

1) Particles in SM:

+4

fermions:  $\begin{pmatrix} \nu_e \\ e \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}$

← all spin 1/2  
+ antiparticles

$\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$

← 3 colors of each,  
all spin 1/2

bosons:  $Z, W^{\pm}, \gamma, 8 \text{ gluons}, H$  ← all spin-1  
H ← spin 0

2) The coupling constant w/ a factor  $\alpha$  is different for each force. Also, weak force bosons are massive.

+1

3)  $\text{Br}(W^{\pm} \rightarrow l\nu\nu)$

Strong force?

$$= \{e, \bar{\nu}_e, \nu_e, \bar{e}\} + \{\mu, \bar{\nu}_\mu, \nu_\mu, \bar{\mu}\} + \{\tau, \bar{\nu}_\tau, \nu_\tau, \bar{\tau}\}$$

$$= \text{Br}(W \rightarrow e\nu) + \text{Br}(W \rightarrow \mu\nu) + \text{Br}(W \rightarrow \tau\nu)$$

and  $\text{Br}(Z \rightarrow \nu\nu)$  and  $\text{Br}(Z \rightarrow e\bar{e})$

$\frac{1}{3 + 2 \cdot 3}$

no color

re sorry, see  
next page

$$3) \text{Br}(\omega Z \rightarrow (e/\mu + \nu\nu))$$

$$= \text{Br}(\omega \rightarrow e/\mu + \nu) \cdot \text{Br}(Z \rightarrow \nu\nu)$$

$$\left( \frac{\cancel{8} + \{e, \bar{\nu}_e\} + \{\mu, \bar{\nu}_\mu\}}{3 \text{ leptons} + 3 \cdot 2 \text{ quarks}} \right) \cdot \left( \frac{3 \text{ neutrinos}}{6 \text{ leptons} + 3 \text{ quarks} = 5 \text{ quarks}} \right)$$

$$= \left( \frac{2}{9} \cdot \frac{3}{21} \right)$$

$$= \frac{2}{3} \cdot \frac{1}{21} = \frac{2}{63}$$

+ 6

# 4) Higgs Decay:

+ 5

$WW$  ( $Br \sim 20\%$ )

$ZZ$  ( $Br \sim 3\%$ )

$\gamma\gamma$  ( $Br \sim 0.2\%$ )

boxon decays to  $e$  or  $\mu^-$

$$Br(WW \rightarrow l\nu l'\nu') = \left( \frac{\text{either } e \text{ or } \mu^-}{(3\nu) + (3\nu + 2 \text{ neutrinos of gluons})} \right)^2 = \left( \frac{2}{9} \right)^2 = \boxed{\frac{4}{81} = Br(WW \rightarrow l\nu l'\nu')}$$

no top gluon decays

$$Br(ZZ \rightarrow ll' ll') = \left( \frac{\text{either } e\bar{e}^+ \text{ or } \mu\bar{\mu}^+}{6 \text{ leptons} + 3 \times 5 \text{ possible gluons}} \right)^2 = \left( \frac{2}{21} \right)^2$$

$$\Rightarrow \boxed{Br(ZZ \rightarrow ll' ll') = \left( \frac{2}{21} \right)^2} \sim \frac{1}{100}$$

w/ these factors

$$Br(H \rightarrow WW \rightarrow l\nu l'\nu') \text{ becomes } \frac{1}{5} \times \frac{4}{81} \sim \frac{1}{100}$$

$$Br(H \rightarrow ZZ \rightarrow ll' ll') \text{ becomes } \sim \frac{1}{100} \times \frac{9}{100} = \frac{3}{10000}$$

$$Br(H \rightarrow \gamma\gamma) = 0.2\% = \frac{1}{1000}$$

$\Rightarrow$  rank

(1)	$WW$	$\leftarrow$ most likely
(2)	$\gamma\gamma$	
(3)	$ZZ$	$\leftarrow$ least

## 5) accelerators

(a) the ~~size~~ available power (to create a magnetic field) limits the energy of circular proton accelerators. Also, there will be more synchrotron radiation.

(b) limits to the <sup>linear</sup> electron accelerators are also synchrotron radiation, but this will become more apparent since ~~they need more energy to be accelerated to the same energy~~ (since they will go faster than the protons at some energy (since they have less mass)).

+4

## 6) $e^-e^+$ collisions

(a)  $R(E_{cm}) \equiv \frac{\sigma(ee \rightarrow \text{etc})}{\sigma(ee \rightarrow \mu\mu)}$

$E_{cm}$  must be  $\geq 2m_c$

$$\sigma \sim \frac{1}{E_1 E_2} |M|^2 d\Omega$$

as you go above  $2m_c$ , you will expect more jets to be made, so we expect

$$R(m < 2m_c) \sim \text{goes to } 0$$

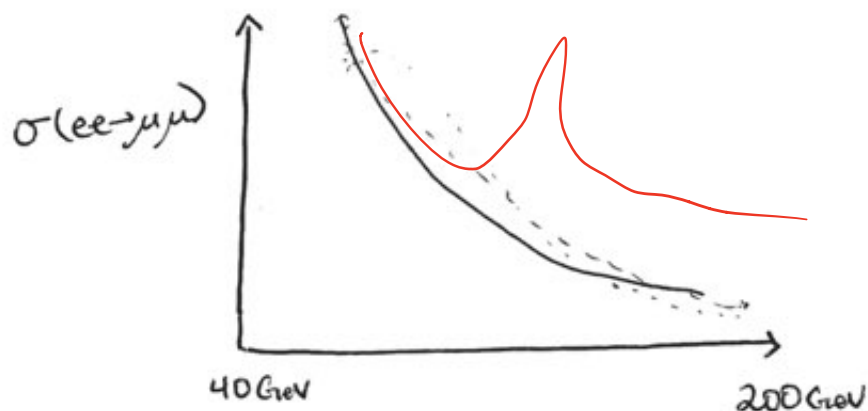
$$R(m > 2m_c) \sim \text{goes to } \infty$$

See solutions

(over)

6 cont'd)

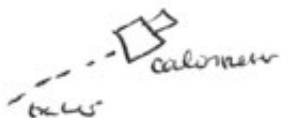
(b) cross section of  $ee \rightarrow \mu\mu$ , fun of  $E_{cm}$



7) Collider Detectors

(a) ~~electrons~~  $\mu^-$ 's behave like heavy  $e^-$ 's as they are 2nd generation, but they have a ~~longer lifetime~~ than  $e^-$  and will go farther in detector.

(b) photons have no charge and do not "curve" in tracker  
lon line



$e^-$ 's do appear in tracker



8) Hadronic showers are more difficult since there are more complex reactions going on than just  $\gamma \rightarrow e^-e^+$  or  $e^- \rightarrow e^- \gamma$  as seen in EM showers.

+ 4



9)  $m_x \sim 2 \text{TeV} = 2000 \text{GeV}$

$x \rightarrow e^+e^-$  or  $x \rightarrow \mu^+\mu^-$ ?

+ 3

Expect to measure mass more precisely w/  $x \rightarrow e^+e^-$  since the  $e^+e^-$ s are stopped in the calorimeters. This gives an accurate reading to their energies which can be used to select for the  $m_x$  4-vectors.

10)  $V_s$  are detected by comparing what was measured in the beginning vs at the end. If their energies are not equal, then  $V_s$  were created.

+ 2

11) (a) decay  
 $H \text{ --- } \langle \begin{matrix} \text{ring} \\ \text{ring} \end{matrix} \rangle$

production:



+ 3

(b) Higgs show up in the "looped" Feynman diagrams which appear less often than the  $W$  or  $Z$  bosons.



12) Higgs decays to pairs of leptons:

$$H \rightarrow (ll)$$

+1

The decay mode of  $H \rightarrow ee^+$  would be best to study couplings to Higgs fields. The  $ee^+$  can easily be studied (w/ accurate reading of energy) @ the LHC. The coupling could then easily be found.

13) (a) the group symmetry we input (e.g.  $SU(2)_L \times U(1)$ ) +5  
dictates the particle content [i.e. which particles satisfy this symmetry]

(b) electroweak symmetry groups are  
 $SU(2)_L$  [left-handed] and  $U(1)$

(c) as a result, only left handed ~~particles~~  $\nu_L$  obey the symmetry. In other words,  $\nu_R$  may as well not exist as they do not interact in the standard model.

(14)

- (a) Spontaneous symmetry breaking is when you have a 'displaced' minima about the 0 in your field ~~there~~. A minima is randomly chosen, thus spontaneously breaking symmetry in order to give a non-negative mass to the particle.
- (b) SSB is responsible for giving mass to particles and is closely related to the Higgs field.
- (c) The discovery of the Higgs gives evidence that SSB is actually occurring.

X

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