

Week 3

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Qual Study

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Solution 2018.6. The partition function for one link is

$$\begin{aligned}
 Z_1 &= \sum_s e^{-\beta E_s} \\
 &= \sum_{n=0}^{\infty} \left(e^{-\beta[(n+1/2)\hbar\omega_a - Fa]} + e^{-\beta[(n+1/2)\hbar\omega_b - Fb]} \right) \\
 &= e^{-\beta\hbar\omega_a/2} e^{\beta Fa} \sum_{n=0}^{\infty} e^{-\beta n\hbar\omega_a} + e^{-\beta\hbar\omega_b/2} e^{\beta Fb} \sum_{n=0}^{\infty} e^{-\beta n\hbar\omega_b} \\
 &= \frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}}
 \end{aligned}$$

a. The average length for one link is given by

$$\begin{aligned}
 \langle \ell \rangle &= \frac{1}{Z_1} \sum_s \ell_s e^{-\beta E_s} \\
 &= \frac{1}{Z_1} \sum_{n=0}^{\infty} \left(a e^{-\beta[(n+1/2)\hbar\omega_a - Fa]} + b e^{-\beta[(n+1/2)\hbar\omega_b - Fb]} \right) \\
 &= \frac{1}{Z_1} \left[a \frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + b \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}} \right] \\
 &= \left[a \frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + b \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}} \right] \left[\frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}} \right]^{-1}
 \end{aligned}$$

The links are non-interacting, so the average total length is just

$$L = N \left[a \frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + b \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}} \right] \left[\frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{1 - e^{-\beta\hbar\omega_a}} + \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{1 - e^{-\beta\hbar\omega_b}} \right]^{-1}$$

b. In the $T \rightarrow 0$ limit, the system should occupy its minimum energy state. Therefore, $\lim_{T \rightarrow 0} L = Na$ if $\frac{1}{2}\hbar\omega_a - Fa < \frac{1}{2}\hbar\omega_b - Fb$, otherwise the limit is Nb . In the $T \rightarrow \infty$ limit, $\beta \rightarrow 0$, and we have

$$\lim_{T \rightarrow \infty} L = N \lim_{T \rightarrow \infty} \left[a \frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{\beta\hbar\omega_a} + b \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{\beta\hbar\omega_b} \right] \left[\frac{e^{-\beta\hbar\omega_a/2} e^{\beta Fa}}{\beta\hbar\omega_a} + \frac{e^{-\beta\hbar\omega_b/2} e^{\beta Fb}}{\beta\hbar\omega_b} \right]^{-1}$$

$$\begin{aligned}
&= N \lim_{T \rightarrow \infty} \left[a \frac{1}{\beta \hbar \omega_a} + b \frac{1}{\beta \hbar \omega_b} \right] \left[\frac{1}{\beta \hbar \omega_a} + \frac{1}{\beta \hbar \omega_b} \right]^{-1} \\
&= N \frac{a/\hbar \omega_a + b/\hbar \omega_b}{1/\hbar \omega_a + 1/\hbar \omega_b}
\end{aligned}$$

Solution 2018.8. a. This is worked in box 6.2 of the P556 notes. I reproduce the solution here

$$\begin{aligned}
\frac{N}{A} &= \int \frac{1}{y^{-1} e^{\beta E} - \eta} \frac{d^2 \mathbf{k}}{(2\pi)^2} \\
&= \int_0^\infty \frac{1}{y^{-1} e^{\beta \hbar^2 k^2 / 2m} - \eta} \frac{k dk}{2\pi} \quad u = e^{-\beta \hbar^2 k^2 / 2m}, \quad du = e^{-\beta \hbar^2 k^2 / 2m} \frac{\beta \hbar^2}{2m} k dk \\
&= \frac{2m}{\beta \hbar^2} \int_0^1 \frac{1}{y^{-1}/u - \eta} \frac{du}{u} \\
&= \frac{m}{\pi \beta \hbar^2} \int_0^1 \frac{du}{y^{-1} - \eta u} \\
&= -\frac{m}{\pi \beta \hbar^2 \eta} [\log(y^{-1} - \eta u)]_0^1 \\
&= -\frac{m}{\pi \beta \hbar^2 \eta} [\log(y^{-1} - \eta) - \log(y^{-1})] \\
&= -\frac{m}{\pi \beta \hbar^2 \eta} [\log(1 - \eta y)] \\
1 - \eta y &= \exp \left(-\frac{\pi \hbar^2 \eta}{m k_B T} \frac{N}{A} \right) \\
y &= \eta - \eta \exp \left(-\frac{\pi \hbar^2 \eta}{m k_B T} \frac{N}{A} \right) \\
\mu &= k_B T \log \left[\eta - \eta \exp \left(-\frac{\pi \hbar^2 \eta}{m k_B T} \frac{N}{A} \right) \right]
\end{aligned}$$

b.

i++i

Solution 2017.7. a. The single particle partition function is

$$\begin{aligned}
Z_1 &= Z_{1,trans} Z_{1,rot} \\
Z_{1,trans} &= \frac{V}{\lambda^3} \\
Z_{1,rot} &= \frac{1}{h^2} \int \exp \left[-\beta \frac{1}{2I} \left(p_\theta^2 + \frac{p_\phi^2}{\sin^2 \theta} \right) \right] d\phi d\theta dp_\phi dp_\theta \\
&= \frac{1}{h^2} (2\pi) \int_0^\pi \sqrt{\frac{\pi}{1/2I}} \sqrt{\frac{\pi}{1/2I \sin^2 \theta}} d\theta \\
&= \frac{4I\pi^2}{h^2} \int_0^\pi \sin \theta d\theta = \frac{2I}{h^2}
\end{aligned}$$

Then, the partition function for N particles is

$$Z(N, \beta) = \frac{1}{N!} Z_1^N$$

since the particles are indistinguishable, and the grand partition function is

$$\begin{aligned} Z_G &= \sum_{N=0}^{\infty} z^N Z(N, \beta) \\ &= \exp(z Z_1) \end{aligned}$$

b. The average number of particles is

$$\begin{aligned} \mathcal{N} &= -\frac{\partial}{\partial \alpha} \log Z_G \\ &= z \frac{\partial}{\partial z} \log Z_G \\ &= z \frac{\partial}{\partial z} (z Z_1) \\ &= z Z_1 \end{aligned}$$

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Solution 2017.16. $\mathfrak{j}+++i$

Solution 2016.6. $\mathfrak{j}+++i$

Solution 2016.8. $\mathfrak{j}+++i$

Solution 2015.7. $\mathfrak{j}+++i$

Solution 2015.16. $\mathfrak{j}+++i$

Solution 2014.6. $\mathfrak{j}+++i$

Solution 2014.8. $\mathfrak{j}+++i$

Solution 2013.7. $\mathfrak{j}+++i$

Solution 2013.16. $\mathfrak{j}+++i$