

# Lecture 12 — Finite Potential Well & Quantum Tunneling

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**Reading:** Engel 4th ed., Chapter 4 (Sections 4.5–4.7)

## Learning Objectives

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- Contrast the finite and infinite potential wells qualitatively and quantitatively
  - Explain the physical meaning of wavefunction penetration into classically forbidden regions
  - Derive and apply the transmission coefficient for rectangular barriers
  - Describe quantum tunneling and its role in chemistry (proton transfer, STM, nuclear fusion)
  - Recognize when tunneling is significant vs. negligible
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## 1. The Finite Potential Well

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### The Potential

$$V(x) = \begin{cases} 0 & |x| \leq L/2 \\ V_0 & |x| > L/2 \end{cases}$$

where  $V_0$  is finite (not infinite).

### Qualitative Differences from the Infinite Well

Feature	Infinite well	Finite well
Wavefunction at walls	Exactly zero	Nonzero (penetrates into barriers)
Number of bound states	Infinite	Finite (depends on $V_0$ and $L$ )
Energies	$E_n = n^2 h^2 / (8ma^2)$	Lower than infinite well values
"Effective" box size	Exactly $L$	Greater than $L$ (particle "leaks" out)

**Inside the Well ( $|x| < L/2, E < V_0$ ):**

$$\psi(x) = A \cos(kx) + B \sin(kx), \quad k = \frac{\sqrt{2mE}}{\hbar}$$

**Outside the Well ( $|x| > L/2$ ):**

$$\psi(x) = Ce^{-\kappa x} + De^{\kappa x}, \quad \kappa = \frac{\sqrt{2m(V_0-E)}}{\hbar}$$

Physical requirement:  $\psi \rightarrow 0$  as  $|x| \rightarrow \infty$  (normalizability).

### Key Result: Wavefunction Penetration

The wavefunction decays exponentially in the classically forbidden region with a **penetration depth**:

$$\delta = \frac{1}{\kappa} = \frac{\hbar}{\sqrt{2m(V_0 - E)}}$$

- Lighter particles penetrate farther (smaller  $m$ )
- Particles with energy close to  $V_0$  penetrate farther
- The probability of finding the particle in the forbidden region is nonzero but small

### Matching Conditions

At  $x = \pm L/2$ , we require:

1.  $\psi$  is continuous
2.  $d\psi/dx$  is continuous

These matching conditions lead to transcendental equations that determine the allowed energies. They must be solved graphically or numerically — no closed-form solution exists.

### Bound State Count

The number of bound states in a 1-D finite well of depth  $V_0$  and width  $L$ :

$$N \approx 1 + \text{floor}\left(\frac{L\sqrt{2mV_0}}{\pi\hbar}\right)$$

There is **always at least one** bound state in a 1-D finite well (not true in 3-D!).

[!NOTE] **Concept Check 12.1** In the finite potential well, as  $V_0 \rightarrow \infty$ , what happens to the penetration depth  $\delta$ ? Show that the energies of the finite well approach the energies of the infinite well in this limit.

## 2. Quantum Tunneling

### The Rectangular Barrier

$$V(x) = \begin{cases} V_0 & 0 \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$

A particle with energy  $E < V_0$  approaches from the left. Classically, it would be completely reflected. Quantum mechanically, there is a nonzero probability of transmission.

### The Transmission Coefficient

For  $E < V_0$ :

$$T \approx e^{-2\kappa a}$$

where  $\kappa = \frac{\sqrt{2m(V_0-E)}}{\hbar}$  and  $a$  is the barrier width.

More precisely:

$$T = \frac{1}{1 + \frac{V_0^2 \sinh^2(\kappa a)}{4E(V_0 - E)}}$$

For thick barriers ( $\kappa a \gg 1$ ), the exponential approximation is excellent.

### What Controls Tunneling?

$$T \propto e^{-2a\sqrt{2m(V_0-E)}/\hbar}$$

Tunneling probability increases when:

- **Barrier is narrow** (small  $a$ )
- **Barrier is low** ( $V_0 - E$  small)
- **Particle is light** (small  $m$ )

[!NOTE] **Concept Check 12.2** If you double the width of a tunneling barrier ( $a \rightarrow 2a$ ), while keeping the height  $V_0$  and the particle's energy  $E$  constant, how does the transmission coefficient  $T$  change? (Assume the barrier is "thick" so the exponential approximation holds.)

System	$m$	Tunneling?
Electron through 1 nm oxide	$m_e$	Significant
Proton through H-bond barrier	$m_p$	Measurable
Deuteron (same barrier)	$2m_p$	Reduced by $e^{-\sqrt{2}}$
Carbon atom	$12m_p$	Negligible

### 3. Chemical Applications of Tunneling

#### Scanning Tunneling Microscope (STM)

Electrons tunnel from a sharp metal tip across a vacuum gap ( $\sim 1 \text{ nm}$ ) to a surface. The tunneling current depends exponentially on the tip-surface distance:

$$I \propto e^{-2\kappa d}$$

This extreme distance sensitivity (factor of  $\sim 10$  per  $0.1 \text{ nm}$ ) enables **atomic-resolution** imaging. Gerd Binnig and Heinrich Rohrer won the 1986 Nobel Prize for the STM.

#### Proton Tunneling in Chemistry

- **Proton transfer reactions:** Tunneling contributes significantly to reaction rates, especially at low temperatures
- **Kinetic isotope effect:** Replacing H with D reduces tunneling  $\rightarrow$  measurable rate decrease. The ratio  $k_H/k_D > 1$  is a signature of tunneling
- **DNA mutations:** Some theories suggest proton tunneling between base pairs could cause rare tautomeric forms

#### Nuclear Fusion

In the Sun's core, protons tunnel through Coulomb barriers to fuse — classical energies are far too low. Without tunneling, stars would not shine.

## Alpha Decay

George Gamow (1928) first explained radioactive alpha decay as tunneling of an alpha particle through the nuclear potential barrier.

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## 4. The WKB Approximation (Brief)

For barriers of arbitrary shape  $V(x)$ , the transmission coefficient can be approximated as:

$$T \approx \exp\left(-\frac{2}{\hbar} \int_{x_1}^{x_2} \sqrt{2m[V(x) - E]} dx\right)$$

where  $x_1$  and  $x_2$  are the classical turning points. This **WKB (Wentzel-Kramers-Brillouin)** approximation is widely used in chemistry for reaction rate calculations.

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## Key Equations Summary

Equation	Expression
Penetration depth	$\delta = \hbar / \sqrt{2m(V_0 - E)}$
Transmission coefficient	$T \approx e^{-2\kappa a}$
Decay constant	$\kappa = \sqrt{2m(V_0 - E)} / \hbar$
WKB approximation	$T \approx \exp\left(-\frac{2}{\hbar} \int_{x_1}^{x_2} \sqrt{2m(V - E)} dx\right)$

## Recent Literature Spotlight

"Recent Developments in Symmetry-Adapted Perturbation Theory" *K. Patkowski*, WIREs Computational Molecular Science, 2020, 10, e1452. [DOI](#)

This comprehensive review describes how quantum mechanical tunneling and finite-barrier effects are handled in symmetry-adapted perturbation theory (SAPT), the premier method for computing intermolecular interaction energies from first principles. The decomposition into electrostatics, induction, dispersion, and exchange — each arising

from different quantum mechanical operators acting on finite-barrier wavefunctions — demonstrates the power of perturbation theory for chemical problems.

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## Practice Problems

1. **Finite well.** An electron is in a finite well of depth  $V_0 = 10$  eV and width  $L = 0.50$  nm.  
(a) Calculate the penetration depth  $\delta$  for the ground state energy  $E \approx 0.6$  eV. (b) How does  $\delta$  change if the particle is a proton instead?
  2. **Tunneling probability.** An electron with  $E = 5.0$  eV encounters a rectangular barrier of height  $V_0 = 10.0$  eV and width  $a = 0.50$  nm. Calculate the transmission coefficient. Repeat for  $a = 1.0$  nm and  $a = 2.0$  nm.
  3. **Isotope effect.** A proton and a deuteron each encounter the same barrier ( $V_0 - E = 0.4$  eV,  $a = 0.05$  nm). Calculate the ratio  $T_H/T_D$ . This is related to the kinetic isotope effect in chemistry.
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Next week: *Particle-in-a-Box Applications; Operators & Entanglement*