

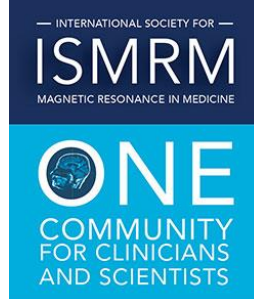
ANISOTROPIC CEREBRAL VASCULATURE CAUSES ORIENTATION DEPENDENCY IN MR SIGNAL



Simulation in the Human Brain

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Declaration of Financial Interests or Relationships

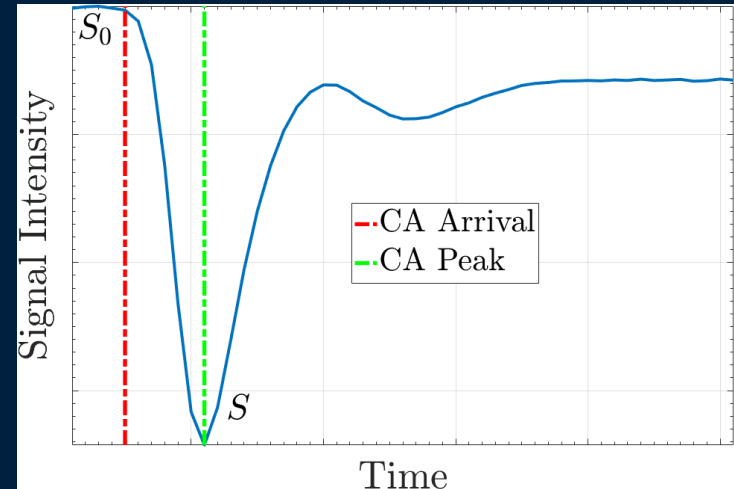
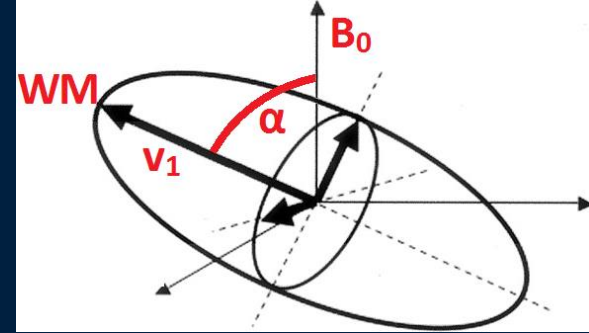
Speaker Name: **Jonathan Doucette**

I have no financial interests or relationships to disclose with regard to the subject matter of this presentation.

INTRODUCTION AND MOTIVATION

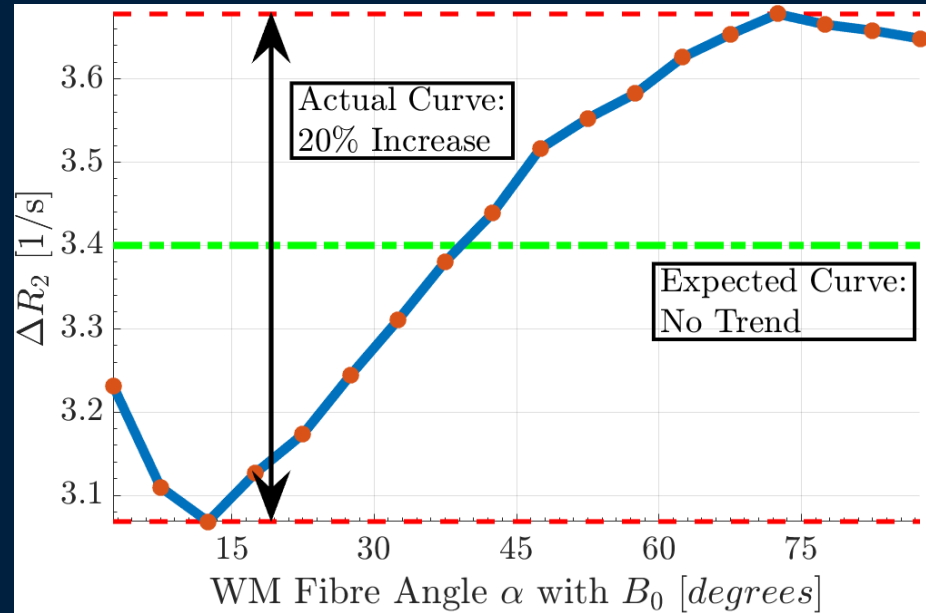
- Subject Data
 - DTI and Spin Echo DSC scans obtained for 19 healthy subjects
 - WM Fibre orientation angle α determined from DTI data, correlated with change in R_2 with and without CA
 - Change in decay rate ΔR_2 at points S_0 and S can be calculated via:

$$\Delta R_2 = -\frac{1}{TE} \ln\left(\frac{S}{S_0}\right)$$



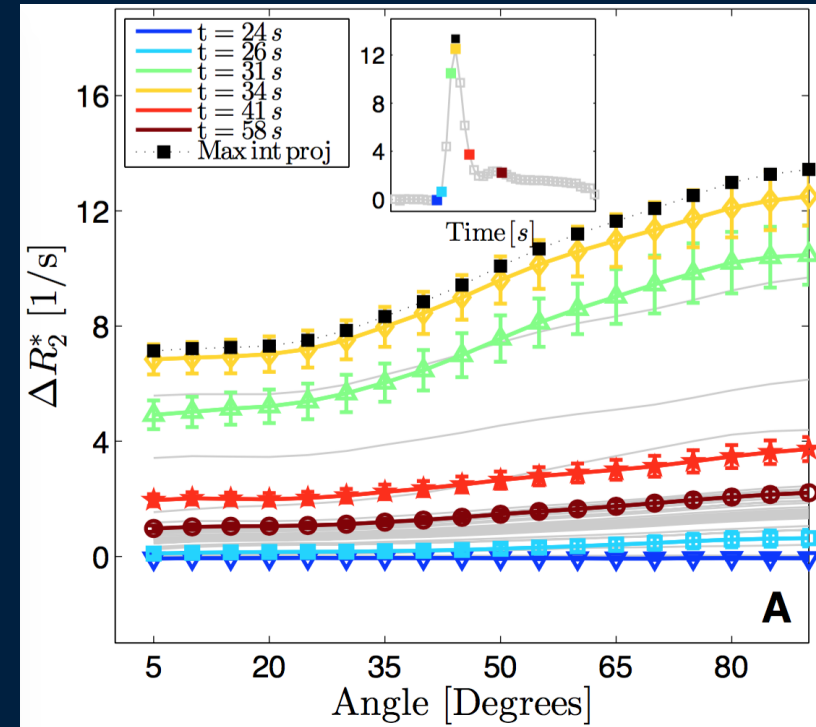
INTRODUCTION AND MOTIVATION

- Expected Spin Echo DSC Signal:
 - SE spins should be refocused by TE
 - Decay should be due only to local (CA dependent) R_2 -value
 - Decoherence should be negligible
- Actual Spin Echo DSC Signal:
 - There is extra signal decay!
 - Even more peculiar: extra decay depends on the local fibre orientation



INTRODUCTION AND MOTIVATION

- Comparison with Gradient Echo DSC case:
 - Our previous study* showed that orientation dependency is strong in GRE-DSC
- This is *expected* for GRE-DSC
 - For larger angles, the extravascular inhomogeneities occupy larger volume fractions of a voxel
- We now aim to explain why this same effect is present in SE-DSC



*Hernandez-Torres, et al. *JCBFM* (2016)

THE BLOCH EQUATIONS FOR SE-DSC WITH INHOMOGENEITIES

- We start by considering the Bloch Equations:

$$\frac{d\vec{M}}{dt} = \gamma \vec{M} \times \vec{B} - \left(\frac{M_x}{T_2}, \frac{M_y}{T_2}, \frac{M_z}{T_1} \right)$$

- In a rotating frame with constant $\vec{B} = B_0 \hat{z}$ and $T_2 \ll T_1$, we have:

$$\frac{d\tilde{M}_{xy}}{dt} = -\frac{1}{T_2} \tilde{M}_{xy}$$

- Where we define the complex magnetization:

$$\tilde{M}_{xy} := M_x + iM_y$$



THE BLOCH EQUATIONS FOR SE-DSC WITH INHOMOGENEITIES

- We can add induced field inhomogeneities through incorporating susceptibility changes $\delta\chi$ and using the unit dipole kernel $G(\vec{r})$:

$$\delta\omega(\vec{r}) = \gamma B_0 \cdot (G(\vec{r}) * \delta\chi(\vec{r}))$$

- The resulting differential equation and solution are then:

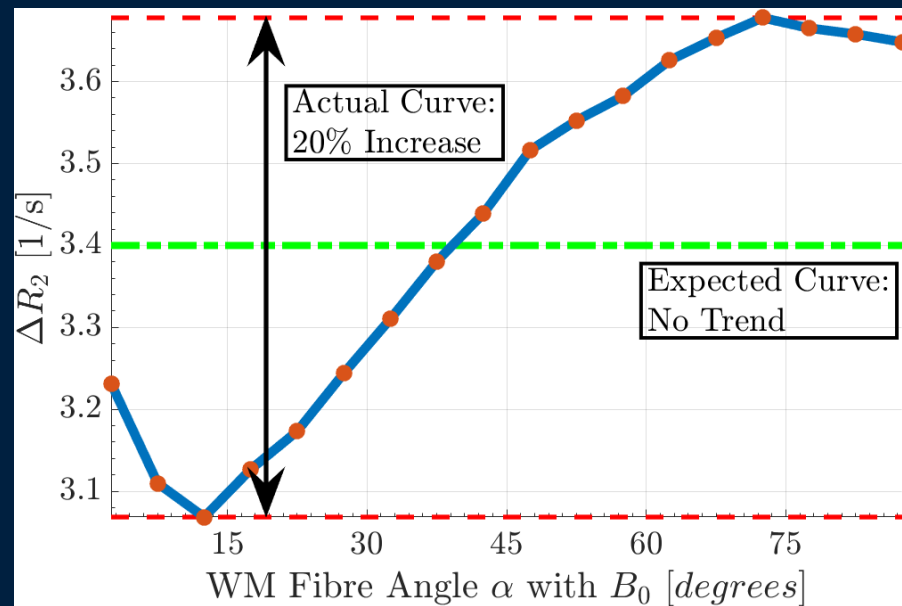
$$\begin{aligned}\frac{d\tilde{M}_{xy}}{dt} &= -(R_2(\vec{r}) + i\delta\omega(\vec{r})) \cdot \tilde{M}_{xy} \\ \Rightarrow \tilde{M}_{xy}(\vec{r}, t) &= \tilde{M}_{xy}(\vec{r}, 0) e^{-(R_2 + i\delta\omega)t}\end{aligned}$$



THE BLOCH EQUATIONS FOR SE-DSC WITH INHOMOGENEITIES

- However, this solution can only produce the Expected Curve, not the Actual Curve
- This is because this solution completely refocuses by time TE (as it should for SE!):

$$\tilde{M}_{xy}(\vec{r}, TE) = \overline{\tilde{M}_{xy}(\vec{r}, 0)} e^{-R_2 TE}$$



THE BLOCH-TORREY EQUATION FOR SE-DSC

- Remedying this model
 - So far, we have implicitly assumed that all spins are stationary
 - If spins are allowed to move, they will see different local inhomogeneities and will not perfectly rephase – there will be extra decay



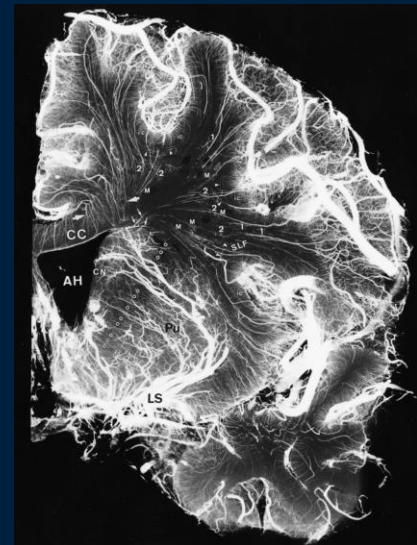
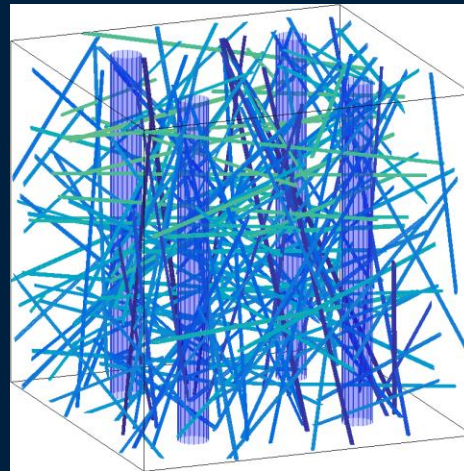
- We add diffusion to the equation:

$$\frac{dM}{dt} = D \cdot \nabla^2 M - (R_2 + i\delta\omega) \cdot M$$

- Where ∇^2 is the 3D Laplacian operator. This is the so-called Bloch-Torrey equation, and has no closed form solution and must be solved numerically

GEOMETRY OF THE SE-DSC PROBLEM

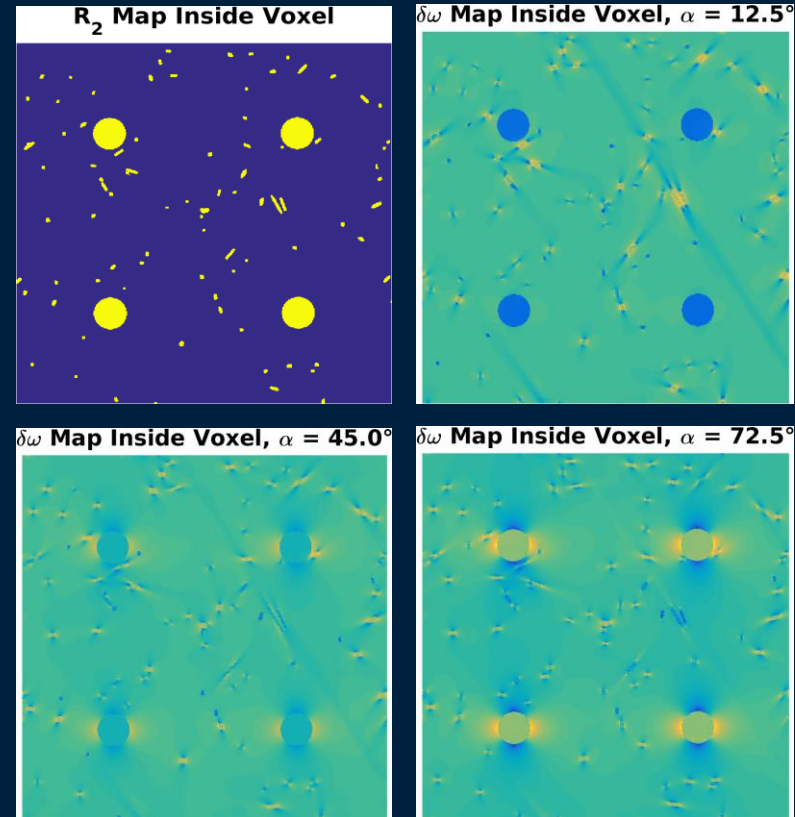
- Now, we consider the geometry of the problem within a single voxel
- The geometry is simulated as follows:
 - Create an isotropic vascular bed of small vessels
 - Add anisotropic large vessels to simulate large vasculature running in parallel to WM tracks*
- These large vessels introduce the possibility of angular dependence



*Okudera, et al.
Neuropathology (1999)

GEOMETRY OF THE SE-DSC PROBLEM

- Given the geometry, we now calculate $R_2(\vec{r})$ and $\delta\chi(\vec{r})$:
 - Each are piecewise constant, taking different values in tissue and in blood
 - Value in blood depends on CA
- Choose a WM fibre angle α relative to \vec{B}_0
 - This represents the orientation of the whole voxel
- Calculate the *dephasing*:
 - $\delta\omega_\alpha = \gamma B_0 \cdot (G_\alpha * \delta\chi)$



FORMALIZING THE SE-DSC PROBLEM

1. *Initialize*

- Choose parameters: BVF, iBVF, CA concentration
- Calculate geometry: $R_2(\vec{r})$ and $\delta\omega_\alpha(\vec{r}) = \gamma B_0 \cdot (G_\alpha * \delta\chi)$
- Initialize magnetization: $M(\vec{r}, 0) = i$ ($\pi/2$ pulse into xy -plane)

2. *Propagate*

- Solve the Bloch-Torrey equation for $M(\vec{r}, TE)$
 - $\frac{dM}{dt} = D \cdot \nabla^2 M - (R_2 + i\delta\omega) \cdot M$

3. *Sum*

- The final signal is given by:
 - $S(TE) = | \iiint M(\vec{r}, TE) d^3r |$



FORMALIZING THE SE-DSC PROBLEM

4. *Repeat*

- The simulation is executed for each angle in the range $\{2.5^\circ, 7.5^\circ, \dots, 87.5^\circ\}$, both **with** and **without** contrast agent

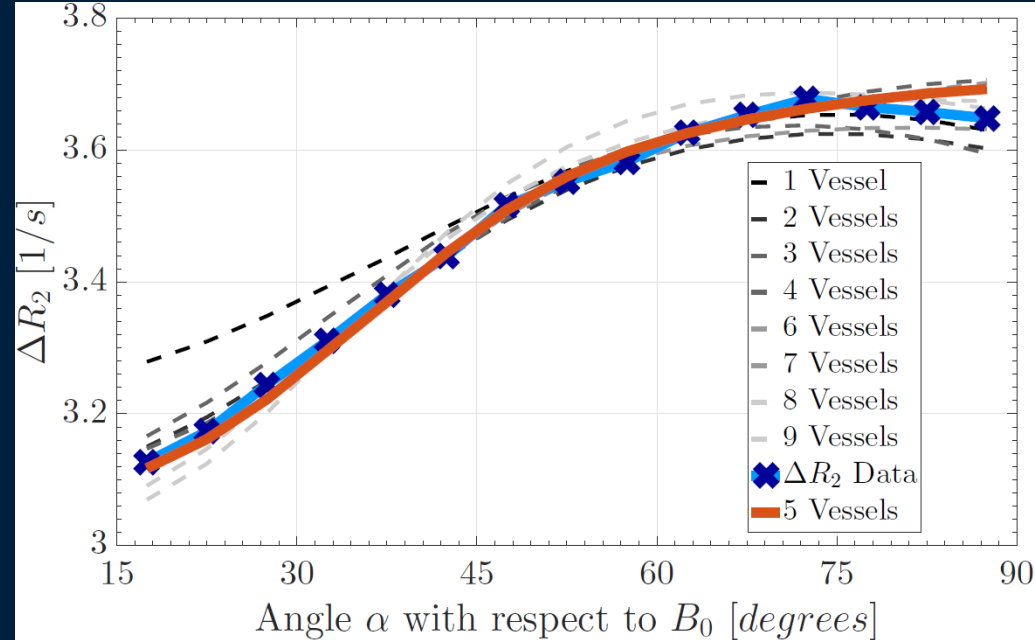
5. *Fit*

- Given $S_0(TE)$ and $S(TE)$, calculate $\Delta R_2 = -\frac{1}{TE} \ln S/S_0$ for each angle and compare with observed data
- If fit is poor, adjust parameters (BVF, iBVF, CA) and go to step 1



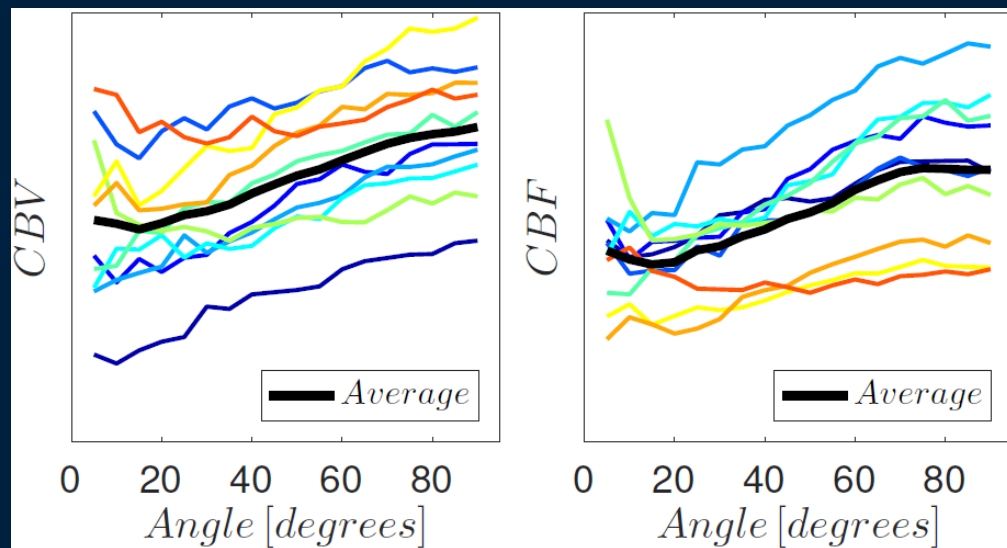
SE-DSC SIMULATION RESULTS AND DISCUSSION

- Simulation matches Actual Curve well
- Simulation was ran for a variable number of large anisotropic vessels
 - Best results occurred for 4-6 anisotropic vessels
 - Corresponding radii of approx. 90 μm , 80 μm , 70 μm
- Resulting Fit Parameters:
 - BVF: 3.49%
 - iBVF: 2.30%
 - CA: 6.23 mM



SE-DSC SIMULATION RESULTS AND DISCUSSION

- Quantities that are derived from R_2 are affected, too!
- Perfusion parameters such as CBV and CBF exhibit orientation dependencies on the order of 20%, similar to R_2



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