

## Reading Notes

### Intro

- Oscillatory phenomena can be found in a lot of different areas of physics, and one common case is a pendulum.
- Pendulums exhibit a harmonic force.

### 3.1

- A simple pendulum can be connected to a mass on a rigid support, where the parallel forces all add up to zero.
- We then get the equation:
- $F(\theta) = -mg\sin(\theta)$
- $g$  = acceleration due to gravity
- this is negative because force is opposite of the displacement
- With Newton's second law representing  $F = ma$ , we can assume that the second derivative of  $\theta$  ( $d^2\theta/dt^2 = -(g/l)\theta$ ), which is the equation of simple harmonic motion.
- We can also take a numerical approach to this problem; with the second-order DE that we made, it can be rewritten into two first-order DE's:
- $d\omega/dt = -(g/l)\theta$
- $d\theta/dt = \omega$
- There is also a pseudo code that uses Euler's method to calculate values of  $\theta$  and  $\omega$ .
- The total energy equation can be defined as:
- $E = \frac{1}{2}ml^2\omega^2 + mgl(1 - \cos(\theta))$
- $KE = \frac{1}{2}mv^2$

### 3.2

- We saw before that the equation of motion uses a frictionless pendulum.
- They incorporate a frictional force  $-q(d\theta/dt)$ , resulting in our equation of motion to be:
- $(d^2\theta/dt^2 = -(g/l)\theta - q d\theta/dt)$
- The pendulum can be listed in three separate variants: underdamped, overdamped, and critically damped.
- Resonance occurs when our amplitude ( $\theta_0$ ) can be large even with a small friction.
- With the equation of motion listed above in this section, this has a nonlinear pendulum without friction and a driving force. (according to Eq. 3.17)

- This means that there is no extra adding/removing energy to our system, and that the period cannot be independent of the amplitude.

### 3.3

- When we incorporate the sinusoidal driving force  $F_d \sin(\Omega_d t)$ , friction of the form  $-q(d(\theta)/dt)$ , and no expansion of  $\sin(\theta)$ , we get the following equation of motion, which is for a **nonlinear damped driven pendulum**:

- $d^2(\theta)/dt^2 = -(g/l)*\theta - q d(\theta)/dt + F_d \sin(\Omega_d t)$
- In order to understand how  $\theta$  acts as a function of time in this equation, we need to make a program to generate a numerical solution for this.

- By rewriting this into 2 differential equations:

- $d(\omega)/dt = -(g/l)*\theta - q d(\theta)/dt + F_d \sin(\Omega_d t)$
- $d(\theta)/dt = \omega$

- There are a lot of different  $\theta$  vs. time graphs in Figure 3.6 that all show different behaviors for  $\theta$ .

- The first one shows vertical leaps in  $\theta$  when it is reset within the range of  $\pm \pi$ .

- the second behavior of  $\theta(t)$  for  $F_d = 1.2$  has activity within a couple of oscillations for the decay. After this, the pendulum settles into a steady oscillation.

- The third graph shows behavior that corresponds to the angular velocity of the pendulum.

- When we have a larger value of  $F_d$ , we see that the changes in  $\theta$  increase rapidly and irregularly with  $t$ . This can be indicated with a logarithmic scale to denote how fast the change in  $\theta$  can be with time, and this also corresponds to a log relation which is similar to a value of  $\text{Lyapunov} * \text{time}$ :

- it implies that the change in  $\theta$  is similar to  $e^{\text{Lyapunov} * \text{time}}$ .
- When there is a small driving force in a trajectory of  $\omega$ , the phase space gets easier to understand due to how the pendulum can react. On Figure 3.8, we see two separate graphs that show the behavior of the absolute value of  $\theta = \pi$  as well as how  $\omega$  can behave as a function of  $\theta$  for a pendulum.