

3D

Ch 7

Generalization to 3D straight forward

$$i\hbar \frac{\partial \psi}{\partial t} = H \psi$$

H - total energy
"Hamiltonian"

3D

$$H \rightarrow \frac{1}{2m} (p_x^2 + p_y^2 + p_z^2) + V(\vec{x})$$

$$p_x \rightarrow \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\vec{p} \rightarrow \frac{\hbar}{i} \vec{\nabla}$$

Sch Eq

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V \psi$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Note ψ & V now functions of \vec{x} & t

Just as before, we will consider solutions

$$\Psi(\vec{x}, t) = \psi(\vec{x}) e^{-iEt/\hbar}$$

where $\psi(x)$ satisfies

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi = E\psi$$

Questions

Infinite Cube

$$V(x, y, z) = \begin{cases} 0 & \text{if } x, y, z \text{ in } 0-a \\ \infty & \text{otherwise} \end{cases}$$

Outside of cube $\psi(x, y, z) = 0$
(Same logic as in 1D)

Inside the cube

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) = E \psi$$

Assume separation of variables

$$\psi(x, y, z) = X(x) Y(y) Z(z)$$

then (substitute & dividing by ψ)

$$\underbrace{\frac{1}{X} \frac{d^2 X}{dx^2}}_{f(x)} + \underbrace{\frac{1}{Y} \frac{d^2 Y}{dy^2}}_{f(y)} + \underbrace{\frac{1}{Z} \frac{d^2 Z}{dz^2}}_{f(z)} = -\frac{2m}{\hbar^2} E$$

constant

Only possible if the different $f()$'s are equal to constants (k_x^2, k_y^2, k_z^2)

$$\text{So, } \frac{d^2 X}{dx^2} = -k_x^2, \quad \frac{d^2 Y}{dy^2} = -k_y^2, \quad \frac{d^2 Z}{dz^2} = -k_z^2$$

$$\text{And } E = \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2)$$

$$\rightarrow (p = \hbar k, E = \frac{p_x^2 + p_y^2 + p_z^2}{2m})$$

As expected...

Now easy to solve the separated eqn's

$$X(x) = A_x \sin k_x x + B_x \cos k_x x$$



$$k_x = \frac{n_x \pi}{a} \quad (\text{from } \psi\text{-cont at } a)$$

Also $Y(y), Z(z)$

$$\psi(x, y, z) = A \sin\left(\frac{n_x \pi}{a}\right) \sin\left(\frac{n_y \pi}{a}\right) \sin\left(\frac{n_z \pi}{a}\right)$$

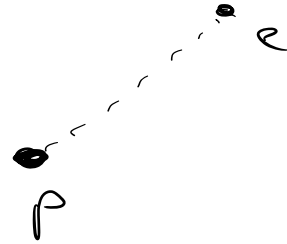
So $\psi_{n_x, n_y, n_z}(x, y, z) \quad E_{n_x, n_y, n_z} = \frac{\hbar^2 \pi^2}{2mL^2} (n_x^2 + n_y^2 + n_z^2)$

Atoms

However most physical problems in 3D
have $V(\vec{r})$

Appropriate to use

r, θ, ϕ not x, y, z .



All the complications (fancy math we'll need)
arises from going to spherical coordinates
not inherent in QM or in 3D Sch.

I will give a high-level view of
what's going on and how I think
about it.

Skip most of the math (all of the
diff eq's)

What I want you to pay attention to:

-) Where the separation constants come from and how they are constrained.
-) Physical meaning of separation constants
-) Possible (Allowed) Quantum #'s (expect 3)
-) Qualitative understanding of ψ

Ok, need to change basis

$$(x, y, z) \rightarrow (r, \theta, \phi)$$

$$\frac{\partial^2 \psi}{\partial x^2} \rightarrow \frac{\partial^2 \psi}{\partial r^2} \frac{\partial r}{\partial x} + \frac{\partial^2 \psi}{\partial \theta^2} \frac{\partial \theta}{\partial x} + \frac{\partial^2 \psi}{\partial \phi^2} \frac{\partial \phi}{\partial x}$$

+ others (Then need to "square" $\frac{\partial^2 \psi}{\partial x^2}$)

Ton of algebra, not enlightening (you do in ch 3)

Answer:

$$\frac{\nabla^2 \psi}{2x^2} + \frac{\nabla^2 \psi}{2y^2} + \frac{\nabla^2 \psi}{2z^2} \longrightarrow$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial^2 \psi}{\partial \phi^2} \right)$$

Start by looking for separable solutions

$$\psi(r, \theta, \phi) = R(r) Y(\theta, \phi)$$

Putting this in and dividing by ψ (like always)

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E]$$

Only depends on r

$$+ \frac{1}{Y} \left\{ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial Y}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 Y}{\partial \phi^2} \right\} = 0$$

Only cares about θ, ϕ

First Separation Constant C_1

Note the equation for $\theta \neq \phi$, V independent

\Rightarrow Angular Solutions universal (like $\phi(t)$)

Angular E_ϕ well known "Laplace's E_ϕ "

It itself is separable

$$Y(\theta, \phi) = T(\theta) P(\phi)$$

Second Separation Constant. C_2

$$\frac{1}{P} \frac{d^2 P}{d\phi^2} = -C_2$$

Easy Posing $P(\phi) = e^{i\sqrt{C_2}\phi}$

Remember ϕ periodic $\phi = \phi + 2\pi$

$$\Rightarrow e^{i\sqrt{C_2}2\pi} = 1$$

S_0 , $\sqrt{C_2}$ must be a real integer

$$C_2 \equiv m^2 \quad \text{w/ } m = 0, \pm 1, \pm 2, \dots$$

Ok, one down. Now the Θ eq

Turns out this is a special diff eq
that 19th century math figured out

Solution

$$T(\theta) = A P_l^m(\cos \theta)$$

where $C_1 = l(l+1)$ and l integer

$$\text{where } P_l^m(x) \equiv (1-x^2)^{|m|/2} \left(\frac{d}{dx} \right)^{|m|} P_l(x)$$

$$\text{and } P_l \equiv \frac{1}{2^l l!} \left(\frac{d}{dx} \right)^l (x^2-1)^l$$

where l is $0, 1, 2, \dots$

(Note: you will not need these details)

Whats important:

P_l - polynomial in x of degree l .

$$\Rightarrow P_m^l \sim \left(\frac{d}{dx} \right)^{|m|} O(x^l) = 0$$

if $|m| > l$

This is the
take away

$$\Rightarrow l: 0, 1, 2, \dots$$

$$m: -l, -l+1, \dots, -1, 0, 1, \dots, l-1, l$$

Back to the Radial eq

Can't go further until we specify $V(r)$

Here we will only study hydrogen potential.

Most other atoms can only be solved numerically

$$V(r) \sim -\frac{\alpha}{r}$$

Again you end up w/ complicated diff eq.

Note that we expect:

-) Quantized states given by n , E_n

-) Link between n & l

Solutions $R_{nl}(r) = A e^{-r/a_0 n} r^l L_n^l\left(\frac{r}{a_0 n}\right)$

"Laguerre Polynomials"

($n > l$ from same logic as $|m| < l$, $R=0$ otherwise)

Turns out

$$E_n = - \left(\frac{kZe^2}{h} \right)^2 \frac{m}{2n^2} = - \frac{Z^2 E_1}{n^2}$$

$$E_1 = 13.6 \text{ eV}$$

Same energy levels as those in Bohr model!!

- All upside of getting E levels right
- Now have a solid theory behind it.

$$n = 1, 2, 3, \dots \quad \text{But } n > 0$$

Summary of Quantum Numbers

3 as expected

$$n = 1, 2, 3, \dots \quad (r)$$

$$l = 0, 1, 2, \dots, (n-1) \quad (\theta)$$

$$m = -l, (-l+1), \dots, 0, 1, \dots, l \quad (\phi)$$

Note E_n - only depends on n .
(special to $\frac{1}{n}$ law)

- 1
- 2
- 3

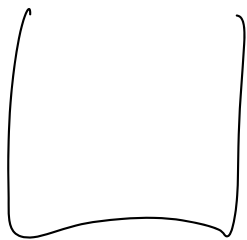
P shell

boron 5 B 10.811	carbon 6 C 12.011	nitrogen 7 N 14.007	oxygen 8 O 15.999	fluorine 9 F 18.998	neon 10 Ne 20.180
aluminum 13 Al 26.982	silicon 14 Si 28.086	phosphorus 15 P 30.974	sulfur 16 S 32.065	chlorine 17 Cl 35.453	argon 18 Ar 39.948

lanthanum 57 La 138.91	cerium 58 Ce 140.12	praseodymium 59 Pr 140.91	neodymium 60 Nd 144.24	promethium 61 Pm [145]	samarium 62 Sm 151.96	europium 63 Eu 157.25	gadolinium 64 Gd 158.93	terbium 65 Tb 162.50	dysprosium 66 Dy 164.93	holmium 67 Ho 167.26	erbium 68 Er 168.93	thulium 69 Tm 168.93	ytterbium 70 Yb 173.04
actinium 89 Ac [227]	thorium 90 Th 232.04	protactinium 91 Pa 231.04	uranium 92 U 238.03	neptunium 93 Np [237]	plutonium 94 Pu [244]	americium 95 Am [243]	curium 96 Cm [247]	berkelium 97 Bk [247]	californium 98 Cf [251]	einsteinium 99 Es [252]	fermium 100 Fm [257]	mendelevium 101 Md [258]	nobelium 102 No [259]

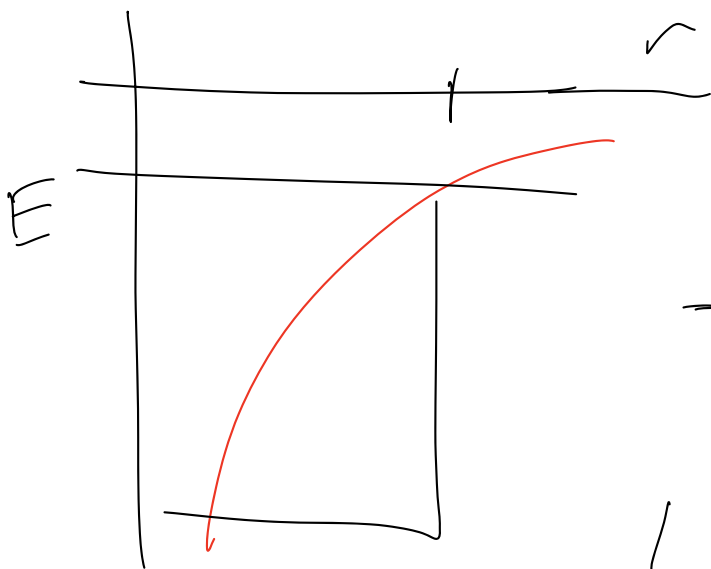
eg) similar elements in same column.

Chem. of C, Si + Ge similar



L

$$E = \frac{\hbar^2 \pi^2 n^2}{2mL^2}$$



$$E \Rightarrow \alpha$$

$$L = \frac{\alpha}{E}$$

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2m\alpha^2} E_n^2$$

$$r_e \sim \frac{1}{m\alpha}$$

$$\Rightarrow \frac{2m\alpha^2}{\hbar^2 \pi^2} \frac{1}{n^2}$$

$$E \sim m\alpha^2$$