# Lecture 6

#### **Review Quantum Mechanics**

QM Linear Algebra in a complex vector space.

State of a system is a vector (ray) in the Complex Vector Space.

 $|\alpha\rangle$  - state vector

### **Linear Superpositions**

$$|\psi\rangle = c_1 |\alpha_1\rangle + c_2 |\alpha_2\rangle$$

If the  $|\alpha' s\rangle$  are vectors and the c's are complex numbers, then  $|\psi\rangle$  is another vector in the space.

#### **Dual Space:**

For every vector  $|\alpha\rangle$  ("ket") there is another vector  $\langle\alpha|$  ("bra") in a "dual" space.

Dual space is a mirro image of the ket space.

If,  $|\psi\rangle = c_1 |\alpha_1\rangle + c_2 |\alpha_2\rangle$ , then  $\langle \psi | = c_1^* \langle \alpha_1 | + c_2^* \langle \alpha_2 |$  where the  $c^*$  is the complex conjugate.

**Inner Product:** Dual space allows us to define an inner product between to vectors.

Given 2 vectors can get a "c#"

 $\langle \alpha | \beta \rangle$ 

# **Properties**

- 1.  $\langle \beta | \alpha \rangle = \langle \alpha | \beta \rangle^*$
- 2.  $\langle \alpha | \alpha \rangle \geq 0$
- 3.  $\langle \beta | (c_1 | \alpha_1 \rangle + c_2 | \alpha_2 \rangle) = c_1 \langle \beta | \alpha_1 \rangle + c_2 \langle \beta | \alpha_2 \rangle$

 $|\alpha\rangle$  and  $|\beta\rangle$  are orthogonal if  $\langle\alpha|\beta\rangle = 0$ .

States can be normalized  $\langle \alpha | \alpha \rangle = 1$ .

Operators: Thigs that act on a given state, and return another state

$$X|\alpha\rangle = |\alpha'\rangle$$

Product:

$$(YX)|\alpha\rangle = Y(X|\alpha\rangle$$

In general, not commutative:  $XY \neq YX$ , but they are associative.

$$\langle \alpha | = \langle \alpha' | X^{\dagger}$$

In general,  $X^{\dagger} \neq X$ , if so X is said to be "Hermitian".

System characheterized by single observable A

eg: position / Momentum / energy

Measuremnet of A gives possible values

$$a_1, a_2, ...$$

 $|\alpha\rangle$  = State for which A has a value of a

$$\sum_{\alpha'} |\alpha'\rangle\langle\alpha'| = 1$$

$$\langle \alpha | \alpha' \rangle = \delta_{aa'}$$

Any physical observable corrisponds to an operator like

$$A = \sum_{\alpha'} \alpha' |\alpha'\rangle \langle \alpha'|$$

eg:

$$A |\alpha\rangle = \left(\sum_{a'} a' |\alpha'\rangle\langle\alpha'|\right) |\alpha\rangle = a |\alpha\rangle$$

Physical observables are real numbers, therefore physical operators A are <u>Hermitian</u> proof:

$$A^{\dagger} = \sum_{a} a^{*} |\alpha\rangle\langle\alpha| = \sum_{a} a |\alpha\rangle\langle\alpha| = A$$

#### **Probabilities:**

Condsider a filter  $M(a) = |a\rangle\langle a|$  on a general state  $|s\rangle$ .

$$M(a)|s\rangle = |a\rangle\langle a|s\rangle$$
  
=  $\langle a|s\rangle|a\rangle$ 

where  $\langle a|s\rangle$  is a c# that tells you something about what fraction of the time you get through.

 $\langle a|s\rangle$  is related to the pass fraction

\*But, a) not real, b) not normalized

However, we know that  $|\langle a|s\rangle|^2 = \langle a|s\rangle\langle s|a\rangle = \langle s|a\rangle\langle a|s\rangle$  is both real and normalized.

$$\sum_{a} \langle s|a\rangle \langle a|s\rangle = \langle s|s\rangle = 1$$

### Interpretation

 $|\langle a|s\rangle|^2$  = Probability that a system prepared in state  $|s\rangle$  will be found in a state  $|a\rangle$  with value a for observable A after measurement.

Comments on the measurement problem....

# **Position Operator**

$$\vec{X} | \vec{x} \rangle = \vec{x} | \vec{x} \rangle$$

where,  $\vec{X}$  is position operator and  $\vec{x}$  is position eigenvalue.

Position of course is a continous observable so "sums go to intergals" etc. eg: (Completnees and Orthogonality)

$$\sum_{\alpha'} |\alpha'\rangle\langle\alpha'| = 1 \Rightarrow \int d^3x |x\rangle\langle x| = 1$$

$$\langle \alpha' | \alpha' \rangle = \delta_{\alpha,\alpha'} \Rightarrow \langle \vec{x} | \vec{x'} \rangle = \delta^3 (\vec{x} - \vec{x'})$$

#### Wave function

$$|\psi\rangle = \int d^3x |x\rangle \langle x|\psi\rangle = \int d^3x \psi(x) |x\rangle$$

where  $\psi(x) = \langle x | \psi \rangle$  is called the Position-space wavefunction.

$$\langle \psi | \psi \rangle = 1 \Rightarrow \int d^3x \psi(x) \langle \psi | x \rangle = \int d^3x \psi^*(x) \psi(x) = 1$$

# Interpretation

 $|\psi(x)|^2 d^3x$  is the probability to find the particle in volume  $d^3x$  around  $\vec{x}$ .

**Translation Operator** "The operator that moves you over"

$$T(\vec{a})|\vec{x}\rangle = |\vec{x} + \vec{a}\rangle$$

What is  $T^{\dagger}(\vec{a})$ ???

Well,

$$\langle \vec{x'} | \left( T(\vec{a}) \, | \vec{x} \rangle \right) = \delta^3((\vec{x} + \vec{a}) - \vec{x'})$$

or

$$\left(\langle \vec{x'}|T(\vec{a})\right)|\vec{x}\rangle = \delta^3(\vec{x} - (\vec{x'} - \vec{a}))$$

$$\Rightarrow \langle x' | T(a) = \langle x - a |$$

which says that  $T^{\dagger}(a) | \vec{x} \rangle = | \vec{x} - \vec{a} \rangle$ 

So,

$$T^{\dagger}(\vec{a}) = T(-\vec{a}) = T^{-1}(\vec{a})$$

Properties of T

- 1. Unitary  $T^{\dagger}T = 1$
- 2.  $T(\vec{a})T(\vec{b}) = T(\vec{a} + \vec{b}) = T(\vec{b})T(\vec{a})$ , Translations commute  $([T(\vec{a}), T(\vec{b})] = 0)$
- 3. T(0) = 1

# **Infintesimal Translations**

consider  $\vec{a} = N\vec{\epsilon}$