

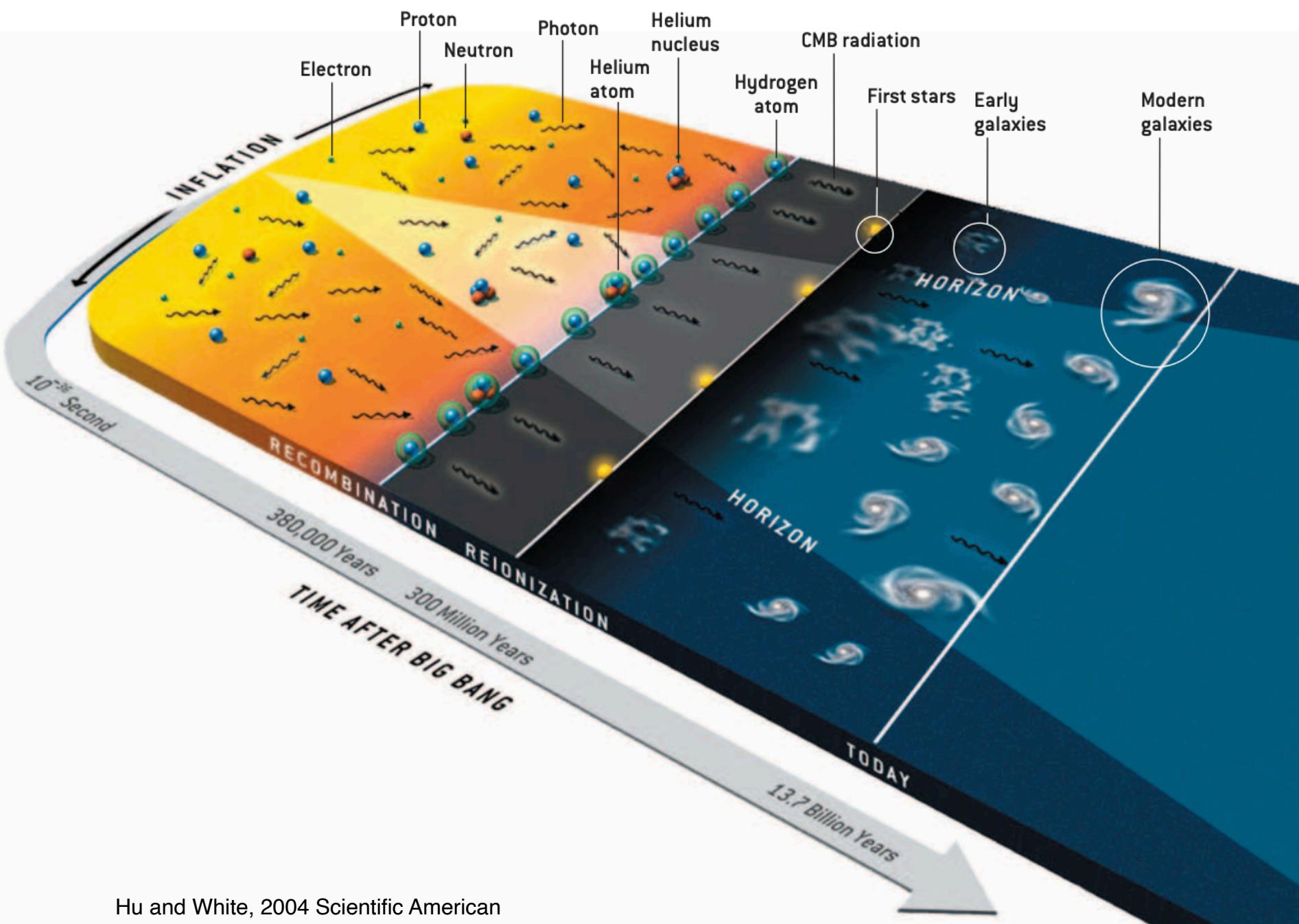
CMB Probes of LSS: Lensing & SZ

Gil Holder



Outline

- the “surface of last scattering” is actually not the final word for lots of photons
 - ★ Thomson scattering
 - ★ lensing
 - ★ extragalactic foregrounds

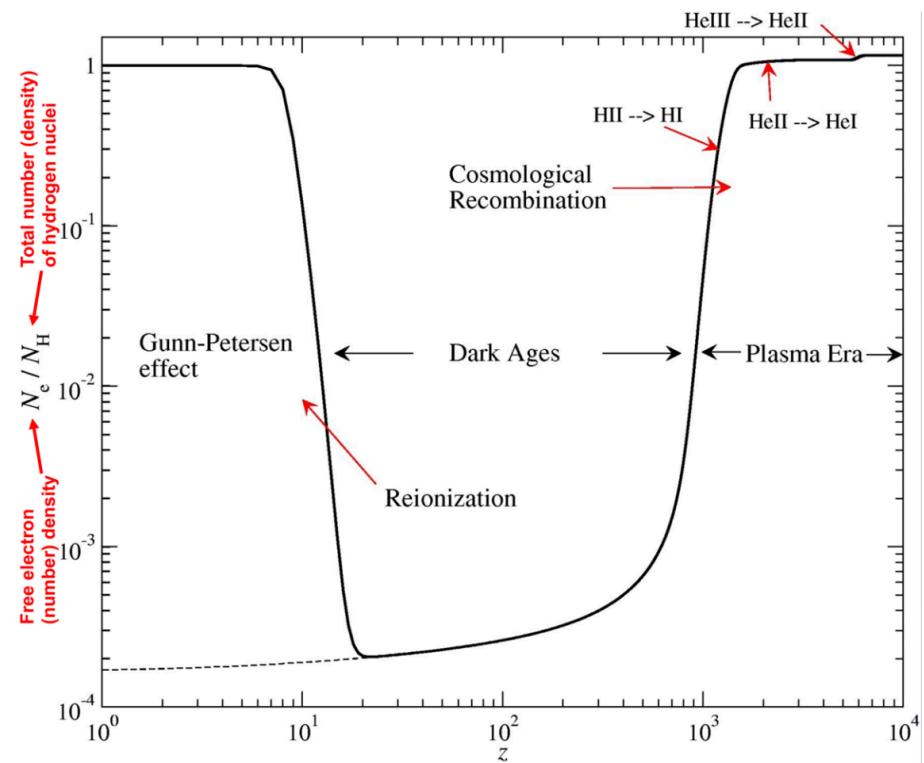


Ionization non-equilibrium

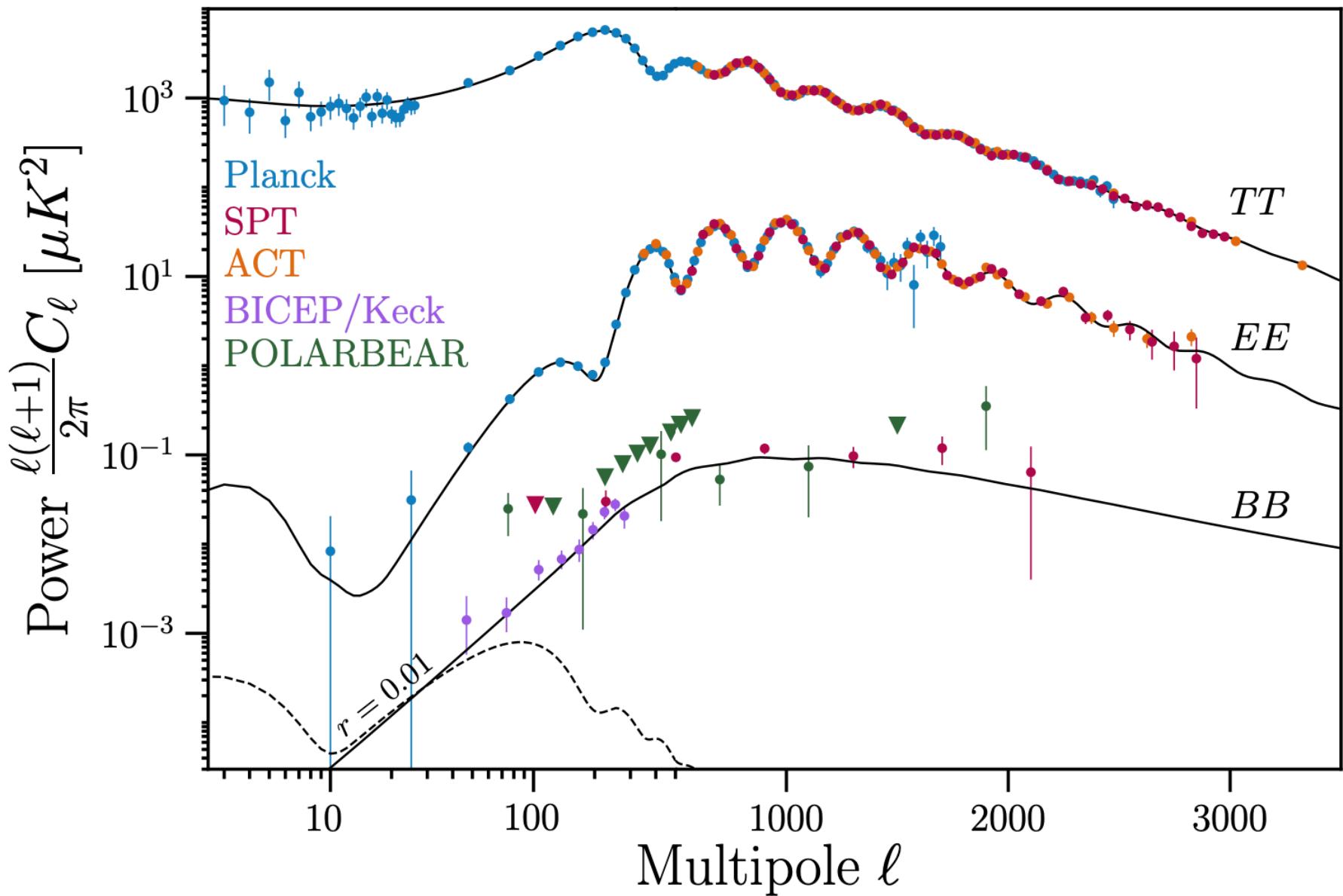
Hubble expansion causes recombinations to “freeze out” as e- and p+ can’t find each other in the dilute universe

small residual ionization keeps gas and CMB thermally coupled for a surprisingly long time

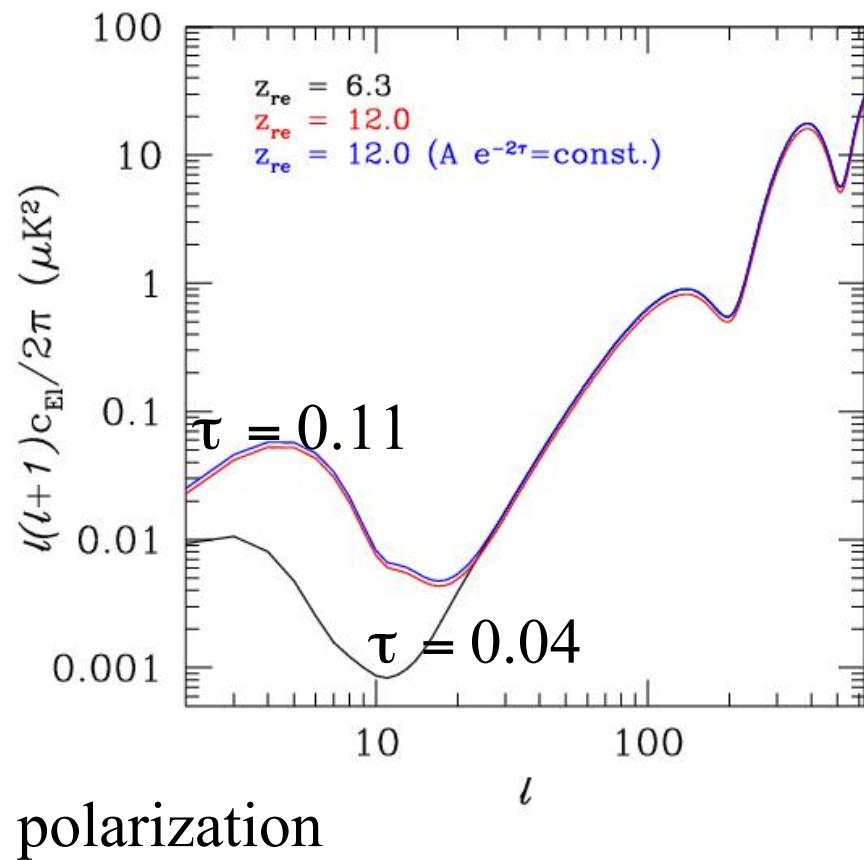
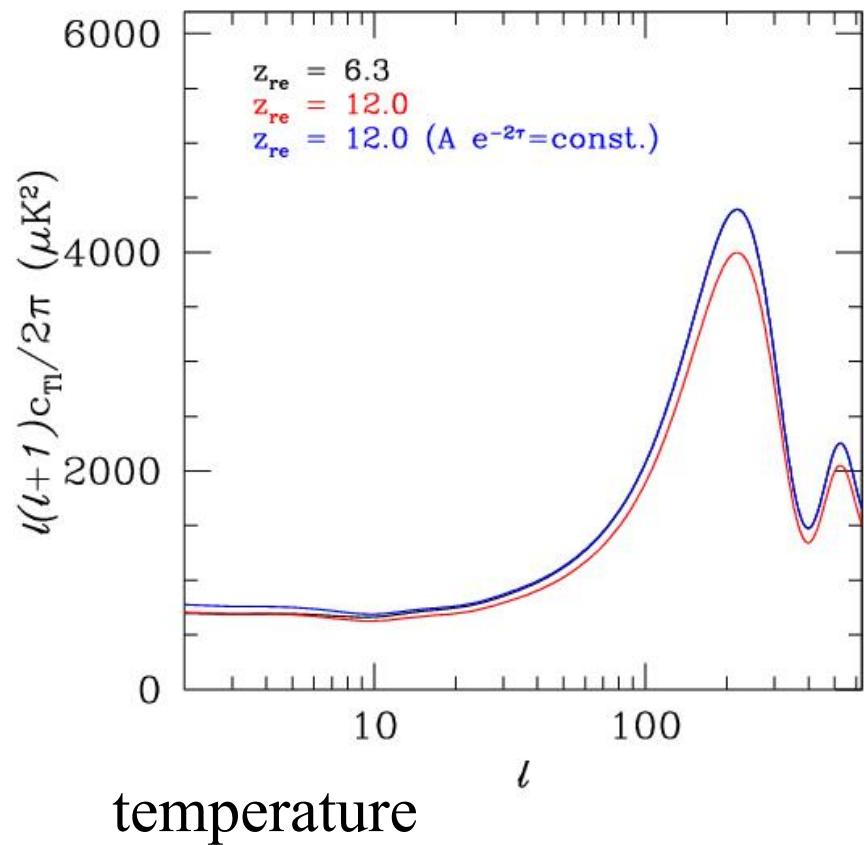
reionization leads to unbinding of electrons from H atoms due to UV background ionizing field



Sunyaev & Chluba 2009

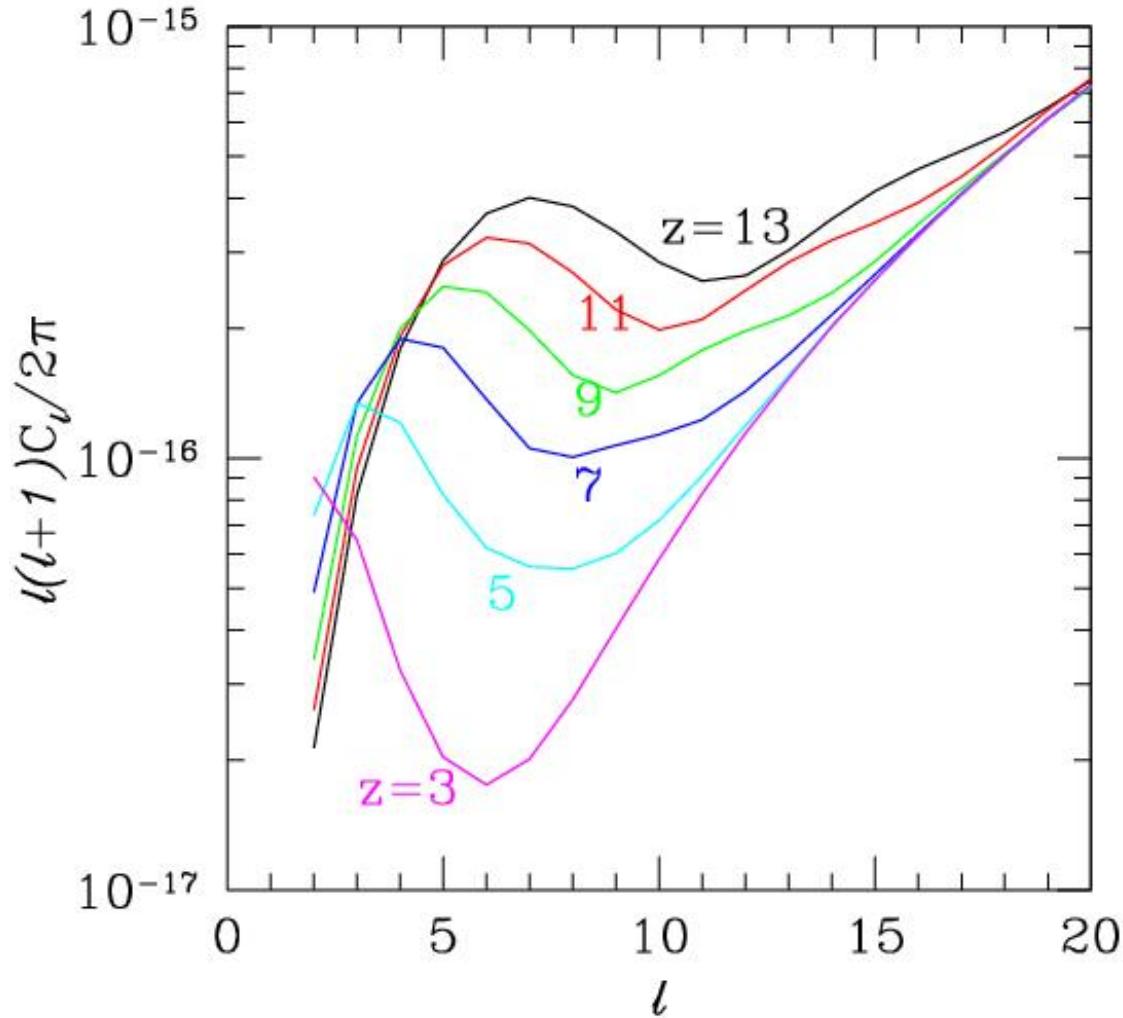


WMAP: +/- 0.015 ; Planck: +/-0.005 ; ????: +/-0.002



Ionization and CMB Polarization

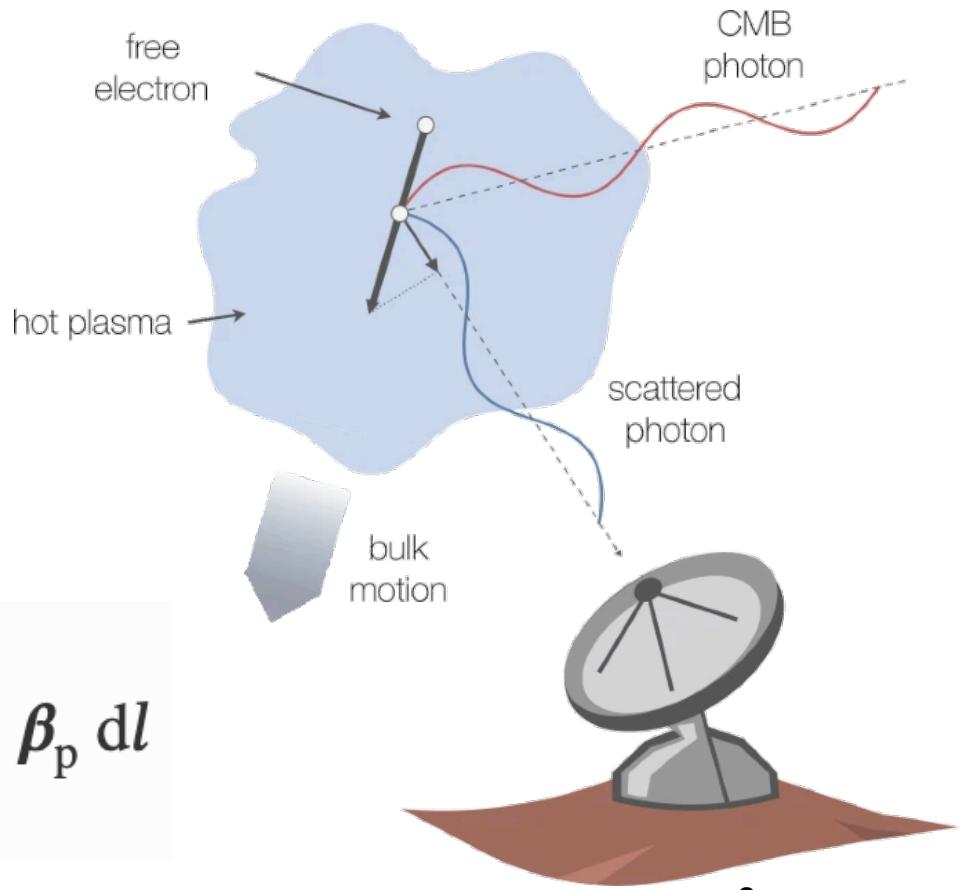
“Pulse” of
ionization
 $dz=1$



Scattering on moving electrons: kSZ

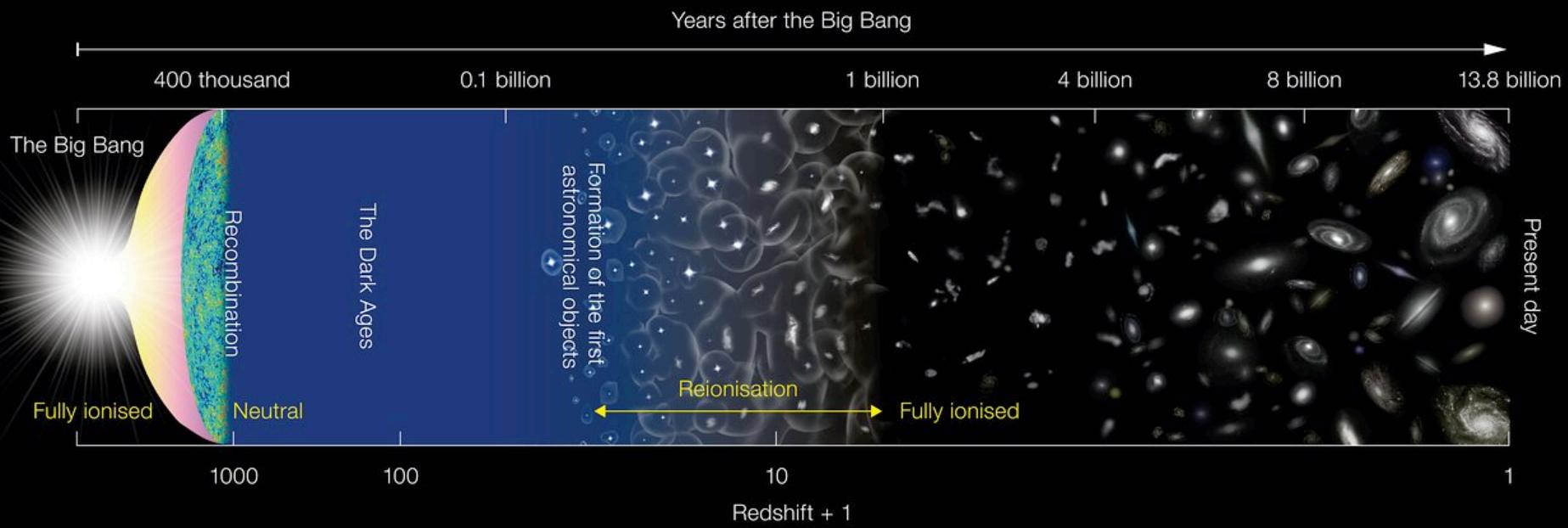
- kinetic Sunyaev-Zeldovich effect: Thomson scattering by bulk flow of electrons

$$\frac{\Delta T_{\text{CMB}}}{T_{\text{CMB}}} \approx - \int \sigma_T n_e \mathbf{n} \cdot \boldsymbol{\beta}_p dl$$



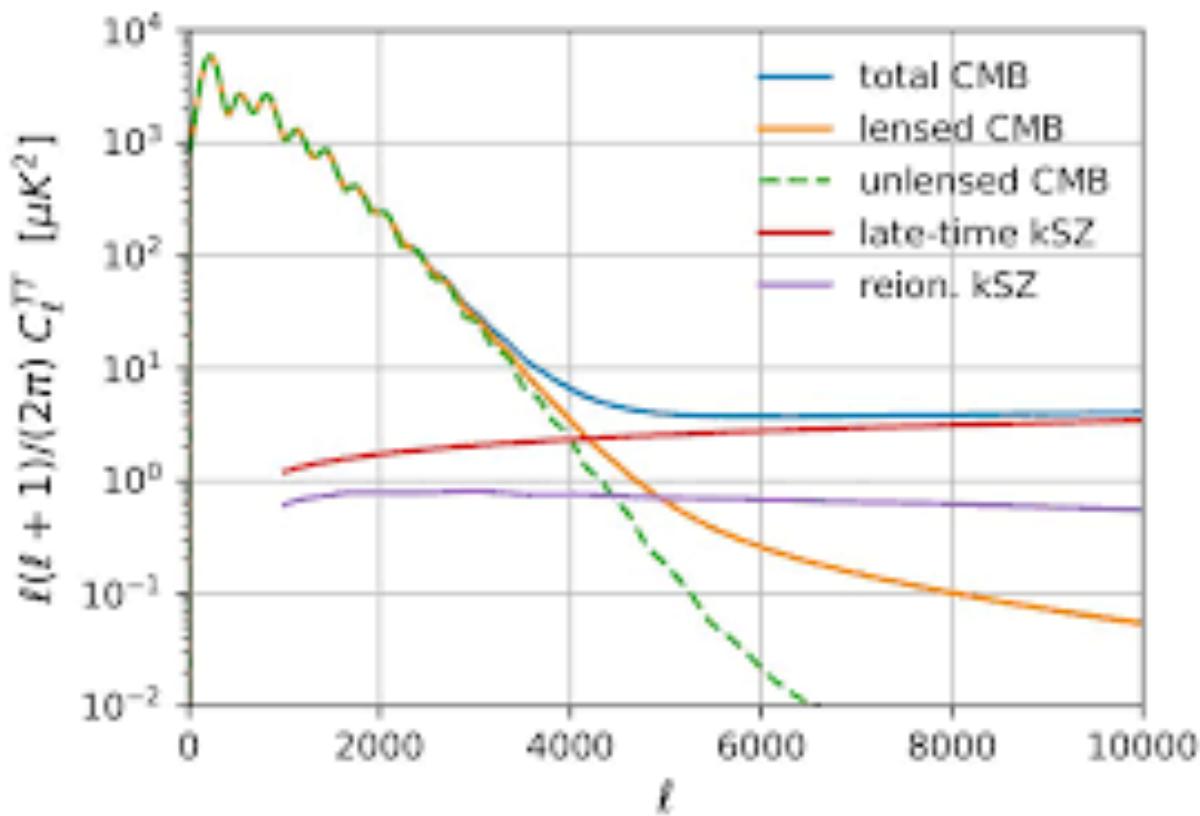
Scattering on moving electrons: kSZ

clumps of moving electrons at reionization, and at late times



ESA

Scattering on moving electrons: kSZ



Madhavacheril

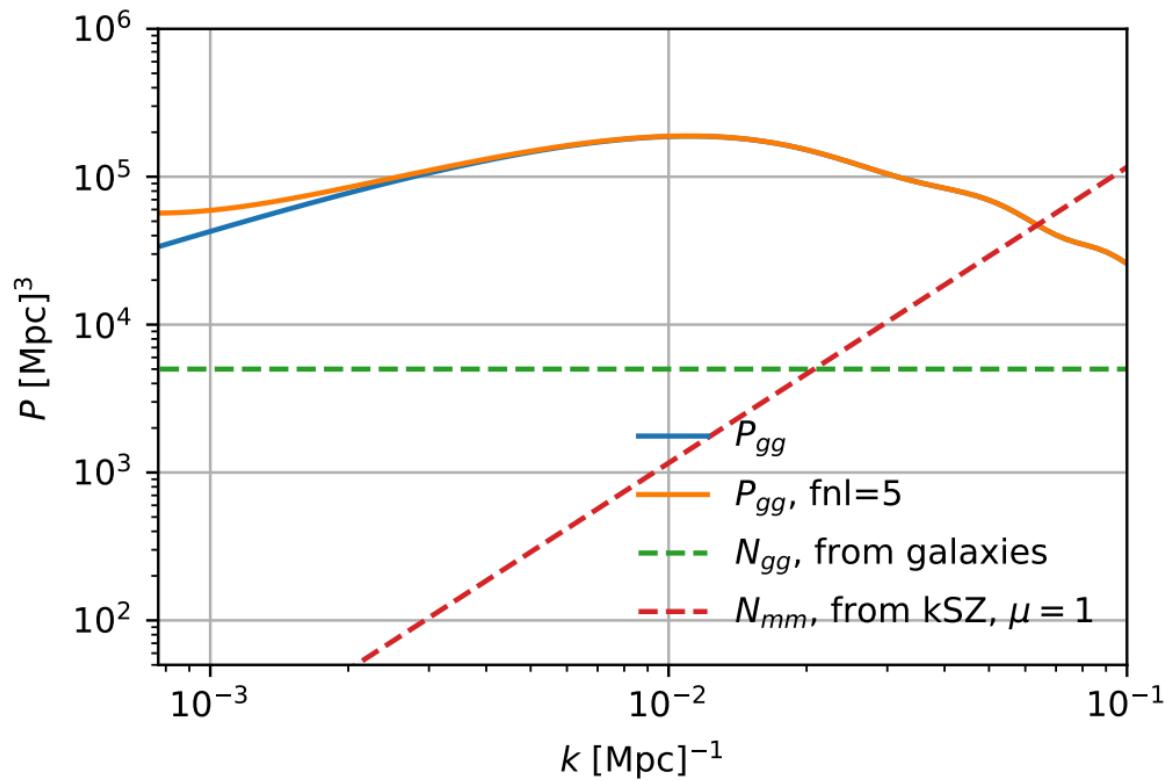
Scattering on moving electrons: kSZ

current status:

detected in cross-correlation with galaxies/clusters

forecast:

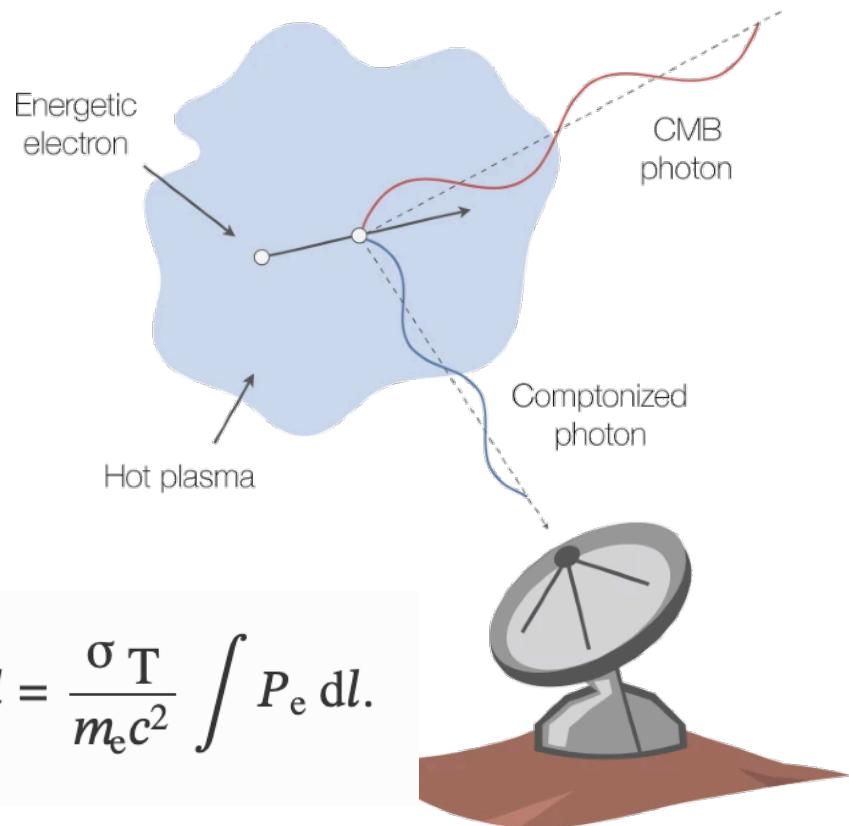
soon to be detected in auto-spectrum, higher order correlations could be very powerful for largest scales



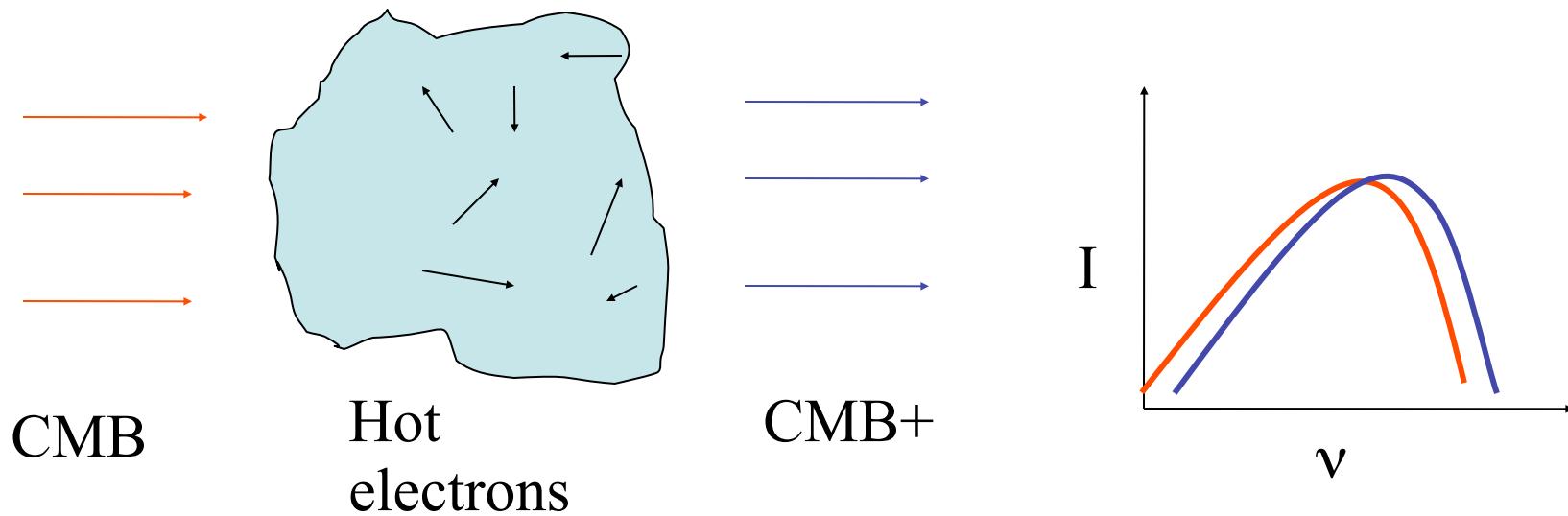
Scattering on moving electrons: tSZ

- thermal Sunyaev-Zeldovich effect: Thomson scattering by thermal motions of electrons

$$y \equiv \int \frac{k_B T_e}{m_e c^2} d\tau_e = \int \frac{k_B T_e}{m_e c^2} n_e \sigma_T dl = \frac{\sigma_T}{m_e c^2} \int P_e dl.$$



Thermal Sunyaev-Zel'dovich Effect



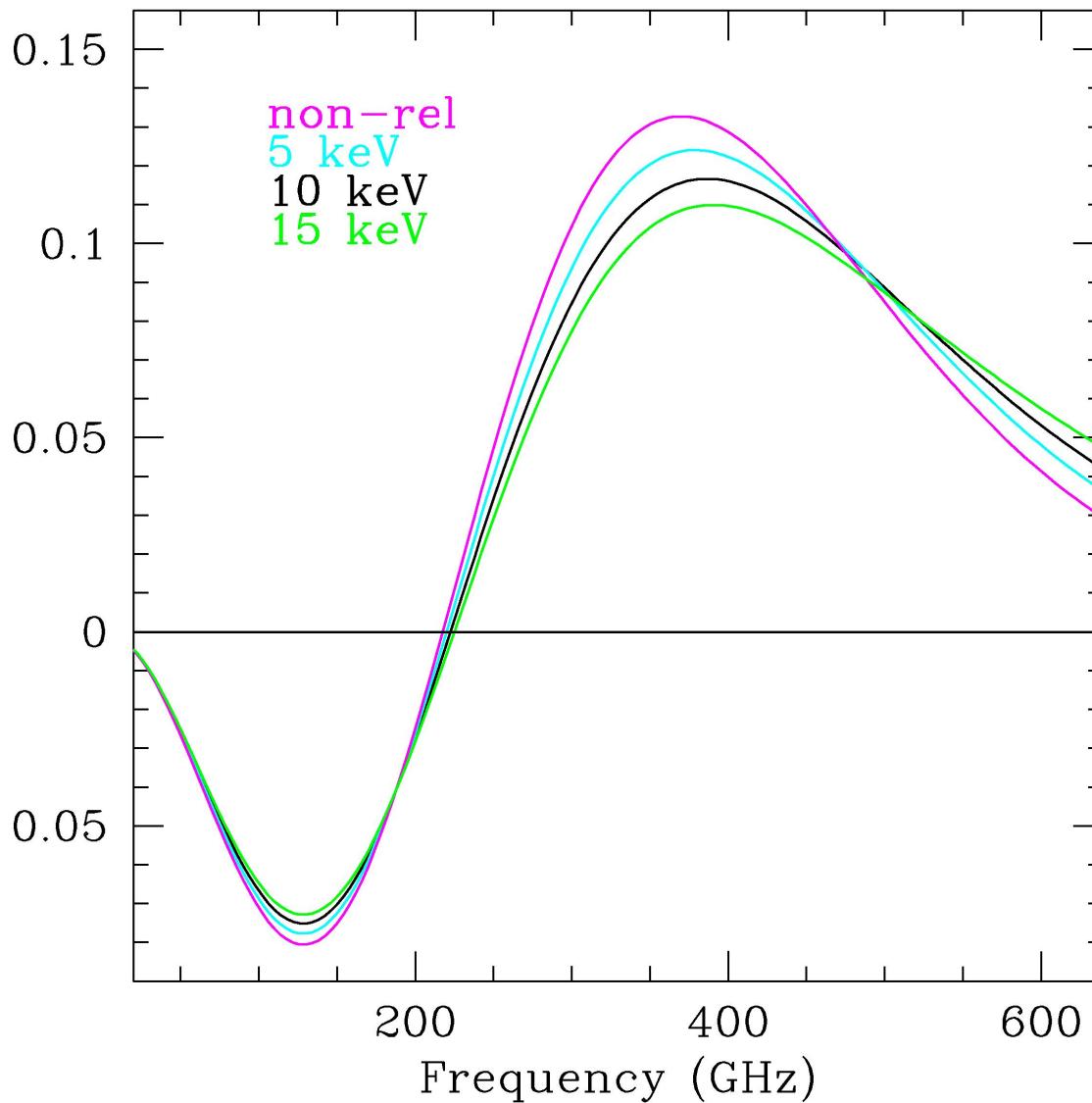
Optical depth: $\tau \sim 0.01$

Fractional energy gain per scatter: $\frac{kT}{m_e c^2} \sim 0.01$

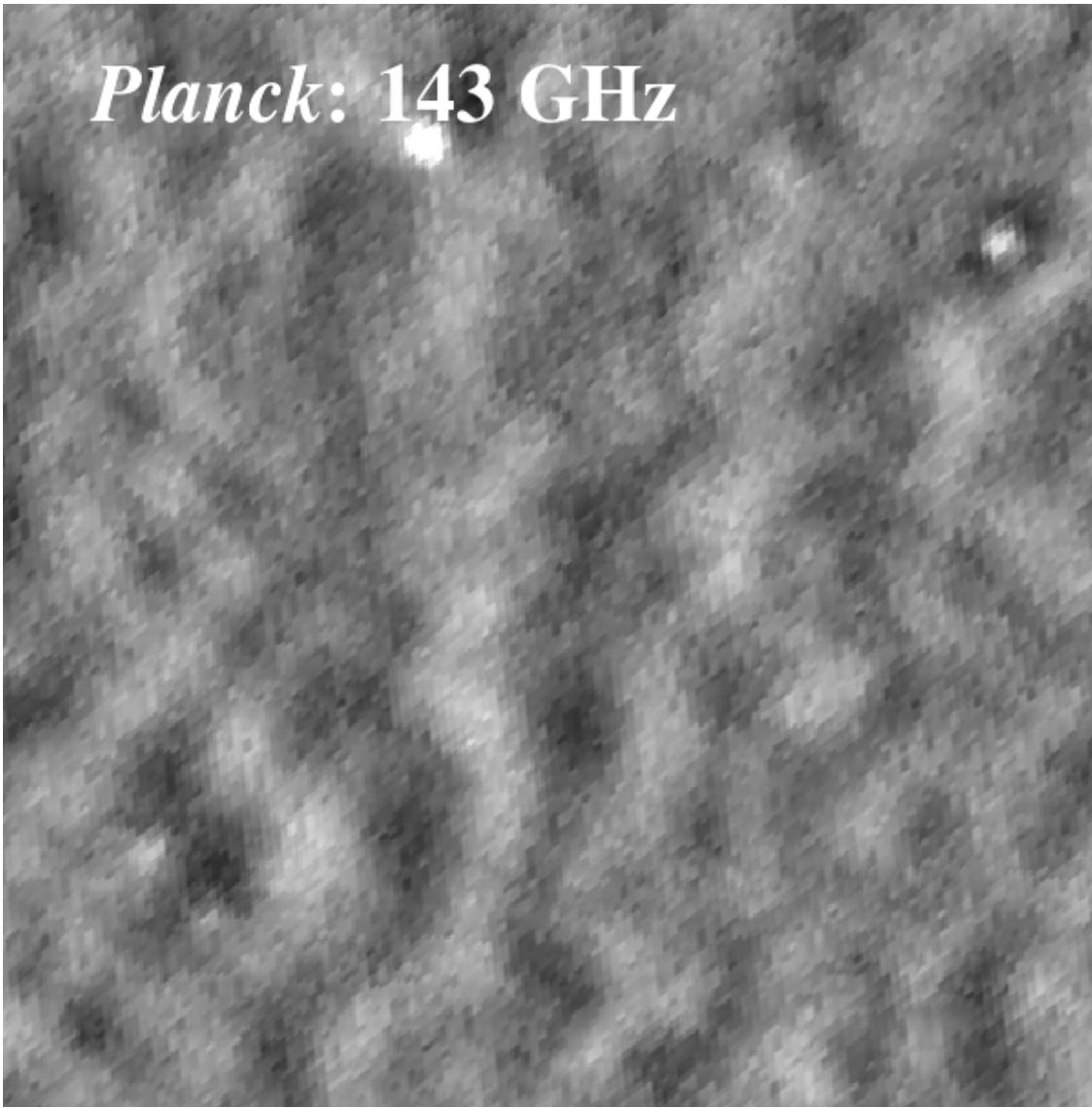
Typical cluster signal: $\sim 500 \mu K$

Thermal SZ Effect (and relativistic corrections)

uK imaging
would allow
1 kev
accuracy in
SZ
temperature



3 degrees

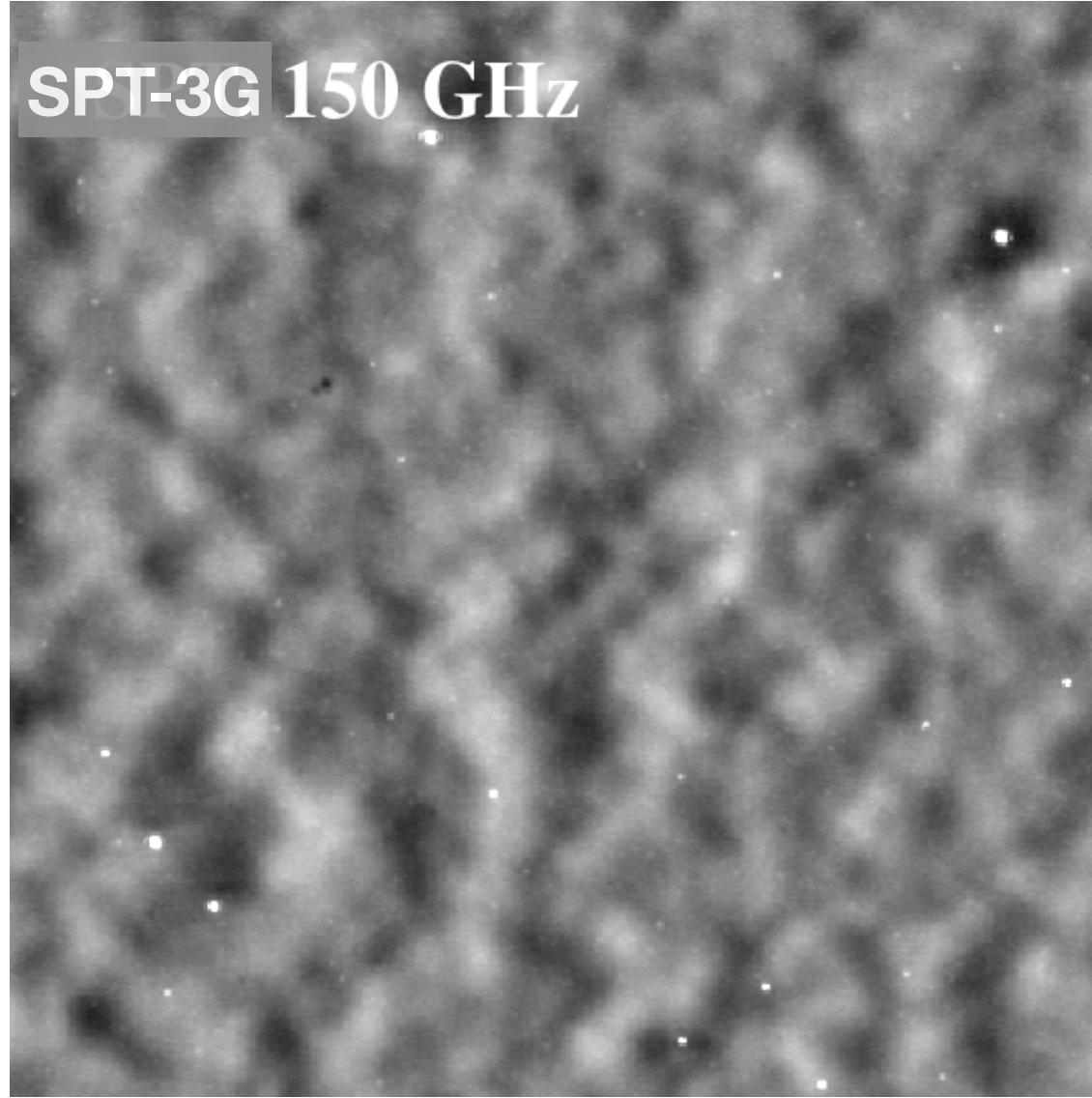


Planck: 143 GHz

3 degrees

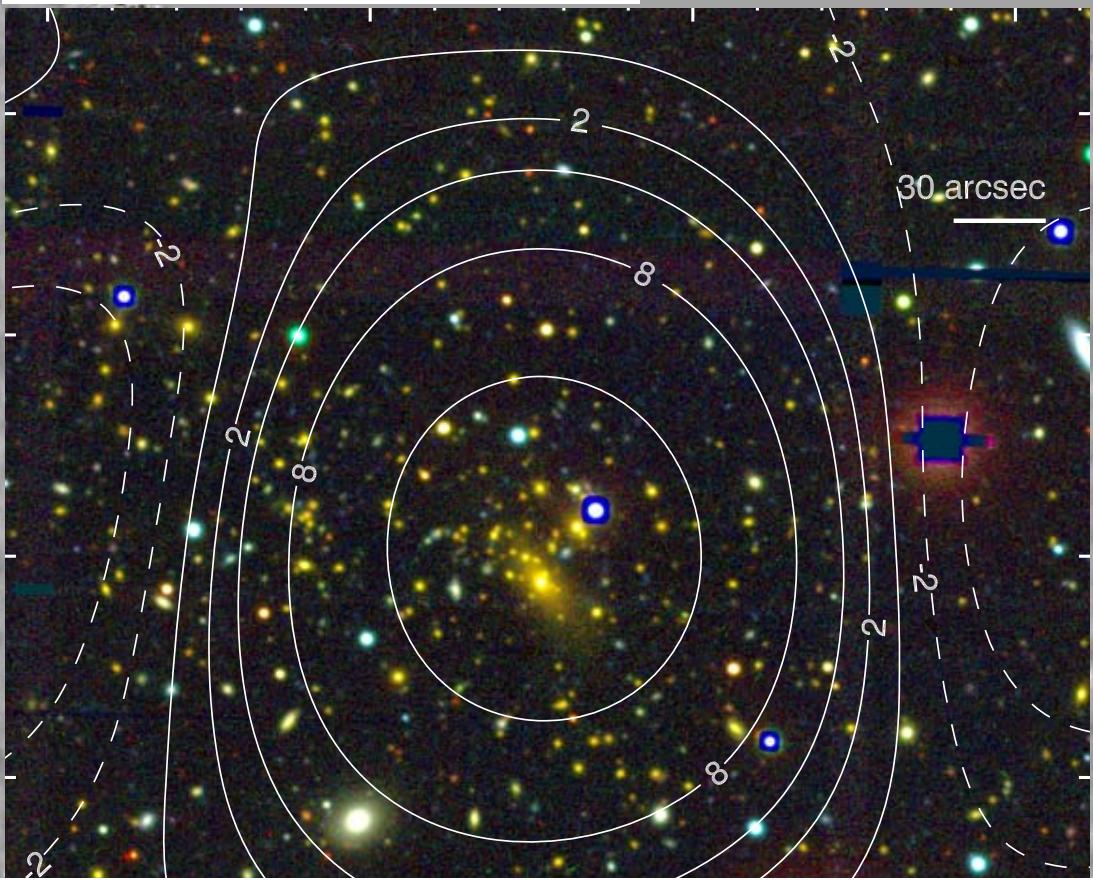
Srini Raghunathan

3 degrees



3 degrees

Srini Raghunathan



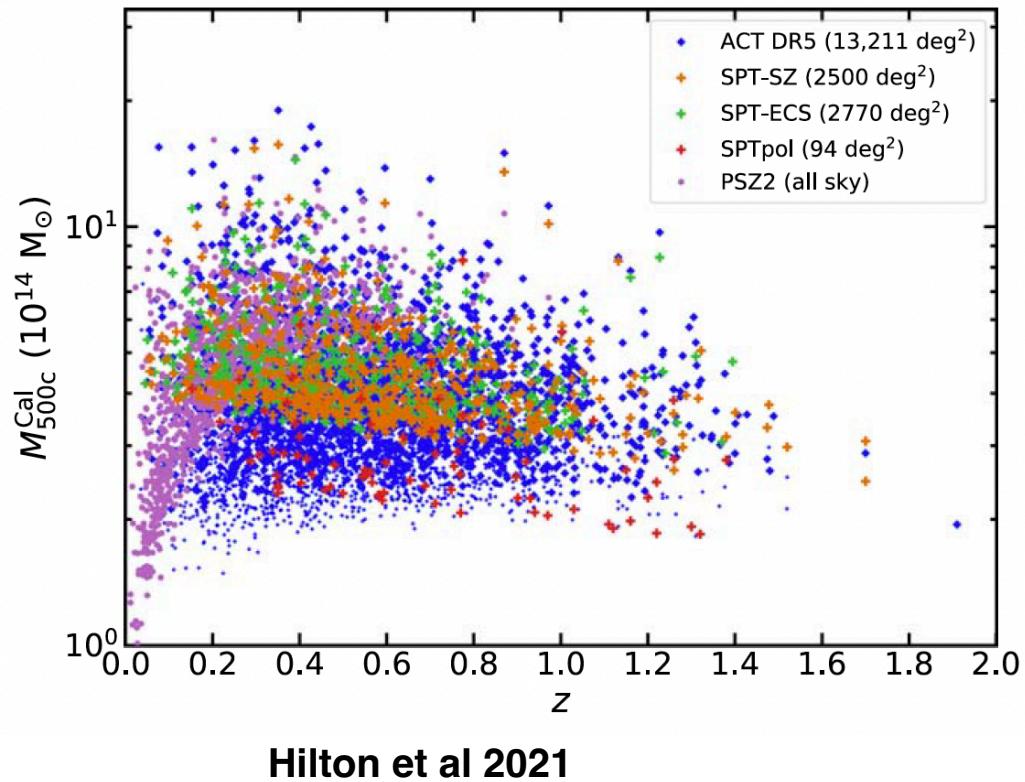
*One of the heaviest objects in the universe
 $>10^{15}$ solar masses*

patch of
isolated cosmic
fog

1 degree

tSZ-selected Galaxy Clusters

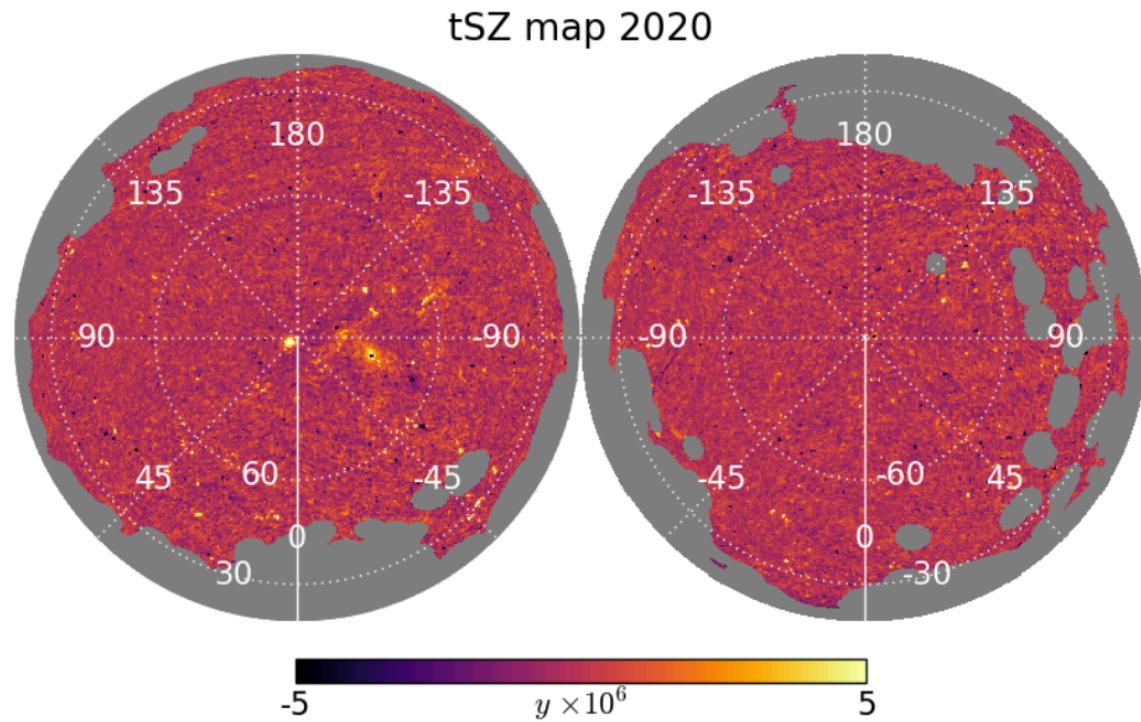
- now many thousands of galaxy clusters have been discovered by their CMB signatures



Hilton et al 2021

Compton γ maps

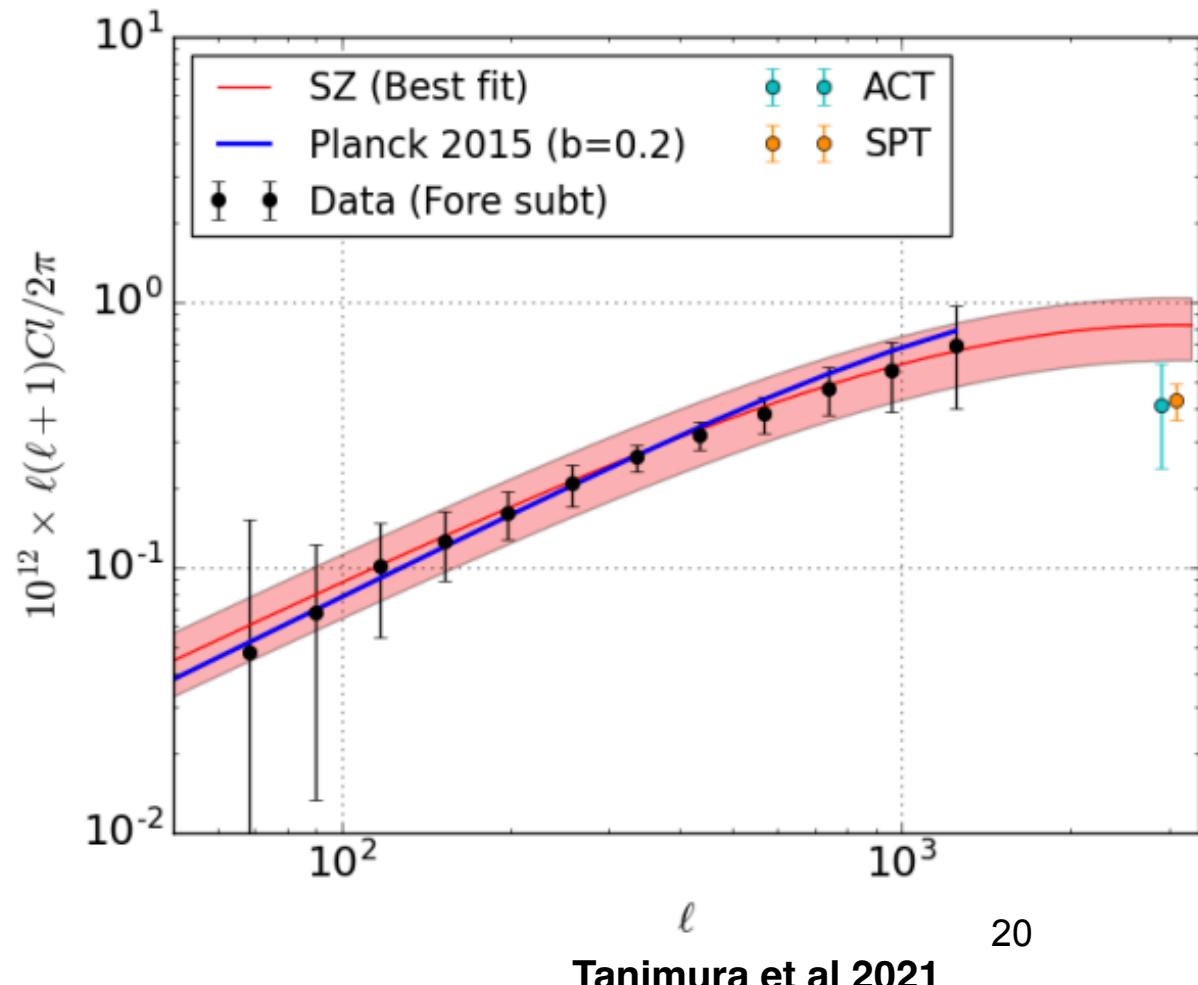
Tanimura et al.



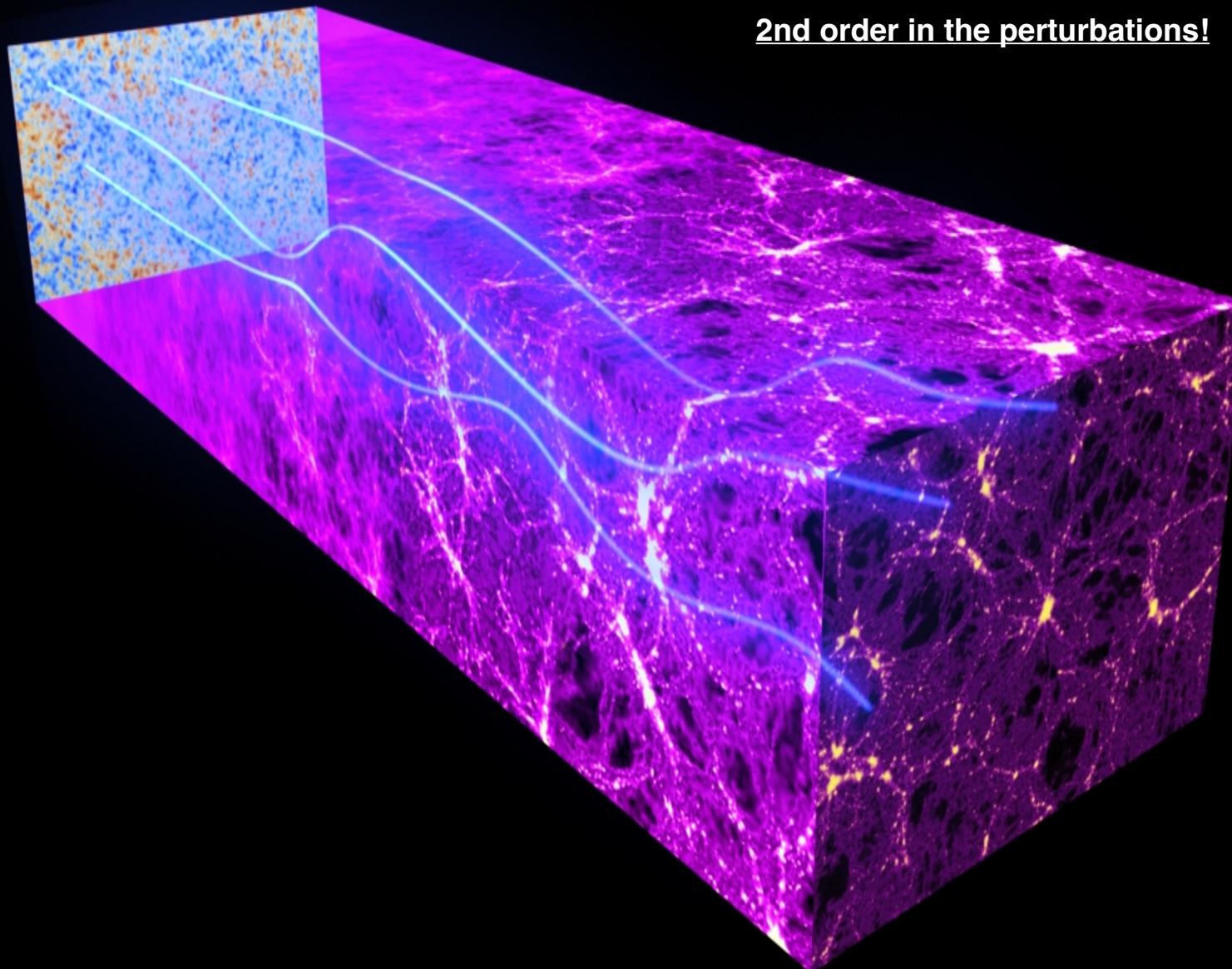
Compton γ power spectrum

Hint that maybe tSZ power is low at high ell

Almost entirely just l-halo term



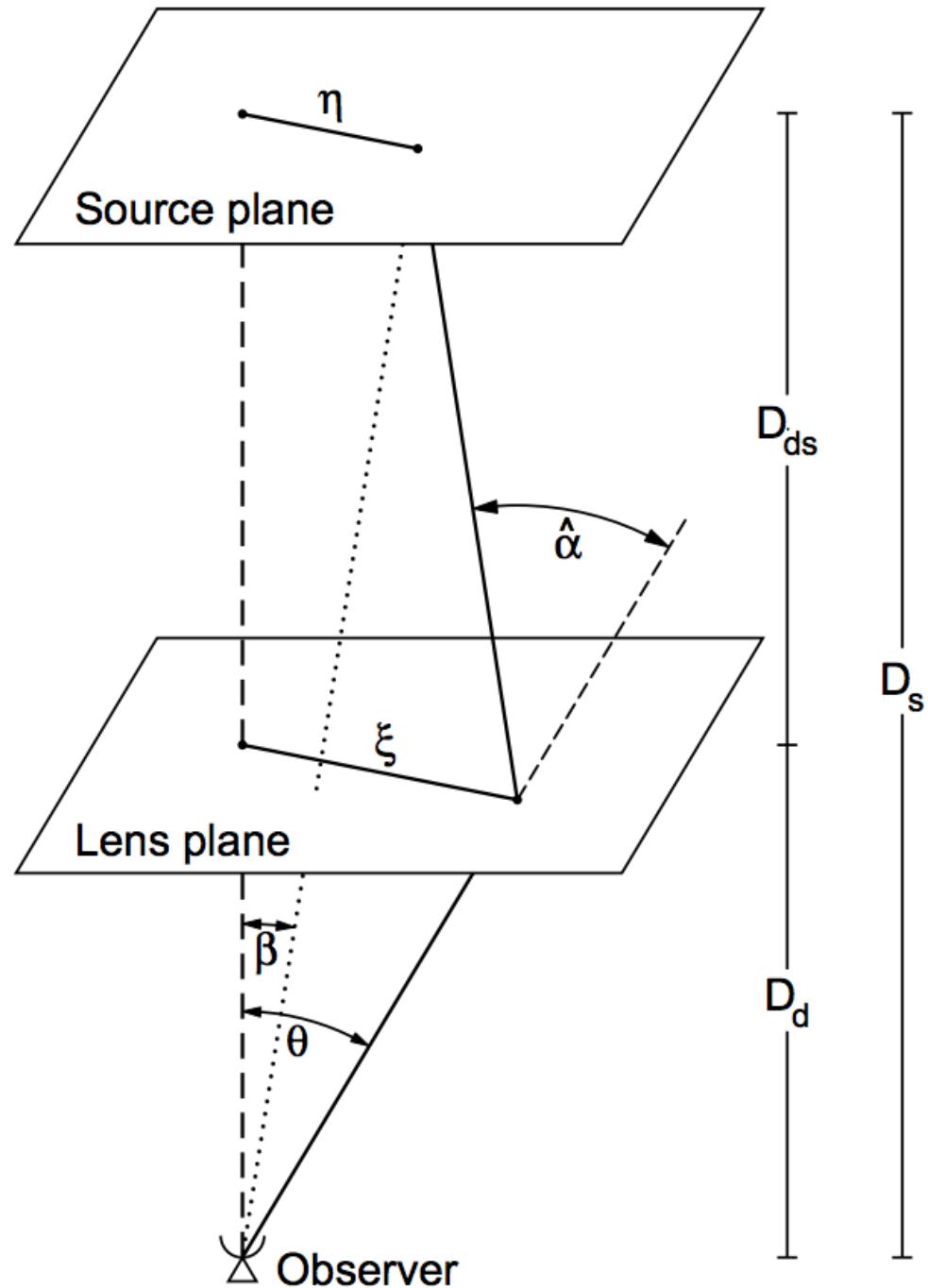
lensing of primordial fluctuations by intervening fluctuations



ESA and the Planck Collaboration

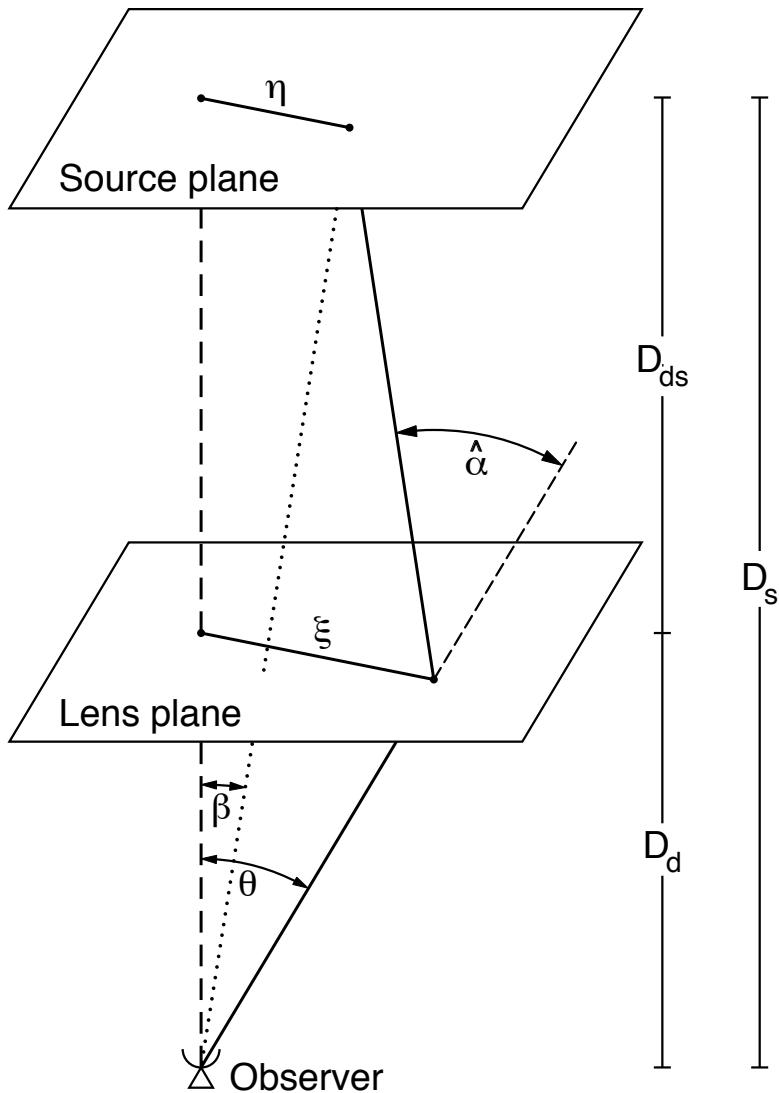
Gravitational lensing

$$\hat{\alpha} = \frac{4GM}{c^2 \xi}$$



Bartelmann & Schneider

<http://arxiv.org/abs/astro-ph/9912508>



source
position

image
position

deflection
angle

$$\vec{\beta} = \vec{\theta} - \frac{D_{ds}}{D_s} \hat{\vec{\alpha}}(D_d \vec{\theta}) \equiv \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

$$\vec{\alpha}(\vec{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\vec{\theta}') \frac{\vec{\theta} - \vec{\theta}'}{|\vec{\theta} - \vec{\theta}'|^2}$$

$$\psi(\vec{\theta}) = \frac{1}{\pi} \int_{\mathbb{R}^2} d^2\theta' \kappa(\vec{\theta}') \ln |\vec{\theta} - \vec{\theta}'|$$

lensing
potential

convergence

$$\mathcal{A}(\vec{\theta}) = \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \left(\delta_{ij} - \frac{\partial^2 \psi(\vec{\theta})}{\partial \theta_i \partial \theta_j} \right) = \begin{pmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix},$$

where we have introduced the components of the shear $\gamma \equiv \gamma_1 + i\gamma_2 = |\gamma|e^{2i\varphi}$

$$\gamma_1 = \frac{1}{2}(\psi_{,11} - \psi_{,22}), \quad \gamma_2 = \psi_{,12},$$

$$\mathcal{A} = (1 - \kappa) \begin{pmatrix} 1 - g_1 & -g_2 \\ -g_2 & 1 + g_1 \end{pmatrix}$$

distortion has overall magnification
image gets bigger (or smaller),
not brighter (dimmer)

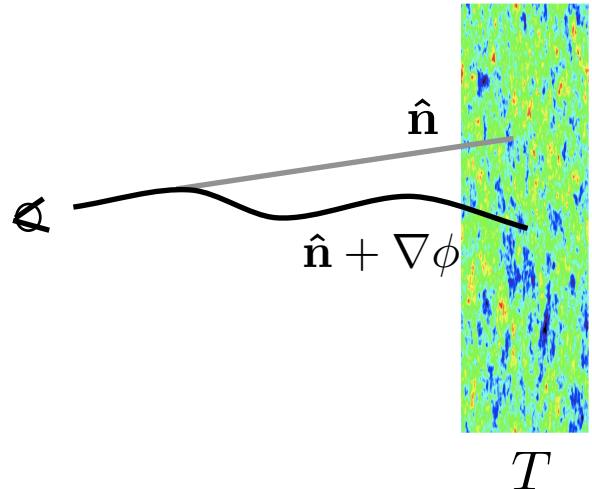
$$g(\vec{\theta}) \equiv \frac{\gamma(\vec{\theta})}{1 - \kappa(\vec{\theta})}$$

reduced shear

CMB Lensing

Photons get shifted

$$T^L(\hat{\mathbf{n}}) = T^U(\hat{\mathbf{n}} + \nabla\phi(\hat{\mathbf{n}}))$$



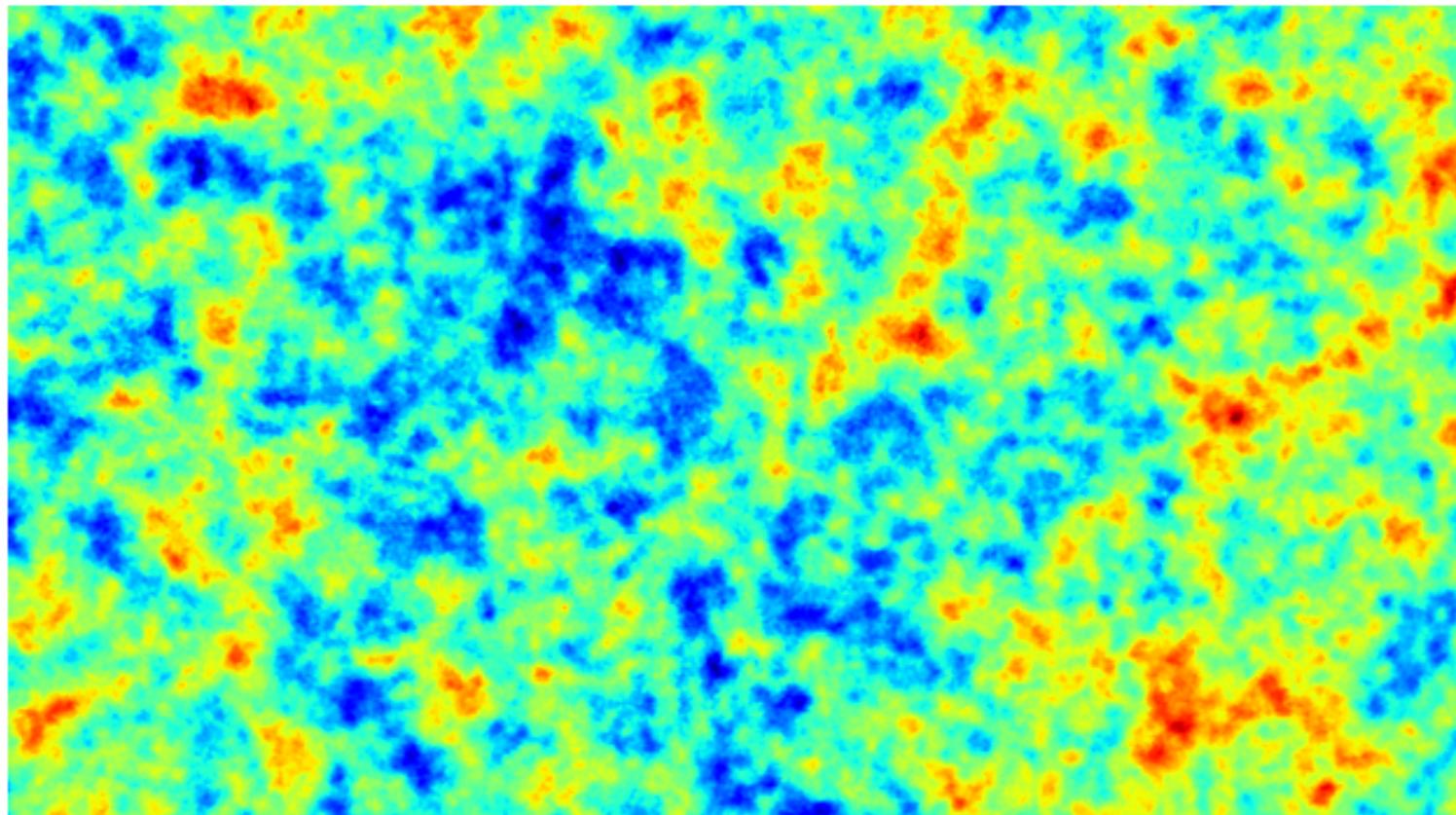
In WL limit, add many deflections along line of sight

$$\nabla\phi(\hat{\mathbf{n}}) = -2 \int_0^{\chi_*} d\chi \frac{\chi_* - \chi}{\chi_* \chi} \nabla_{\perp} \Phi(\chi \hat{\mathbf{n}}, \chi)$$

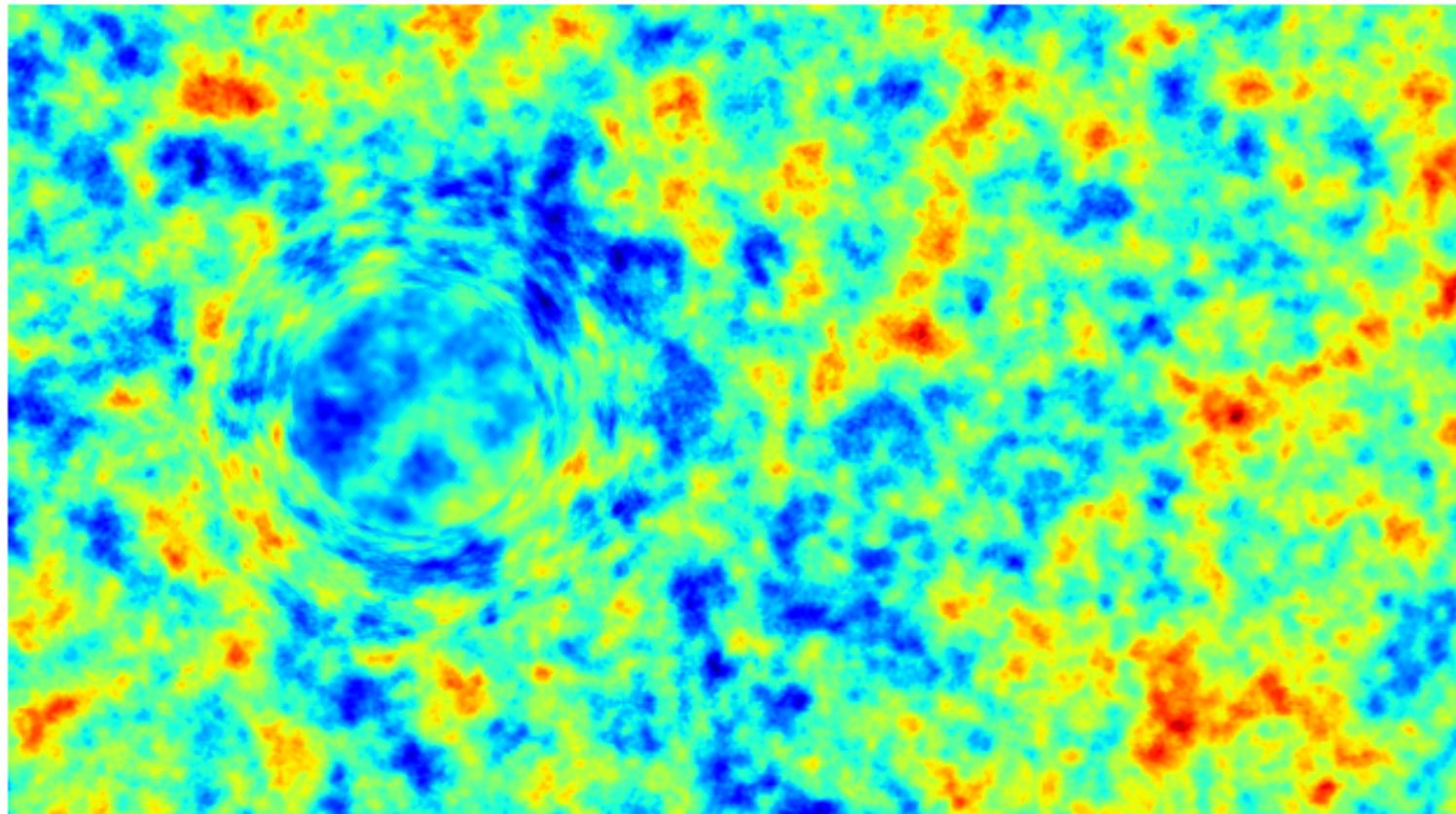
Broad kernel, peaks at $z \sim 2$

- CMB is a unique source for lensing
 - Gaussian, with well-understood power spectrum (contains all info)
 - At redshift which is (a) unique, (b) known, and (c) highest

patch of sky (the North pole) as seen by Planck (17x10 degrees)

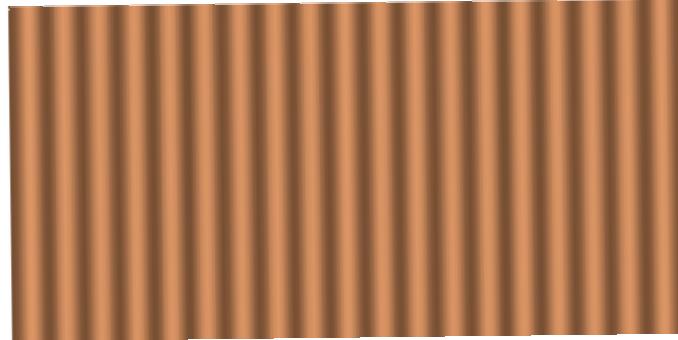


SIMULATED lensing effect (20x larger than typical)

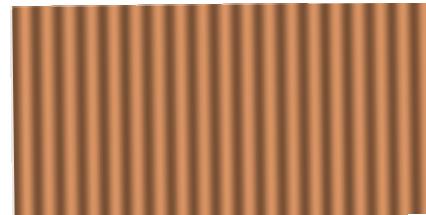


Lensing simplified

- gravitational potentials distort images by stretching, squeezing, shearing

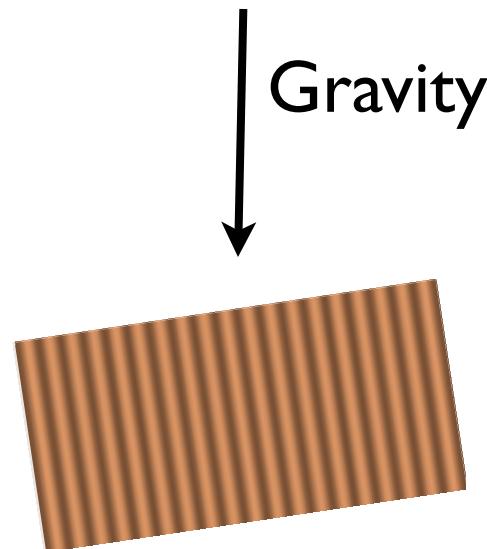
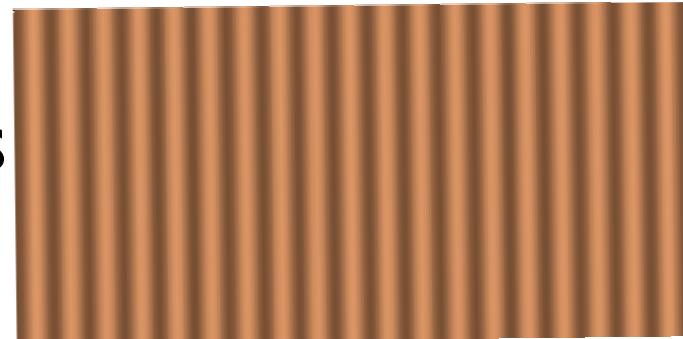


Gravity
↓

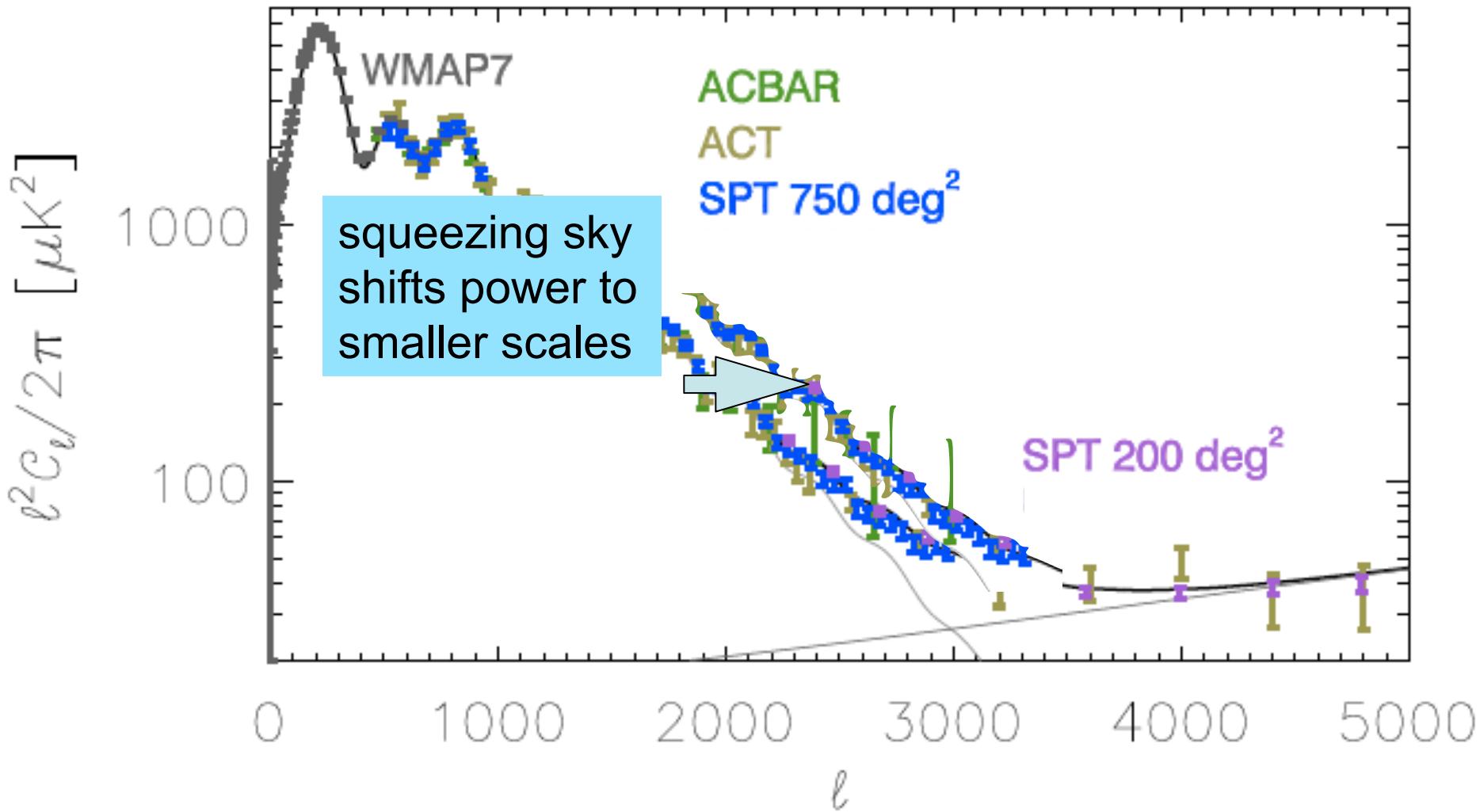
A black downward-pointing arrow is positioned to the right of the top image, pointing towards the bottom image. To the right of the arrow, the word "Gravity" is written in a black sans-serif font.

Lensing simplified

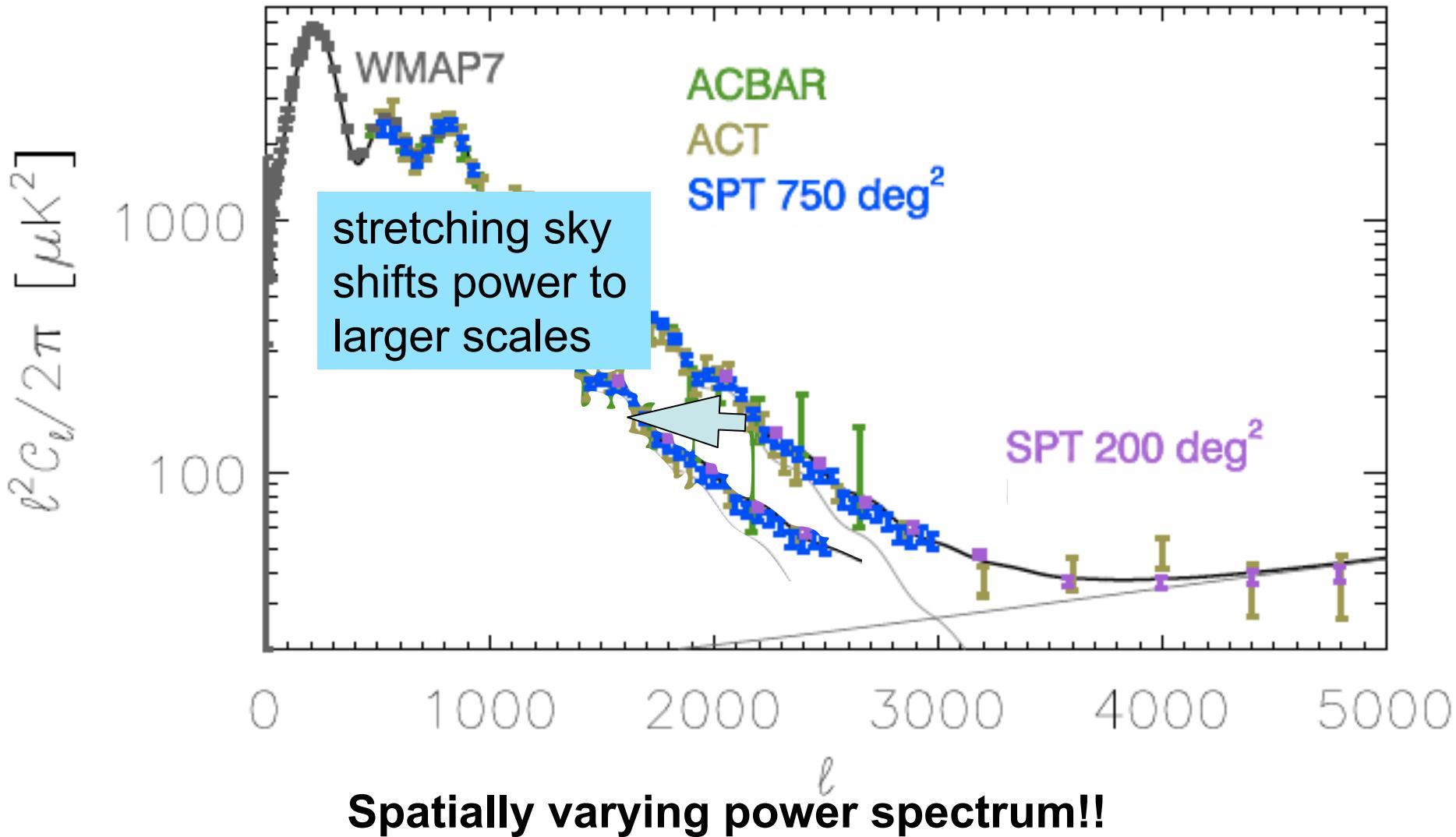
- where gravity stretches, gradients become smaller
- where gravity compresses, gradients are larger
- shear changes ***direction***



CMB Power Spectrum

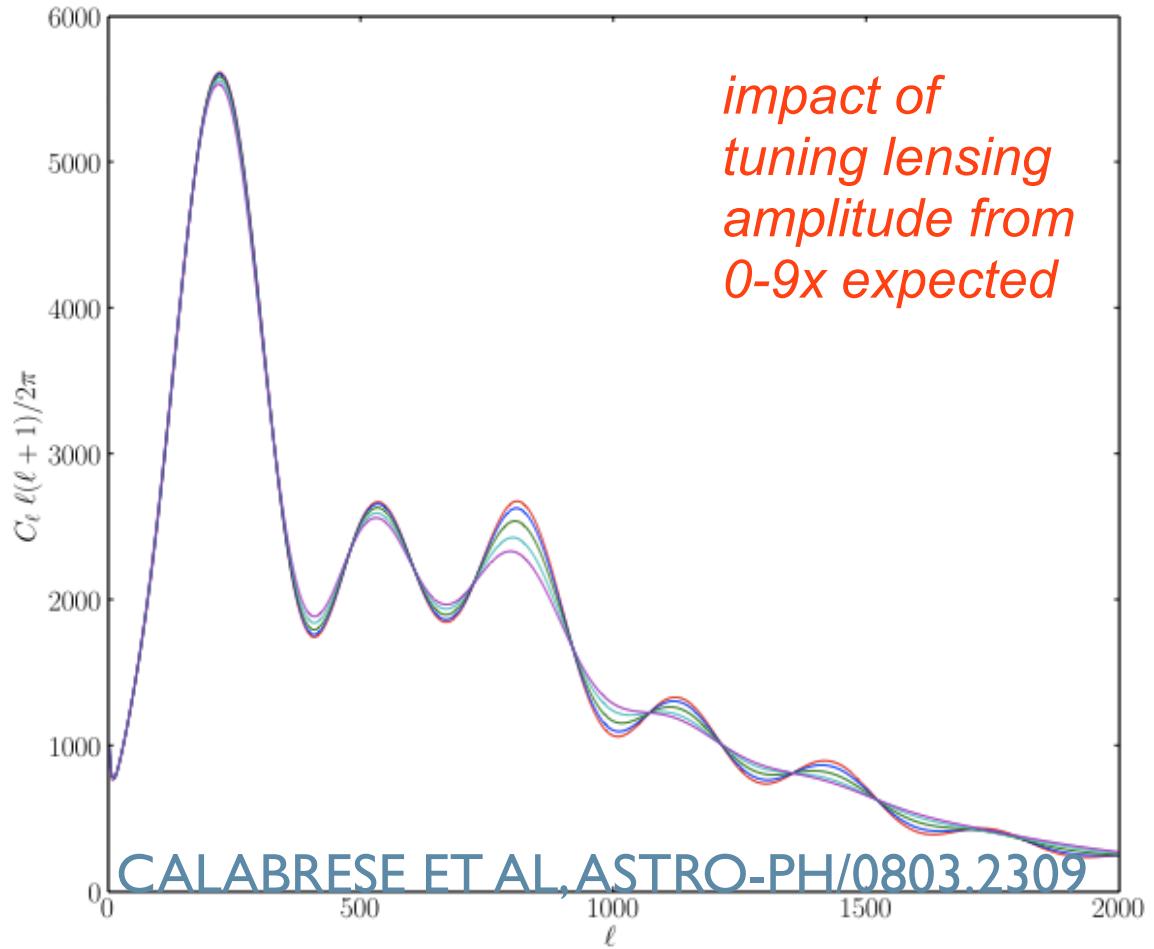


CMB Power Spectrum



Effect on CMB Power Spectrum

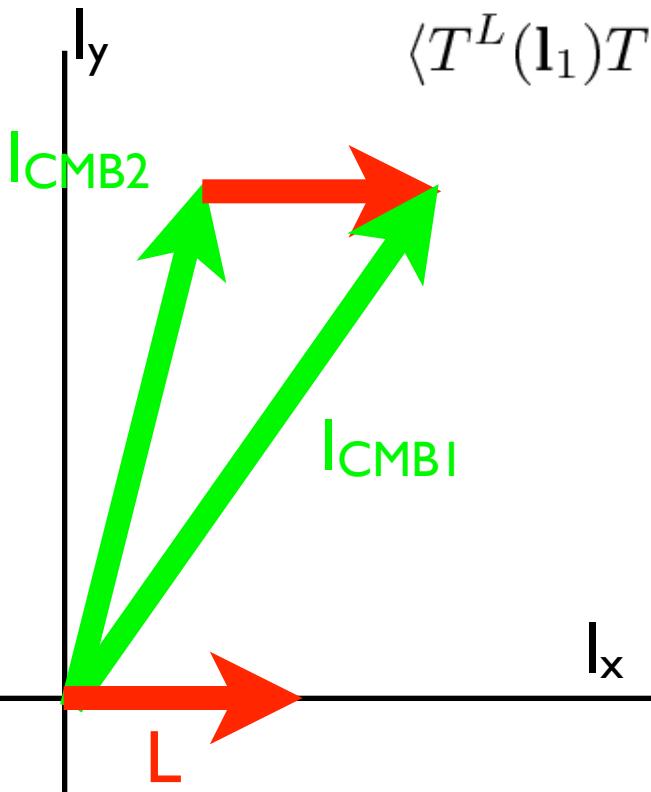
- mixing of power leads to smoothing of acoustic peaks
- small effect but data is really good



Mode Coupling from Lensing

$$\begin{aligned} T^L(\hat{\mathbf{n}}) &= T^U(\hat{\mathbf{n}} + \nabla\phi(\hat{\mathbf{n}})) \\ &= T^U(\hat{\mathbf{n}}) + \nabla T^U(\hat{\mathbf{n}}) \cdot \nabla\phi(\hat{\mathbf{n}}) + O(\phi^2), \end{aligned}$$

- Non-gaussian mode coupling for $\mathbf{l}_1 \neq -\mathbf{l}_2$:



$$\langle T^L(\mathbf{l}_1)T^L(\mathbf{l}_2) \rangle = \mathbf{L} \cdot (\mathbf{l}_1 C_{l_1}^T + \mathbf{l}_2 C_{l_2}^T) \phi(\mathbf{L}) + O(\phi^2)$$
$$\mathbf{L} = \mathbf{l}_1 + \mathbf{l}_2$$

- We extract ϕ by taking a suitable average over CMB multipoles separated by a distance L
- We use the standard Hu quadratic estimator.

E-modes and B-modes

$$Q(\mathbf{l}) = [E(\mathbf{l}) \cos(2\phi_{\mathbf{l}}) - B(\mathbf{l}) \sin(2\phi_{\mathbf{l}})]$$
$$U(\mathbf{l}) = [E(\mathbf{l}) \sin(2\phi_{\mathbf{l}}) + B(\mathbf{l}) \cos(2\phi_{\mathbf{l}})].$$

- E/B is a different way to express polarization field
- easy to understand in flat-sky limit (i.e. Fourier modes)

E-modes/B-modes

- E-modes vary spatially parallel or perpendicular to polarization direction
- B-modes vary spatially at 45 degrees
- CMB
 - scalar perturbations only generate **only** E
- ***Lensing of CMB is much more obvious in polarization!***

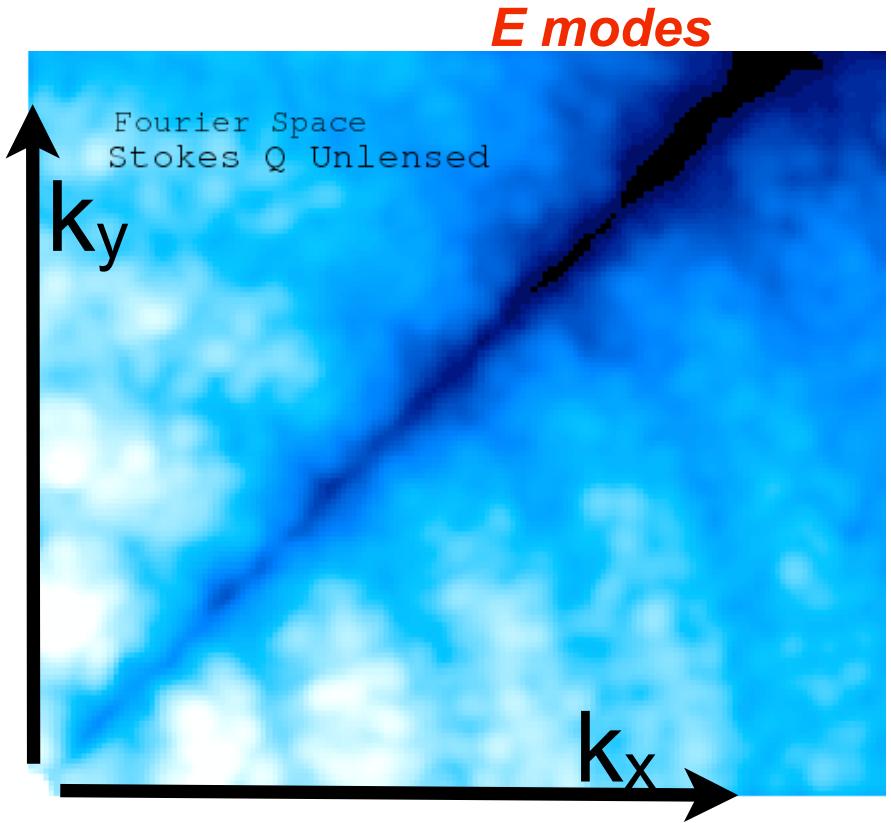


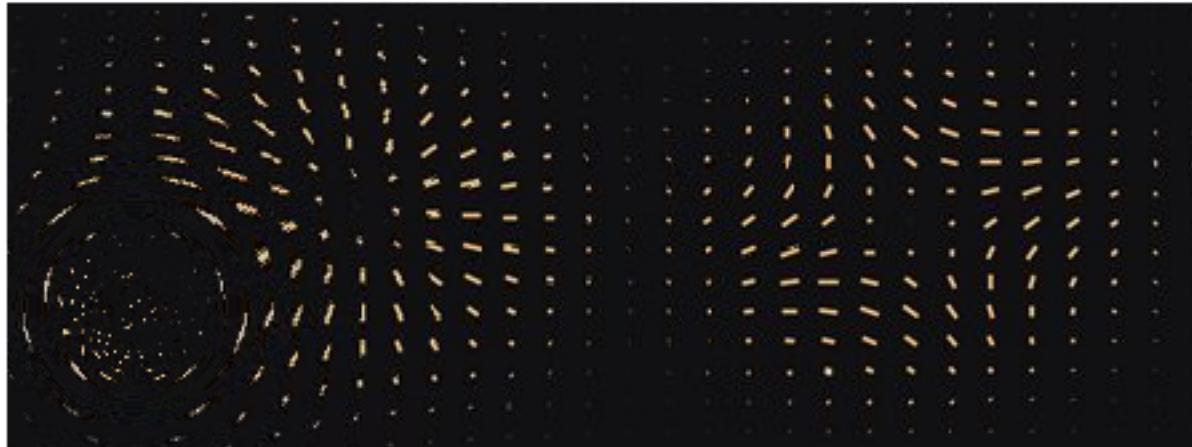
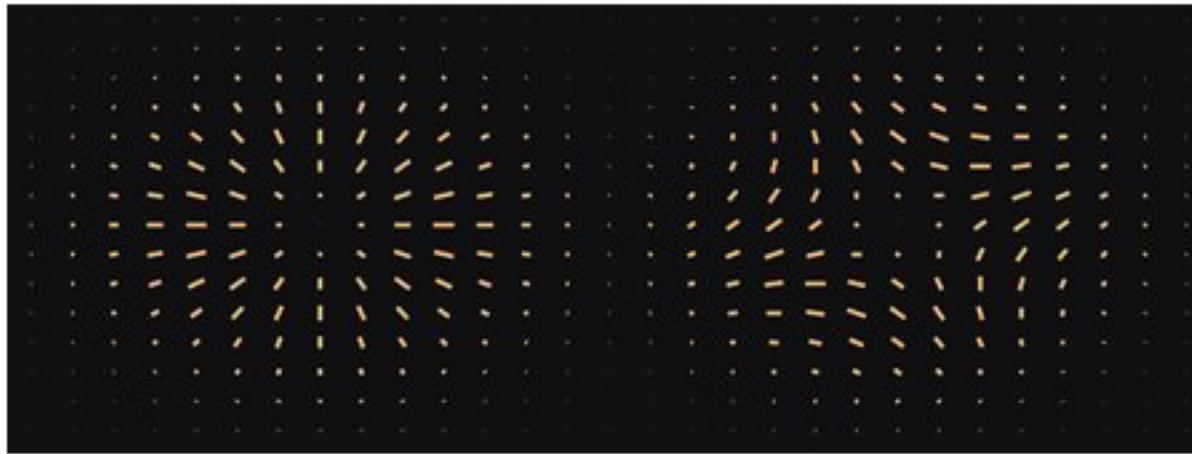
Image of positive k_x /positive k_y Fourier transform of a 10x10 deg chunk of Stokes Q CMB map [simulated; nothing clever done to it]

B Modes from E Modes

Before: pure E mode (left) and pure B mode (right)

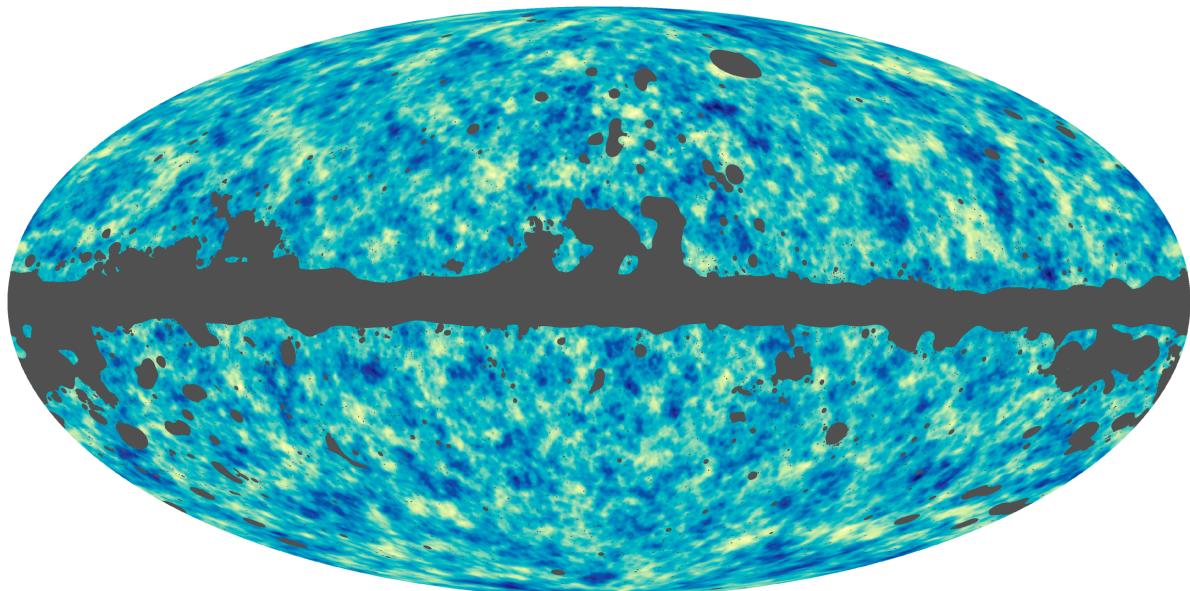
From B-pol.org

After: large point mass lenses image

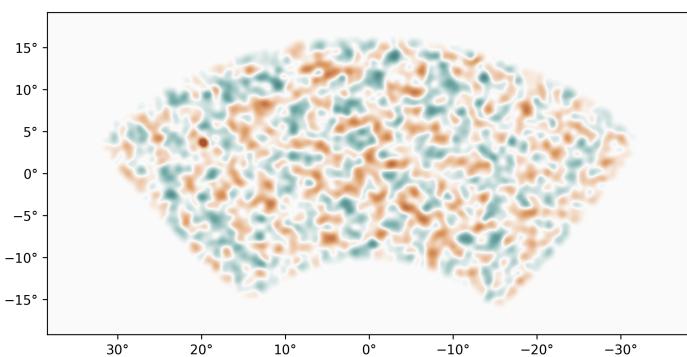


Lensing done with “Lens an astrophysicist”

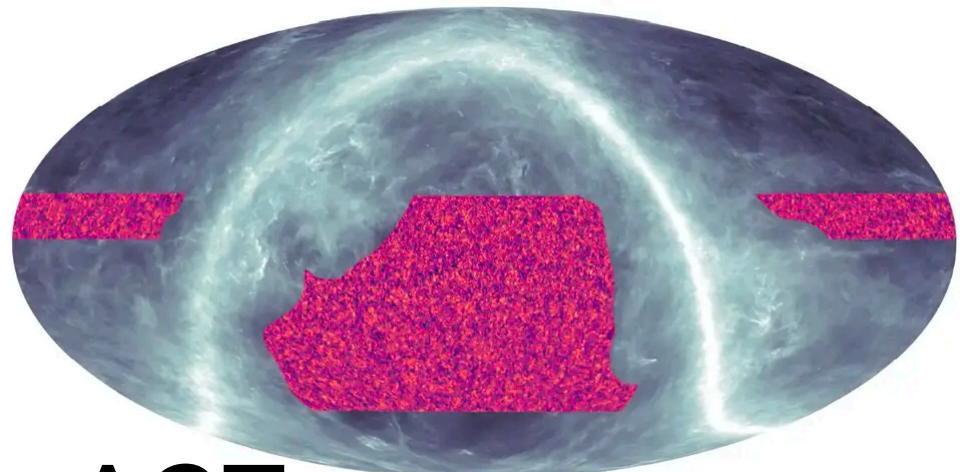
<http://theory2.phys.cwru.edu/~pete/GravitationalLens/>



Planck
(~all-sky)



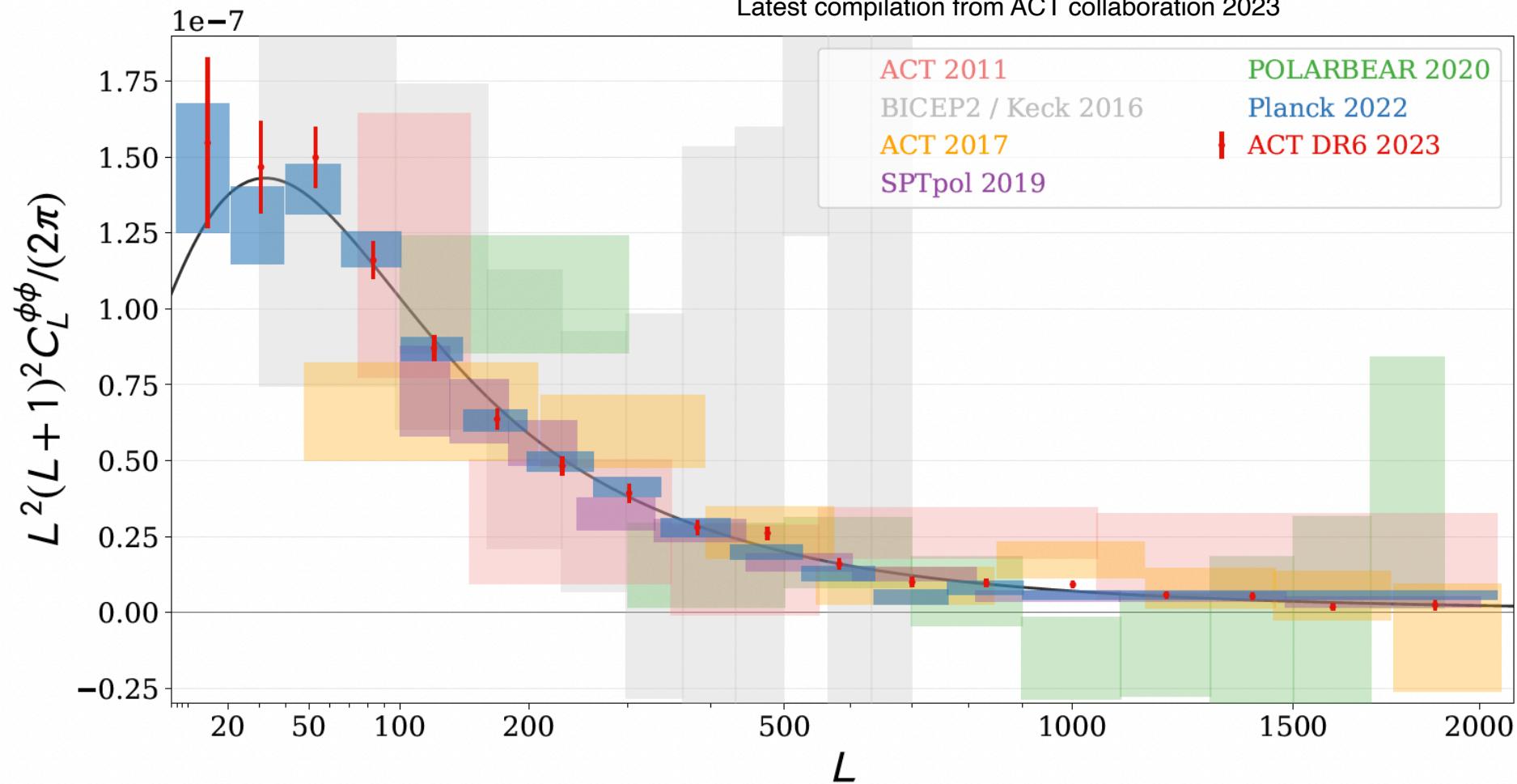
SPT-3G
(1500 square degrees)



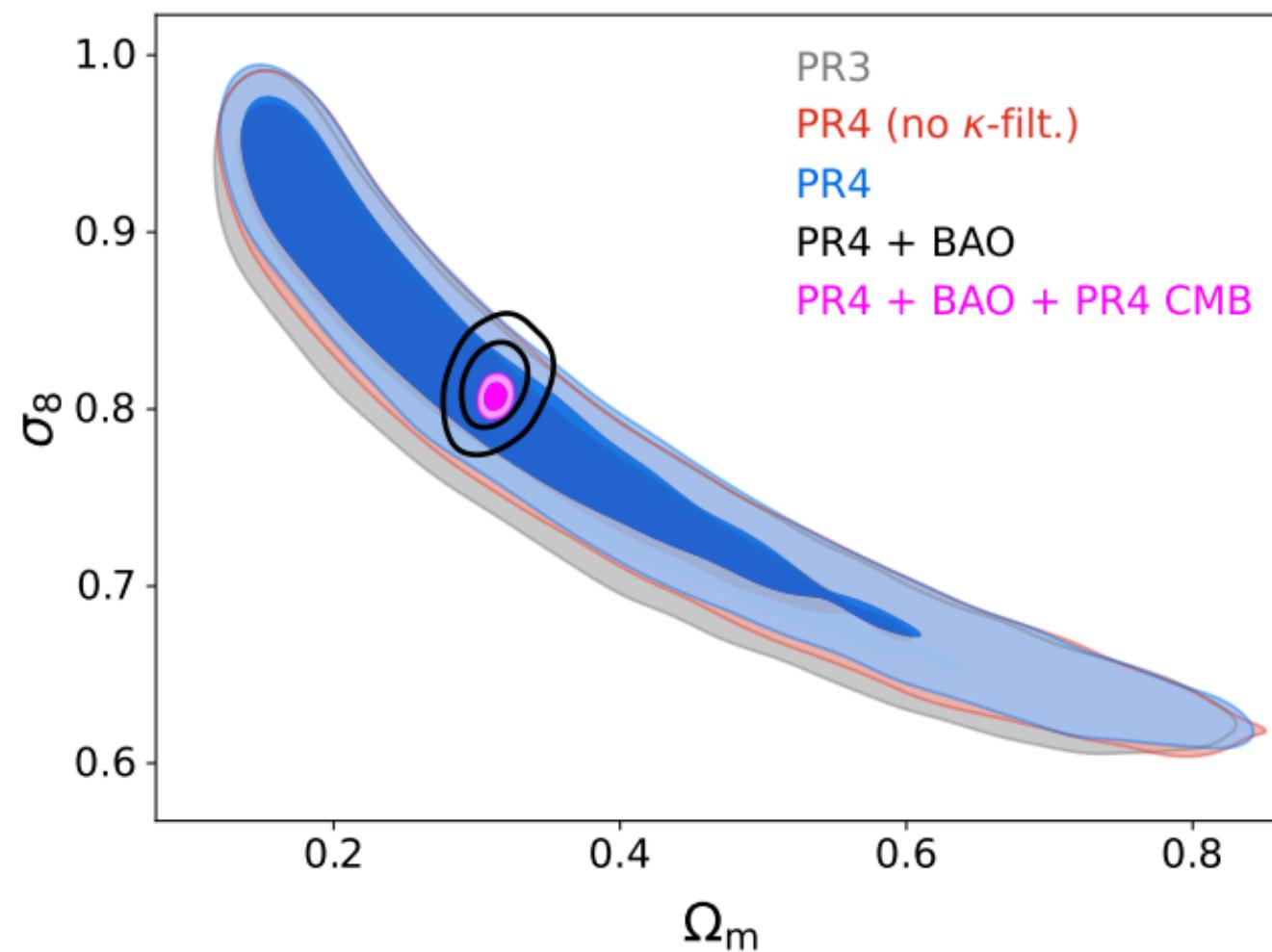
ACT
(9400 square degrees)

CMB Lensing Power Spectra

Latest compilation from ACT collaboration 2023



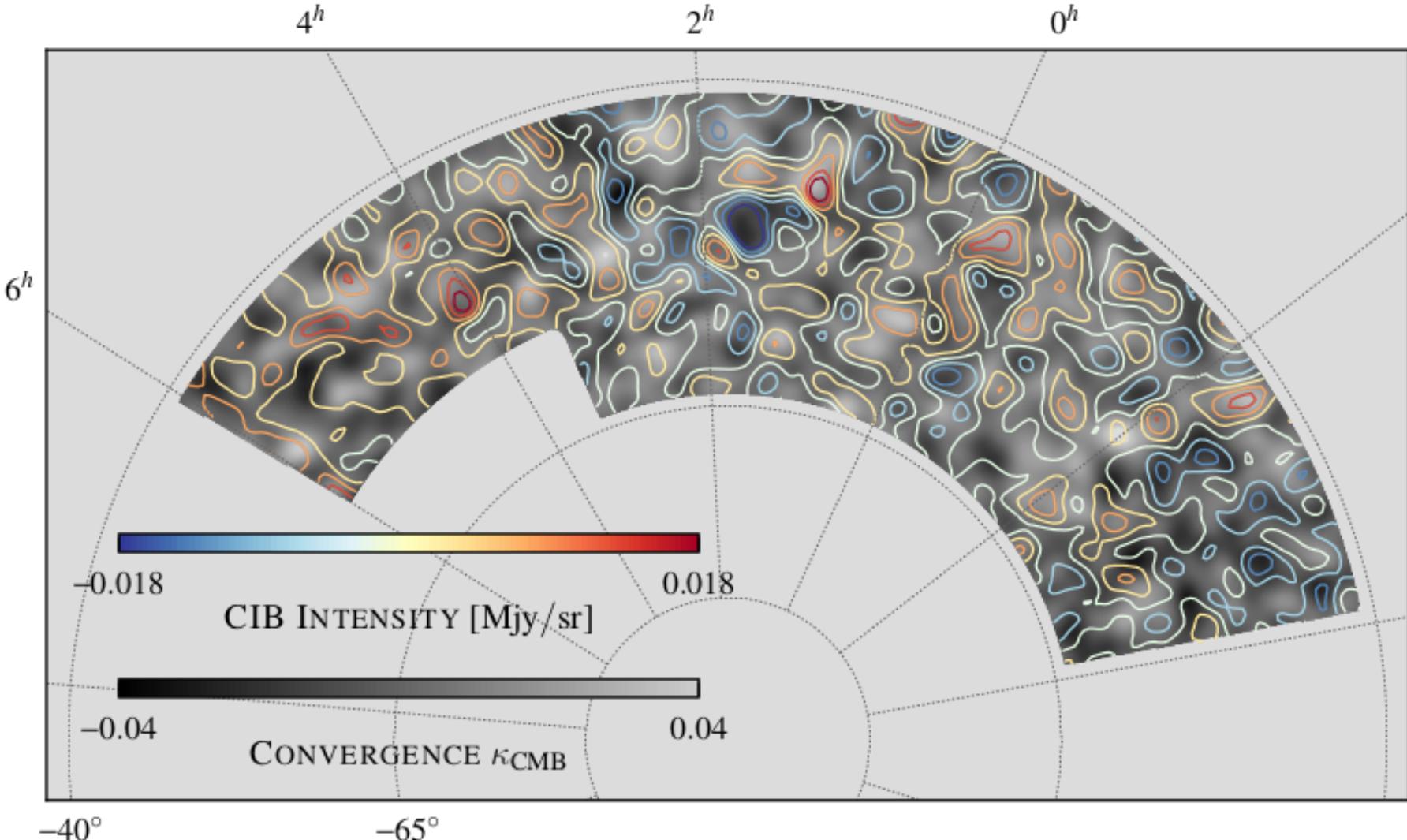
Cosmological constraints on structure formation



Planck: Carron 2022

see also Madhavacheril 2023

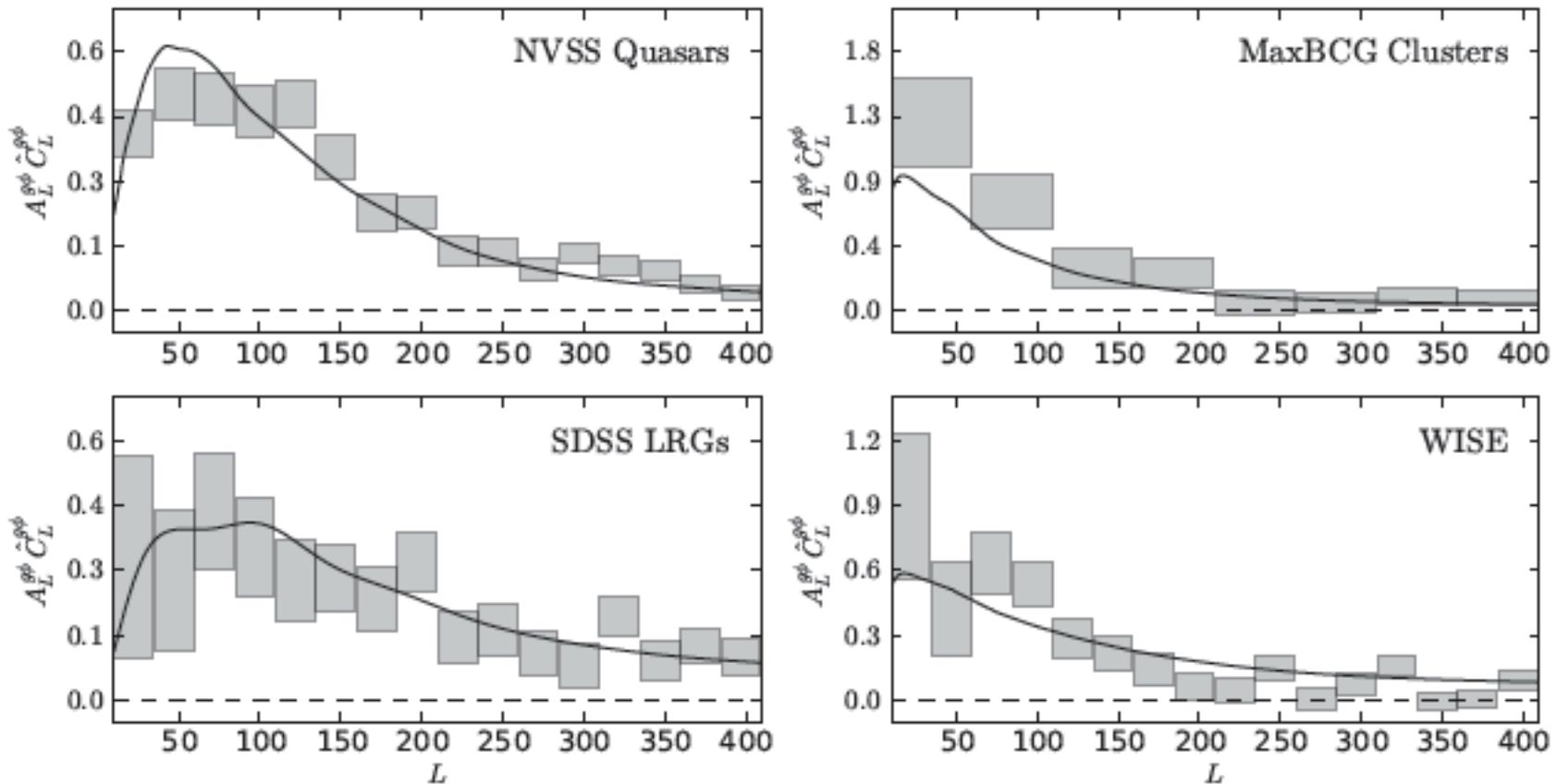
CMB-LSS cross-correlation: CIB



CIB map from Planck GNLIC 545 GHz

Omori, Chown, Simard, KTS, et. al (arXiv:1705.00743)

Planck X Galaxies, etc.



Angular Clustering

Angular power spectrum of power spectrum between two maps X & Y (could be same map!)

$$C_\ell^{XY} = \frac{2}{\pi} \int_0^\infty d\chi_1 d\chi_2 W^X(\chi_1) W^Y(\chi_2) \int_0^\infty k^2 dk P_{XY}(k; z_1, z_2) j_\ell(k\chi_1) j_\ell(k\chi_2)$$

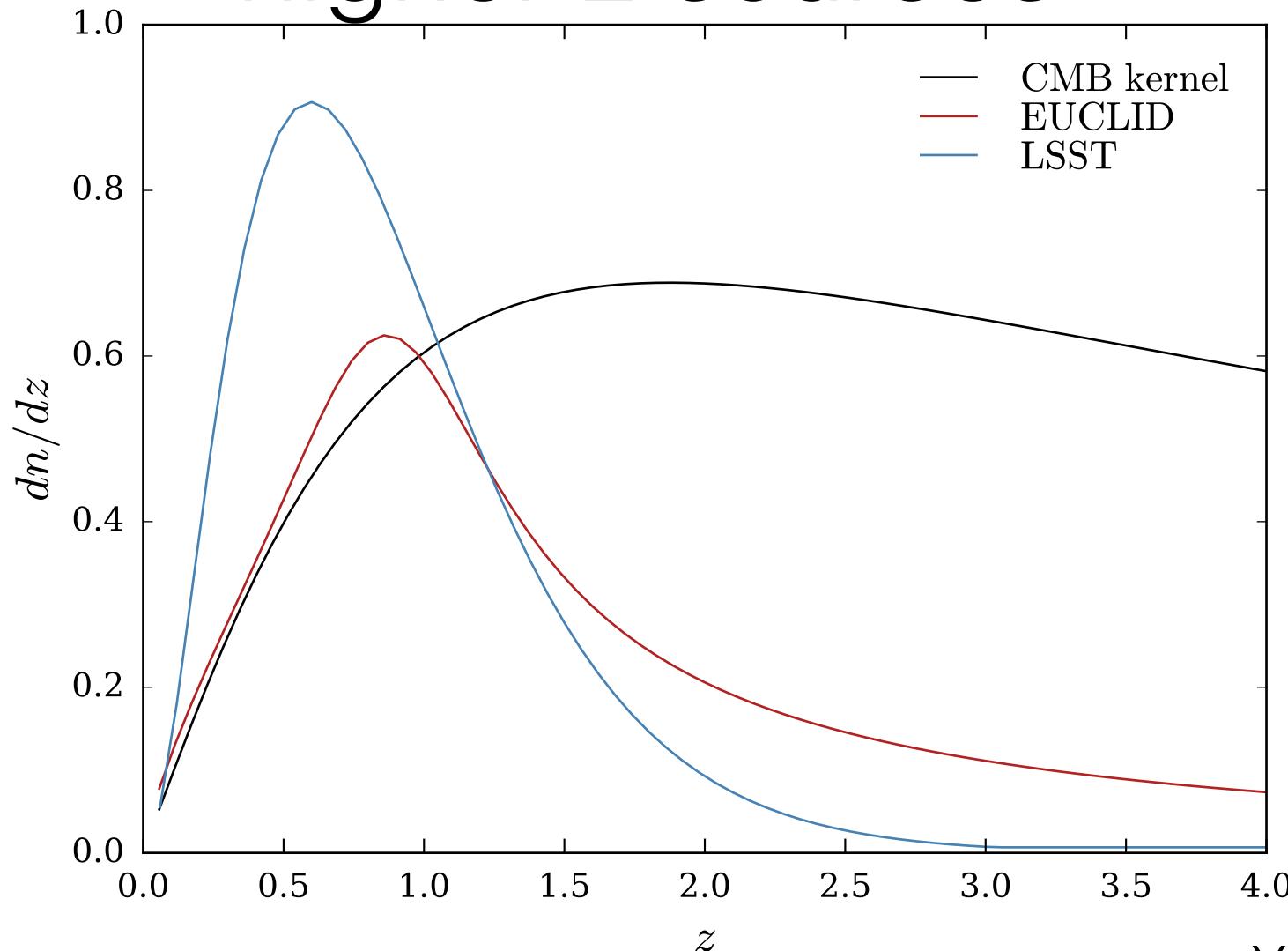
Limber approximation, which generally works pretty well except for really large scales

$$C_\ell^{XY} = \int d\chi \frac{W^X(\chi) W^Y(\chi)}{\chi^2} P_{XY} \left(k_\perp = \frac{\ell + 1/2}{\chi}, k_z = 0 \right)$$

weights for CMB lensing or some galaxy tracer

$$W^\kappa(\chi) = \frac{3}{2} (\Omega_m + \Omega_\nu) H_0^2 (1+z) \frac{\chi(\chi_\star - \chi)}{\chi_\star} , \quad W^g(\chi) = b(z) H(z) \frac{dN}{dz}$$

CMB lensing is sensitive to higher z sources



Angular Clustering

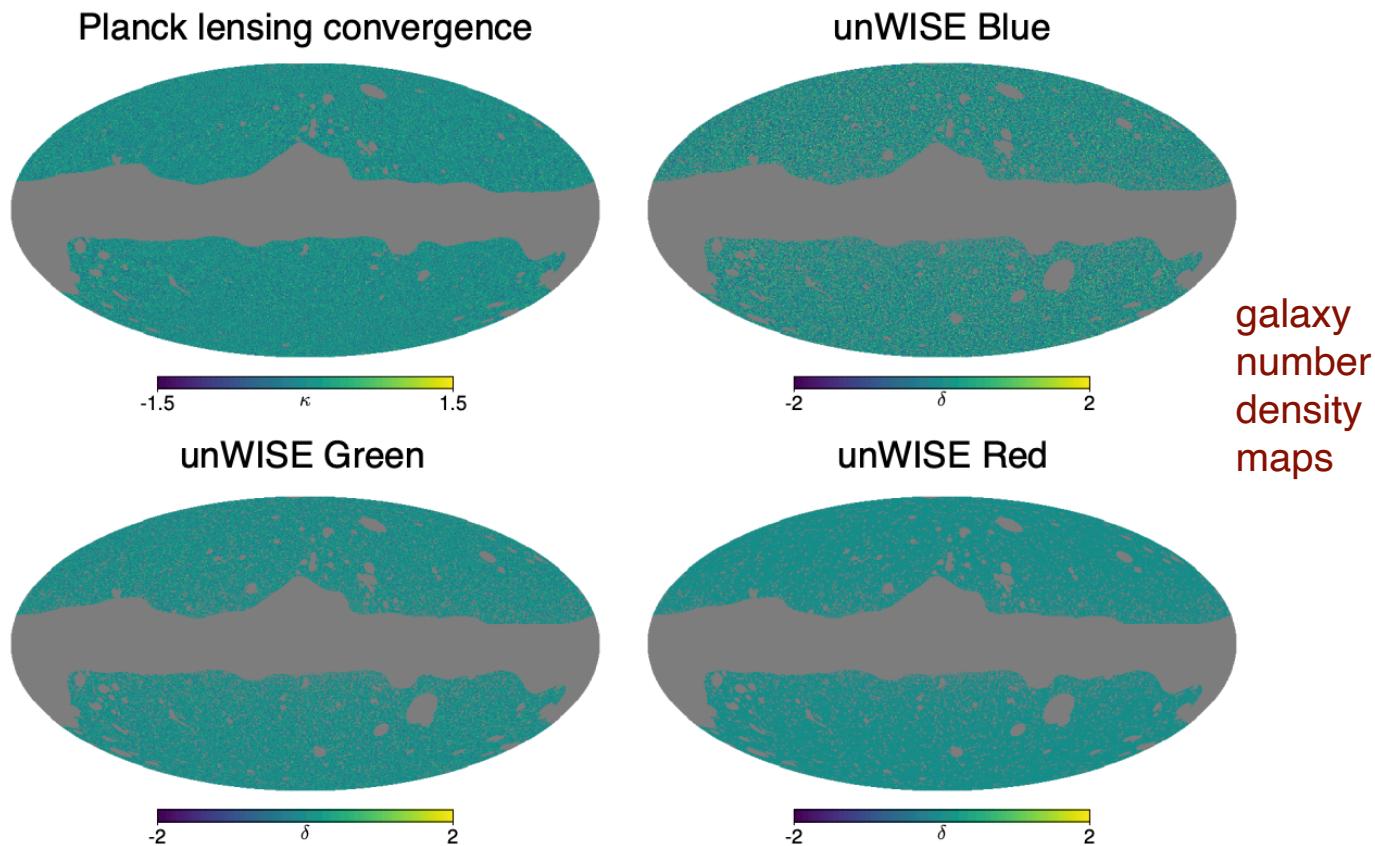
$$C_\ell^{\kappa g} = b^{\text{eff}} \int d\chi \frac{W^\kappa(\chi)}{z} H(z) \left[f(z) \frac{dN_p}{dz} \right] P(k\chi = \ell + 1/2)$$

$$C_\ell^{gg} = (b^{\text{eff}})^2 \int d\chi \frac{1}{\chi^2} H(z)^2 \left[f(z) \frac{dN_p}{dz} \right]^2 P(k\chi = \ell + 1/2)$$

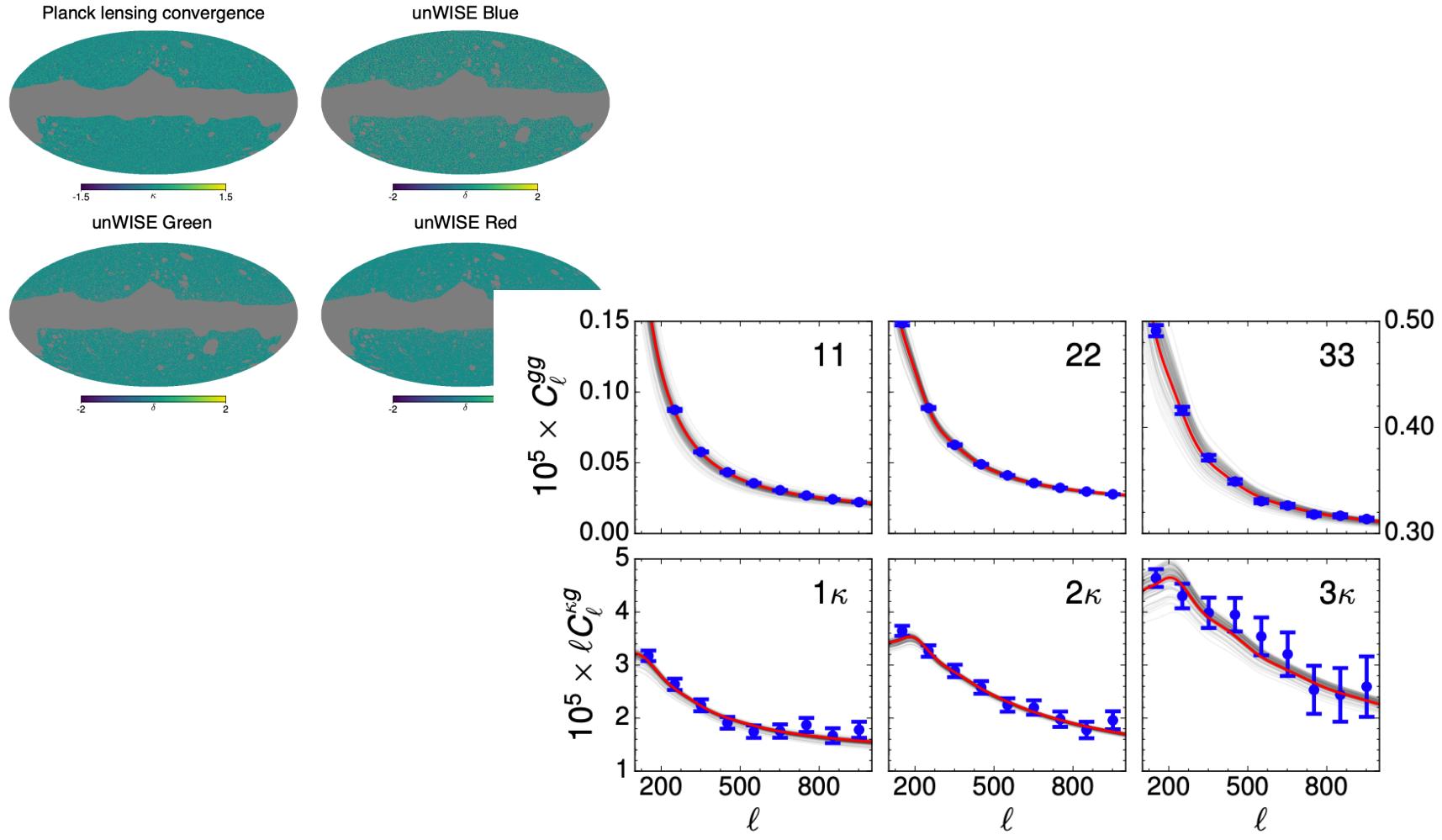
$$W^\kappa(\chi) = \frac{3}{2}(\Omega_m + \Omega_\nu)H_0^2(1+z) \frac{\chi(\chi_* - \chi)}{\chi_*} , \quad W^g(\chi) = b(z)H(z) \frac{dN}{dz}$$

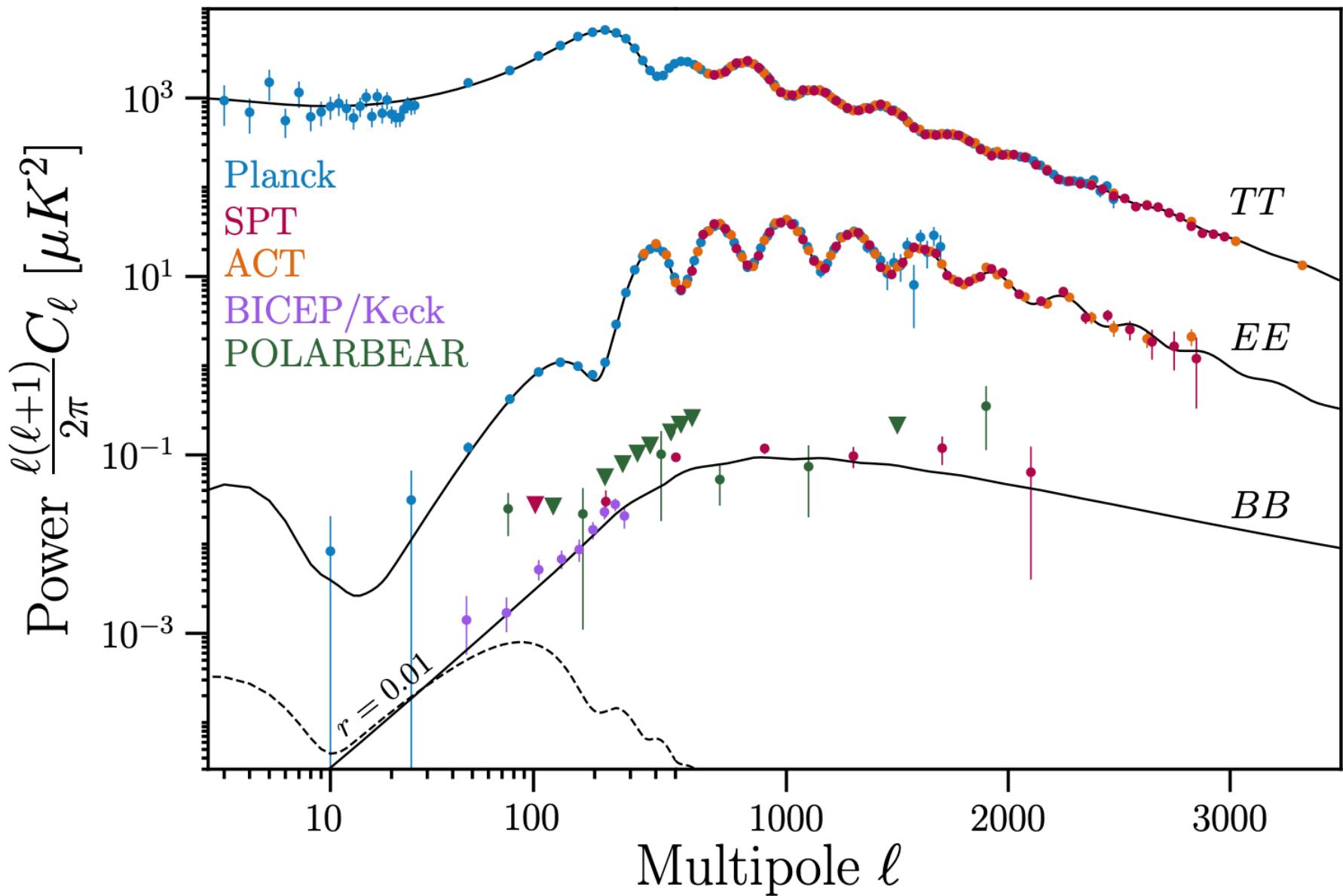
- CMB lensing power measures projected power of all matter (**no b**)
- galaxy clustering measures projected power of biased tracers (**b dN/dz**)²
- CMB lensing X galaxies measures projected power in common (**b dN/dz**)

Example: WISE X Planck lensing



Example: WISE X Planck lensing





Power spectrum Uncertainties

- fundamentally limited by number of independent measurements, noise
- $C_{l;\text{meas}} = C_{l;\text{true}} + C_{l;\text{noise}}$ *in any single map you can't tell the difference*
- $\text{Var}(C_l) \sim (2/n_{\text{meas}}) C_l^2$ **“sample variance”**
- more modes means better measurement of $C_{l;\text{true}} + C_{l;\text{noise}}$
- lower noise gives better measure of $C_{l;\text{true}}$

Delensing lowers sample variance for B-mode searches

SPT-3G + external tracers (galaxies+CIB) can remove 80% of lensing power

BICEP/Keck is signal-dominated, so delensing directly reduces the error bar for constraints on tensors (also true for SPT-3G, for however low in ℓ can be reached)

