

PHYS440 - Exercise on basis states

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2.3. Given basis states $\{|0\rangle, |1\rangle\}$ for measurements on the Z axis and $\{|+\rangle, |-\rangle\}$ for measurements on the X axis find suitable basis states for measurements on the Y axis.

We'll follow the outline provided in exercise 2.3 in Susskind and Friedman[1].

We start with the states $|0\rangle$ and $|1\rangle$ for the Z axis.

For the X axis, we have:

$$|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$
$$|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

Rearranging, we also have:

$$|0\rangle = \frac{1}{\sqrt{2}}(|+\rangle + |-\rangle)$$
$$|1\rangle = \frac{1}{\sqrt{2}}(|+\rangle - |-\rangle)$$

And also these facts:

$$\begin{aligned}
\langle 0|0\rangle &= \langle 1|1\rangle = 1 \\
\langle 0|1\rangle &= \langle 1|0\rangle = 0 \\
\langle +|+\rangle &= \langle -|-\rangle = 1 \\
\langle +|-\rangle &= \langle -|+\rangle = 0 \\
|\langle 0|+\rangle|^2 &= \langle 0|+\rangle \langle +|0\rangle = \frac{1}{2} \\
|\langle 0|-\rangle|^2 &= \langle 0|-\rangle \langle -|0\rangle = \frac{1}{2} \\
|\langle 1|+\rangle|^2 &= \langle 1|+\rangle \langle +|1\rangle = \frac{1}{2} \\
|\langle 1|-\rangle|^2 &= \langle 1|-\rangle \langle -|1\rangle = \frac{1}{2}
\end{aligned}$$

We will denote the basis states for the Y axis as $|a\rangle$ and $|b\rangle$ for now, and we will define them as a superposition of the Z axis basis states:

$$\begin{aligned}
|a\rangle &= \alpha|0\rangle + \beta|1\rangle \\
|b\rangle &= \gamma|0\rangle + \delta|1\rangle
\end{aligned}$$

where

$$\begin{aligned}
\alpha &= \langle 0|a\rangle & \bar{\alpha} &= \langle a|0\rangle \\
\beta &= \langle 1|a\rangle & \bar{\beta} &= \langle a|1\rangle \\
\gamma &= \langle 0|b\rangle & \bar{\gamma} &= \langle b|0\rangle \\
\delta &= \langle 1|b\rangle & \bar{\delta} &= \langle b|1\rangle
\end{aligned}$$

For symmetry between the Z and Y axes, we have these relationships:

$$\begin{aligned}
|\langle 0|a\rangle|^2 &= \langle 0|a\rangle \langle a|0\rangle = \alpha\bar{\alpha} = \frac{1}{2} \\
|\langle 1|a\rangle|^2 &= \langle 1|a\rangle \langle a|1\rangle = \beta\bar{\beta} = \frac{1}{2} \\
|\langle 0|b\rangle|^2 &= \langle 0|b\rangle \langle b|0\rangle = \gamma\bar{\gamma} = \frac{1}{2} \\
|\langle 1|b\rangle|^2 &= \langle 1|b\rangle \langle b|1\rangle = \delta\bar{\delta} = \frac{1}{2}
\end{aligned}$$

We also have similar relationships between the X and Y axes, in terms of $|+\rangle$

and $|-\rangle$, which we can rewrite in terms of $|0\rangle$ and $|1\rangle$. We can use $\langle +|a\rangle$ as follows:

$$\begin{aligned}
 |\langle +|a\rangle|^2 &= \langle +|a\rangle \langle a|+\rangle = \frac{1}{2} \\
 \frac{1}{\sqrt{2}}(\langle 0|a\rangle + \langle 1|a\rangle) \frac{1}{\sqrt{2}}(\langle a|0\rangle + \langle a|1\rangle) &= \frac{1}{2} \\
 \frac{1}{2}(\langle 0|a\rangle \langle a|0\rangle + \langle 0|a\rangle \langle a|1\rangle + \langle 1|a\rangle \langle a|0\rangle + \langle 1|a\rangle \langle a|1\rangle) &= \frac{1}{2} \\
 \frac{1}{2} + \langle 0|a\rangle \langle a|1\rangle + \langle 1|a\rangle \langle a|0\rangle + \frac{1}{2} &= 1 \\
 \langle 0|a\rangle \langle a|1\rangle + \langle 1|a\rangle \langle a|0\rangle &= 0 \\
 \alpha\bar{\beta} + \beta\bar{\alpha} &= 0 \\
 \alpha\bar{\beta} &= -(\alpha\bar{\beta})
 \end{aligned}$$

For $z = -\bar{z}$, we have z purely imaginary, so $\alpha\bar{\beta}$ is purely imaginary, and α and β cannot both be real.

In the same way, we can find that $\gamma\bar{\delta}$ is purely imaginary, and so γ and δ cannot both be real.

We can set

$$\begin{aligned}
 \alpha &= \frac{1}{\sqrt{2}} \\
 \beta &= \frac{i}{\sqrt{2}} \\
 \gamma &= \frac{1}{\sqrt{2}} \\
 \delta &= -\frac{i}{\sqrt{2}}
 \end{aligned}$$

Then we have

$$\begin{aligned}
 |a\rangle &= \frac{1}{\sqrt{2}}|0\rangle + \frac{i}{\sqrt{2}}|1\rangle \\
 |b\rangle &= \frac{1}{\sqrt{2}}|0\rangle - \frac{i}{\sqrt{2}}|1\rangle
 \end{aligned}$$

Checking this, we have

$$\begin{aligned}
 \langle a|a \rangle &= \left(\frac{1}{\sqrt{2}} \langle 0| - \frac{i}{\sqrt{2}} \langle 1| \right) \left(\frac{1}{\sqrt{2}} |0\rangle + \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle + \frac{i}{2} \langle 0|1 \rangle - \frac{i}{2} \langle 1|0 \rangle + \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} + 0 - 0 + \frac{1}{2} \\
 &= 1
 \end{aligned}$$

$$\begin{aligned}
 \langle a|b \rangle &= \left(\frac{1}{\sqrt{2}} \langle 0| - \frac{i}{\sqrt{2}} \langle 1| \right) \left(\frac{1}{\sqrt{2}} |0\rangle - \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle - \frac{i}{2} \langle 0|1 \rangle - \frac{i}{2} \langle 1|0 \rangle - \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} - 0 - 0 - \frac{1}{2} \\
 &= 0
 \end{aligned}$$

$$\begin{aligned}
 \langle b|b \rangle &= \left(\frac{1}{\sqrt{2}} \langle 0| + \frac{i}{\sqrt{2}} \langle 1| \right) \left(\frac{1}{\sqrt{2}} |0\rangle - \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle - \frac{i}{2} \langle 0|1 \rangle + \frac{i}{2} \langle 1|0 \rangle + \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} - 0 + 0 + \frac{1}{2} \\
 &= 1
 \end{aligned}$$

$$\begin{aligned}
 \langle b|a \rangle &= \left(\frac{1}{\sqrt{2}} \langle 0| + \frac{i}{\sqrt{2}} \langle 1| \right) \left(\frac{1}{\sqrt{2}} |0\rangle + \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle + \frac{i}{2} \langle 0|1 \rangle + \frac{i}{2} \langle 1|0 \rangle - \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} + 0 + 0 - \frac{1}{2} \\
 &= 0
 \end{aligned}$$

We have shown that the basis states for the Y axis require complex numbers. But the particular $|a\rangle$ and $|b\rangle$ we have chosen are not unique. We can introduce an arbitrary phase factor $e^{i\theta}$ ($\theta \in \mathbb{R}$) to $|a\rangle$ and $|b\rangle$, and they will still be valid basis states for the Y axis.

Let

$$\begin{aligned}
 |c\rangle &= e^{i\theta} |a\rangle = \frac{e^{i\theta}}{\sqrt{2}} (|0\rangle + i|1\rangle) \\
 |d\rangle &= e^{i\theta} |b\rangle = \frac{e^{i\theta}}{\sqrt{2}} (|0\rangle - i|1\rangle)
 \end{aligned}$$

The magnitudes of the amplitudes are unchanged, so the Born rule requirement is still met.

We can also check the inner products of these states:

$$\begin{aligned}
 \langle c|c \rangle &= \frac{e^{-i\theta}}{\sqrt{2}} \left(\langle 0| - \frac{i}{\sqrt{2}} \langle 1| \right) \frac{e^{i\theta}}{\sqrt{2}} \left(|0\rangle + \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle + \frac{i}{2} \langle 0|1 \rangle - \frac{i}{2} \langle 1|0 \rangle + \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} + 0 - 0 + \frac{1}{2} \\
 &= 1
 \end{aligned}$$

$$\begin{aligned}
 \langle c|d \rangle &= \frac{e^{-i\theta}}{\sqrt{2}} \left(\langle 0| - \frac{i}{\sqrt{2}} \langle 1| \right) \frac{e^{i\theta}}{\sqrt{2}} \left(|0\rangle - \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle - \frac{i}{2} \langle 0|1 \rangle - \frac{i}{2} \langle 1|0 \rangle - \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} - 0 - 0 - \frac{1}{2} \\
 &= 0
 \end{aligned}$$

$$\begin{aligned}
 \langle d|d \rangle &= \frac{e^{-i\theta}}{\sqrt{2}} \left(\langle 0| + \frac{i}{\sqrt{2}} \langle 1| \right) \frac{e^{i\theta}}{\sqrt{2}} \left(|0\rangle - \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle - \frac{i}{2} \langle 0|1 \rangle + \frac{i}{2} \langle 1|0 \rangle + \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} - 0 + 0 + \frac{1}{2} \\
 &= 1
 \end{aligned}$$

$$\begin{aligned}
 \langle d|c \rangle &= \frac{e^{-i\theta}}{\sqrt{2}} \left(\langle 0| + \frac{i}{\sqrt{2}} \langle 1| \right) \frac{e^{i\theta}}{\sqrt{2}} \left(|0\rangle + \frac{i}{\sqrt{2}} |1\rangle \right) \\
 &= \frac{1}{2} \langle 0|0 \rangle + \frac{i}{2} \langle 0|1 \rangle + \frac{i}{2} \langle 1|0 \rangle - \frac{1}{2} \langle 1|1 \rangle \\
 &= \frac{1}{2} + 0 + 0 - \frac{1}{2} \\
 &= 0
 \end{aligned}$$

The $|a\rangle$ and $|b\rangle$ above are the simplest of infinitely many possible choices. Following the nice convention in Wong[2], let's relabel these:

$$\begin{aligned}
 |i\rangle &= |a\rangle = \frac{1}{\sqrt{2}} |0\rangle + \frac{i}{\sqrt{2}} |1\rangle \\
 |-i\rangle &= |b\rangle = \frac{1}{\sqrt{2}} |0\rangle - \frac{i}{\sqrt{2}} |1\rangle
 \end{aligned}$$

References

- [1] Art Friedman Leonard Susskind. *Quantum Mechanics: The Theoretical Minimum*. 1st ed. Allen Lane, 2014. ISBN: 978-0-241-00344-1.
- [2] Thomas G. Wong. *Introduction to Classical and Quantum Computing*. 1st ed. Rooted Grove, 2022. ISBN: 979-8-9855931-0-5.