

Explicit Hydrogen-like Momentum Wave Functions

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Correcting for the Levchuk effect i.e. the presence of non-zero target energy due to the Fermi motion of the bound electrons requires knowing the momentum distribution of the electrons in the target material. The procedure applied in the Swartz *et al.* paper [1] used screened hydrogen-like wave functions to obtain the momentum distribution, by explicit summing over the wavefunctions for each shell weighted by its occupancy. The Levchuk correction code based on the procedure in this paper was written two decades ago and this document is a check of the calculations in the code. I compute all the pieces required to get the probability distributions given in Eq. 15 of the Swartz paper [1]. Here I explicitly write the momentum wave functions $\phi_{n,l}(p)$ required for an Fe foil, where p is in atomic units of $\hbar/a_0 = m_e\alpha$ or for larger Z , $\tilde{Z}m_e\alpha$ (\tilde{Z} is adjusted to account for screening). Since the electron configuration for atomic Fe is $[\text{Ar}]3d^64s^2$, at a minimum, the following wave functions are required: $\phi_{1,0}$, $\phi_{2,0}$, $\phi_{2,1}$, $\phi_{3,0}$, $\phi_{3,1}$, $\phi_{3,2}$, $\phi_{4,0}$. The momentum wave functions come from a reference written by Bethe [3], where equation 8.8 gives the following form for the radial momentum:

$$\phi_{n,l}(p) = \left(\frac{2(n-l-1)!}{\pi(n+l)!} \right)^{\frac{1}{2}} n^2 2^{2(l+1)} l! \frac{n^l p^l}{(n^2 p^2 + 1)^{l+2}} C_{n-l-1}^{l+1} \left(\frac{n^2 p^2 - 1}{n^2 p^2 + 1} \right), \quad (1)$$

where $C_\beta^\alpha(x)$ are the Gegenbauer polynomials. The Gegenbauer functions are given explicitly by

$$\begin{aligned} C_0^\alpha(x) &= 1 \\ C_1^\alpha(x) &= 2\alpha x \\ C_0^\alpha(x) &= \frac{1}{n} \left[2x(\beta + \alpha - 1)C_{\beta-1}^\alpha(x) - (\beta + 2\alpha - 2)C_{\beta-2}^\alpha(x) \right]. \end{aligned}$$

¹Linus Pauling argues in [2] that a better model for metallic Fe is $[\text{Ar}]3d^54s^14p^2$. He suggests we can view Fe as having about 5.8 valence electrons in hybridized wave functions which are linear combinations of the 3d, 4s and 4p functions. This would require a further wave function $\phi_{4,2}$ beyond that required for the simple atomic electron configuration.

The required Gegenbauer polynomials are:

$$\begin{aligned}
C_0^1(x) &= 1 \\
C_1^1(x) &= 2x \\
C_2^1(x) &= 4x^2 - 1 \\
C_3^1(x) &= 8x^3 - 4x \\
C_0^2(x) &= 1 \\
C_1^2(x) &= 4x \\
C_0^3(x) &= 1.
\end{aligned}$$

Substituting these polynomials into Eq. 1 gives the momentum wave functions:

$$\phi_{1,0} = \frac{8}{\sqrt{2\pi}(p^2 + 1)^2} \quad (2)$$

$$\phi_{2,0} = \frac{32(4p^2 - 1)}{\sqrt{\pi}(4p^2 + 1)^3} \quad (3)$$

$$\phi_{2,1} = \frac{128p}{\sqrt{3\pi}(4p^2 + 1)^3} \quad (4)$$

$$\phi_{3,0} = \frac{72}{\sqrt{6\pi}(9p^2 + 1)^2} \left(4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^2} - 1 \right) \quad (5)$$

$$\phi_{3,1} = \frac{864p(9p^2 - 1)}{\sqrt{3\pi}(9p^2 + 1)^4} \quad (6)$$

$$\phi_{3,2} = \frac{5184}{\sqrt{15\pi}} \frac{p^2}{(9p^2 + 1)^4} \quad (7)$$

$$\phi_{4,0} = \frac{256}{\sqrt{2\pi}} \frac{(16p^2 - 1)}{(16p^2 + 1)^3} \left(2 \frac{(16p^2 - 1)}{(16p^2 + 1)} - 1 \right) \quad (8)$$

And since we need probability distributions let's just go ahead and square them all.

$$|\phi_{1,0}|^2 = \frac{32}{\pi(p^2 + 1)^4} = 10.185916 \times \frac{1}{(p^2 + 1)^4} \quad (9)$$

$$|\phi_{2,0}|^2 = \frac{1024(4p^2 - 1)^2}{\pi(4p^2 + 1)^6} = 325.94932 \times \frac{(4p^2 - 1)^2}{(4p^2 + 1)^6} \quad (10)$$

$$|\phi_{2,1}|^2 = \frac{16384p^2}{3\pi(4p^2 + 1)^6} = 1738.3964 \times \frac{p^2}{(4p^2 + 1)^6} \quad (11)$$

$$|\phi_{3,0}|^2 = \frac{864}{\pi(9p^2 + 1)^4} \left(4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^2} - 1 \right)^2 \quad (12)$$

$$= 275.01974 \times \frac{1}{(9p^2 + 1)^4} \left(4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^2} - 1 \right)^2 \quad (13)$$

$$|\phi_{3,1}|^2 = \frac{248832p^2}{\pi} \left(\frac{(9p^2 - 1)}{(9p^2 + 1)^4} \right)^2 = 79205.686 \times p^2 \left(\frac{(9p^2 - 1)}{(9p^2 + 1)^4} \right)^2 \quad (14)$$

$$|\phi_{3,2}|^2 = \frac{8957952}{5\pi} \frac{p^4}{(9p^2 + 1)^8} = 570280.94 \times \frac{p^4}{(9p^2 + 1)^8} \quad (15)$$

$$|\phi_{4,0}|^2 = \frac{32768}{\pi} \left(\frac{(16p^2 - 1)}{(16p^2 + 1)^3} \right)^2 \left(2 \frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 1 \right)^2 \quad (16)$$

$$= 10430.378 \times \left(\frac{(16p^2 - 1)}{(16p^2 + 1)^3} \right)^2 \left(2 \frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 1 \right)^2 \quad (17)$$

$$\text{or alternately} \quad (18)$$

$$= 651.89865 \times \left(\frac{(16p^2 - 1)}{(16p^2 + 1)^3} \right)^2 \left(8 \frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 4 \right)^2 \quad (19)$$

Using the procedure in the Swartz paper[1], the polarized and unpolarized momentum distributions for a pure elemental target foil are given by

$$f_{unp} = \sum_{n,l} \frac{C_{n,l}}{P_n} \left(\frac{p}{P_n} \right)^2 \left| \phi_{n,l} \left(\frac{p}{P_n} \right) \right|^2, \quad (20)$$

where p is the radial momentum, $C_{n,l}$ is the fraction of the total unpolarized electron population in the n, l orbital and $P_n = Z_n m_e \alpha$ is an atomic momentum scale associated with the given orbital. The term Z_n is adjusted to account for screening assuming that a given electron is fully screened by electrons in inner orbitals and by half of its neighboring electrons in the same orbital. It is given explicitly by

$$Z_n = Z - \frac{N_n - 1}{2} - \sum_i^{n-1} N_i,$$

where N_i is the number of electrons in the i^{th} shell. The calculated P_n values for Fe, Co and Ni are given in Table 1. If the polarized distribution is assumed to come solely from D-wave,

Table 1: Atomic momentum scale P_n in keV/c for the first four shells of the three ferromagnetic elements.

Element	n=1	n=2	n=3	n=4
Fe	95.113	76.464	35.434	5.595
Co	98.843	80.193	37.299	5.595
Ni	102.573	83.923	39.164	5.595

M-shell electrons, this polarized distribution is given by

$$f_{pol} = \sum_{n,l} \frac{1}{P_n} \left(\frac{p}{P_3} \right)^2 \left| \phi_{3,2} \left(\frac{p}{P_3} \right) \right|^2.$$

References

- [1] M. Swartz, H.R. Band, F.J. Decker, P. Emma, M.J. Fero, R. Frey, R. King, A. Lath, T. Limberg, R. Prepost, P.C. Rowson, B.A. Schumm, M. Woods, and M. Zolotarev. Observation of target electron momentum effects in single-arm m ller polarimetry. *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, 363(3):526 – 537, 1995.
- [2] Linus Pauling. A theory of ferromagnetism. *Proceedings of the National Academy of Sciences*, 39(6):551–560, 1953.
- [3] Hans A. Bethe and Edwin E. Salpeter. *Quantum Mechanics of One- and Two-Electron Atoms*. Springer Berlin Heidelberg, 1957.