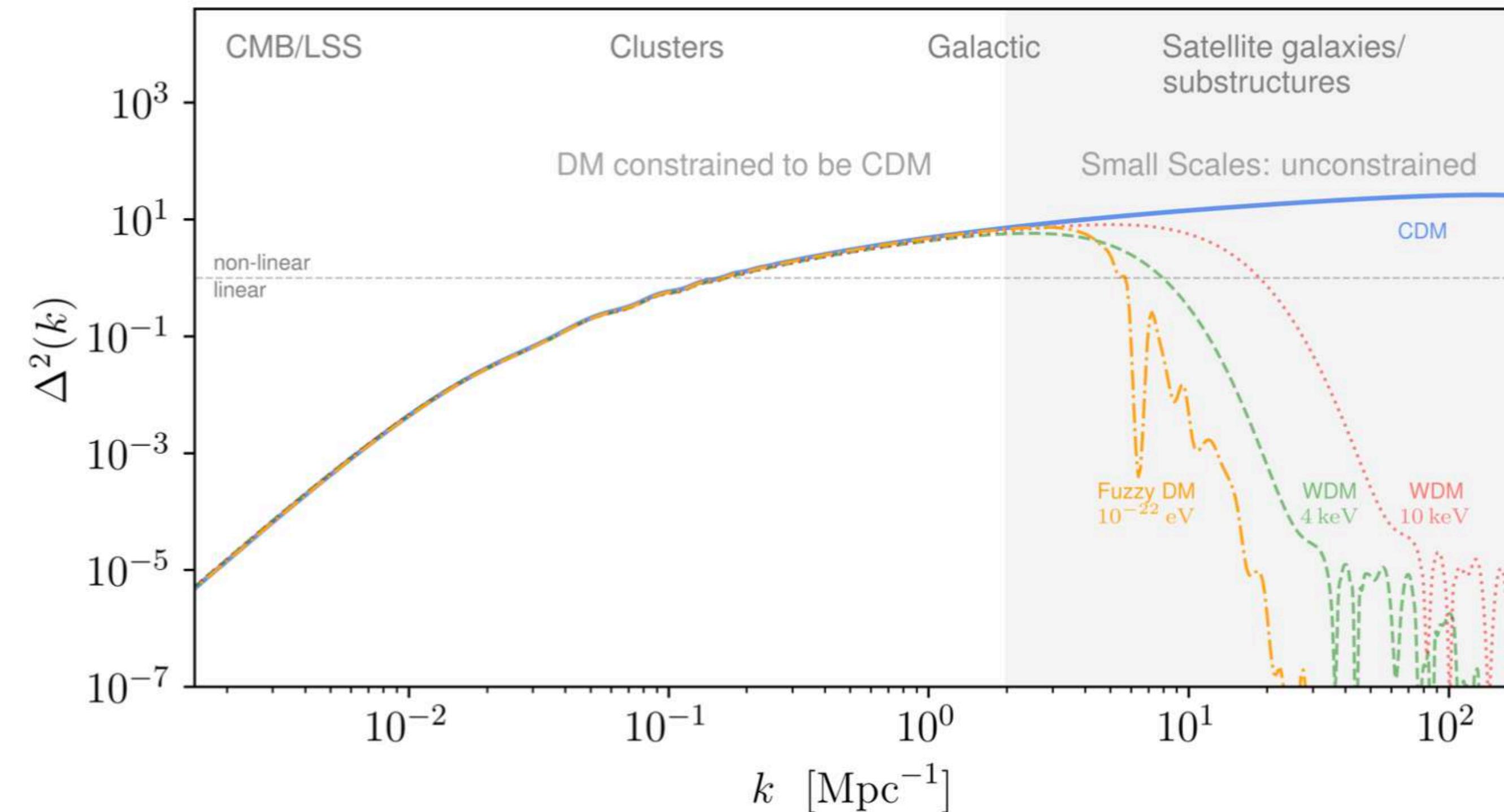


New probes

Sub-galactic power spectrum

New probes

Using gravitational probes, **strong lensing** and **stellar streams**, to describe substructures



Sub-galactic power spectrum

New probes

Using gravitational probes, **strong lensing** and **stellar streams**, to describe substructures

A. Diaz Rivero, et al. (2017); Diaz Rivero, et al. , (2018)

Substructure convergence power spectrum

Develop a formalism to compute the substructure convergence power spectrum for different populations of dark matter subhalos.

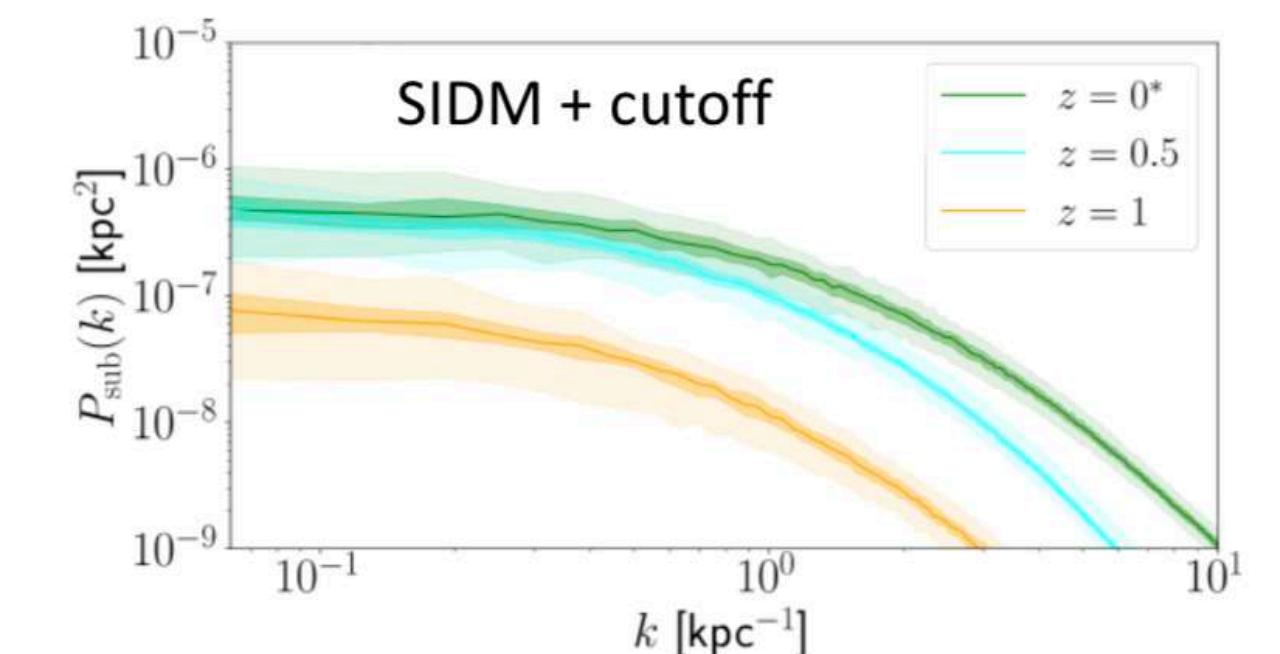
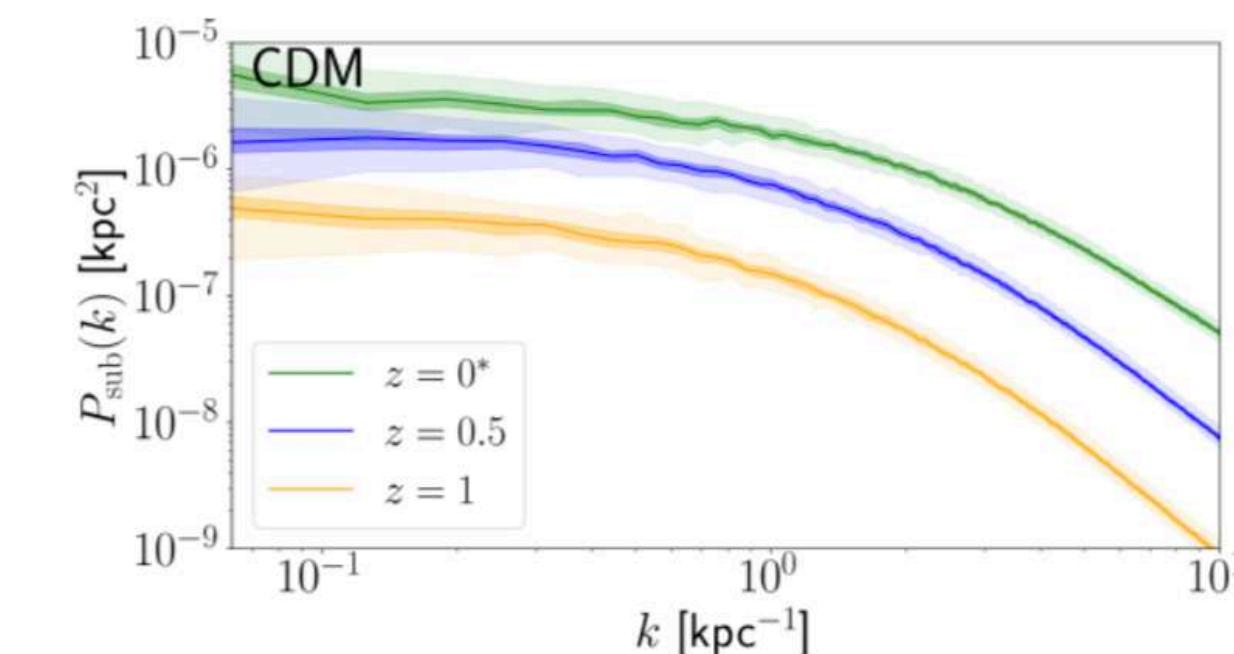
Bayer et al. (2018) ; Auger et al. 2009
FDM: Kawai et al. (2021)

Hezaveh et al. (2016) (projected mPS by using strong lensing)

Change of language: instead of talking about lensing perturbations in terms of individual subhalos, look at the correlation function of the projected density field.

(based on Dvorkin's slide)

$$P_{\text{sub}}(k) = P_{1\text{sh}}(k) + P_{2\text{sh}}(k)$$



Sub-galactic power spectrum

Using gravitational probes, **strong lensing** and **stellar streams**, to describe substructures

Substructure convergence power spectrum

Sten Delos and Fabian Schmidt (2021)

Stellar streams: perturbed by passing substructure. Good gravitational probe, since given their low dynamical temperature and negligible self-interaction, it retains the memory of those encounters.

THIS WORK: Fully analytical understanding of the stream perturbations!

Power spectrum of a stream's stellar density is analytically related to that of the substructure background:

$$\boxed{P_*(k, t)} = \chi_* \left(k\sigma_0 t, \frac{D}{k\sigma_0^3} \right) \frac{k^2 t^2}{3} P_{\Delta v}(k, t)$$

Stream power Substructure power

$$P_{\Delta v}(k, t) = 16\pi^4 G^2 \bar{\rho}^2 k^2 t \int_k^\infty \frac{dq}{q} \frac{\mathcal{P}(q)}{q^6} \int d^3 u \frac{f(u)}{u} \theta_H(qu - kv)$$

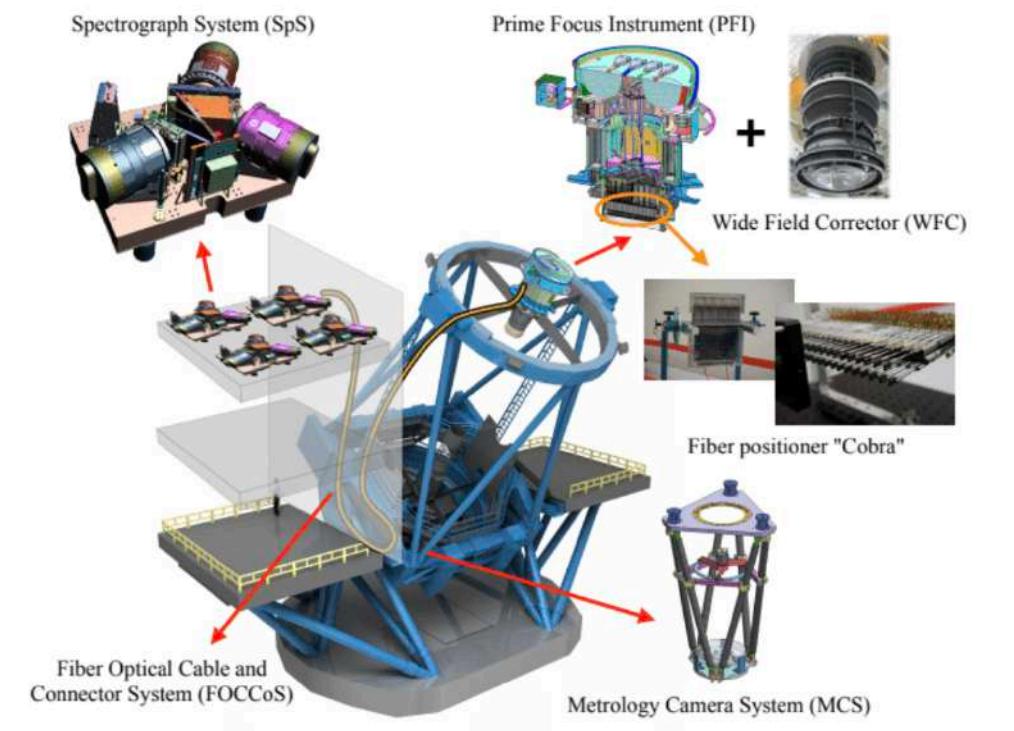
- Previous:
- Mostly numerical
 - Perturbations → sub-halo mass function

Relates the stellar stream perturbation to the surrounding matter distribution, from dark and luminous substructure

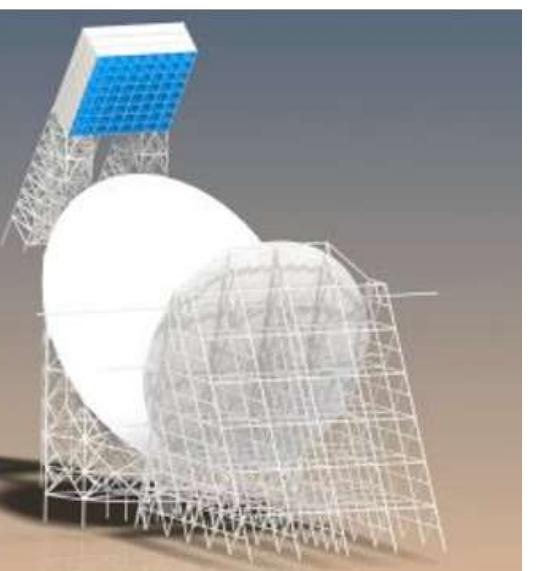
Future

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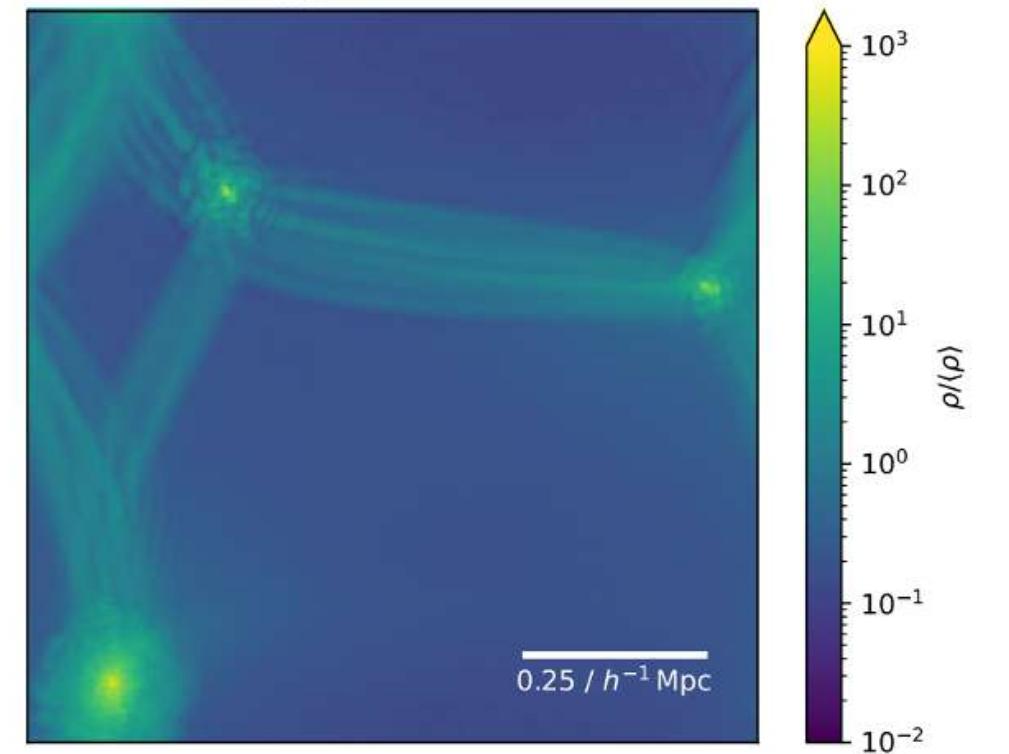


Vera Rubin observatory (LSST)



Simulations

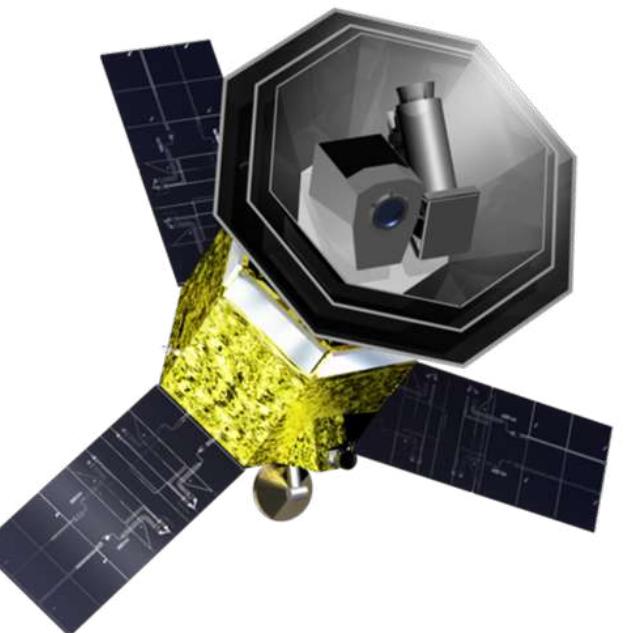
FDM: 256^3 , $mc^2 = 1.75 \times 10^{-23}$ eV, $z = 0.00$
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CMB-S4



LiteBIRD



New probes

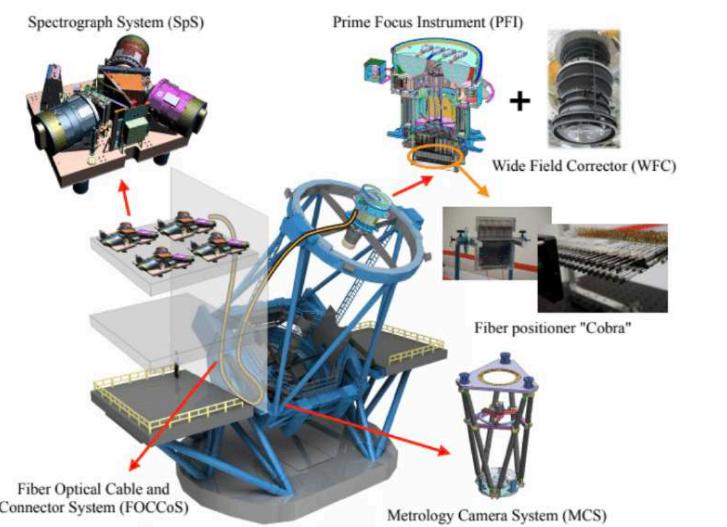
Future - signals in cosmology

Observations

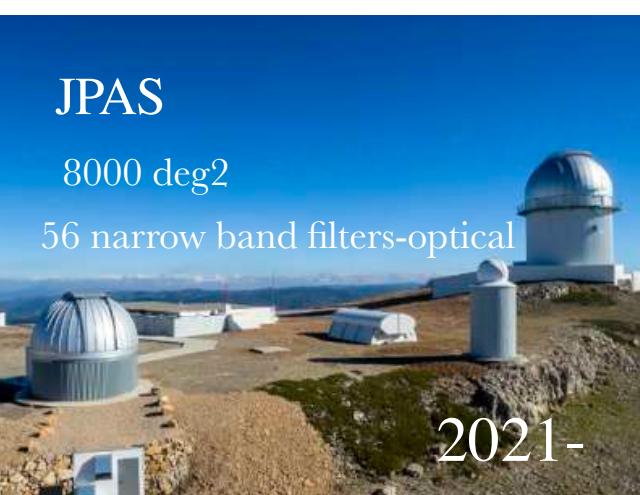
Photometric and spectroscopic surveys



Prime Focus Spectrograph (PFS)



21cm



GWs



CMB

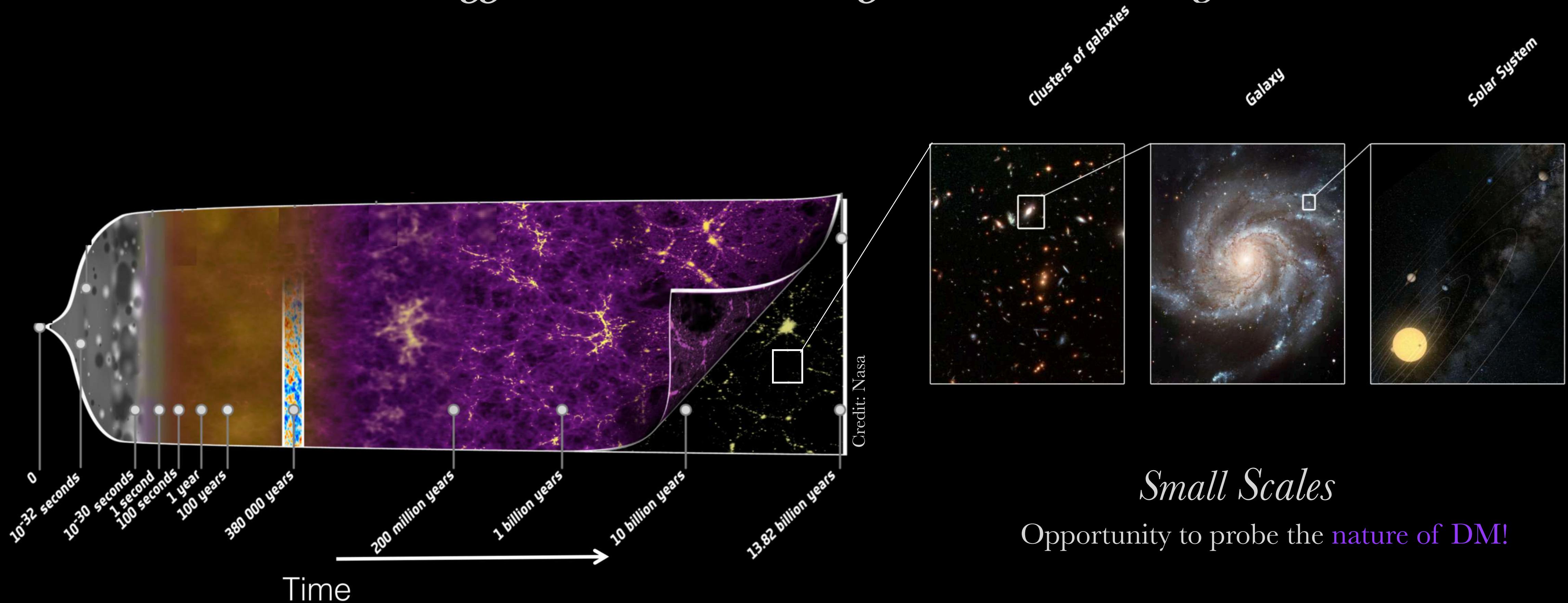


CMB-S4
Next Generation CMB Experiment



Modified from Jia Liu

*Small scales can offer some **hints** of the nature of DM*



Small Scales

Opportunity to probe the **nature of DM!**



Summary

DM builder's guide

What we learned from observations

- Cold or warm Thermal candidate: $m_{dm} \geq$ keV
Or produced cold by a non-thermal mechanism
 - Reproduce large and small scale distribution
Clusters like CDM on large scales $k \lesssim 10 \text{ Mpc}^{-1}$
Clustering on scales smaller than $k \gtrsim 10 \text{ Mpc}^{-1}$ highly unconstrained

- Non-interacting or weakly interacting

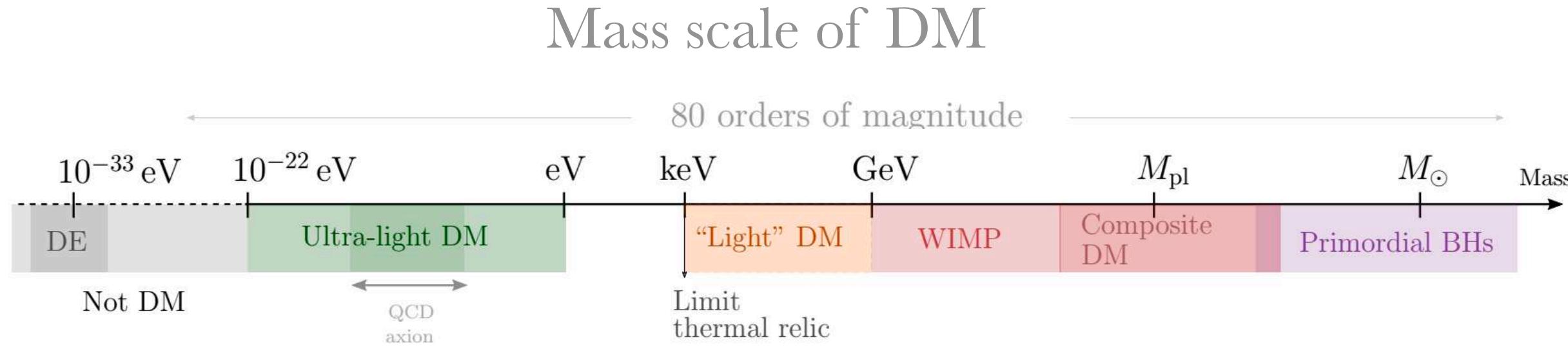
Can have a small **self interaction**.

Can interact via the *weak force*

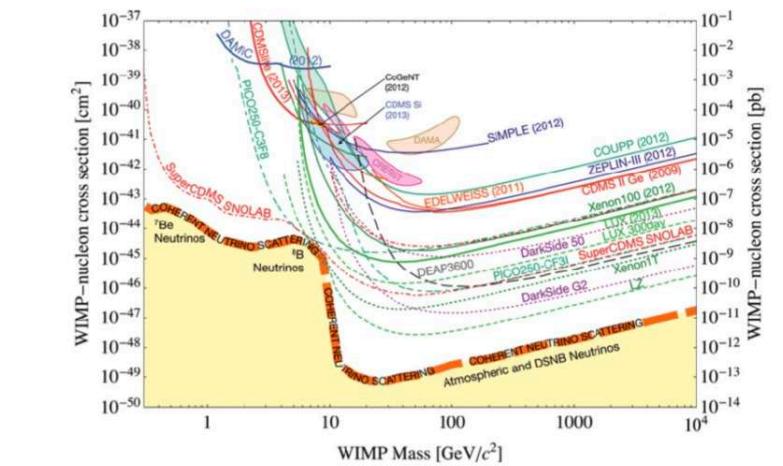
You can interact via the [weasyprint](#) command-line tool.

- Abundance $\Omega_m = 0.308 \pm 0.012$
 - Stable

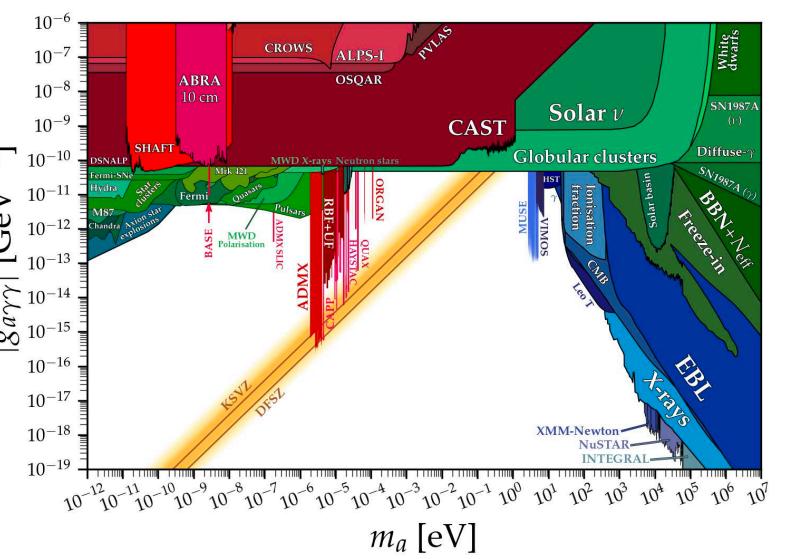
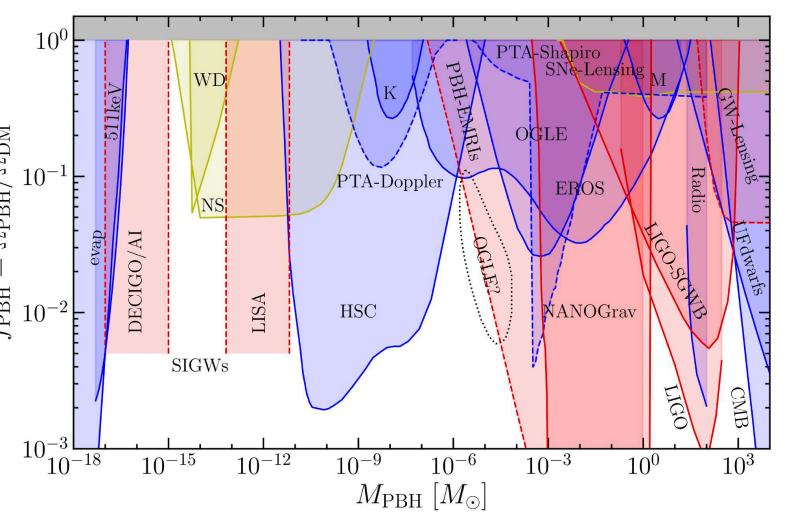
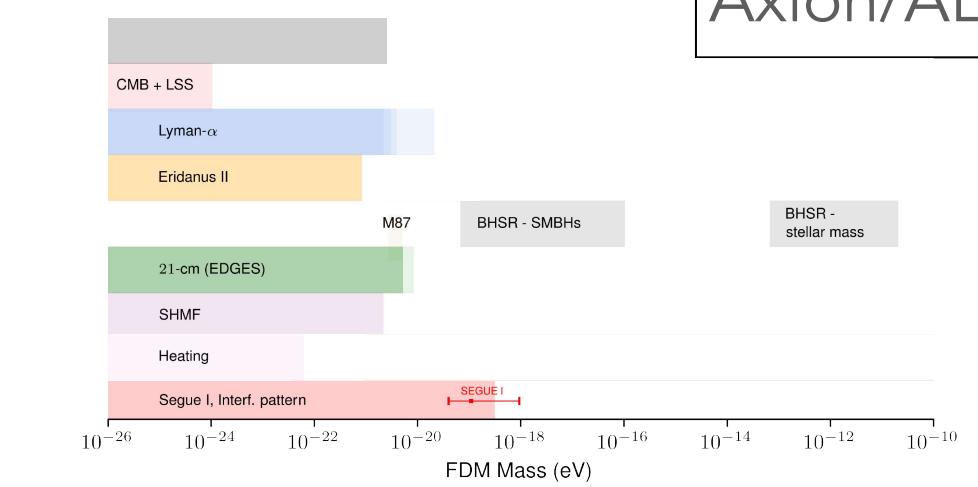
(Planck 2018)



Search for DM



Axion/ALP





Thank you very much!