The TOV Equation and the Mass of Stars

Jonas Pleyer

March 3, 2021

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- Questions to answer:
 - Well defined radius?
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- Mathematically interesting: show that a solution of the differential equation has a zerovalue without knowing the solution

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- 1. General Relativity
- 2. Thermodynamics
- 3. Numerical Solutions
- 4. Exact Results
- 5. Outlook

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- 1.1 Concepts
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Figure: Third servicing mission of the Hubble Telescope 1993 [NAS99].

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$$G_{\mu\nu} + \Lambda g_{\mu\nu} = R_{\mu\nu} + \left(\frac{1}{2}R + \Lambda\right)g_{\mu\nu} = 8\pi T_{\mu\nu}$$

• Heavy objects create curvature in space

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Energy-Momentum Tensor of perfect fluid

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Solve Einstein Equations (without cosm. constant)

$$G_{\mu\nu} = R_{\mu\nu} + \frac{1}{2} R g_{\mu\nu} = 8\pi T_{\mu\nu} \tag{3}$$

$$-8\pi T_0^0 = 8\pi \rho = \frac{\lambda' e^{-\lambda}}{r} + \frac{1 - e^{-\lambda}}{r^2} \tag{4}$$

$$8\pi T_1^1 = 8\pi p = \nu' \frac{e^{-\lambda}}{r} - \frac{1 - e^{-\lambda}}{r^2}$$
 (5)

$$8\pi T_2^2 = 8\pi p = \frac{e^{-\lambda}}{2} \left[\nu'' + \left(\frac{\nu'}{2} + \frac{1}{r} \right) (\nu' - \lambda') \right]$$
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Use equation 4 and identify Mass m(r)

$$e^{-\lambda} = 1 - \frac{2}{r} \int_{0}^{r} 4\pi \rho(r') r'^{2} dr' =: 1 - \frac{2m(r)}{r}$$
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Divergence of Energy-Momentum Tensor is $\nabla_{\mu}T^{\mu\nu}=0$. Then obtain

$$\frac{\partial p}{\partial r} = -\frac{p+\rho}{2} \frac{\partial \nu}{\partial r} \tag{8}$$

$$\frac{\partial m}{\partial r} = 4\pi \rho r^2 \tag{9}$$

$$\frac{\partial p}{\partial r} = -\frac{Gm\rho}{r^2} \left(1 + \frac{p}{\rho c^2} \right) \left(\frac{4\pi r^3 p}{mc^2} + 1 \right) \left(1 - \frac{2Gm}{rc^2} \right)^{-1} \tag{10}$$

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- Needs equation of state (EoS) $f(\rho, p, r) = 0$ to be solvable

Non-relativistic Limit of 2nd TOV equation 10 is

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Obtain Lane-Emden equation [Lan70; Emd07]

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• Some exact solutions are known [Cha58]

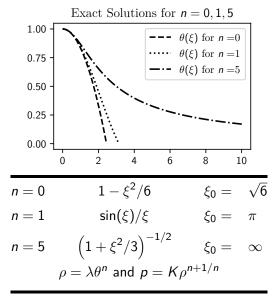


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Calculating an EoS

Statistical theory of manyparticle systems



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- Partition function contains all information about N particles with position $x_i \in M$ and momentum $p_i \in T_{x_i}M$ with V = vol(M)

$$\mathcal{Z}(T, V, N) = \int_{TM^N} \exp\left(-\frac{H(x_1, \dots, p_N)}{k_B T}\right) \frac{dx_1 \dots dp_N}{N! h^{3N}}$$
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ullet Calculate equation of state via internal Energy ${\cal U}$ and

$$p = k_B T \frac{1}{Z} \frac{\partial Z}{\partial V} \qquad \rho = \frac{U}{V} = \frac{k_B T^2}{V} \frac{\partial Z}{\partial T}$$
 (15)

• Partition Function for $H = \sqrt{m^2 + p^2}$

$$\mathcal{Z} = \frac{1}{N!} \left(8\pi V \left(\frac{k_B T}{hc} \right)^3 \frac{\alpha^2 K_2(\alpha)}{2} \right)^N \tag{16}$$

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$$p = \frac{Nk_BT}{V} = CNmc^2 \frac{1}{K_2(\alpha)\alpha^2} \exp\left(-\alpha \frac{K_1(\alpha) + K_3(\alpha)}{2K_2(\alpha)}\right)$$
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The function $p(\alpha)$ above is a bijection.

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Theorem

The function $p(\alpha)$ above is a bijection.

- ullet Sketch of proof: take differential, use properties/representations of $K_
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- Use inverse and previous equations 15, 16 and 17 to obtain EOS

$$\rho = \frac{\mathcal{U}}{V} = p \left(1 + \alpha(p) \frac{K_1(\alpha(p)) + K_3(\alpha(p))}{2K_2(\alpha(p))} \right)$$
(18)

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- 3.2 Verifying LE Results
- 3.3 Zero Values of TOV and LE equation
- 3.4 TOV Hypothesis
- 4. Exact Results
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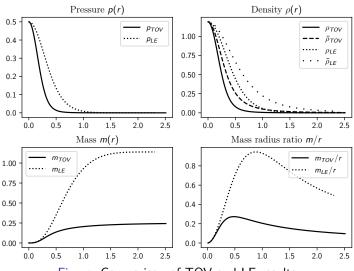


Figure: Comparison of TOV and LE results

Verifying LE Results

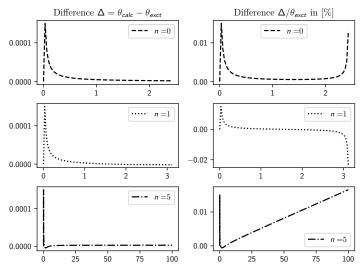


Figure: Validation of numerically calculated Lane-Emden results

Zero Values of TOV and LE equation

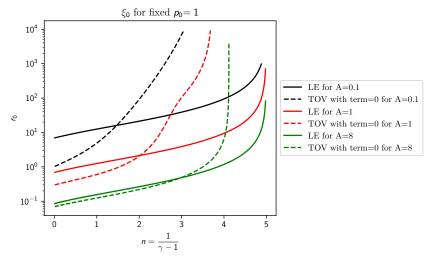


Figure: TOV and LE results for varying A parameter of $\rho = Ap^{n/(n+1)}$

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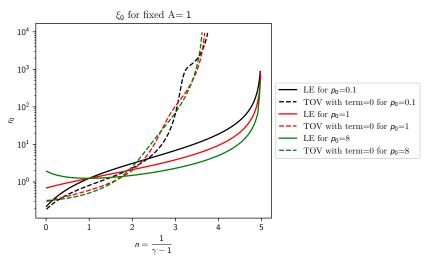


Figure: TOV and LE solutions for varying parameter p_0 .

Intersection: $r = \beta \xi$ with $4\pi \beta A^{n/(n+1)} = (n+1)\lambda^{1-1/n}$ is independent of $\lambda = \rho(p_0)$.

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TOV Hypothesis

Hypothesis

Given the TOV differential equation with $\rho = Ap^{\frac{n}{n+1}}$ and $p_0, A > 0$

$$\begin{split} &\frac{\partial m}{\partial r} = 4\pi \rho r^2 \\ &\frac{\partial p}{\partial r} = -\frac{m\rho}{r^2} \left(1 + \frac{p}{\rho}\right) \left(\frac{4\pi r^3 p}{m} + 1\right) \left(1 - \frac{2m}{r}\right)^{-1} \end{split}$$

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- (i) There exists a solution for some $n_0 \ge 0$
- (ii) All solutions with same parameters A, p_0 and smaller exponent $n < n_0$ have a $p(r_0)$ for some $r_0 > 0$.

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- 4.1 New Exact Solution for LE at n=2
- 4.2 Hypothesis
- 4.3 Limiting Case TOV
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• Found new solution for n=2 by using simple power-series $\theta=\sum a_m\xi^m$

$$\frac{1}{\xi^2} \frac{\partial}{\partial \xi} \left(\xi^2 \frac{\partial \theta}{\partial \xi} \right) + \theta^2 = 0 \tag{19}$$

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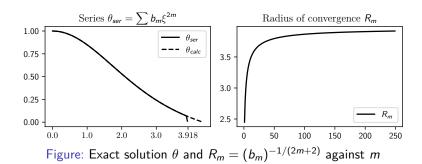
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The series $\theta = \sum_{m=0}^{\infty} b_m \xi^{2m}$ converges for $\xi < 1$.

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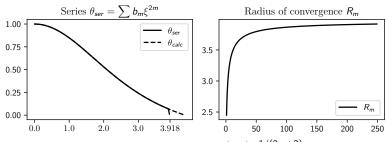


Figure: Exact solution θ and $R_m = (b_m)^{-1/(2m+2)}$ against m

Hypothesis

The radius of convergence R > 1 and $\exists \xi_0 \geq R$ such that $p(\xi_0) = 0$.

Hypothesis

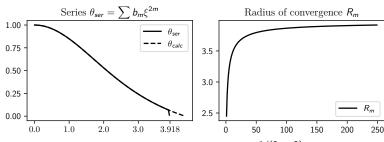


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Hypothesis

Let $n \ge 0$ and $\theta_n = \sum a_m \xi^m$ be a LE solution. Then $a_{2m+1} = 0$ for $m \in \mathbb{N}$.

Limiting Case TOV

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Theorem

Let p_A be a solution of the TOV equation with $\rho = Ap^{1/\gamma}$. Then

$$\lim_{A \to 0} p_A = \frac{p_0}{2\pi p_0 r^2 + 1} \tag{21}$$

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Proof.

With $\partial m/\partial r = 4\pi A p^{1/\gamma} r^2$, define v := m/A

$$\frac{\partial p}{\partial r} = -\frac{p^{1/\gamma}}{r^2} \left(A + p^{1-1/\gamma} \right) \left(4\pi r^3 p + vA \right) \left(1 - \frac{2vA}{r} \right)^{-1} \tag{22}$$



Proof (Cont ...)

Then for A = 0, we have

$$\frac{\partial p}{\partial r} = -4\pi r p^2 \tag{23}$$

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- Problem: Assume that TOV equation is solvable
- Simple transformation tricks for 1D singular ODEs not working.

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- Solvability of TOV equation is complicated and subject of current research [Mar+19; BVW07]

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