

Modèle [Tordieu & Seyfried, PRE (2014)]

- Modèle de Pipes (1953)

$$\dot{x}_i(t) = V(\Delta x_i(t))$$

- Modèle de Newell (1961)

$$\dot{x}_i(t+\tau) = V(\Delta x_i(t))$$

- Modèle de Bando (1995)

we take the model of Newell 1961 and expand the left hand side =

$$\left. \begin{aligned} v_i(t+\tau) &\stackrel{\text{Taylor}}{=} v_i(t) + \tau \frac{dv_i}{dt}(t) \\ &\stackrel{\text{model}}{=} V(\Delta x_i) \end{aligned} \right\} \Rightarrow \frac{dv_i}{dt}(t) = \frac{1}{\tau} [V(\Delta x_i) - v_i]$$

2nd order (acceleration).

Modèle [Tordieu & Seyfried (2014)]:

Instead of expanding the l.h.s, we expand the r.h.s in:

$$\dot{x}_i(t) = V(\Delta x_i(t-\tau))$$

$$= V[\Delta x_i(t) - \tau \underbrace{\dot{\Delta x}_i(t)}_{\dot{x}_{i-1} - \dot{x}_i}]$$

$$\dot{x}_i(t) = V[\Delta x_i(t) - \tau [V(\Delta x_{i-1}(t)) - V(\Delta x_i(t))]]$$

implicit in speed ; 2 predecessors ; first order (velocity and not acceleration)
sometimes called generalized optimal velocity model (GOV).

[Tordoux & Seyfried (2014)]

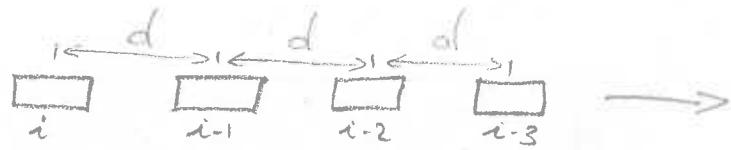
$$\dot{x}_i(t) = V \left[\Delta x_i(t) - \tau \left[V(\Delta x_{i-1}(t)) - V(\Delta x_i(t)) \right] \right]$$

Local stability

- Uniform solⁿ

$$\Delta x_k^u = d \quad \forall k$$

$$v_k^u = V(d) \quad \forall k$$



- Small perturbation applied on car i

$$x_i(t) = x_i^u(t) + y(t)$$

$$= x_{i-1}^u(t) - d + y(t)$$

$$\begin{aligned} \Delta x_i(t) &\equiv x_{i-1}(t) - x_i(t) \\ &= d - y(t) \end{aligned}$$

↓ derivⁿ

$$\dot{x}_i(t) = V(d) + \dot{y}(t)$$

We put this in the model:

$$V(d) + \dot{y}(t) = V \left[d - y(t) - \tau \left[V(d) - V(d - y(t)) \right] \right]$$

$$= V \left[d - y - \tau \underbrace{V(d) - V(d - y)}_{\text{Taylor expansion (y small)}} \right]$$

$$\stackrel{\approx}{=} V(d) - y(1 + \tau V'(d)) V'(d)$$

$$\Leftrightarrow \dot{y}(t) = - \underbrace{y(1 + \tau V'(d)) V'(d)}_{\equiv \alpha}$$

$$\Rightarrow y = y_0 e^{-\alpha t}$$

stable if $\alpha > 0$

→ always stable if $V'(d) > 0$

(the solⁿs $V'(d) < -\frac{1}{\tau}$ are not really physical).

Exponential convergence with no oscillaⁿ, even for $\tau > 0$ large

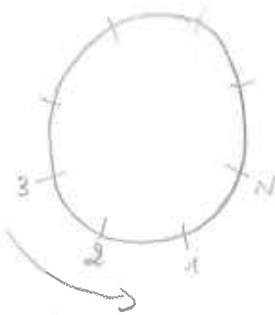
→ collision free if $V'(d) > 0$.

On a ring =

Uniform solve Θ U:

$$d_i = d = \frac{L}{N}$$

$$v_i = v = V(d)$$



$$\Delta x_i^u = x_{i-1}^u(t) - x_i^u(t) = d \quad \forall i$$

$$\Delta x_1^u = x_N^u - x_1^u + L \quad (\text{up to periodic BCs})$$

Small perturbation Θ : in all cases

$$y_i(t) = x_i(t) - x_i^u(t) \rightarrow \Delta x_i(t) = x_{i-1} - x_i = (y_{i-1} + x_{i-1}^u) - (y_i + x_i^u) = \Delta y_i + d$$

$$\dot{y}_i(t) = \dot{x}_i(t) - v = V[\Delta x_i(t) - \tau[V(\Delta x_{i-1}(t)) - V(\Delta x_i(t))]] - v$$

$$\dot{y}_i + v = V[\Delta y_i + d - \tau[V(d + \Delta y_{i-1}) - V(d + \Delta y_i)]]$$

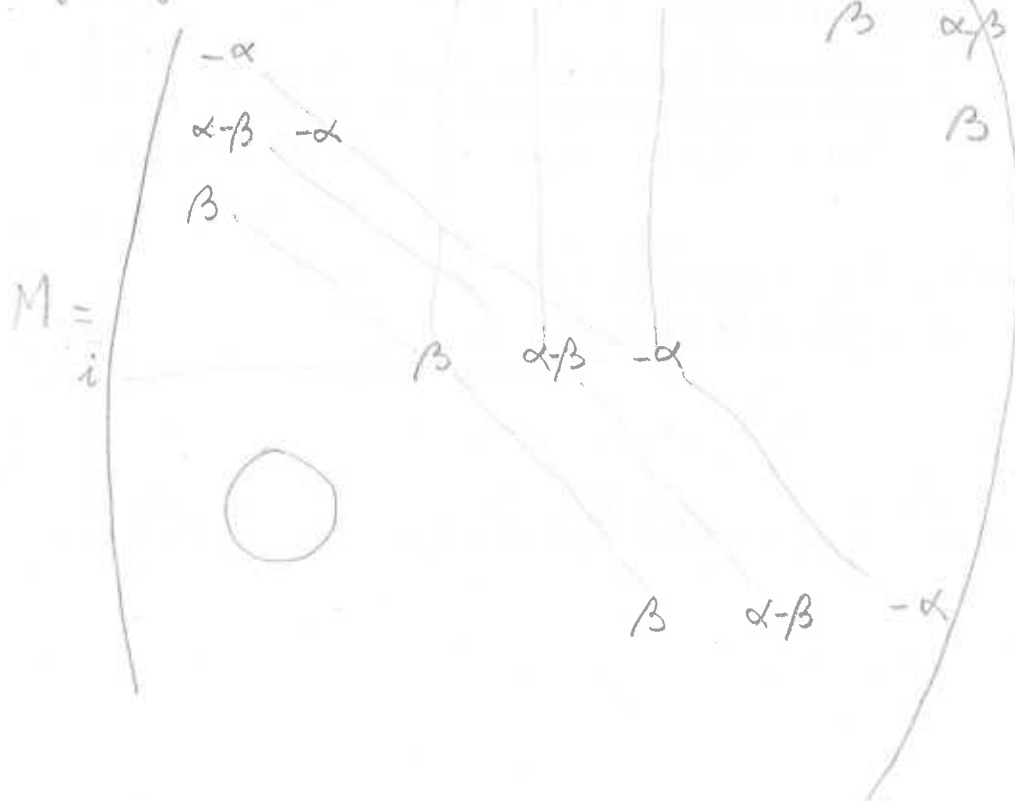
and Δy_i small

$$\hookrightarrow \dot{y}_i + v = V[d + \Delta y_i - \tau V(d) + \tau V(d) - \tau \Delta y_{i-1} V'(d) + \tau \Delta y_i V'(d)]$$

$$= \underbrace{V(d)}_v + [\Delta y_i - \tau \Delta y_{i-1} V'(d) + \tau \Delta y_i V'(d)] V'(d)$$

$$\dot{y}_i = [1 + \tau V'(d)] V'(d) \Delta y_i - \tau [V'(d)]^2 \Delta y_{i-1}$$

$$\dot{y} = M y$$



Eigenvalues:

Preliminary results =

$$\text{Let } K = \begin{pmatrix} 0 & & & 1 \\ 1 & & & \\ & \ddots & & \\ 0 & & 1 & 0 \end{pmatrix}$$

$N \times N$ matrix

$$K^2 = \begin{pmatrix} 0 & & & 2 & 0 \\ 0 & 1 & & & \\ & \ddots & \ddots & & \\ 0 & & & 1 & 0 \\ & & & & \ddots \end{pmatrix}$$

$$K^3 = \begin{pmatrix} 0 & & & 3 & 0 & 0 \\ 0 & 1 & & & & \\ & \ddots & \ddots & & & \\ 0 & & & 1 & 0 & \\ & & & & \ddots & \end{pmatrix}$$

• $K^N = \text{Id}$

Eigenvalues for K : $\mu_k = \mu^k$ with $\mu = e^{\frac{2i\pi}{N}}$

• Each circulant matrix can be expressed as a polynomial of matrix K .

• Here: $M = -\alpha \text{Id} + (\alpha - \beta)K + \beta K^2$

Eigenvalues of M :

$$\begin{aligned} \lambda_k &= -\alpha + (\alpha - \beta)\mu_k + \beta\mu_k^2 \\ &= \alpha(\mu_k - 1) + \beta\mu_k(\mu_k - 1) \end{aligned}$$

$\lambda_0 = 0$ ($\mu_0 = 1$) (associated to the stationary state)

$\text{Re}(\lambda_k)$? $\forall k = 1$ to $N-1$

$$\text{Re}(\lambda_k) = \alpha \left(\cos \frac{2k\pi}{N} - 1 \right) + \beta \left(\cos^2 \frac{2k\pi}{N} - \sin^2 \frac{2k\pi}{N} - \cos \frac{2k\pi}{N} \right)$$

$$= \alpha (C_k - 1) + \beta (C_k^2 - 1 + C_k^2 - C_k)$$

$$= [1 + \tau V'(d)] V'(d) (C_k - 1) - \tau [V'(d)]^2 (2C_k^2 - 1 - C_k)$$

$$= V'(d) (C_k - 1) + \tau [V'(d)]^2 (2C_k - 2C_k^2)$$

$$= V'(d) (1 - C_k) [-1 + 2\tau (V'(d))^2 C_k]$$

$$Re(\lambda_k) = V'(d) \overbrace{(1 - C_k)}^{\geq 0} (-1 + \tau V'(d) 2 C_k)$$

• If $V'(d) > 0$, the sign of $Re(\lambda_k)$ is the same as

It should be so!

$$Re(\lambda_k) \stackrel{\text{sign}(\tau V'(d) 2 C_k - 1)}{\geq 0} \quad \forall k$$

$$\Leftrightarrow \tau V'(d) 2 C_k - 1 < 0 \quad \forall k$$

$$\Leftrightarrow \tau V'(d) C_k < \frac{1}{2} \quad \forall k$$

This is ^{always} the case if $\tau V'(d) < \frac{1}{2}$

• If $V'(d) \leq 0$, $Re(\lambda_k) \stackrel{?}{\leq} 0 \quad \forall k$

$$\Leftrightarrow \tau V'(d) 2 C_k \geq 1 \quad \forall k$$

$$\Leftrightarrow \underbrace{\tau V'(d)}_{> 0} \underbrace{C_k}_{< 0} \geq \frac{1}{2} \quad \forall k$$

Not possible to ensure for all $k / C_k \geq 0$, i.e. if $N > 3$

→ global stability for long lines of vehicles cannot be achieved.



CCL: for global stability we must have

$$\boxed{\begin{array}{l} V'(d) > 0 \\ \text{and} \\ \tau V'(d) < \frac{1}{2} \end{array}}$$

If $V'(d) > 0$:

If the condition does not hold, the largest N for which we have stability is such that

$$\cos \frac{2\pi}{N} < \frac{1}{2\tau V'(d)}$$

↘ \cos^{-1} is decreasing

$$\frac{2\pi}{N} > \cos^{-1}\left(\frac{1}{2\tau V'(d)}\right)$$

$$\boxed{N < \frac{2\pi}{\cos^{-1}\left(\frac{1}{2\tau V'(d)}\right)}}$$

