

Relativistic Wave Eqs

①

Sch eq : $i \frac{d}{dt} \psi = \left(-\frac{\nabla^2}{2m} + V \right) \psi$

Problems - Conservation of non-relativistic Energy.

$$E \leftrightarrow i \frac{d}{dt}$$

$$p \leftrightarrow -i \nabla$$

$$\Rightarrow \text{Sch eq} \Rightarrow E = \frac{p^2}{2m} + V$$

Not L.I.

- time / position not on equal footing.

Start w/ relativistic E/p relation!

$$E^2 - p^2 - m^2 = 0$$

Klein-Gordon Eq $\left(-\frac{\partial^2}{\partial t^2} + \nabla^2 - m^2 \right) \phi(x,t) = 0$

Solution

$$\phi(x,t) = e^{i p \cdot x}$$

"on shell solutions"

every body is L.I.

Note two solutions
 $E > 0 + E < 0$

Alternative Standard Powerful Formulation of Physics (2)

via Lagrangians

$$KG \quad \underline{eq} \quad \frac{d^2}{dt^2} \phi = \nabla^2 \phi - m^2 \phi$$

$\phi(x,t)$ - permeates time & space

looks like $\frac{d^2}{dt^2} x(t) = - \frac{2U}{2X} \leftarrow \text{Potential Energy}$

$\frac{d^2}{dt^2} \phi$ - looks like the "acceleration" of the field

$-m^2 \phi$ - effective "force" from the mass.
(harmonic oscillator)

$\nabla^2 \phi$ - "shear force" shearing fields takes force

Now have some intuition about what KG telling us

Integrate force to get the potential

See from finite diff method

$$-\frac{2U}{2\phi} = \nabla^2 \phi - m^2 \phi \Rightarrow U = +\frac{1}{2}(\nabla\phi)(\nabla\phi) + \frac{m^2}{2}\phi^2$$

↖ Potential energy

K.E. from generalization of $\frac{1}{2}\dot{x}^2$ to fields

$$K = \frac{1}{2} \left(\frac{\partial}{\partial t} \phi \right)^2$$

$$K+U = \frac{1}{2} \left(\frac{\partial}{\partial t} \phi \right)^2 + \frac{1}{2}(\nabla\phi) \cdot (\nabla\phi) + \frac{m^2}{2}\phi^2 \leftarrow \text{at the point } (x,t)$$

Total energy from integrating

$$H(t) = \int d^3x \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 + \frac{1}{2} (\nabla \phi)(\nabla \phi) + \frac{m^2}{2} \phi^2 \right]$$

→ with this you could go back to 3 lectures ago
and replace $\frac{1}{2} \frac{\partial \phi^\dagger \partial \phi}{m}$ and to everything again
(We won't bother here ...)

However we can use this to formulate the field eq
in a totally different, ultimately more useful way...

think of classical mechanics of point particles
can all be formulated with Lagrangian & principle of
least action.

↙ difference
KE & Potential

$$L = \frac{1}{2} \dot{x}^2 - U(x)$$

The action $S[x(t)] = \int dt L$

↙ you could take this
as the starting point
& derive Newton's 2nd law
by minimizing the action

$$\int_{x(t)}^{\phi} S = \int_{x(t)}^{\phi} \left(\frac{1}{2} \dot{x}^2 - U(x) \right) dt$$

Do the same for relativistic fields

$$L = K - U = \int d^3x \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \frac{1}{2} (\nabla \phi)(\nabla \phi) - \frac{m^2}{2} \phi^2 \right]$$

$$S[\phi(t)] = \int d^4x \left[\begin{array}{c} \nearrow \end{array} \right]$$

K.G. can be found from
minimizing action wrt ϕ

$$z_\mu = (\frac{z}{z_t}, -\vec{z})$$

$$S[\phi] = \int d^4x \left[\frac{1}{2} (z_\mu \phi)(z^\mu \phi) - \frac{m^2}{2} \phi^2 \right]$$

manifestly L. I.

This tells us how to construct general L I descriptions of Relativistic QM systems.

- Need L I Lagrangian and we are "done"

Substantially $z^\mu \phi \longleftrightarrow z_\mu \phi$

all of particle physics follows from this.

Two more relativistic systems important for S.M.

Spin 1 "EM" KG Spin 0

Spin 1/2 Dirac equation

KG 2nd order in x^μ . Contains no info about intrinsic angular momentum. (No free Lorentz indices)

ϕ from KG is Spin-0 field

Need relativistic eq. with free Lorentz indices

Can we find 1st order relativistic eq of motion?

Assume there exists a wave eq linear in $t + x$

$$\left(\alpha \frac{z}{z_t} + \vec{\beta} \cdot \vec{\nabla} \right) \psi = m \psi$$

for some $\alpha, \vec{\beta}, m$

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If rel invariant, must imply the KG

Assume that it is the square root

$$\left(\alpha \frac{d}{dt} + \vec{\beta} \cdot \vec{\nabla}\right) \left(\alpha \frac{\partial}{\partial t} + \vec{\beta} \cdot \vec{\nabla}\right) \psi = m^2 \psi$$

$$\Rightarrow \left[\alpha^2 \frac{\partial^2}{\partial t^2} + \alpha \vec{\beta} \cdot \vec{\nabla} \frac{\partial}{\partial t} + \vec{\beta} \cdot \vec{\nabla} \alpha \frac{\partial}{\partial t} + (\vec{\beta} \cdot \vec{\nabla})^2 \right] \psi = m^2 \psi$$

$$\text{to get } \left(-\frac{\partial^2}{\partial t^2} + \nabla^2\right) \psi = m^2 \psi$$

$$\Rightarrow \alpha^2 = -1 \quad \beta_i \beta_j = \delta_{ij} \quad \alpha \beta_i + \beta_i \alpha = 0$$

\nearrow
C#

\nearrow

Cannot be #'s

Lets keep going...

$$\alpha =: \gamma_0 \quad \beta_i =: \gamma_i$$

$$\gamma_0 \gamma_0 = 1$$

$$\gamma_i \gamma_j = -\delta_{ij}$$

$$\gamma_0 \gamma_i + \gamma_i \gamma_0 = 0$$

$$\gamma_\mu \gamma_\nu + \gamma_\nu \gamma_\mu = \begin{pmatrix} 2 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 \\ 0 & 0 & -2 & 0 \\ 0 & 0 & 0 & -2 \end{pmatrix} = 2 \eta_{\mu\nu}$$

$\{\gamma_\mu, \gamma_\nu\}$

$$(\gamma_\mu \partial^\mu - m) \psi = 0 \quad \leftarrow \text{Dirac equation}$$

γ matrices 4×4

$$\gamma_\mu = \begin{pmatrix} 0 & \sigma_\mu \\ \bar{\sigma}_\mu & 0 \end{pmatrix}$$

"Weyl basis"

$$\sigma = (1, \sigma_i)$$

$$\bar{\sigma} = (1, -\sigma_i)$$

Other choices you will
work out...

Dirac eq describes spin $\frac{1}{2}$ particles (e)

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B/c γ_μ are matrices

↳ has the Pauli spin matrices

the solution ψ will be a 4-component vector.
"Spinor"

Solution to Dirac Eq

Assume of the form $\psi(x) = u(p) e^{-ip \cdot x}$

$$(i \gamma_\mu \partial^\mu - m) \psi = (i \gamma_\mu (-ip^\mu) - m) u e^{-ip \cdot x} = 0$$

$$\Rightarrow \text{matrix eq } (\gamma \cdot p - m) u = 0$$

Massive Particle

$$p = (E, 0, 0, 0)$$

$$\begin{pmatrix} -mI & EI \\ EI & -mI \end{pmatrix} u = 0$$

↳ 2×2 sub blocks

$$u = \begin{pmatrix} C \\ D \end{pmatrix}$$

↖ C & D
2 component
spinors

$$\Rightarrow -mC + ED = 0 \quad \& \quad EC - mD = 0$$

$$E^2 = m^2 \Rightarrow E = \pm m \quad \leftarrow \text{Same as seen in K.G.}$$

$$E = +m \Rightarrow C = D \quad E = -m \Rightarrow C = -D$$

$$\psi_+ = \begin{pmatrix} C \\ C \end{pmatrix} e^{-imt}$$

$$\psi_- = \begin{pmatrix} C \\ -C \end{pmatrix} e^{imt}$$

C - two components represent spin up or down