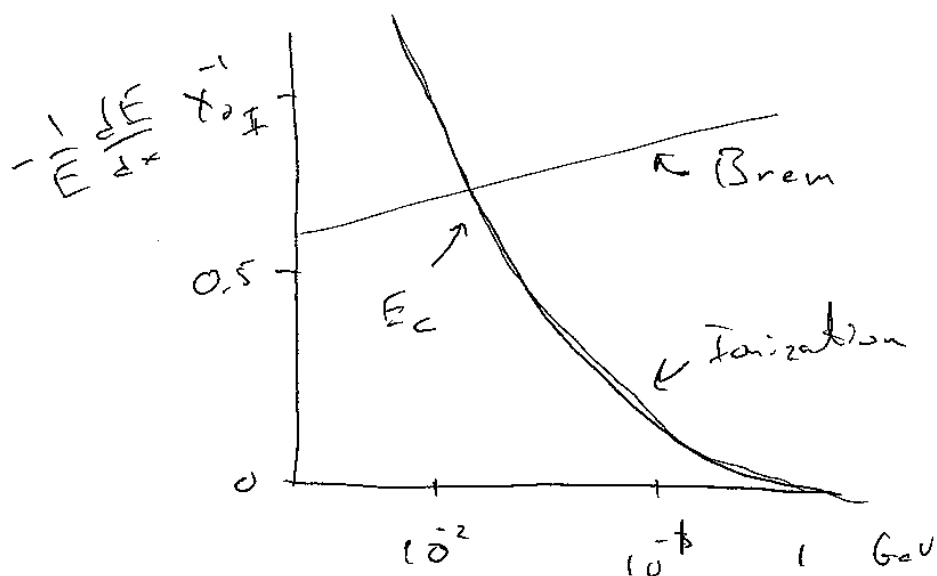
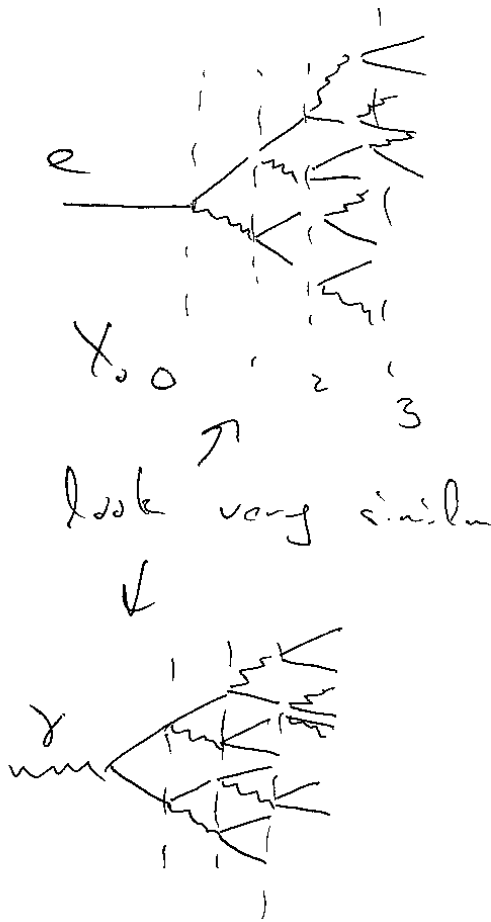


Lecture 24



Similar story for photons.

E_c - critical energy $O(10^{-2} - 10^{-1} \text{ GeV})$



“Electromagnetic Shower”

-) For each e (or γ) with $E > E_c$ travels $\sim 1x_0$ then gives up 1/2 energy to γ (or ee).

-) e 's, γ 's with energy $< E_c$ get absorbed via ionization

If initial energy $E_0 \gg E_c$ then after t -radiation lengths there will be 2^t particles. Approximately equal e 's and γ 's each with energy $E(t) \sim \frac{E_0}{2^t}$.

Shower will stop growing when

$$E(t) \simeq E_c \equiv E(t_{max})$$

t_{max} is point in shower with max particles

$$t_{max} = \frac{\ln \frac{E_0}{E_c}}{\ln 2}$$

\Rightarrow shower depth grows as \ln

N_{max} given by E_0/E_c

For heavier particles, Brem does not kick in until much higher energies.

eg: $E_c \sim 3000 \text{ GeV}$ for a muon in lead.

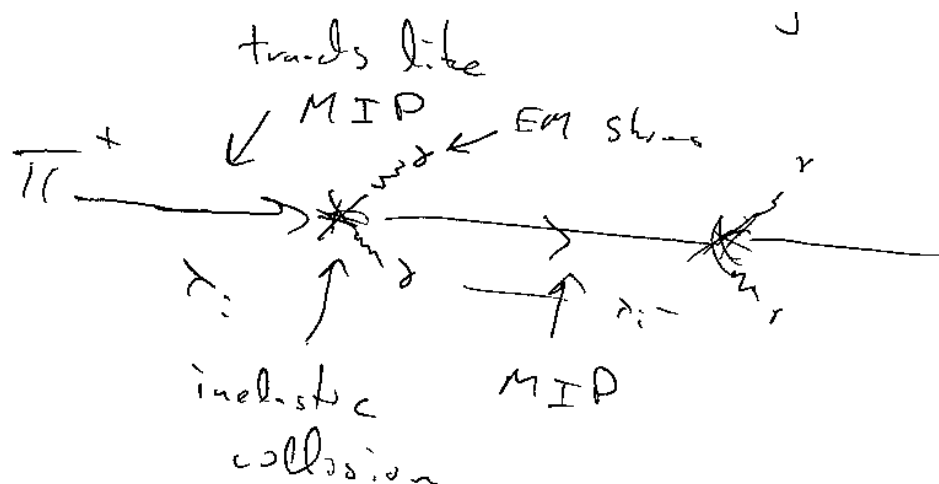
Nuclear Interactions

If the particle traversing the medium is a hadron, can interact via the strong interaction.

eg: π^\pm moving through detector material will suffer an inelastic collision in distance λ_i

This collision takes energy from π and converts it to additional hadrons (charged and neutral)

Leads to what is called “Hadronic Shower” - same basic idea as an EM shower, but much more complex. Involves a variety of processes at different length scales



At the inelastic collision, typically produce

$\pi^+ \sim 1/3$ of the time. Travels like a MIP w/scale λ_i

$\pi^0 \sim 1/3 \leftarrow Br(\gamma\gamma) \sim 100\%$ travels $10\text{nm} \Rightarrow$ EM shower w/ scale X_0

$\pi^- \sim 1/3$

Both π^\pm travel like MIPs of order λ_i . \Rightarrow Hadronic showers develop over larger distances and contain much more fluctuations.

Trackers

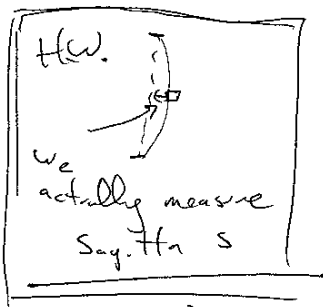
Ionization detectors in a magnetic field.

Charged particles traveling in a magnetic field move in circles.

$$\frac{1}{R[\text{m}]} = Q \cdot B[\text{T}] \cdot \frac{0.3}{p_T[\text{GeV}]}$$

Particle positions are measured by finely etched silicon sensors.
 $\sim 10\mu\text{m}$ resolution

We measure the curvature $\sim \frac{1}{R} \sim \frac{1}{p_T}$



$$s = \frac{qBL^2}{8p_T}$$

(More important to have bigger length than larger B)

$$\Rightarrow \Delta \frac{1}{p_T} = \frac{\Delta p_T}{p_T^2}$$

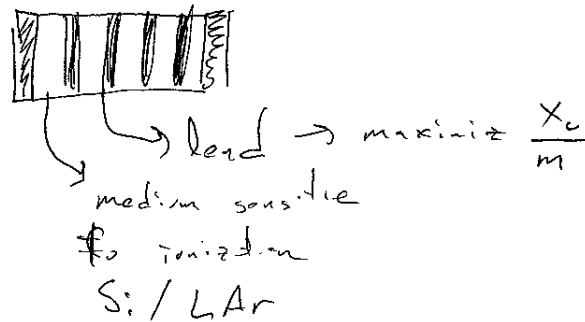
So the relative uncertainty in p_T degrades with p_T .

LHC typical performance

$$\frac{\Delta p_T}{p_T} \sim (\text{few } \%) \frac{p_T}{100 \text{ GeV}}$$

EM Calorimeter

Contains and measures e^- 's γ 's and $\pi^0 \rightarrow \gamma\gamma$.



Total depth $\sim 20X_0$ (Enough to contain high p_T showers)



Only energy deposited in Si/LAr is measured.
Rest is lost in uninstrumented lead.

So only a small fraction of the total EM shower energy is collected (measured).

$E \sim N_c$ amount collected scales like the number of electrons produced

$$\frac{\Delta E}{E} \sim \frac{\Delta N_c}{N_c} \sim \frac{1}{\sqrt{N_c}} \sim \frac{1}{\sqrt{E}}$$

At the LHC typical performance

$$\frac{\Delta E}{E} \sim \frac{10\%}{\sqrt{E}}$$

Unlike trackers the relative energy measurement improves with E .

Hadronic Calorimeter

Compare shooting a π^\pm in the detector instead of an electron.

- π will travel further because it has to interact via the strong force (shorter range than EM)
- π will make other π^\pm and π^0
- Compared to electrons, make less particles (fewer at fixed E) b/c π mass is bigger.
- More varied showers depending on what particles are produced π^\pm vs π^0 .

\Rightarrow hadronic calorimeters need to be bigger, and will measure showers that have more fluctuations.

At the LHC typical performance worse than for EM calorimeters

$$\frac{\Delta E}{E} \sim \frac{50\%}{\sqrt{E}}$$

Relative energy measurement also improves with E.

Muon Detectors

Put tracking detectors outside of the calorimeters.

If any charged particle makes it through $\sim 3\text{m}$ of lead it has to be a muon.

μ | no strong interaction
| only rarely radiate γ s (below $\sim 3\text{ TeV}$)