Steinmann Relations and the Two-Loop MHV Amplitude in Eight-Particle Kinematics

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ABSTRACT: We present the full functional form of the two-loop eight-point MHV amplitude in the planar limit of maximally supersymmetric Yang-Mills theory, in terms of cluster polylogarithms. We also compute the two BDS-like ansätze that can be formulated in eight-particle kinematics, and find that ...

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1 Introduction

2 Promoting $R_8^{(2)}$ from Symbol to Function

2.1 Cluster Polylogarithms Review

Briefly describe the method outlined in [1] for upgrading the n-point two-loop MHV symbol to a function.

2.2 The A_5 Function

- describe the A_5 cluster algebra

- what makes "the" A_5 function unique?

The A_5 cluster algebra is generated from the seed cluster

$$x_1 \to x_2 \to x_3 \to x_4 \to x_5. \tag{2.1}$$

The full A_5 algebra contains 132 clusters with 140 distinct $\mathcal{X}$-coordinates.

Define:

$$x_{i_1 \ldots i_k} = \sum_{a=1}^{k} \prod_{b=1}^{a} x_{i_b} = x_{i_1} + x_{i_1} x_{i_2} + \ldots + x_{i_1} \cdots x_{i_k}. \tag{2.2}$$

The A_5 cluster algebra has an eight-fold cyclic symmetry, which we denote by σ:

$$\sigma: \quad x_1 \mapsto \frac{x_2}{1 + x_{12}}, \quad x_2 \mapsto \frac{x_3 \left(1 + x_1\right)}{1 + x_{123}}, \quad x_3 \mapsto \frac{x_4 \left(1 + x_{12}\right)}{1 + x_{1234}},$$
$$x_4 \mapsto \frac{x_5 \left(1 + x_{123}\right)}{1 + x_{12345}}, \quad x_5 \mapsto \frac{1 + x_{1234}}{x_1 x_2 x_3 x_4 x_5}. \tag{2.3}$$

A_5 also has a two-fold flip symmetry, which we denote by τ:

$$\tau: \quad x_i \mapsto \frac{1}{x_{6-i}}. \tag{2.4}$$

There are 56 distinct A_2 subalgebras in A_5 ($56 = \binom{8}{5} =$ number of distinct pentagons inside an octagon), they can be parameterized by:

$$\left\{ x_1 \to x_2, \quad x_2 \to x_3 \left(1 + x_4\right), \quad x_2 \left(1 + x_3\right) \to \frac{x_3 x_4}{1 + x_3} \right\} + \sigma,$$
$$\left\{ x_2 \to x_3, \quad x_1 \left(1 + x_2\right) \to \frac{x_2 x_3}{1 + x_2} \right\} + \sigma + \tau \tag{2.5}$$

where by "$+ \sigma$" and "$+ \sigma + \tau$" I mean "+ cyclic copies" and "+ cyclic and flip copies," respectively. This correspond to the following topologies of pentagons inside an octagon:

(8 of each)

(16 of each)
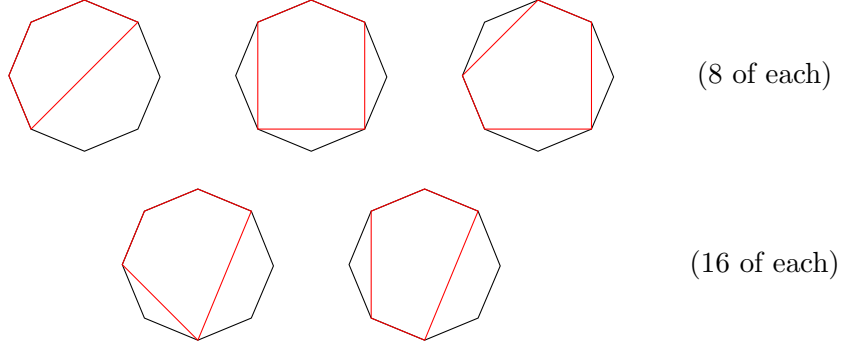
2.3   Identifying A_5 subalgebras in $\mathrm{Gr}(4, 8)$

Comment on infinite nature of $\mathrm{Gr}(4, 8)$, define "good" (in the two-loop MHV sense) subalgebras, describe algorithm for finding "good", and describe the 56 good A_5s in $\mathrm{Gr}(4, 8)$.

2.4   Representing $R_8^{(2)}$

- A_5 contribution

- Li_4 contribution

- $\mathrm{Li}_2 \mathrm{Li}_2$ contribution

- $\mathrm{Li}_3 \mathrm{Li}_1$ contribution

- Li_2Li_1^2 contribution

- Li_1^4 contribution

- $\text{Li}_2\pi^2$ contribution

- $\text{Li}_1^2\pi^2$ contribution

- π^4 constribution

2.5   Analytic Properties of $R_8^{(2)}$

- some plots

- agrees with numerics

3   The Steinmann Relations for Eight Particles

When the number of gluons n is not a multiple of 4, the BDS-like ansatz is unique because there exists only a single decomposition of the BDS ansatz

$$\mathcal{A}_n^{\text{BDS}} = \mathcal{A}_n^{\text{BDS-like}}(\{s_{i,i+1}\}) \exp\left[\frac{\Gamma_{\text{cusp}}}{4}Y_n(\{u_i\})\right], \quad n \neq 4K, \tag{3.1}$$

such that the kinematic dependence of $A_n^{\text{BDS-like}}$ involves only two-particle Mandelstam invariants while Y_n is a weight-two function depending only on dual conformal invariant cross ratios [2]. When n is a multiple of 4, no decomposition of this type exists, and we are forced to consider multiple BDS-like ansätze if we want to transparently expose all Steinmann relations between higher-particle Mandelstam invariants.

In eight particle kinematics, there are two natural BDS-like normalization choices we might consider. Namely, we can let our BDS-like ansatz depend on either three- or four-particle Mandelstam invariants in addition to two-particle invariants. In this spirit, let us define a pair of Bose-symmetric BDS-like ansätze, respectively satisfying

$$\mathcal{A}_8^{\text{BDS}} = {}^3\mathcal{A}_8^{\text{BDS-like}}(\{s_{i,i+1}\}, \{s_{i,i+1,i+2,i+3}\}) \exp\left[\frac{\Gamma_{\text{cusp}}}{4}\,{}^3Y_8(\{u_i\})\right], \tag{3.2}$$

$$\mathcal{A}_8^{\text{BDS}} = {}^4\mathcal{A}_8^{\text{BDS-like}}(\{s_{i,i+1}\}, \{s_{i,i+1,i+2}\}) \exp\left[\frac{\Gamma_{\text{cusp}}}{4}\,{}^4Y_8(\{u_i\})\right]. \tag{3.3}$$

In fact, the functions ${}^3A_8^{\text{BDS-like}}$ and ${}^4A_8^{\text{BDS-like}}$ are not uniquely fixed by these decomposition choices; each admits a family of Bose-symmetric solutions.

Any choice for ${}^3A_8^{\text{BDS-like}}$ or ${}^4A_8^{\text{BDS-like}}$ that satisfy eqns. (3.2) or (3.3) gives rise to a BDS-like normalized amplitude that manifestly exhibits a subset of the Steinmann relations. In particular, defining

$$^X\mathcal{E}_8 \equiv \frac{\mathcal{A}_8^{\text{MHV}}}{^X\mathcal{A}_8^{\text{BDS-like}}} = \exp\left[R_8 - \frac{\Gamma_{\text{cusp}}}{4}\,{}^XY_8\right] \tag{3.4}$$

for any label X, we expect that $^3\mathcal{E}_8$ should satisfy Steinmann relations between all partially overlapping pairs of three-particle invariants, while $^4\mathcal{E}_8$ should satisfy Steinmann relations between all partially overlapping pairs of four-particle invariants. That is, $^3\mathcal{E}_8$ is expected to satisfy the relations

$$\text{Disc}_{s_{i+1,i+2,i+3}}\left[\text{Disc}_{s_{i,i+1,i+2}}(^3\mathcal{E}_8)\right] = 0, \tag{3.5}$$

$$\text{Disc}_{s_{i+2,i+3,i+4}}\left[\text{Disc}_{s_{i,i+1,i+2}}(^3\mathcal{E}_8)\right] = 0, \tag{3.6}$$

for all i, while $^4\mathcal{E}_8$ is expected to satisfy

$$\text{Disc}_{s_{i+1,i+2,i+3,i+4}}\left[\text{Disc}_{s_{i,i+1,i+2,i+3}}(^4\mathcal{E}_8)\right] = 0, \tag{3.7}$$

$$\text{Disc}_{s_{i+2,i+3,i+4,i+5}}\left[\text{Disc}_{s_{i,i+1,i+2,i+3}}(^4\mathcal{E}_8)\right] = 0, \tag{3.8}$$

$$\text{Disc}_{s_{i+3,i+4,i+5,i+6}}\left[\text{Disc}_{s_{i,i+1,i+2,i+3}}(^4\mathcal{E}_8)\right] = 0. \tag{3.9}$$

However, conditions (3.5) through (3.9) don't exhaust the set of Steinmann relations obeyed by generic eight-particle amplitudes—there are also Steinmann relations between partially overlapping three- and four-particle invariants. If we want to make these additional relations manifest, we can use the fact that it proves possible to define a BDS-like ansatz that depends on all but one of the four-particle invariants (and on no three-particle invariants). That is, we decompose the BDS ansatz as

$$\mathcal{A}_8^{\text{BDS}} = {}^{3,j}\mathcal{A}_8^{\text{BDS-like}}(\{s_{i,i+1}\}, \{s_{i,i+1,i+2,i+3} \neq s_{j,j+1,j+2,j+3}\})\exp\left[\frac{\Gamma_{\text{cusp}}}{4}\,{}^{3,j}Y_8(\{u_i\})\right], \tag{3.10}$$

and in so doing define a BDS-like normalized amplitude that satisfies

$$\text{Disc}_{s_{j+2,j+3,j+4}}\left[\text{Disc}_{s_{j,j+1,j+2,j+3}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.11}$$

$$\text{Disc}_{s_{j+3,j+4,j+5}}\left[\text{Disc}_{s_{j,j+1,j+2,j+3}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.12}$$

$$\text{Disc}_{s_{j-1,j,j+1}}\left[\text{Disc}_{s_{j,j+1,j+2,j+3}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.13}$$

$$\text{Disc}_{s_{j-2,j-1,j}}\left[\text{Disc}_{s_{j,j+1,j+2,j+3}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.14}$$

$$\text{Disc}_{s_{i+1,i+2,i+3}}\left[\text{Disc}_{s_{i,i+1,i+2}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.15}$$

$$\text{Disc}_{s_{i+2,i+3,i+4}}\left[\text{Disc}_{s_{i,i+1,i+2}}(^{3,j}\mathcal{E}_8)\right] = 0, \tag{3.16}$$

where the relations (3.15) and (3.16) are the Steinmann relations satisfied by $^3\mathcal{E}_8$. Clearly it is not possible to find a solution to eq. (3.10) that is Bose-symmetric, but we can require that it is invariant under the dihedral flip that maps $s_{i\ldots l} \to s_{9-i\ldots 9-l}$. There is again a family of solutions that satisfy these constraints.

Momentum conservation implies that $^{3,j+4}\mathcal{A}_8^{\text{BDS-like}} = {}^{3,j}\mathcal{A}_8^{\text{BDS-like}}$, and thus every eight-point Steinmann relation is made manifest in at least one of the amplitudes

$$\{^4\mathcal{E}_8, {}^{3,1}\mathcal{E}_8, {}^{3,2}\mathcal{E}_8, {}^{3,3}\mathcal{E}_8, {}^{3,4}\mathcal{E}_8\}.$$

However, in practice it may prove convenient to retain $^3\mathcal{E}_8$ in one's toolbox, since it can be chosen to respect Bose symmetry. It may additionally prove possible to chose these

amplitudes to satisfy other desirable properties, since $^3Y_8$, $^4Y_8$, and $^{3,j}Y_8$ are not uniquely picked out by eqns. (3.2), (3.3), and (3.10). In fact there is a 10-dimensional (3-dimensional) space of (Bose-symmetric) solutions for $^3Y_8$, a 36-dimensional (4-dimensional) space of (Bose-symmetric) solutions for $^4Y_8$, and a 5-dimensional (3-dimensional) space of (dihedral flip-invariant) solutions for $^{3,j}Y_8$.

To take full advantage of the Steinmann relations, it is convenient to work in terms of symbol letters that isolate different Mandelstam invariants. There are twelve independent dual conformally invariant cross ratios that can appear in these symbols

$$u_1 = \frac{s_{12}s_{4567}}{s_{123}s_{812}}, \quad \text{and cyclic (8-orbit)} \tag{3.17}$$

$$u_9 = \frac{s_{123}s_{567}}{s_{1234}s_{4567}}, \quad \text{and cyclic (4-orbit).} \tag{3.18}$$

It is not possible to isolate all three- and four-particle Mandelstam invariants simultaneously into twelve different symbol letters. (More than twelve symbol letters will appear in these amplitudes, but we here restrict our attention to the twelve that will appear in the first entry.) However, different choices of letters can be made such that either all the four-particle invariants, or all the three-particle invariants, are isolated.

One choice that isolates the four-particle invariants is

$$^4d_1 = u_2\,u_6 = \frac{s_{23}\,s_{67}\,(s_{1234})^2}{s_{123}\,s_{234}\,s_{567}\,s_{678}}, \quad \text{and cyclic (4-orbit)} \tag{3.19}$$

$$^4d_5 = u_2/u_6 = \frac{s_{23}\,s_{567}\,s_{678}}{s_{67}\,s_{123}\,s_{234}}, \quad \text{and cyclic (4-orbit)} \tag{3.20}$$

$$^4d_9 = u_1\,u_2\,u_5\,u_6\,u_9^2 = \frac{s_{12}\,s_{23}\,s_{56}\,s_{67}}{s_{234}\,s_{456}\,s_{678}\,s_{812}}, \quad \text{and cyclic (4-orbit).} \tag{3.21}$$

In this alphabet $^4d_1, {}^4d_2, {}^4d_3$, and $^4d_4$ each contain a different four-particle Mandelstam invariant, while the other letters only involve two- and three-particle invariants. The extended Steinmann relations then tell us that $^4d_1, {}^4d_2, {}^4d_3$, and $^4d_4$ can never appear next to each other in the symbol of $^4A_8^{\text{BDS-like}}$ (but each can still appear next to themselves).

Similarly, we can isolate the three-particle invariants by choosing

$$^3d_1 = \frac{u_1\,u_2\,u_4\,u_7}{u_3\,u_5\,u_6\,u_8\,u_9^2} = \frac{s_{12}\,s_{23}\,s_{45}\,s_{78}\,(s_{1234})^2\,(s_{4567})^2}{s_{34}\,s_{56}\,s_{67}\,s_{81}\,(s_{123})^2}, \quad \text{and cyclic (8-orbit)} \tag{3.22}$$

$$^3d_9^4 = u_1\,u_5\,u_9\,u_{12} = \frac{s_{12}\,s_{56}}{s_{1234}\,s_{3456}}, \quad \text{and cyclic (4-orbit),} \tag{3.23}$$

in which case $^3d_1$ through $^3d_8$ each contain a different three-particle Mandelstam invariant, as well as four-particle Mandelstams that they don't partially overlap with. The remaining four letters only contain two- and four-particle invariants. In these letters, conditions (3.15) and (3.16) tell us that $^3d_7, {}^3d_8, {}^3d_2$, and $^3d_3$ can never appear next to $^3d_1$ in the symbols of $^3\mathcal{E}_8$ or $^{3,j}\mathcal{E}_8$ (plus the cyclic images of this statement). Moreover, conditions (3.11) through (3.14) give us the additional restrictions that none of $^3d_1, {}^3d_5, {}^3d_9$ and $^3d_{10}$ can ever appear next to

$^3d_3, {}^3d_4, {}^3d_7$, or $^3d_8$ in the symbol of $^{3,1}\mathcal{E}_8$ (analogous relations hold for the other $^{3,j}\mathcal{E}_8$). These are the restrictions given by the Steinmann relations involving s_{1234} and one of $s_{781}, s_{812}, s_{345}$, or s_{456}. The other Steinmann relations between three- and four-particle invariants will not be respected by $^{3,1}\mathcal{E}_8$, since $^{3,j}\mathcal{A}_8^{\text{BDS-like}}$ depends on s_{2345}, s_{3456}, and s_{4567}.

4 Conclusion

A A_5 Representation of $R_7^{(2)}$

References

[1] J. Golden and M. Spradlin, "An analytic result for the two-loop seven-point MHV amplitude in $\mathcal{N} = 4$ SYM," *JHEP*, vol. 1408, p. 154, 2014.

[2] G. Yang, "Scattering amplitudes at strong coupling for 4K gluons," *JHEP*, vol. 12, p. 082, 2010.