# A NOTE ON THE GAUSSIAN MAXIMAL FUNCTION

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ABSTRACT. This note presents a proof that the non-tangential maximal function of the Ornstein-Uhlenbeck semigroup is bounded almost surely by the Gaussian Hardy-Littlewood maximal function. In particular this entails improvement on a result by Pineda and Urbina [2] who proved a similar result for a 'trunctated' version of the non-tangential maximal function. We actually obtain boundedness of the maximal function on non-tangential cones of arbitrary aperture.

#### 1. Introduction

Maximal functions are among the most studied objects in harmonic analysis. It is well known that the classical non-tangential maximal function associated with the heat semigroup is bounded almost everywhere by the Hardy-Littlewood maximal function,

function, 
$$\sup_{\substack{(y,t) \in \mathbf{R}_{+}^{d+1} \\ |x-y| < t}} |\mathrm{e}^{-t\Delta} u(y)| \lesssim \sup_{r>0} \frac{1}{|B_r(x)|} \int_{B_r(x)} |u| \, \mathrm{d}\lambda.$$

Here the action of heat semigroup  $e^{-t\Delta}u = \rho_t * u$  is given by a convolution of u with the heat kernel

$$\rho_t(s) := \frac{e^{-|s|^2/4t}}{(4\pi t)^{\frac{d}{2}}}.$$

In this note we are interested in its Gaussian counterpart. The change from Lebesgue measure to the Gaussian measure

(2) 
$$d\gamma(x) := \pi^{-\frac{d}{2}} e^{-|x|^2} dx$$

introduces quite some intricate technical and conceptual difficulties which are caused by the fact that the Gaussian measure is non-doubling. The Gaussian analogue to the Laplacian is the  $Ornstein-Uhlenbeck\ operator\ L$ ,

(3) 
$$L := -\frac{1}{2}\Delta + \langle x, \nabla \rangle = \frac{1}{2}\nabla_{\gamma}^* \nabla_{\gamma},$$

where  $\nabla_{\gamma}$  denotes the realisation of the gradient in  $L^2(\mathbf{R}^d, \gamma)$ . Our main result, to be proved in Theorem 1, is the following Gaussian analogue of (1):

(4) 
$$\sup_{(y,t) \in \Gamma_x^{(A,a)}} |e^{-t^2 L} u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \, d\gamma.$$

Here.

(5) 
$$\Gamma_x^{(A,a)} := \Gamma_x^{(A,a)}(\gamma) := \{(y,t) \in \mathbf{R}_+^{d+1} \, : \, |x-y| < At \text{ and } t \leqslant am(x)\}$$

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is the Gaussian cone with aperture A and cut-off parameter a, and

$$(6) \hspace{3.1em} m(x) := \min \bigg\{ 1, \frac{1}{|x|} \bigg\}.$$

A slighly weaker version of the inequality (4) has been proved by Pineda and Urbina [2] who showed that

$$\sup_{(y,t)\in\widetilde{\Gamma}_x} |\mathrm{e}^{-t^2L} u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \,\mathrm{d}\gamma,$$

where

$$\widetilde{\Gamma}_x(x) = \{(y, t) \in \mathbf{R}^d_+ : |x - y| < t \leqslant \widetilde{m}(x)\}$$

is the 'reduced' Gaussian cone corresponding to the function

$$\widetilde{m}(x) = \min\left\{\frac{1}{2}, \frac{1}{|x|}\right\}.$$

Their proof does not seem to easily generalize the range of t from  $\frac{1}{2}$  up to 1. Our proof of (4) is different and, we believe, more transparent than the one presented in [2]. It has the further advantage of allowing the extension to cones with arbitrary aperture A>0 and cut-off parameter a>0 without any additional technicalities. This additional generality is very important and has already been used by Portal (cf. the claim made in [3, discussion preceding Lemma 2.3]) to prove the  $H^1$ -boundedness of the Riesz transform associated with L.

To save writing, let us fix some notation. The number d is a positive integer. To avoid possible confusion, we define the *positive integers* as the set  $\mathbf{Z}_{+} = \{1, 2, 3, \ldots\}$ .

### 2. The Mehler Kernel

The Mehler kernel (see e.g., [4]) is the Schwartz kernel associated to the Ornstein-Uhlenbeck semigroup  $(e^{-tL})_{t\geqslant 0}$ , that is,

(7) 
$$e^{-tL}u(x) = \int_{\mathbf{R}^d} M_t(x,\cdot)u \, d\gamma.$$

There is an abundance of literature on the Mehler kernel and its properties. We shall only use the fact, proved e.g. in the review paper [4], that the Mehler kernel is given explicitly by

$$M_t(x,y) = \frac{\exp\left(-\frac{|e^{-t}x - y|^2}{1 - e^{-2t}}\right)}{(1 - e^{-2t})^{\frac{d}{2}}} e^{|y|^2}.$$

Note that  $M_t(x, y)$  is symmetric in x and y. A formula for  $M_t$  which honors this observation is:

(8) 
$$M_t(x,y) = \frac{\exp\left(-e^{-2t} \frac{|x-y|^2}{1 - e^{-2t}}\right)}{(1 - e^{-t})^{\frac{d}{2}}} \frac{\exp\left(2e^{-t} \frac{\langle x, y \rangle}{1 + e^{-t}}\right)}{(1 + e^{-t})^{\frac{d}{2}}}.$$

#### 3

#### 3. Some Lemmata

We use m as defined in (6) in our next lemma, which is taken from [1].

- 1. **Lemma.** Let a, A be strictly positive real numbers and t > 0. We have for  $x, y \in \mathbf{R}^d$  that:
  - (1) If |x y| < At and  $t \le am(x)$ , then  $t \le (1 + aA)m(y)$ ;
  - (2) If |x y| < Am(x), then  $m(x) \le (1 + A)m(y)$  and  $m(y) \le 2(1 + A)m(x)$ .

The next lemma will come useful when we want to cancel exponential growth in one variable with exponential decay in the other as long both variables are in a Gaussian cone.

2. **Lemma.** Let  $\alpha > 0$  and  $|x - y| \leq \alpha m(x)$ . Then:

$$e^{-\alpha^2 - 2\alpha} e^{|y|^2} \le e^{|x|^2} \le e^{\alpha^2 (1+\alpha)^2 + 2\alpha(1+\alpha)} e^{|y|^2}.$$

*Proof.* By the inverse triangle inequality and  $m(x)|x| \leq 1$  we get,

$$|y|^2 \le (\alpha m(x) + |x|)^2 \le \alpha^2 + 2\alpha + |x|^2$$
.

This gives the first inequality. For the second we use Lemma 1 to infer  $m(x) \leq (1+\alpha)m(y)$ . Proceeding as before we obtain:

$$|x|^2 \le \alpha^2 (1+\alpha)^2 + 2\alpha(1+\alpha) + |y|^2$$
.

As required.

- 3.1. An estimate on Gaussian balls. Let  $B = B_t(x)$  be the open Euclidean ball with radius t and center x and let  $\gamma$  be the Gaussian measure as defined by (2). We shall denote by  $S_d$  the surface area of the unit sphere in  $\mathbf{R}^d$ .
- 3. **Lemma.** For all  $x \in \mathbf{R}^d$  and t > 0 we have the inequality:

(9) 
$$\gamma(B_t(x)) \leqslant \frac{1}{2} S_d t^d e^{2t|x|} e^{-|x|^2}.$$

*Proof.* Remark that

$$\int_{B} e^{-|\xi|^{2}} d\xi = e^{-|x|^{2}} \int_{B} e^{-|\xi - x|^{2}} e^{-2\langle x, \xi - x \rangle} d\xi 
\leqslant e^{-|x|^{2}} \int_{B} e^{-|\xi - x|^{2}} e^{2|x||\xi - x|} d\xi 
\leqslant e^{-|x|^{2}} e^{2t|x|} \int_{B} e^{-|\xi - x|^{2}} d\xi 
= \pi^{\frac{d}{2}} e^{-|x|^{2}} e^{2t|x|} \gamma(B_{t}(0)).$$

So, there holds that

(10) 
$$\gamma(B_t(x)) \leqslant e^{-|x|^2} e^{2t|x|} \gamma(B_t(0)).$$

We will estimate the Gaussian volume of the ball  $B_t(0)$ . Using polar coordinates we proceed for  $d \ge 2$  by:

$$\gamma(B_t(0)) = \pi^{-\frac{d}{2}} \int_{B_t(0)} e^{-|\xi|^2} d\xi$$

$$= \pi^{-\frac{d}{2}} S_d \int_0^t e^{-r^2} r^{d-1} dr$$

$$\leq \frac{1}{2} \pi^{-\frac{d}{2}} S_d t^{d-2} \int_0^t 2r e^{-r^2} dr$$

$$= \frac{1}{2} \pi^{-\frac{d}{2}} S_d t^{d-2} (1 - e^{-t^2})$$

$$\leq \frac{1}{2} \pi^{-\frac{d}{2}} S_d t^d,$$

where the last step uses  $1 - e^{-x} \le x$  for  $x \ge 0$ . The case for d = 1 follows by a simplified argument. Upon combining this result with (10) we obtain (9).

3.2. On-diagonal kernel estimates on annuli. As is common in harmonic analysis, we often wish to decompose  $\mathbf{R}^d$  into sets on which certain phenomena are easier to handle. Here we will decompose the space into disjoint annuli.

Throughout this subsection we fix  $x \in \mathbf{R}^d$ , constants A, a > 0, a pair  $(y, t) \in \Gamma_x^{(A,a)}$ , and we put  $\alpha := Aa$ . We use the notation rB to mean the ball obtained from the ball B by multiplying its radius by r.

The annuli  $C_k$  are given by:

(11) 
$$C_k := (2^{k+1} - 1)B_{\alpha t}(x) \setminus (2^k - 1)B_{\alpha t}(x) \text{ with } k \geqslant 0.$$

Note that  $C_0 = B_{\alpha t}(x)$ . Whenever  $\xi$  is in  $C_k$ , we get for  $k \ge 0$ :

$$(12) (2^k - 1)\alpha t < |y - \xi| \le (2^{k+1} - 1)\alpha t.$$

On  $C_k$  we have the following bound for  $M_{t^2}(y,\cdot)$ :

4. **Lemma.** For all  $\xi \in C_k$  we have:

(13) 
$$M_{t^2}(y,\xi) \leqslant \frac{e^{|y|^2}}{(1-e^{-2t^2})^{\frac{d}{2}}} \exp((2^{k+1}-1)\alpha t|y|) \exp(-\frac{\alpha^2}{2e^{2a^2}}(2^k-1)^2),$$

*Proof.* Considering the first exponential which occurs in the Mehler kernel (8) together with (12) gives for  $k \ge 0$ :

$$\exp\left(-e^{-2t^2} \frac{|y-\xi|^2}{1-e^{-2t^2}}\right) \leqslant \exp\left(-e^{-2t^2} \frac{(2^k-1)^2 \alpha^2 t^2}{1-e^{-2t^2}}\right)$$

$$\stackrel{(\dagger)}{\leqslant} \exp\left(-\frac{\alpha^2}{2e^{2t^2}} (2^k-1)^2\right) \stackrel{(\dagger\dagger)}{\leqslant} \exp\left(-\frac{\alpha^2}{2e^{2a^2}} (2^k-1)^2\right),$$

where (†) follows from  $1 - e^{-x} \le x$  for  $x \ge 0$ , and (††) uses that  $t \le am(x) \le a$ . Using the estimate  $1 + x \ge 2x$  for  $0 \le x \le 1$ , for the second exponential in the Mehler kernel (8) we obtain, by (12):

$$\exp\left(2e^{-t^2}\frac{\langle y,\xi\rangle}{1+e^{-t^2}}\right) \leqslant \exp(|\langle y,\xi\rangle|)$$
$$\leqslant \exp(|\langle y,\xi-y\rangle|)e^{|y|^2}$$
$$\leqslant \exp((2^{k+1}-1)\alpha t|y|)e^{|y|^2}.$$

Combining things, we obtain the estimate in the formulation of the lemma.

### 4. The main result

Our theorem is a modification of [2, lemma 1.1] with a new proof.

1. **Theorem.** Let A, a > 0. For all  $x \in \mathbf{R}^d$  and all  $u \in L^2(\mathbf{R}^d, \gamma)$  we have

(14) 
$$\sup_{(y,t)\in\Gamma_x^{(A,a)}} |\mathrm{e}^{-t^2L}u(y)| \lesssim \sup_{r>0} \frac{1}{\gamma(B_r(x))} \int_{B_r(x)} |u| \, \mathrm{d}\gamma.$$

*Proof.* We fix  $x \in \mathbf{R}^d$  and  $(y,t) \in \Gamma_x^{(A,a)}$ . Set  $\alpha = aA$ . We will prove (14) by splitting up the integration domain into the annuli  $C_k$  as defined by (11):

(15) 
$$|e^{-t^2L}u(y)| \le \sum_{k=0}^{\infty} I_k(y)$$
, where  $I_k(y) := \int_{C_k} M_{t^2}(y, \cdot) |u(\cdot)| d\gamma$ .

From  $t \leq am(x)$  we get  $t|x| \leq a$ , and by Lemma 1 we have  $t|y| \leq 1 + \alpha$ . From this and Lemma 4 we infer, for  $\xi \in C_k$ , that:

$$(16) \quad M_{t^2}(y,\xi) \leqslant \frac{e^{|y|^2}}{(1 - e^{-2t^2})^{\frac{d}{2}}} \exp((2^{k+1} - 1)\alpha(1 + \alpha)) \exp\left(-\frac{\alpha^2}{2e^{2a^2}}(2^k - 1)^2\right).$$

Combining (16) and Lemma 2, we obtain

$$M_{t^2}(y,\xi) \leqslant c_k \frac{e^{|x|^2}}{(1 - e^{-2t^2})^{\frac{d}{2}}}$$

where  $c_k$  depends only on A, a, and t. Also, by (12),

$$|x - \xi| \le |x - y| + |\xi - y| \le (A + 2^{k+1}\alpha)t.$$

It follows that  $C_k$  is contained in  $D_k := B_{(A+2^{k+1}\alpha)t}(x)$ .

Let us denote the supremum on right-hand side of (14) by  $M_{\gamma}u(x)$ . Using (16), we can bound the integral on the right-hand side of (15) by

$$\int_{C_{k}} M_{t^{2}}(y,\cdot)|u(\cdot)| \, d\gamma \leqslant c_{k} \frac{e^{|x|^{2}}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} \int_{C_{k}} |u| \, d\gamma$$

$$\leqslant c_{k} \frac{e^{|x|^{2}}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} \int_{D_{k}} |u| \, d\gamma$$

$$\leqslant c_{k} \frac{e^{|x|^{2}}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} \gamma(D_{k}) M_{\gamma} u(x)$$

$$\stackrel{(\dagger)}{\lesssim}_{A,a,d} c_{k} \frac{t^{d}}{(1 - e^{-2t^{2}})^{\frac{d}{2}}} e^{2((A + 2^{k+1}\alpha)t)|x|} M_{\gamma} u(x)$$

$$\stackrel{(\dagger\dagger)}{\lesssim}_{A,a,d} c_{k} e^{2^{k+2}\alpha a} M_{\gamma} u(x),$$

where (†) uses Lemma 3 applied to  $D_k$  and (††) uses that  $t \leq am(x)$  implies  $t|x| \leq a$  and  $t \leq a$ , the latter implying

$$\left(\frac{t^2}{1 - \mathrm{e}^{-t^2}}\right)^{\frac{d}{2}} \leqslant \left(\frac{a^2}{1 - \mathrm{e}^{-a^2}}\right)^{\frac{d}{2}}$$

(note that  $x/(1-e^{-x})$  is increasing).

Inserting the dependency of  $c_k$  upon k as coming from (16) and using that  $t \leq a$ , we can then bound the maximal function as follows:

$$|\mathrm{e}^{-t^2L}u(y)| = \sum_{k=0}^{\infty} I_k \lesssim_{A,a,d} M_{\gamma}u(x) \sum_{k=0}^{\infty} 2^{kd} \mathrm{e}^{\alpha(1+\alpha)2^{k+1}} \mathrm{e}^{-\frac{\alpha^2}{2\mathrm{e}^{2a^2}}4^k} \mathrm{e}^{2^{k+2}\alpha a}.$$

Evidently the sum on the right-hand side converges.

## References

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