

Fermion-Statistics

Recall from last class, we discussed how we could take the many bodied wave function $\Psi(\vec{r}_1, \vec{r}_2, \dots \vec{r}_n) =$ slater determinant that ensures that we have the correct anti-symmetric combination of orbitals. These orbitals $\phi_i(\vec{r}_j)$ are the energy eigenstates of the single particle system. The *trick* we are going to be using is to take each of the orbitals and treat them as if they were a separate systems... weird.

Now think of our $\phi_i(\vec{r})$ as having an occupancy of 1 or 0.



Figure 1: general idea of energy levels for our single particle

We have that our Gibbs sum for a ϕ_i From Pauli's exclusion principle, we can only have one particle :

$$\begin{aligned}
 \mathbb{Z} &= \sum_i e^{-\beta(E_i - \mu N_i)} \\
 &= 1 + e^{-\beta(\epsilon - \mu)} \\
 \Rightarrow \langle E \rangle &= \epsilon \frac{e^{-\beta(\epsilon - \mu)}}{1 + e^{-\beta(\epsilon - \mu)}} \\
 &= \epsilon \frac{1}{1 + e^{\beta(\epsilon - \mu)}} \\
 \langle N \rangle &= \frac{e^{-\beta(\epsilon - \mu)}}{1 + e^{-\beta(\epsilon - \mu)}} \\
 &= \frac{1}{1 + e^{\beta(\epsilon - \mu)}} = \text{Fermi-Dirac function}
 \end{aligned}$$

Definitions The *Fermi-Level* is synonymous with the chemical potential μ . The *Fermi-energy* is $\mu(T = 0)$.

Boson-statistics

We do not have the anti-symmetry requirement and so we can have many combinations of orbitals to make a combined state. Now our orbitals ϕ_i with energy ϵ and occupancy $N = 0, 1, 2, 3, \dots \infty$.

$$\begin{aligned}\mathbb{Z} &= 1 + \sum_{j=1}^{\infty} \left(e^{-\beta(\epsilon - \mu j)} \right) \\ &= 1 + \sum_{j=1}^{\infty} \left(e^{-\beta(\epsilon - \mu)} \right)^j \\ &= \frac{1}{1 - e^{-\beta(\epsilon - \mu)}} \\ \langle N \rangle &= \sum_j \frac{j e^{-\beta(\epsilon - \mu)j}}{\mathbb{Z}} \\ &= \frac{1}{\mathbb{Z}} \frac{\partial \mathbb{Z}}{\partial(\mu\beta)} \\ &= \frac{e^{-\beta(\epsilon - \mu)}}{1 - e^{-\beta(\epsilon - \mu)}} = \frac{1}{e^{\beta(\epsilon - \mu)} - 1} = f_{BE}(\epsilon)\end{aligned}$$