Stress, Strain and Deformation **Bending Stresses and Strains**

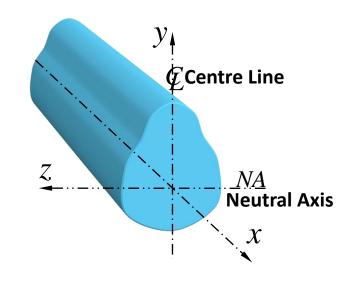
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- Beams are straight
- Cross-sections are constant
- X
- Material response is linear elastic
- Material properties are 'constant':
 - Isotropic
 - Homogeneous
- Loading in the plane of symmetry
 - i.e. always in the plane x-y
 - Also called symmetric bending

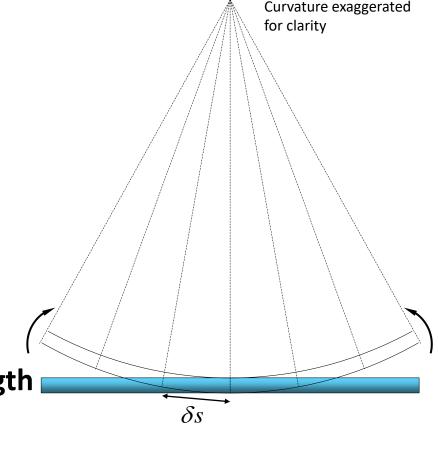


Section A-A

Note: choice of axes may be different in different textbooks!



- Small deformations
- Pure bending
 - Neglecting shear deformation
 - Deformation in shape of circular arc
- Cross-sections remain plane
 - and rotate about the neutral plane (i.e. no shear deformation)
- Consider a small element of arc-length δs under a **bending moment** M
 - The moment is equivalent to compression on one 'side' and tension on the other, e.g. a 'couple'
 - So the length decreases on one side and increases on the other forming a **curvature**, with radius R and curvature $1/R = \kappa \rightarrow \text{(kappa)}$

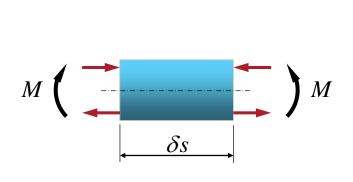


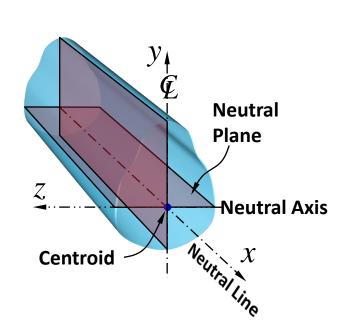
 δs



Neutral Plane, Neutral Line and Neutral Axis

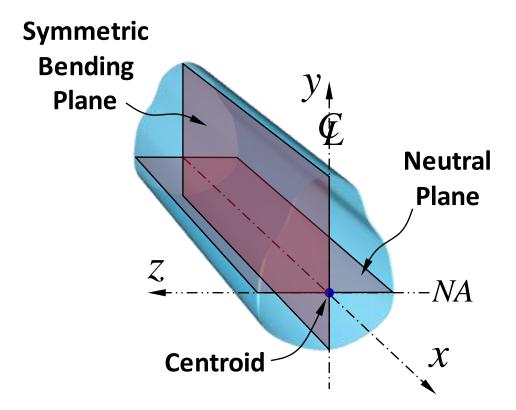
- Somewhere in the middle the length will be unchanged; here the normal stress will be zero - this is the <u>neutral plane</u>
- The intersection of the neutral plane with the loading plane is called the <u>neutral line</u>
- The intersection of the neutral plane with the cross section plane is called the <u>neutral axis</u>







- Bending acts in a section-symmetric plane
- Neutral axis perpendicular to symmetric bending plane and through section centroid
- (This is not the case for non-symmetric bending as we will see in StM2)





- Consider the beam element ds at a distance y from the neutral line,
 where R = <u>radius of curvature</u>
- After bending, assuming small displacements:

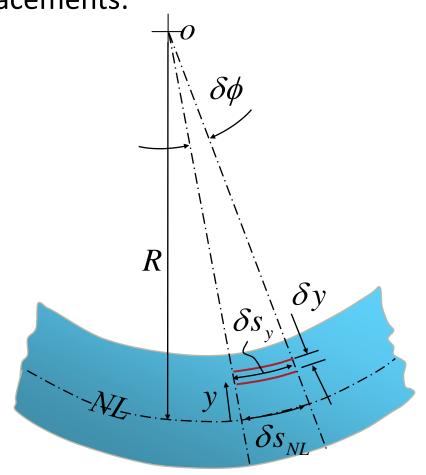
$$\delta \phi \cong \tan \delta \phi \cong \frac{\delta s_{NL}}{R}$$

$$\delta s_{NL} = R \, \delta \phi$$

And also:

$$\delta\phi \cong \tan\delta\phi \cong \frac{\delta s_y}{R - y}$$

$$\delta s_{y} = (R - y)\delta \phi$$





• But before bending (i.e. when straight)

$$\delta s_y = \delta s_{NL}$$

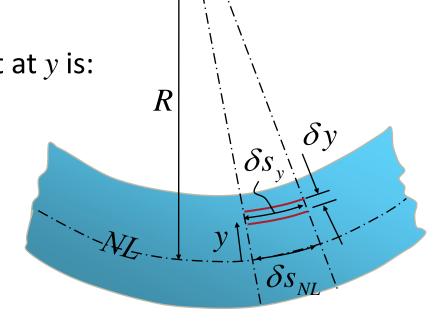
• So the change in length of the 'element' at y is:

$$(\delta s_y - \delta s_{NL}) = (R - y) \delta \phi - R \delta \phi = -y \delta \phi$$

• Therefore the strain of the element at y is:

$$\varepsilon = \frac{\delta s_y - \delta s_{NL}}{\delta s_{NL}} = \frac{-y \, \delta \phi}{R \, \delta \phi} = -\frac{1}{R} \, y$$

$$\mathcal{E} = -\frac{1}{R} y \qquad \Rightarrow \text{KEY POINT} \leftarrow$$

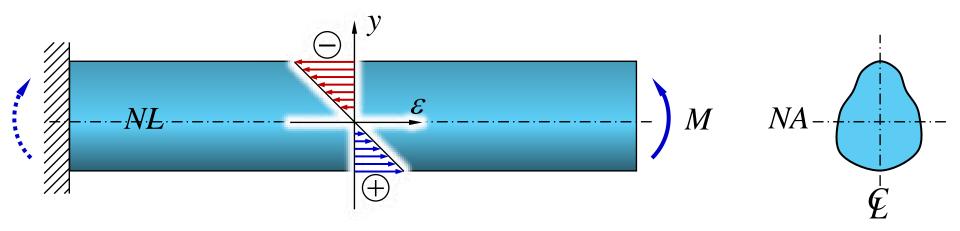




• The derived strain distribution is:

$$\varepsilon = -\frac{1}{R} y$$

 Note the linear variation of strains with distance from the neutral axis NA:

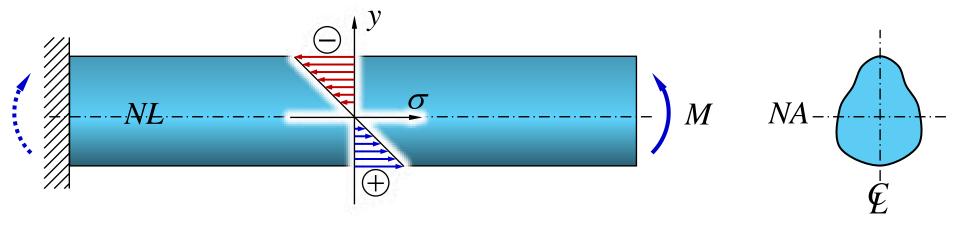




For linear elastic stress-strain behaviour (i.e. applying Hooke's law)
 the direct stress (normal to section) is:

$$\varepsilon = \frac{-y}{R}$$
 , $\sigma = E \varepsilon$: $\sigma = -\frac{E}{R} y$

 Note the linear variation of stresses with distance from the neutral axis NA:

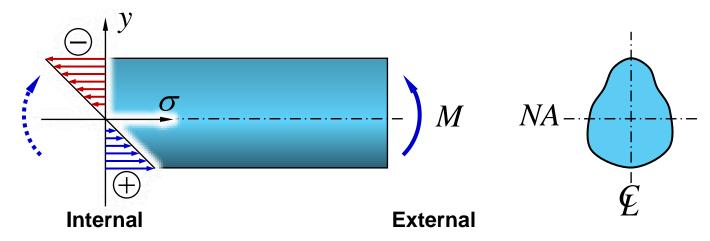




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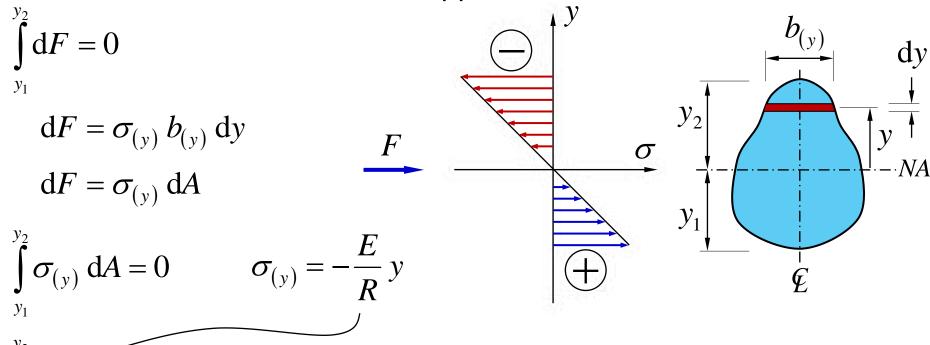
 Note the linear variation of stresses with distance from the neutral axis NA:



 External forces and moments are reacted by internal forces and moments which are the result of internal stresses



For pure bending the total internal resultant force must be zero,
 since no external axial force is applied



$$\int_{y}^{y_2} -\frac{E}{R} y \, dA = 0$$

However E and R are constant with respect to y, therefore:

 $\int_{y_1}^{y_2} y \, \mathrm{d}A = 0$

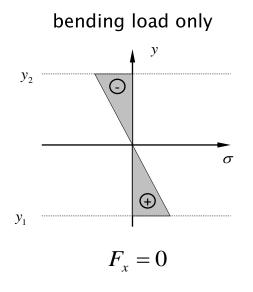
- \rightarrow The **1**st moment of area is zero about NA
- \rightarrow NA must pass through the **centroid** of the cross-section

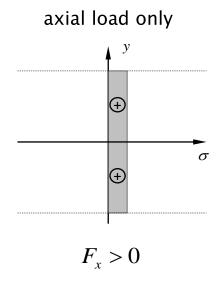
Internal Axial Forces

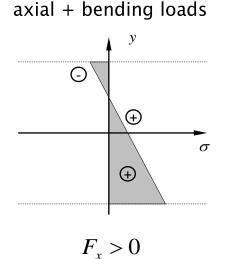
• For a constant width b, the resultant internal axial force F_x can be interpreted as the 'summation' of the 'areas' under the curves of stress distribution through the thickness

 $F_x = b \int_{y_1}^{y_2} \sigma_{(y)} \, \mathrm{d}y$

(Note that these areas can have different 'signs')









• For pure bending the internal resultant moment must be nonzero, to balance out the external moment M

$$\int_{y_1}^{y_2} y \, dF = -M$$

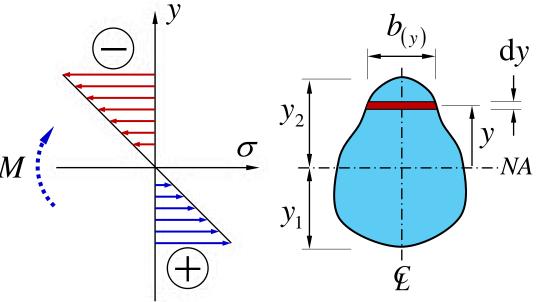
$$\int_{y_1}^{y_2} y \, \sigma_{(y)} \, dA = -M$$

$$\int_{y_1}^{y_2} y \, \sigma_{(y)} \, dA = -M$$

$$\int_{y_2}^{y_2} \frac{E}{R} y^2 \, dA = M$$

$$\frac{\int_{y_1} \overline{R} y \, dA = M}{E \int_{y_1}^{y_2} y^2 \, dA = M} \qquad \int_{y_1}^{y_2} y^2 \, dA = I$$

$$\frac{1}{R}EI = \kappa \ EI = M$$



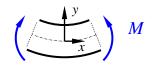
- \rightarrow The **2**nd **moment of area** of the cross-section is called $I_{(ref. axis)}$
- \rightarrow The product EI is called the 'flexural modulus' of the beam



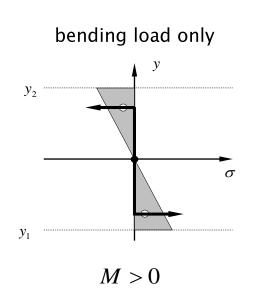
Bending Moment

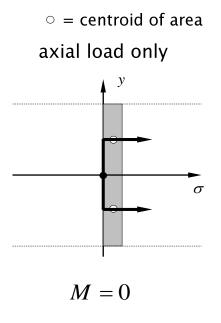
• The internal bending moment is given by the integral:

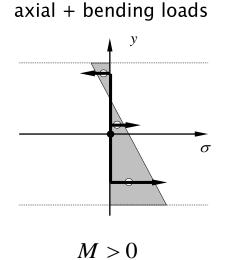
$$M = -b \int_{y_1}^{y_2} \sigma_{(y)} y \, \mathrm{d}y$$



• This can be interpreted as the 'moment of area' of the stress distribution:









We can now re-arrange the equations:

$$\frac{EI}{R} = M$$
 \therefore $\frac{M}{I} = \frac{E}{R}$

$$\sigma = -\frac{E}{R}y$$
 : $-\frac{\sigma}{y} = \frac{E}{R}$



$$-\frac{\sigma}{y} = \frac{M}{I} = \frac{E}{R}$$



Stress, Strain and Deformation **Second Moment of Areas**

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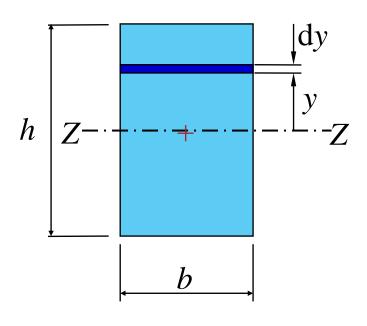
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Rectangular and Square Cross-Sections

- Consider a **rectangular cross-section** of height h and width (or breadth) b
- The second moment of area about a centroidal Z-Z axis is:



$$I = \int_{-h/2}^{h/2} y^2 \, \mathrm{d}A$$

$$I = \int_{-h/2}^{h/2} y^2 \ b \ \mathrm{d}y$$

$$I = \left[\frac{y^3}{3}b\right]_{-h/2}^{h/2}$$

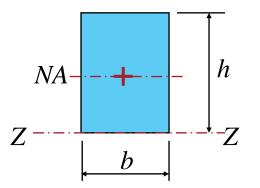
$$I = \frac{b h^3}{12}$$

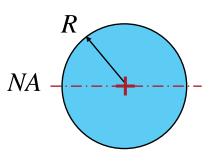
• For square cross-sections b=h, therefore

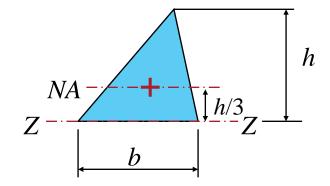
$$I = \frac{h^4}{12}$$



Rectangular







$$I_{NA} = \frac{b h^3}{12}$$

$$I_{NA} = \frac{\pi R^4}{4}$$

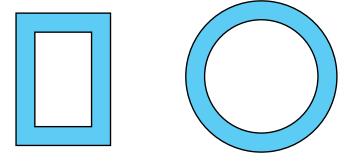
$$I_{NA} = \frac{b h^3}{36}$$

$$I_{ZZ} = \frac{b h^3}{3}$$

$$I_{ZZ} = \frac{b h^3}{12}$$



• For concentric hollow cross-sections, e.g.



The second moment of area is simply:

$$I = I_{\text{outer}} - I_{\text{inner}}$$

