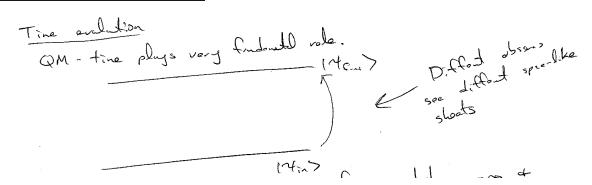
Lecture 10

QFT Continued...

Summary From Last Time



Only have a hope of Lorentz invariance if we start at $-\infty$ and go to $+\infty$.

Throw particles in from ∞ let them scatter & go back out to ∞ .

Define S-Matrix

$$\underbrace{|p_1\sigma_1,...p_n\sigma_n\rangle}_{t=-\infty} \to \underbrace{\mathcal{S}|p_1\sigma_1,...p_n\sigma_n\rangle}_{t=+\infty}$$

S might be (at least a hope) Lorentz Invariant.

Big Picture: The plan is to Figure out what S is in a totally generic theory, then see what it would take to make it Lorentz Invariant.

Sure doesnt look like ti will be L.I. \mathcal{S} is the only object that you could even have a hope to make L.I.

We will see that for <u>very special</u> choices of the interaction it will barely be possible for it to be Lorentz Invariant. These choices force on us anti-particle and the connection between spin and statistics.

Something annoying that we should get rid of right away. Free evolution, just evolves w/phase. Totally irrelevant part.

Standard way of removing the free evolution "Interaction Representation".

$$H = H_{\text{free}} + H_{\text{Int}}$$

$$i\frac{d|\psi\rangle}{dt} = (H_{\text{free}} + H_{\text{Int}})|\psi\rangle$$

For $H_{\text{Int}} = 0$

$$|\psi\rangle = e^{-iH_f t} |\psi_{in}\rangle$$

Now, we dont have a free theory, but if the interaction is small going to be pretty close to evolving like this.

$$|\psi\rangle = e^{-iH_f t} \underbrace{|\psi_I\rangle}_{\text{definition}}$$
 (1)

If $H_{\text{Int}} = 0$, $|\psi_{\text{Int}}\rangle$ doesnt evolve at all. Bc there is H_{Int} , $|\psi_{\text{Int}}\rangle$ will evolve.

$$i\frac{d}{dt}|\psi\rangle = H_{f}|\psi\rangle + e^{-iH_{f}t}i\frac{d}{dt}|\psi_{I}\rangle$$
$$= (H_{f} + H_{int})e^{-iH_{f}t}|\psi_{I}\rangle$$

Note: first line from derivitive of 1, the second from the Schrodinger Equation. The RHSs imply,

$$i\frac{d}{dt}|\psi_I\rangle = \underbrace{e^{-iH_ft}H_{\rm Int}e^{-iH_ft}}_{\text{Interaction Hamiltonian in the interaction representation}}|\psi_I\rangle$$

So,

$$i\frac{d}{dt}|\psi_I\rangle = H_I|\psi_I\rangle$$

where H_I can be time dependent.

Lets formally sovle this

Just intergrating gives,

$$|\psi_I(t_2)\rangle = |\psi_I(t_1)\rangle - i \int_{t_1}^{t_2} dt H_I |\psi_I(t)\rangle$$

Now we can keep iterating,

$$= |\psi_I(t_1)\rangle - i \int_{t_1}^{t_2} dt H_I(t) \left(|\psi_I(t_1)\rangle - i \int_{t_1}^t dt' H_I(t') |\psi_I(t')\rangle \right)$$

or

$$= |\psi_I(t_1)\rangle - i \int_{t_1}^{t_2} dt H_I \, |\psi_I(t_1)\rangle + (-i)^2 \int_{t_1}^{t_2} dt \int_{t_1}^t dt' H_I(t) H_I(t') \, |\psi_I(t')\rangle$$

Pattern is clear, can keep going...

$$\begin{split} |\psi_I(t_2)\rangle &= [\quad 1 \quad + (-i) \int_{t_1}^{t_2} dt H_I(t) \\ &+ \quad (-i)^2 \int_{t_1}^{t_2} dt \int_{t_1}^t dt' H_I(t) H_I(t') \\ &+ \quad (-i)^3 \int_{t_1}^{t_2} dt \int_{t_1}^t dt' \int_t^{t'} dt'' H_I(t) H_I(t') H_I(t''') \\ &+ \quad ...] |\psi_I(t_1)\rangle \end{split}$$