

Shor algorithm

QFT → "Phase estimation"

In the algorithm we only worked with real numbers. The QFT adds a phase to the problem. One can formulate the problem as follows.

You are given a state on n qubits as follows.

$$\begin{aligned} \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} e^{\frac{2i\pi}{2^n} kx} |k\rangle &= \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} \left(e^{\frac{2i\pi}{2^n} x} \right)^k |k\rangle \\ &= \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} \omega^k |k\rangle \end{aligned} \quad \left(\omega = e^{\frac{2i\pi}{2^n} x} \right)$$

Can you get the value of x , seen as a number between 0 and $2^n - 1$?

The answer is yes: we need a circuit computing

$$\frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} e^{\frac{2i\pi}{2^n} kx} |k\rangle \mapsto |x\rangle \quad (\text{beware: there will be re-indexing of the register } x)$$

Let us try with $n = 1$ and $x = 0$ ou 1. Then

$$\sum_{k=0}^1 e^{\frac{2i\pi}{2} kx} |k\rangle = |0\rangle + (-1)^x |1\rangle$$

and we want to get back $|x\rangle$... Hadamard is enough.

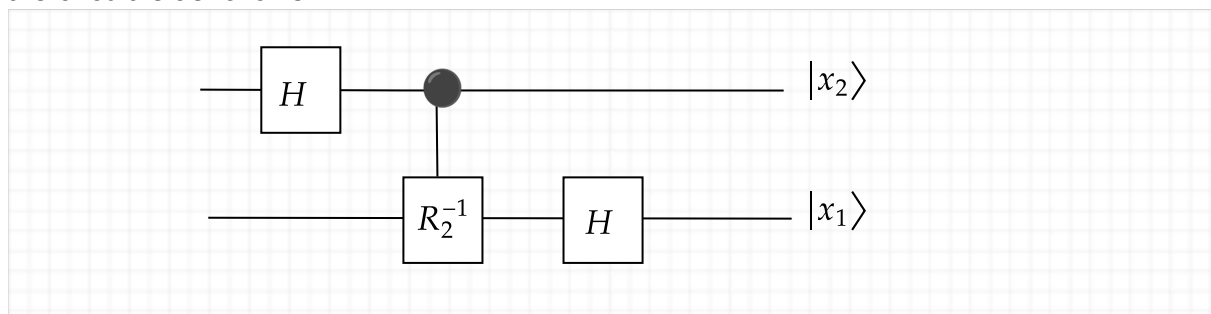
Let us try with $n = 2$ and $x = 0, 1, 2$ or 3 : $[x]_2 = x_1x_2$, i.e. $x = 2x_1 + x_2$

We have

$$\begin{aligned} \sum_{k=0}^{2^n-1} e^{\frac{2i\pi}{2^n} kx} |k\rangle &= \sum_{k=0}^3 e^{2i\pi \left(\frac{x_1}{2} + \frac{x_2}{4} \right) k} |k\rangle = \\ &= \sum_{k=0}^3 e^{2i\pi \frac{x_1}{2} k} e^{2i\pi \frac{x_2}{4} k} |k\rangle = \sum_{k=0}^3 e^{i\pi x_1 k} e^{i\pi \frac{x_2}{2} k} |k\rangle \\ &= \frac{1}{2} (|00\rangle + e^{i\pi x_1} e^{i\pi \frac{x_2}{2}} |01\rangle + e^{2i\pi x_1} e^{i\pi x_2} |10\rangle + \\ &\quad e^{3i\pi x_1} e^{3i\pi \frac{x_2}{2}} |11\rangle) \\ &= \end{aligned}$$

$$\begin{aligned}
& \frac{1}{2} \left(|00\rangle + e^{i\pi x_1} e^{i\pi \frac{x_2}{2}} |01\rangle + e^{i\pi x_2} |10\rangle + e^{i\pi x_1} e^{3i\pi \frac{x_2}{2}} |11\rangle \right) \\
&= \\
& \frac{1}{2} (|0\rangle + e^{i\pi x_2} |1\rangle) \otimes (|0\rangle + e^{i\pi x_1} e^{i\pi \frac{x_2}{2}} |1\rangle) \\
&= \\
& \frac{1}{2} (|0\rangle + (-1)^{x_2} |1\rangle) \otimes (|0\rangle + (-1)^{x_1} e^{i\pi \frac{x_2}{2}} |1\rangle)
\end{aligned}$$

so to get back $|x_2 x_1\rangle$ (!!! BEWARE OF THE RE-INDEXING !!!)
the circuit is as follows



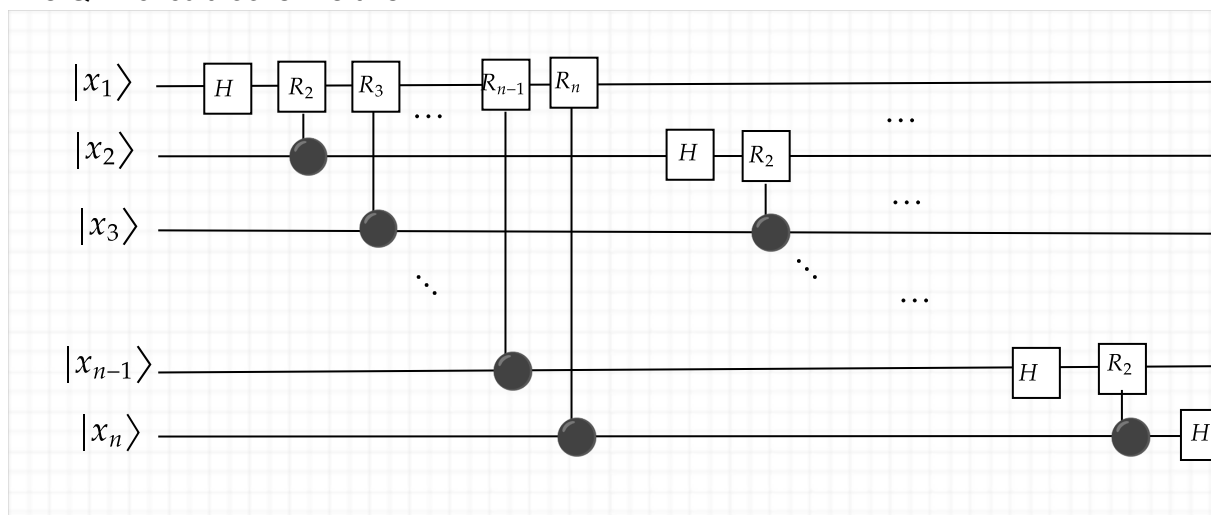
with $R_k = \begin{pmatrix} 1 & 0 \\ 0 & e^{2i\pi/2^k} \end{pmatrix}$

This generalizes to n qubits, and the circuit is call "QFT inverse" (for "Quantum Fourier Transform"). The name comes from the fact that in the other direction, we compute

$$|x\rangle \mapsto \sum_{k=0}^{2^n-1} e^{\frac{2i\pi}{2^n} kx} |k\rangle \quad (\text{modulo the } x \text{ reindexing})$$

which is very close to a discrete Fourier transform.

The QFT circuit looks like this :



with

$$R_k = \begin{pmatrix} 1 & 0 \\ 0 & e^{2i\pi/2^k} \end{pmatrix}$$

Quantum Phase Estimation (QPE)

An very useful algorithm makes it possible to find the eigenvalues of a unitary map U . For such a map the eigenvalues are necessary of the form $e^{2i\pi\omega}$, and, without loss of generality, ω is a real number between 0 and 1.

Recall:

If $U|\psi\rangle = \lambda|\psi\rangle$ we say that $|\psi\rangle$ is an **eigenvector** of U and λ an **eigenvalue**

In the case where U is unitary, λ is of the form $e^{2i\pi\omega}$ because U preserves the norm... so $|\lambda| = 1$.

The algorithm QPE ("Quantum Phase Estimation") makes it possible to find ω .

To understand how this work, let us consider ω set to $0.x_1x_2$ in binary form.

$$\text{So } \omega = \frac{x_1}{2} + \frac{x_2}{4}.$$

Let $|\psi\rangle$ be the corresponding eigenvector.

$$\text{We then have } U|\psi\rangle = e^{2i\pi\omega}|\psi\rangle.$$

Computing C-U :

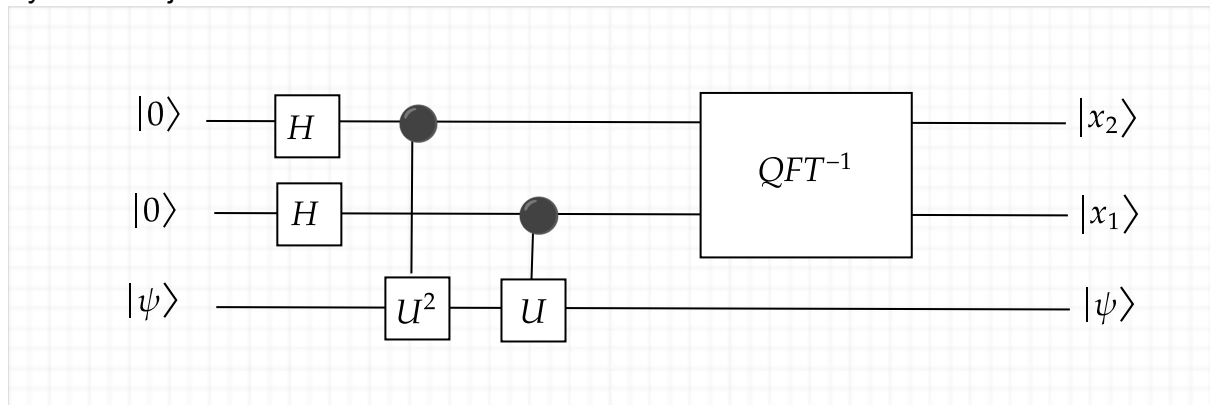
$$\begin{aligned} & \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes |\psi\rangle \\ &= \\ & \frac{1}{\sqrt{2}}(|0\rangle \otimes |\psi\rangle + |1\rangle \otimes |\psi\rangle) \\ & \xrightarrow{C-U} \\ & \frac{1}{\sqrt{2}}(|0\rangle \otimes |\psi\rangle + |1\rangle \otimes (U|\psi\rangle)) \\ &= \\ & \frac{1}{\sqrt{2}}(|0\rangle \otimes |\psi\rangle + |1\rangle \otimes (e^{2i\pi\omega}|\psi\rangle)) \\ &= \end{aligned}$$

$$\begin{aligned} & \frac{1}{\sqrt{2}}(|0\rangle \otimes |\psi\rangle + e^{2i\pi\omega} |1\rangle \otimes |\psi\rangle) \\ &= \\ & \frac{1}{\sqrt{2}}(|0\rangle + e^{2i\pi\omega} |1\rangle) \otimes |\psi\rangle \end{aligned}$$

So

$$\begin{aligned} C-U : & \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes |\psi\rangle \\ \mapsto & \frac{1}{\sqrt{2}} \left(|0\rangle + e^{2i\pi \left(\frac{x_1}{2} + \frac{x_2}{4} \right)} |1\rangle \right) \otimes |\psi\rangle \\ (C-U)^2 : & \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes |\psi\rangle \\ \mapsto & \frac{1}{\sqrt{2}} \left(|0\rangle + e^{2i\pi \left(\frac{x_2}{2} \right)} |1\rangle \right) \otimes |\psi\rangle \end{aligned}$$

By what we just saw with the inverse QFT:



(Note how the indices got swapped)

Once again, this generalizes. One can show that this is also working (albeit probabilistically) if ω is not writable on precisely n bits.

Typical use-case: Shor's algorithm

Problem: **FACTORIZATION**

Input : N a composite number: it admits a non-trivial factor.

Output : a divisor of N

Problem: **ORDER-FINDING**

Input : two integers N and a , co-primes

Output : The period r of a , i.e. the smallest $r > 0$ such that $a^r \equiv 1 \pmod{N}$

Problem: **PHASE-ESTIMATION**

Input : A unitary U and an eigenvector $|\phi\rangle$

Output : the corresponding eigenvalue

We can reduce FACTORIZATION to ORDER-FINDING, that can itself be reduced to PHASE-ESTIMATION...

1st step : reduction of FACTORIZATION to ORDER-FINDING.

Meaning: "If I know how to solve ORDER-FINDING I can easily factor a N "

It is purely a math problem.

So, suppose that I can solve ORDER-FINDING. I am given N to factor.

The idea is to realize that if N is the product of two co-prime integers greater than 2, one can derive from the chinese remainder theorem the existence of at least 4 numbers b such that $b^2 \equiv 1 \pmod{N}$. There is then at least one such b distinct from 1 and $-1 \pmod{N}$ such that $b^2 - 1 \equiv 0 \pmod{N}$, i.e. such that N divides $(b-1)(b+1)$.

→ $b \not\equiv 1 \pmod{N}$ so $b-1 \not\equiv 0 \pmod{N}$ so N does not divide $b-1$

→ $b \not\equiv -1 \pmod{N}$ so N does not divide $b+1$

So $\gcd(N, b+1)$ and $\gcd(N, b-1)$ are non-trivial: we got factors of N (which are furthermore efficiently computable!)

An algorithm for FACTORIZATION then ultimately consists in efficiently finding such a b . Shor's algorithm proceeds as follows.

1) We randomly select a number $1 < a < N$. If it is not co-prime with N , we are done: we have a non-trivial factor.

2) Otherwise, it is co-prime with N : we then invoke our algorithm for ORDER-FINDING: it outputs a number r such that $a^r = 1 \pmod{N}$. Because of maths properties, the number r is odd or even with probability $\frac{1}{2}$. We want an even r : we start over to step 1 until we get one.

3) We now have an even r : $r = 2*r'$. So $(a^{r'})^2 \equiv 1 \pmod{N}$. (Note : This is the number b we were looking for!)

So $(a^{r'})^2 - 1 = 0 \pmod{N}$

So $(a^{r'} - 1)(a^{r'} + 1) = 0 \pmod{N}$ and then N divides $(a^{r'} - 1)(a^{r'} + 1)$.

For density reasons, and invoking the chinese remainder theorem, we can assume

that $a^{r'}$ is neither 1 nor -1 modulo N .

4) A factor of N is then for instance $\gcd(N, a^{r'} - 1)$.

2nd step : solving ORDER-FINDING using PHASE-ESTIMATION (and then using a quantum co-processor !)

Note that if a and N are co-primes, then $x \mapsto a*x \bmod N$ is a reversible function (it is a permutation of $\{0 \dots N-1\}$).

The unitary we can choose is simply $U_a : |x\rangle \mapsto |a*x \bmod N\rangle$ (multiplication modulo N)

Beware: this does not say how to implement it efficiently... One possibility is to use the technique we saw with V_f 's and U_f 's, but there are better alternatives, see e.g. <https://arxiv.org/abs/quant-ph/0205095>.

What is an eigenvector for U_a ?

Let us try with successive approximations.

Start with $|1\rangle_n = |0\dots 01\rangle$ (the binary encoding of 1 on n qubits)

If we apply U_a : we get $|a\rangle$

What about $\frac{1}{\sqrt{2}}(|1\rangle + |a \bmod N\rangle)$:

applying U_a , we get $\frac{1}{\sqrt{2}}(|a \bmod N\rangle + |a^2 \bmod N\rangle)$

What about $\frac{1}{\sqrt{3}}(|1\rangle + |a \bmod N\rangle + |a^2 \bmod N\rangle)$:

applying U_a , we get $\frac{1}{\sqrt{3}}(|a \bmod N\rangle + |a^2 \bmod N\rangle + |a^3 \bmod N\rangle)$

... We can continue like that. Eventually, the power of a will reach the order r , and $|a^r \bmod N\rangle$ is then $|1\rangle$.

So, summarizing:

$$\begin{aligned} \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} |a^k \bmod N\rangle &\xrightarrow{U_a} \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} |a^{(k+1)} \bmod N\rangle = \frac{1}{\sqrt{r}} \sum_{k=1}^r |a^k \bmod N\rangle \\ &= \frac{1}{\sqrt{r}} \left(|a^r \bmod N\rangle + \sum_{k=1}^{r-1} |a^k \bmod N\rangle \right) \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{\sqrt{r}} \left(|1 \bmod N\rangle + \sum_{k=1}^{r-1} |a^k \bmod N\rangle \right) \\
&= \frac{1}{\sqrt{r}} \left(|a^0 \bmod N\rangle + \sum_{k=1}^{r-1} |a^k \bmod N\rangle \right) \\
&= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} |a^k \bmod N\rangle
\end{aligned}$$

And we have an eigenvector of U_a , with eigenvalue 1.

Can we find more of them ?

Using the same technique, we can define (indexed with s)

$$|\phi_s\rangle = \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2i\pi \frac{s \cdot k}{r}} |a^k \bmod N\rangle$$

Let's compute:

$$\begin{aligned}
U_a |\phi_s\rangle &= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2i\pi \frac{s \cdot k}{r}} U_a |a^k \bmod N\rangle \\
&= \\
&= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2i\pi \frac{s \cdot k}{r}} |a^{(k+1)} \bmod N\rangle \\
&= \\
&= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2i\pi \frac{s \cdot (k-1)}{r}} |a^k \bmod N\rangle \\
&= \\
&= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{2i\pi \frac{s}{r}} e^{-2i\pi \frac{s \cdot k}{r}} |a^k \bmod N\rangle \\
&= \\
&= e^{2i\pi \frac{s}{r}} |\phi_s\rangle
\end{aligned}$$

This eigenvalue is nice, since it contains a phase parametrized by r ... As QPE gives us the phase, with some luck we can retrieve r out.

U_a have r such eigenvectors, one for each value of s between 0 and $r-1$.

Is it over ? We would just have to use the fact that

$$|0\dots 0\rangle \otimes |\phi_s\rangle \xrightarrow{QPE(U_a)} |x_1\dots x_n\rangle \otimes |\phi_s\rangle$$

where $x_1\dots x_n$ is a binary representation of the phase of the eigenvalue corresponding to $|\phi_s\rangle$, from which ---maybe--- with some processing one can infer r .

The problem is that one cannot directly use these eigenvectors, since we would need to know r .

However, what we can do is use the fact that the QPE is a linear map, so we can place them in superposition:

$$|0\dots 0\rangle \otimes (\alpha|\phi_s\rangle + \beta|\phi_{s'}\rangle) \xrightarrow{QPE(U_a)} \alpha|x_1\dots x_n\rangle \otimes |\phi_s\rangle + \beta|y_1\dots y_n\rangle \otimes |\phi_{s'}\rangle$$

(where $y_1\dots y_n$ is a binary representation of the phase of the eigenvalue corresponding to $|\phi_{s'}\rangle$)

The trick consists in realizing that if we place them all in (equal) superposition, as follows:

$$\begin{aligned} \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |\phi_s\rangle &= \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-2i\pi \frac{s*k}{r}} |a^k \bmod N\rangle \\ &= \\ \frac{1}{r} \sum_{k=0}^{r-1} \sum_{s=0}^{r-1} e^{-2i\pi \frac{s*k}{r}} |a^k \bmod N\rangle \\ &= \\ \frac{1}{r} \sum_{k=0}^{r-1} \left(\sum_{s=0}^{r-1} e^{-2i\pi \frac{s*k}{r}} \right) |a^k \bmod N\rangle \quad (***) \end{aligned}$$

The inner (red) sum is equal to

- r if $k = 0$
- 0 otherwise (because it is a sum of all of the roots of unity).

Therefore, in the sum (***) all the terms are nul except when $k = 0$: we get

$$\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |\phi_s\rangle = \frac{1}{r} (r|a^0 \bmod N\rangle) = |1\rangle_n \text{ (the encoding of 1 on } n \text{ qubits: } |0\dots 01\rangle)$$

If we run $QPE(U_a)$ on this input, we "compute" all of the phases at once. Consider two registers, the first one for retrieving the ω of the eigenvalue and the second one for the eigenvector. So

$$QPE(U_a)(|0\dots 0\rangle \otimes |\phi_s\rangle) = |s/r\rangle \otimes |\phi_s\rangle$$

where by $|s/r\rangle$ we mean the approximation of s/r over the corresponding number of qubits.

Then

$$\begin{aligned} QPE(U_a) \left(|0\dots 0\rangle \otimes \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |\phi_s\rangle \right) &= \\ \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} QPE(U_a) (|0\dots 0\rangle \otimes |\phi_s\rangle) &= \\ \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |s/r\rangle \otimes |\phi_s\rangle \end{aligned}$$

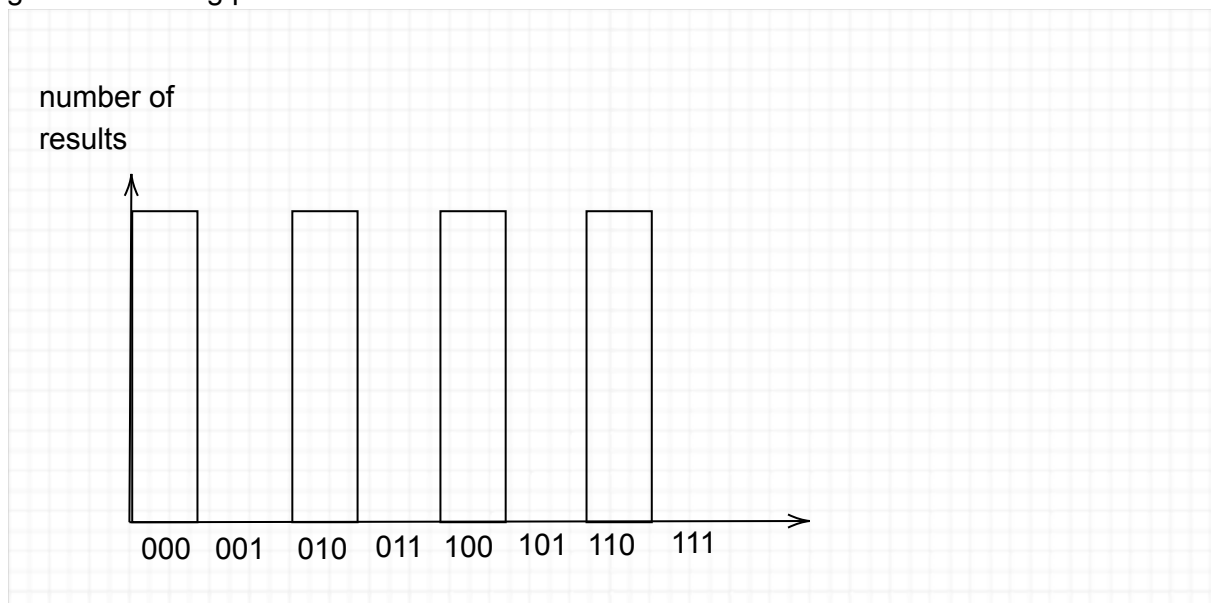
→ measuring the first register, we get one of the $\frac{s}{r}$ (for the sake of the discussion, assume that the decomposition on n bits is exact)

For instance, if $r = 4$, there are exactly 4 elements in the sum: measuring, we get $0/4$, $1/4$, $2/4$ and $3/4$.

→ On 2 bits, this is 00, 01, 10 and 11.

→ On 3 bits, this is 000, 010, 100 and 110 (since $0.x_1x_2x_3 = \frac{x_1}{2} + \frac{x_2}{4} + \frac{x_3}{8}$)

If we were to perform many measurements (for 3 bits) and collecting the results, we would get the following plot



with equiprobable results 000, 010, 100 and 110.

If the decomposition were not exact (for instance when $r = 3$), we would instead get a less precise plot with 3 peaks but not as sharp. Possibly then 3 bits would not be enough to distinguish them, and we would need to get to 5 or 6 bits of precision.

In any case, when the precision is high enough, one can "read out" the period r of $a \bmod N$ from the plot.

→ this is what we shall do in the lab session !

However, to conclude, how to get out the r out of an estimate of s/r for some s ? This can be done with the algorithm of continued fractions. See e.g.
https://en.wikipedia.org/wiki/Continued_fraction#Best_rational_approximations