### Jonas Teuwen

Delft Institute of Applied Mathematics, Delft University of Technology, P.O. Box 5031, 2600 GA Delft, The Netherlands

#### Abstract

This note presents a proof that the non-tangential maximal function of the Ornstein-Uhlenbeck semigroup is bounded almost surely by the Gaussian Hardy-Littlewood maximal function. In particular this entails improvement on a result by Pineda and Urbina [1] who proved a similar result for a 'trunctated' version of the non-tangential maximal function. We actually obtain boundedness of the maximal function on non-tangential cones of arbitrary aperture.

Keywords: ornstein-uhlenbeck, ornstein-uhlenbeck semigroup

#### 1. Introduction

Maximal functions are among the most studied objects in harmonic analysis. It is well-known that the classical non-tangential maximal function associated with the heat semigroup is bounded almost everywhere by the Hardy-Littlewood maximal function,

$$\sup_{\substack{(y,t)\in\mathbf{R}_{+}^{d+1}\\|x-y|< t}} |e^{-t\Delta}u(y)| \lesssim \sup_{r>0} \frac{1}{|B_{r}(x)|} \int_{B_{r}(x)} |u| \,d\lambda. \tag{1}$$

Here the action of heat semigroup  $e^{-t\Delta}u = \rho_t * u$  is given by a convolution of u with the heat kernel

$$\rho_t(\xi) := \frac{e^{-|\xi|^2/4t}}{(4\pi t)^{\frac{d}{2}}}, \text{ with } t > 0 \text{ and } \xi \in \mathbf{R}^d.$$

In this note we are interested in its Gaussian counterpart. The change from Lebesgue measure to the Gaussian measure

$$d\gamma(x) := \pi^{-\frac{d}{2}} e^{-|x|^2} dx \tag{2}$$

Preprint submitted to Indagationes Mathematicae

November 20, 2013

introduces quite some intricate technical and conceptual difficulties which are due to its non-doubling nature. Instead of the Laplacian, will use its Gaussian analogue, the  $Ornstein-Uhlenbeck\ operator\ L$  which is given by,

$$L := -\frac{1}{2}\Delta + \langle x, \nabla \rangle = \frac{1}{2}\nabla_{\gamma}^* \nabla_{\gamma}, \tag{3}$$

where  $\nabla_{\gamma}$  denotes the realisation of the gradient in  $L^2(\mathbf{R}^d, \gamma)$ . Our main result, to be proved in Theorem 1, is the following Gaussian analogue of (1):

$$\sup_{(y,t)\in\Gamma_x^{(A,a)}} |e^{-t^2L}u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \, d\gamma. \tag{4}$$

Here,

$$\Gamma_x^{(A,a)} := \Gamma_x^{(A,a)}(\gamma) := \{ (y,t) \in \mathbf{R}_+^{d+1} : |x-y| < At \text{ and } t \leqslant am(x) \}$$
 (5)

is the  $Gaussian\ cone$  with aperture A and cut-off parameter a, and

$$m(x) := \min\left\{1, \frac{1}{|x|}\right\}. \tag{6}$$

A slighly weaker version of the inequality (4) has been proved by Pineda and Urbina [1] who showed that

$$\sup_{(y,t)\in\widetilde{\Gamma}_x} |e^{-t^2L}u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \,d\gamma,$$

where

$$\widetilde{\Gamma}_x(x) = \{(y, t) \in \mathbf{R}^d_+ : |x - y| < t \leqslant \widetilde{m}(x)\}$$

is the 'reduced' Gaussian cone corresponding to the function

$$\widetilde{m}(x) = \min\left\{\frac{1}{2}, \frac{1}{|x|}\right\}.$$

Our proof of (4) is much shorter and, we believe, more transparent than the one presented in [1]. It has the further advantage of allowing the extension to cones with arbitrary aperture A > 0 and cut-off parameter a > 0 without any additional technicalities. This additional generality is very important and has already been used by Portal (cf. the claim made in [2, discussion preceding Lemma 2.3]) to prove the  $H^1$ -boundedness of the Riesz transform associated with L.

#### 2. The Mehler kernel

The Mehler kernel (see e.g., [3]) is the Schwartz kernel associated to the Ornstein-Uhlenbeck semigroup  $(e^{-tL})_{t\geqslant 0}$ , that is,

$$e^{-tL}u(x) = \int_{\mathbf{R}^d} M_t(x, \cdot)u \, d\gamma. \tag{7}$$

There is an abundance of literature on the Mehler kernel and its properties. We shall only use the fact, proved e.g. in the review paper [3], that it is given explicitly by

$$M_t(x,y) = \frac{\exp\left(-\frac{|e^{-t}x - y|^2}{1 - e^{-2t}}\right)}{(1 - e^{-2t})^{\frac{d}{2}}} e^{|y|^2}.$$
 (8)

Note that  $M_t(x, y)$  is symmetric in x and y. A formula for (8) honoring this observation is:

$$M_t(x,y) = \frac{\exp\left(-e^{-2t} \frac{|x-y|^2}{1 - e^{-2t}}\right) \exp\left(2e^{-t} \frac{\langle x,y \rangle}{1 + e^{-t}}\right)}{(1 - e^{-t})^{\frac{d}{2}}} \frac{\exp\left(2e^{-t} \frac{\langle x,y \rangle}{1 + e^{-t}}\right)}{(1 + e^{-t})^{\frac{d}{2}}}.$$
 (9)

# 3. Some lemmata

We use m as defined in (6) in our next lemma, which is taken from [4].

**1 Lemma.** Let a, A be strictly positive real numbers and t > 0. We have for  $x, y \in \mathbf{R}^d$  that:

- 1. If |x y| < At and  $t \leq am(x)$ , then  $t \leq (1 + aA)m(y)$ ;
- 2. If |x-y| < Am(x), then  $m(x) \leqslant (1+A)m(y)$  and  $m(y) \leqslant 2(1+A)m(x)$ .

The next lemma will come useful when we want to cancel exponential growth in one variable with exponential decay in the other as long both variables are in a Gaussian cone.

**2 Lemma.** Let  $\alpha > 0$  and  $|x - y| \leq \alpha m(x)$ . Then:

$$e^{-\alpha^2 - 2\alpha} e^{|y|^2} \le e^{|x|^2} \le e^{\alpha^2 (1+\alpha)^2 + 2\alpha(1+\alpha)} e^{|y|^2}$$
.

*Proof.* By the inverse triangle inequality and  $m(x)|x| \leq 1$  we get,

$$|y|^2 \le (\alpha m(x) + |x|)^2 \le \alpha^2 + 2\alpha + |x|^2.$$

This gives the first inequality. For the second we use Lemma 1 to infer  $m(x) \leq (1 + \alpha)m(y)$ . Proceeding as before we obtain:

$$|x|^2 \le \alpha^2 (1+\alpha)^2 + 2\alpha(1+\alpha) + |y|^2$$
.

As required.

# 3.1. An estimate on Gaussian balls

Let  $B := B_t(x)$  be the open Euclidean ball with radius t and center x and let  $\gamma$  be the Gaussian measure as defined by (2). We shall denote by  $S_d$  the surface area of the unit sphere in  $\mathbf{R}^d$ .

**3 Lemma.** For all  $x \in \mathbf{R}^d$  and t > 0 we have the inequality:

$$\gamma(B_t(x)) \leqslant \frac{S_d}{2\pi^{\frac{d}{2}}} t^d e^{2t|x|} e^{-|x|^2}.$$
(10)

*Proof.* Remark that, with  $B := B_t(x)$ ,

$$\begin{split} \int_{B} \mathrm{e}^{-|\xi|^{2}} \, \mathrm{d}\xi &= \mathrm{e}^{-|x|^{2}} \int_{B} \mathrm{e}^{-|\xi - x|^{2}} \mathrm{e}^{-2\langle x, \xi - x \rangle} \, \mathrm{d}\xi \\ &\leqslant \mathrm{e}^{-|x|^{2}} \int_{B} \mathrm{e}^{-|\xi - x|^{2}} \mathrm{e}^{2|x||\xi - x|} \, \mathrm{d}\xi \\ &\leqslant \mathrm{e}^{-|x|^{2}} \mathrm{e}^{2t|x|} \int_{B} \mathrm{e}^{-|\xi - x|^{2}} \, \mathrm{d}\xi \\ &= \pi^{\frac{d}{2}} \mathrm{e}^{2t|x|} \mathrm{e}^{-|x|^{2}} \gamma(B_{t}(0)). \end{split}$$

So, there holds that

$$\gamma(B_t(x)) \leqslant e^{2t|x|} e^{-|x|^2} \gamma(B_t(0)).$$
 (11)

We will estimate the Gaussian volume of the ball  $B_t(0)$ . Using polar coordi-

nates we proceed for  $d \ge 2$  by:

$$\gamma(B_t(0)) = \frac{1}{\pi^{\frac{d}{2}}} \int_{B_t(0)} e^{-|\xi|^2} d\xi$$

$$= \frac{S_d}{\pi^{\frac{d}{2}}} \int_0^t e^{-r^2} r^{d-1} dr$$

$$\leqslant \frac{S_d}{2\pi^{\frac{d}{2}}} t^{d-2} \int_0^t 2r e^{-r^2} dr$$

$$= \frac{S_d}{2\pi^{\frac{d}{2}}} t^{d-2} (1 - e^{-t^2})$$

$$\leqslant \frac{S_d}{2\pi^{\frac{d}{2}}} t^d,$$

where the last step uses  $1 - e^{-t} \le t$  for  $t \ge 0$ . The case for d = 1 follows by a simplified argument. Upon combining this result with (11) we obtain (10).

# 3.2. On-diagonal kernel estimates on annuli

As is common in harmonic analysis, we often wish to decompose  $\mathbf{R}^d$  into sets on which certain phenomena are easier to handle. Here we will decompose the space into disjoint annuli.

Throughout this subsection we fix  $x \in \mathbf{R}^d$ , constants  $A, a \ge 1$ , a pair  $(y,t) \in \Gamma_x^{(A,a)}$ . We use the notation rB to mean the ball obtained from the ball B by multiplying its radius by r.

The annuli  $C_k := C_k(B_t(x))$  are given by:

$$C_k := \begin{cases} 2B_t(x), & k = 0, \\ 2^{k+1}B_t(x) \setminus 2^k B_t(x), & k \geqslant 1. \end{cases}$$
 (12)

Whenever  $\xi$  is in  $C_k$ , we get for  $k \geqslant 1$ :

$$2^k t < |y - \xi| \leqslant 2^{k+1} t. \tag{13}$$

On  $C_k$  we have the following bound for  $M_{t^2}(y,\cdot)$ :

**4 Lemma.** For all  $\xi \in C_k$  for  $k \geqslant 1$  we have:

$$M_{t^2}(y,\xi) \leqslant \frac{e^{|y|^2}}{(1-e^{-2t^2})^{\frac{d}{2}}} \exp(2^{k+1}t|y|) \exp\left(-\frac{4^k}{2e^{2t^2}}\right),$$
 (14)

*Proof.* Considering the first exponential which occurs in the Mehler kernel (9) together with (13) gives for  $k \ge 1$ :

$$\exp\left(-e^{-2t^2} \frac{|y-\xi|^2}{1 - e^{-2t^2}}\right) \leqslant \exp\left(-\frac{4^k}{e^{2t^2}} \frac{t^2}{1 - e^{-2t^2}}\right)$$

$$\stackrel{(\dagger)}{\leqslant} \exp\left(-\frac{4^k}{2e^{2t^2}}\right),$$

where (†) follows from  $1 - e^{-s} \le s$  for  $s \ge 0$ . Using the estimate  $1 + s \ge 2s$  for  $0 \le s \le 1$ , for the second exponential in the Mehler kernel (9) we obtain, by (13):

$$\exp\left(2e^{-t^2}\frac{\langle y,\xi\rangle}{1+e^{-t^2}}\right) \leqslant \exp(|\langle y,\xi\rangle|)$$
$$\leqslant \exp(|\langle y,\xi-y\rangle|)e^{|y|^2}$$
$$\leqslant \exp(2^{k+1}t|y|)e^{|y|^2}.$$

Combining things, we obtain the estimate in the formulation of the lemma.

## 4. The main result

In this section we will prove our main theorem for which we have already made the necessary preparations in the previous sections.

**1 Theorem.** Let A, a > 0. For all  $x \in \mathbf{R}^d$  and all  $u \in L^2(\mathbf{R}^d, \gamma)$  we have

$$\sup_{(y,t)\in\Gamma_x^{(A,a)}} |e^{-t^2L}u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \, d\gamma, \tag{15}$$

where the implicit constant only depends on A, a and d.

*Proof.* We fix  $x \in \mathbf{R}^d$  and  $(y,t) \in \Gamma_x^{(A,a)}$ . The proof of (15) is based on splitting the integration domain into the annuli  $C_k$  as defined by (12) and estimating on each annulus. More explicit,

$$|e^{-t^2L}u(y)| \leq \sum_{k=0}^{\infty} I_k(y), \text{ where } I_k(y) := \int_{C_k} M_{t^2}(y,\cdot)|u(\cdot)| d\gamma.$$
 (16)

We have  $t \leq am(x) \leq a$  and, by Lemma 1,  $t|y| \leq 1 + aA$ . Together with Lemma 4 we infer, for  $\xi \in C_k$  and  $k \geq 1$ , that:

$$M_{t^{2}}(y,\xi) \leqslant \frac{e^{|y|^{2}}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} \exp(2^{k+1}(1 + aA)) \exp\left(-\frac{4^{k}}{2e^{2a^{2}}}\right)$$
$$=: \frac{e^{|y|^{2}}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} c_{k}.$$

Combining this with Lemma 2, we obtain

$$M_{t^2}(y,\xi) \lesssim_{A,a} \frac{e^{|x|^2}}{(1-e^{-2t^2})^{\frac{d}{2}}} c_k.$$
 (17)

Also, by (13),

$$|x - \xi| \le |x - y| + |\xi - y| \le (2^{k+1} + A)t$$
.

Let K be the smallest integer such that  $2^{k+1} \ge A$  whenever  $k \ge K$ . Then it follows that  $C_k$  for  $k \ge K$  is contained in  $B_{2^{k+2}t}(x)$  and for k < K is contained in  $B_{2At}(x)$ . We set,

$$D_k := D_k(x) = \begin{cases} B_{2^{k+2}t}(x) & \text{if } k \geqslant K, \\ B_{2At}(x) & \text{elsewhere.} \end{cases}$$

Let us denote the supremum on right-hand side of (15) by  $M_{\gamma}u(x)$ . Using (17), we can bound the integral on the right-hand side of (16) by

$$\int_{C_k} M_{t^2}(y, \cdot) |u(\cdot)| \, d\gamma \lesssim_{A,a} c_k \frac{e^{|x|^2}}{(1 - e^{-2t^2})^{\frac{d}{2}}} \int_{C_k} |u| \, d\gamma 
\leqslant c_k \frac{e^{|x|^2}}{(1 - e^{-2t^2})^{\frac{d}{2}}} \int_{D_k} |u| \, d\gamma 
\leqslant c_k \frac{e^{|x|^2}}{(1 - e^{-2t^2})^{\frac{d}{2}}} \gamma(D_k) M_{\gamma} u(x),$$

where we pause for a moment to compute a suitable bound for  $\gamma(D_k)$ . As above we have both  $t|x| \leq am(x)|x| \leq a$  and  $t \leq a$ . Together with Lemma 3 applied to  $D_k$  for  $k \geq K$  we obtain:

$$\gamma(D_k) e^{|x|^2} \leqslant_A C^d S_d t^d 2^{kd} e^{2^{k+2}t|x|} e^{-|x|^2} e^{|x|^2}$$

$$\lesssim_{A,a,d} t^d 2^{kd} e^{2^{k+2}a}.$$

Similarly, for k < K:

$$\gamma(D_k)e^{|x|^2} \lesssim_{A,a,d} t^d e^{2Aa}$$
.

Using the bound  $t \leq a$ , we can infer that

$$\frac{t^d}{(1 - e^{-2t^2})^{\frac{d}{2}}} \leqslant \frac{a^d}{(1 - e^{-2a^2})^{\frac{d}{2}}} \lesssim_{a,d} 1.$$

(note that  $s/(1 - e^{-s})$  is increasing). Combining these computations with where we left off we get for  $k \ge K$ ,

$$\int_{C_k} M_{t^2}(y,\cdot)|u(\cdot)| \, d\gamma \lesssim_{A,a,d} c_k 2^{kd} e^{2^{k+2}a} M_{\gamma} u(x),$$

while we get for k < K,

$$\int_{C_h} M_{t^2}(y,\cdot)|u(\cdot)| \, d\gamma \lesssim_{A,a,d} c_k M_{\gamma} u(x).$$

Similarly, for  $\xi \in 2B_t(x)$  we obtain:

$$I_0 := \int_{2R} M_{t^2}(y, \cdot) |u(\cdot)| \, \mathrm{d}\gamma \lesssim_{A, a, d} M_{\gamma} u(x).$$

Inserting the dependency of  $c_k$  upon k as coming from (17), we obtain the bound:

$$|e^{-t^{2}L}u(y)| = I_{0} + \sum_{k=1}^{K-1} I_{k} + \sum_{k=K}^{\infty} I_{k}$$

$$\lesssim_{A,a,d} \left[ 1 + \sum_{k=1}^{K-1} c_{k} + \sum_{k=K}^{\infty} c_{k} 2^{kd} e^{2^{k+2}a} \right] M_{\gamma}u(x),$$

$$\lesssim_{A,a,d} \left[ 1 + \sum_{k=1}^{K-1} e^{-\frac{4^{k}}{2e^{2a^{2}}}} + \sum_{k=K}^{\infty} 2^{kd} e^{2^{k+1}(1+2a+aA)} e^{-\frac{4^{k}}{2e^{2a^{2}}}} \right] M_{\gamma}u(x),$$

valid for all  $y \in \Gamma_x^{(A,a)}$ . As the sum on the right-hand side evidently converges, we see that taking the supremum proves (15).

#### References

- [1] E. Pineda, W. R. Urbina, Non Tangential Convergence for the Ornstein-Uhlenbeck Semigroup, Divulgaciones Matemáticas 13 (2) (2008) 1-19. URL http://www.emis.ams.org/journals/DM/v16-1/art7.pdf
- P. Portal, Maximal and quadratic Gaussian Hardy spacesarXiv:1203.1998.
   URL http://arxiv.org/abs/1203.1998
- [3] P. Sjögren, Operators associated with the hermite semigroup A survey, The Journal of Fourier Analysis and Applications 3 (S1) (1997) 813–823. doi:10.1007/BF02656487. URL http://link.springer.com/10.1007/BF02656487
- [4] J. Maas, J. Neerven, P. Portal, Whitney coverings and the tent spaces  $T^{1,q}(\gamma)$  for the Gaussian measure, Arkiv för Matematik 50 (2) (2011) 379–395. doi:10.1007/s11512-010-0143-z.