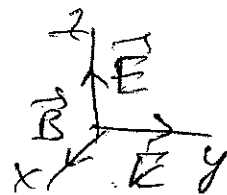


Interaction of an atom with an EM wave (cont.)

Last time:  $H = H_0 + \underbrace{\frac{e}{mc} \vec{P} \cdot \vec{A}(\vec{r}, t)}_{(1)} + \underbrace{\frac{e}{mc} \vec{S} \cdot \vec{B}(\vec{r}, t)}_{(2)} + \underbrace{\frac{e^2}{2mc^2} \vec{A}^2}_{(3)}$

$$V_1 \gg V_2 \gg V_3$$

$$\vec{A} = (A_0 e^{i(ky - \omega t)} + A_0^* e^{-i(ky - \omega t)}) \hat{e}_z$$



$$\vec{E} = E_0 \hat{e}_z \cos(ky - \omega t) \quad ; \quad \frac{E_0}{B_0} = \frac{\omega}{k} = c \quad ; \quad i\omega A_0 = \frac{E_0}{2}$$

$$\vec{B} = B_0 \hat{e}_x \cos(ky - \omega t) \quad ; \quad i k A_0 = \frac{B_0}{2}$$

$$V_1 = \frac{e}{mc} \vec{P} \cdot \vec{A}(\vec{r}, t) = \frac{e}{mc} p_z (A_0 e^{iky}) e^{-i\omega t} + c.c.$$

$$V_0 = V_{0DE} + \dots$$

← electric dipole

Transition rate  $i \rightarrow f \Rightarrow$

$$P_{i \rightarrow f} = \frac{2\pi}{\hbar} \left( \frac{e}{mc} \right)^2 |A_0|^2 |\langle f | e^{iky} p_z | i \rangle|^2$$

$$E_f = E_i + \hbar\omega$$

Electric-dipole approx.:

$$e^{iky} \approx 1 \Rightarrow \langle f | e^{iky} p_z | i \rangle \approx \langle f | p_z | i \rangle =$$

$$= i m \omega_{fi} \langle f | z | i \rangle \Rightarrow$$

selection rules

$$\Delta l = \pm 1$$

$$\Delta m = 0 \quad (\text{or } \pm 1 \text{ if } x\text{- or } y\text{-polar})$$

Now take into account higher-order terms  $\Rightarrow$   
 $e^{iky} \approx 1 + iky + O(ky)^2$

$$V_0 - V_{0,DE} \underset{\substack{\text{neglect} \\ \text{terms} \\ \sim (ky)^2}}{=} \frac{e}{mc} p_z A_0 \cdot (iky) = \frac{e}{mc} \frac{B_0}{2} p_z y =$$

$$= \frac{e}{mc} \frac{B_0}{2} \left[ \frac{1}{2} (p_z y - z p_y) + \frac{1}{2} (p_z y + z p_y) \right] \quad (11.1)$$

$\parallel$   
 $L_x$

$$V_0 \underset{\substack{\text{only} \\ \text{time-independent} \\ \text{part}}}{=} \frac{e}{mc} \vec{S} \cdot \vec{B} \underset{\substack{\text{only} \\ \text{time-independent} \\ \text{part}}}{=} \frac{e}{mc} S_x \frac{B_0}{2} \quad (11.2)$$

$\leftarrow V_0(t) = \frac{e}{mc} S_x B_0 \cos \omega t$   
 $\frac{1}{2} (e^{i\omega t} + e^{-i\omega t})$

Combine (11.1) and (11.2)  $\Rightarrow$

$$V_0 = V_{0,DE} + \frac{e}{mc} \frac{B_0}{2} \cdot \frac{1}{2} (L_x + 2S_x) + \frac{e}{mc} \frac{B_0}{2} \cdot \frac{1}{2} (p_z y + z p_y)$$

$\parallel$   
 $V_{0,DM}$

Selection rules?



$V_{0,DM} \leftarrow$  magnetic dipole transitions

$V_{0,QE} \leftarrow$  electric quadrupole transitions

$\longleftrightarrow$   
the same order!

$$\langle f | \vec{V}_{DM} | i \rangle \sim \langle f | L_x + 2S_x | i \rangle = \quad (3)$$

$$= \langle n_f, l_f, m_f, m_{sf} | L_x + 2S_x | n_i, l_i, m_i, m_{si} \rangle \neq 0 \quad \text{if}$$

$$\begin{array}{ccc} \uparrow & \uparrow & \\ \frac{L_+ + L_-}{2} & \frac{S_+ + S_-}{2} & \end{array} \quad \begin{array}{l} m_f = m_i \pm 1 \\ m_{sf} = m_{si} \pm 1 \end{array}$$

$$\text{If } \vec{B} \parallel Oz \Rightarrow \langle f | L_z + 2S_z | i \rangle \neq 0 \quad \text{if}$$

$$\begin{array}{l} \text{Since } \vec{V}_{DM} \text{ doesn't act on } l \\ \Downarrow \Delta l = l_f - l_i = 0 \end{array} \quad \begin{array}{l} m_f = m_i \\ m_{sf} = m_{si} \end{array}$$

So, selection rules for the magnetic dipole transition are  $\Delta l = 0$ ;  $\Delta M = \pm 1$  or  $0$   
 $\Delta m_s = \pm 1$  or  $0$

What about  $\langle f | \vec{V}_{QE} | i \rangle$  ?  $\Rightarrow$

$$\langle f | P_z y + z P_y | i \rangle = \langle f | \frac{im}{\hbar} [H_0, z] y + \frac{im}{\hbar} [H_0, y] z | i \rangle$$

$$= \langle f | \frac{im}{\hbar} [H_0, zy] | i \rangle \quad \left( \begin{array}{c} \uparrow \\ P_z = \frac{im}{\hbar} [H_0, z] \\ y \end{array} \right)$$

$$\textcircled{=} \frac{im}{\hbar} \omega_{fi} \langle f | zy | i \rangle$$

↖ component of quadrupole moment

$$yz \sim r^2 (A Y_2^1 + B Y_2^{-1}) \Rightarrow$$

$$\langle f | yz | i \rangle \sim \int Y_{\ell_f}^{m_f*} Y_2^{\pm 1} Y_{\ell_i}^{m_i} d\Omega \Rightarrow 0 \text{ unless}$$

also take into account parity of  $Y_{\ell}^m$

$$\rightarrow \Delta \ell = 0, \pm 2$$

$$\Delta m = 0, \pm 1, \pm 2$$

↑  
take into account all  $\langle f | xy | i \rangle, \langle f | xz | i \rangle$

Further expansion: electric octupole, magnetic quadrupole

Analysis . Because of selection rules

$V_{DM}$  and  $V_{QE}$  never compete with  $V_{ED}$

•  $V_{DM}$  &  $V_{QE}$  can be separated by observing  $\Delta \ell = \pm 2$  transitions  $\Rightarrow$  e.g. 557.7 nm line of atomic oxygen

Back to electric dipole approximation and absorption,  
define an absorption cross-section  $\Rightarrow \sigma_{abs} \Rightarrow$

$$\sigma_{abs} = \frac{\text{Energy per unit time, absorbed by the atom}}{\text{Energy flux of the radiation field}} =$$

↑  
Sakurai p. 336

↑ energy per area - per unit time

$$= \frac{\hbar \omega P_{i \rightarrow f}}{c U} = \frac{\hbar \omega \left( \frac{2\pi}{\hbar} \right) \left( \frac{e}{m_e} \right)^2 |A_0|^2 m^2 \omega_{fi}^2 |\langle f | z | i \rangle|^2}{\frac{1}{2\pi} \frac{\omega^2}{c} |A_0|^2} \quad (5)$$

$$c U = \frac{1}{2\pi} \frac{\omega^2}{c} |A_0|^2$$

↑  
energy density

energy absorbed  
by atom,  $E_f = E_i + \hbar \omega$

↑  
Dirac delta  
 $\delta(E_f - E_i - \hbar \omega)$

$$\omega_{fi} \approx \omega$$

$$= 4\pi^2 \frac{e^2}{\hbar c} \omega_{fi} |\langle f | z | i \rangle|^2 \delta(\omega_{fi} - \omega)$$

" $\alpha$ " ← fine-structure constant

Define oscillator strength  $f_{fi} = \frac{2m \omega_{fi}}{\hbar} |\langle f | z | i \rangle|^2$

↑  
determines the strength  
of the transition

$$\sum_f f_{fi} = 1$$

↑  
all possible  
final states

Thomas-Reiche-Kuhn  
sum rule

↑  
show! (HW)

