# Lecture 7

## **Review Quantum Mechanics (Dynamics)**

$$|\alpha, t_0\rangle \rightarrow |\alpha, t\rangle$$

This is what we mean by time evolution.

In QM, then there has to be an operator associated with taking the first state to the second.

Time Evolution Operator:  $U(t, t_0)$ 

$$|\alpha, t\rangle = U(t, t_0) |\alpha, t_0\rangle$$

 $U(t, t_0)$  Properties

- 1.  $U^{\dagger}(t, t_0)U(t, t_0) = 1$  Unitary See this from  $\langle \alpha t_0 | \alpha t_0 \rangle = \langle \alpha t | \alpha t \rangle$
- 2.  $U(t_2, t_0) = U(t_2, t_1)U(t_1, t_0)$  Composition Rule
- 3.  $U(t_0, t_0) = 1$

How can we possibly determine what the time operator is ???? (Should be getting old by now...)

Start with infinitesimal time evolution

$$U(t + \epsilon, t) = 1 - i\Omega\epsilon$$

where  $\Omega$  is a Hermitian operator (b/c) U is unitary So,

$$|\alpha, t + \epsilon\rangle = (1 - i\epsilon\Omega) |\alpha, t\rangle$$

OR,

$$\Omega \left| \alpha, t \right\rangle = i \frac{\left| \alpha, t + \epsilon \right\rangle - \left| \alpha, t \right\rangle}{\epsilon} = i \frac{\partial}{\partial t} \left| \alpha, t \right\rangle$$

in limit  $\epsilon \to 0$ 

### Physical Meaning of $\Omega$ :

As before to get the generate for the finite movement you have to exponetiate

$$U(t) = e^{-i\Omega t}$$

Note that  $\Omega$  has units 1/[time]. Just like energy.

Identify  $\Omega = \frac{1}{\hbar}H$  where H is the Hamiltonian operator.

$$i\frac{\partial}{\partial t}|\psi\rangle = \Omega|\psi\rangle$$

**Schrodinger Equation** 

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = H(t) |\psi\rangle$$

Non-relativistically:  $H \sim \frac{p^2}{2m} = -\frac{1}{2m} \frac{\partial^2}{\partial^2 x}$ 

## **Time Evolution**

$$|\alpha, t\rangle = U(t, t_0) |\alpha, t_0\rangle$$

Case I: H is time independent

$$U(t, t_0) = \lim_{N \to \infty} \prod_{i}^{N} e^{-\frac{i}{\hbar}H\Delta t} = e^{\frac{-iH(t - t_0)}{\hbar}}$$

<u>Case II:</u> H is time dependent, but [H(t), H(t')] = 0

$$U(t,t_0)=e^{\frac{-i}{\hbar}\int_{t_0}^t H(t')dt'}$$

Case III: H is time dependent, and  $[H(t), H(t')] \neq 0$ 

(Note: in general  $e^A e^B \neq e^{A+B}$  if  $[A, B] \neq 0$ )

$$U(t,t_0) = T \left[ e^{\frac{-i}{\hbar} \int_{t_0}^t H(t')dt'} \right]$$

"Time-ordered product" - Power series expansion with earlier terms on the right (Will come back to this later)

## **Example**

For a time independent Hamiltonian.

$$U(t,0) = e^{-iHt}$$

Choose a basis of eigenstates of H

$$H|n\rangle = E_n|n\rangle$$

and

$$H = \sum_{n} E_n |n\rangle \langle n|$$

Then

$$U(t,0) = \sum_{n} e^{-iE_{n}t} |n\rangle \langle n|$$

Now some arbitrary state:

$$\begin{aligned} |\psi(t)\rangle &= U(t,0) \, |\psi(0)\rangle = & \sum_{n} U(t,0) \, |n\rangle \, \langle n|\psi\rangle \\ &= & \sum_{n} |n\rangle \, e^{-iE_{n}t} \, \underbrace{\langle n|\psi\rangle}_{\text{time independent}} \end{aligned}$$

Now lets talk about the expectation values of an observable and how they change with time....

$$\langle A \rangle (t) = \langle \psi(t) | A | \psi(t) \rangle = \qquad \langle \psi(0) | U^{\dagger}(t,0) A U(t,0) | \psi(0) \rangle$$

$$= \underbrace{\left( \langle \psi(0) | U^{\dagger}(t,0) \right) A \left( U(t,0) | \psi(0) \rangle \right)}_{\text{"Schrodinger Picture"}}$$

$$= \underbrace{\left\langle \psi(0) | \left( U^{\dagger}(t,0) A U(t,0) \right) | \psi(0) \right\rangle}_{\text{"Heisenberg Picture"}}$$

#### Schrodinger Picture

- $|\psi(t)\rangle$ s move through Hilbert Space guided by U(t)
- Operators are independent of time
- Basis kets (eigenstates of observables) eg:  $|x\rangle$  and  $|p\rangle$  are independent of time.

## Heisenberg Picture

- $|\psi(t)\rangle = |\psi\rangle_H$  is fixed and independent of time
- Operators in Heisenberg picture are time dependent

$$A_H(t) = U^{\dagger}(t)A_S U(t) = e^{iHt}A_S e^{-iHt}$$

# Time dependent perturbation theory

$$H(t) = H_0 + V(t)$$

where V(t) is small (this will always be the case for us)

$$H_0|n\rangle = E_n|n\rangle$$

Now, a general state at t=0

$$|\psi(0)\rangle = \sum_{n} c_n |n\rangle$$

For V = 0

$$|\psi(t)\rangle = \sum_{n} c_n e^{-iE_n t} |n\rangle$$

For  $V \neq 0$ 

$$|\psi(t)\rangle = \sum_{n} c_n(t)e^{-iE_nt}|n\rangle$$

where time dependence in  $c_n(t)$  due to only V.

### **Interaction Picture**

Define...

$$|\psi(t)\rangle_I \equiv e^{+iH_0t} |\psi(t)\rangle_S$$

$$A_I \equiv e^{+iH_0t} A_S e^{-iH_0t}$$

From these definitions, its clear that

$$_{I}\langle\psi(t)|A_{I}|\psi(t)\rangle_{I} =_{S} \langle\psi(t)|A_{S}|\psi(t)\rangle_{S}$$

When V=0, Heisenberg and Interaction Picture Coincide.

Ok, here's why we care about the interaction picture....

$$i\frac{d}{dt}|\psi(t)\rangle_{I} = i\frac{d}{dt}e^{iH_{0}t}|\psi(t)\rangle_{S} = e^{iH_{0}t}\left(-H_{0}|\psi(t)\rangle_{S} + i\frac{d}{dt}|\psi(t)\rangle_{S}\right)$$

$$= e^{iH_{0}t}\left(-H_{0} + (H_{0} + V_{S}))|\psi(t)\rangle_{S}$$

$$= e^{iH_{0}t}V_{S}e^{-iH_{0}t}e^{iH_{0}t}|\psi(t)\rangle_{S}$$

$$= V_{I}(t)|\psi(t)\rangle_{I}$$

The interaction picture is a hybrid of the Schrodinger and Heisenberg pictures.

Time evolution of state kets and operators depend on different parts of H.

$$|\psi(t)\rangle_S = \sum_n c_n(t)e^{-iE_nt}|n\rangle = e^{-iH_0t}\sum_n c_n(t)|n\rangle$$

$$|\psi(t)\rangle_I = e^{iH_0t} |\psi(t)\rangle_S = \sum_n c_n(t) |n\rangle$$

where,

$$c_n(t) = \langle n | \psi(t) \rangle_I$$

so once we have  $\langle \psi(t) \rangle_I$  we are done.

Solve the "Schrodinger Eq" iteratively

$$i\frac{d}{dt}|\psi(t)\rangle_I = V_I(t)|\psi(t)\rangle_I$$

integrate and get

$$|\psi(t)\rangle_I = |\psi(t_0)\rangle + \int_{t_0}^t dt' \left[ -iV_I(t') |\psi(t')\rangle \right]$$

where the second term is of order  $V_0$  which is small.

Now we keep iterating

$$= |\psi(t_0)\rangle - i \int_{t_0}^t dt' V_I(t') \left[ |\psi(t_0)\rangle - i \int_{t_0}^{t'} dt'' V_I(t'') |\psi(t'')\rangle \right]$$

at 3rd order (iterate again...)

$$= |\psi(t_0)\rangle - i \int_{t_0}^t dt' V_I(t') |\psi(t_0)\rangle + (-i)^2 \int_{t_0}^t dt' \int_{t_0}^{t'} dt'' V(t') V(t'') |\psi(t_0)\rangle$$

$$+ (-i)^3 \int \int \int V(t') V(t'') V(t''') |\psi(t_0)\rangle$$

$$|\psi(t)\rangle = U_I(t,t_0) |\psi(t_0)\rangle$$

where

$$U_I(t,t_0) = 1 + (-i) \int V(t') dt' + (-i)^2 \int \int V(t') V(t'') + \dots$$

"Dyson Series", can be written in slick form (by doing the sum)

$$U_I(t,t_0) = T \left[ e^{-i \int_{t_0}^t V_I(t')dt'} \right]$$

Where the "T" stands for a "time ordered product"