## Explicit Hydrogen-like Momentum Wave Functions

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Correcting for the Levchuk effect i.e. the presence of non-zero target energy due to the Fermi motion of the bound electrons requires knowing the momentum distribution of the electrons in the target material. The procedure applied in the Swartz et al. paper [1] used screened hydrogen-like wave functions to obtain the momentum distribution, by explicit summing over the wavefunctions for each shell weighted by its occupancy. Here I explicitly write the momentum wave functions  $\phi_{n,l}(p)$  required for an Fe foil, where p is in atomic units of  $\hbar/a_0$ . Since the electron configuration for atomic Fe is [Ar]3 $d^64s^{21}$ , at a minimum, the following wave functions are required:  $\phi_{1,0}$ ,  $\phi_{2,0}$ ,  $\phi_{2,1}$ ,  $\phi_{3,0}$ ,  $\phi_{3,1}$ ,  $\phi_{3,2}$ ,  $\phi_{4,0}$ . The momentum wave functions come from a reference written by Bethe [3], where equation 8.8 gives the following form for the radial momentum:

$$\phi_{n,l}(p) = \left(\frac{2(n-l-1)!}{\pi(n+l)!}\right)^{\frac{1}{2}} n^2 2^{2(l+1)} l! \frac{n^l p^l}{(n^2 p^2 + 1)^{l+2}} C_{n-l-1}^{l+1} \left(\frac{n^2 p^2 - 1}{n^2 p^2 + 1}\right), \tag{1}$$

where  $C^{\alpha}_{\beta}(x)$  is the Gegenbauer function. The Gegenbauer functions are given explicitly by

$$C_0^{\alpha}(x) = 1$$

$$C_1^{\alpha}(x) = 2\alpha x$$

$$C_0^{\alpha}\beta(x) = \frac{1}{n} \left[ 2x(\beta + \alpha - 1)C_{\beta-1}^{\alpha}(x) - (\beta + 2\alpha - 2)C_{\beta-2}^{\alpha}(x) \right].$$

<sup>&</sup>lt;sup>1</sup>Linus Pauling argues in [2] that a better model for metallic Fe is [Ar] $3d^54s^14p^2$ . He suggests we can view Fe as having about 5.8 valence electrons in hybridized wave functions which are linear combinations of the 3d, 4s and 4p functions. This would require a further wave function  $\phi_{4,2}$  beyond that required for the simple atomic electron configuration.

The required Gegenbauer polynomials are:

$$C_0^1(x) = 1$$

$$C_1^1(x) = 2x$$

$$C_2^1(x) = 4x^2 - 1$$

$$C_3^1(x) = 8x^3 - 4x$$

$$C_0^2(x) = 1$$

$$C_1^2(x) = 4x$$

$$C_0^2(x) = 1$$

Substituting these polynomials into Eq. 1 gives the momentum wave functions:

$$\phi_{1,0} = \frac{8}{\sqrt{2\pi}(p^2 + 1)^2} \tag{2}$$

$$\phi_{2,0} = \frac{32(4p^2 - 1)}{\sqrt{\pi}(4p^2 + 1)^3} \tag{3}$$

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$$\phi_{2,1} = \frac{128p}{\sqrt{3\pi}(4p^2 + 1)^3}$$
(3)

$$\phi_{3,0} = \frac{72}{\sqrt{6\pi}} \left( 4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^4} - 1 \right) \tag{5}$$

$$\phi_{3,1} = \frac{864p}{\sqrt{3\pi}} \frac{(9p^2 - 1)}{(9p^2 + 1)^3} \tag{6}$$

$$\phi_{3,2} = \frac{5184}{\sqrt{15\pi}} \frac{p^2}{(9p^2 + 1)^4} \tag{7}$$

$$\phi_{4,0} = \frac{256}{\sqrt{2\pi}} \frac{(16p^2 - 1)}{(16p^2 + 1)^3} \left( 2 \frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 1 \right)$$
(8)

And since we need probability distributions let's just go ahead and square them all.

$$|\phi_{1,0}|^2 = \frac{32}{\pi (p^2 + 1)^4} = 10.185916 \times \frac{1}{(p^2 + 1)^4}$$
 (9)

$$|\phi_{2,0}|^2 = \frac{1024(4p^2 - 1)^2}{\pi(4p^2 + 1)^6} = 325.94932 \times \frac{(4p^2 - 1)^2}{(4p^2 + 1)^6}$$
(10)

$$|\phi_{2,1}|^2 = \frac{16384p^2}{3\pi(4p^2+1)^6} = 1738.3964 \times \frac{p^2}{(4p^2+1)^6}$$
(11)

$$|\phi_{3,0}|^2 = \frac{864}{\pi} \left( 4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^4} - 1 \right)^2 = 275.01974 \times \left( 4 \frac{(9p^2 - 1)^2}{(9p^2 + 1)^4} - 1 \right)^2$$
 (12)

$$|\phi_{3,1}|^2 = \frac{248832p^2}{\pi} \left(\frac{(9p^2 - 1)}{(9p^2 + 1)^3}\right)^2 = 79205.686 \times p^2 \left(\frac{(9p^2 - 1)}{(9p^2 + 1)^3}\right)^2 \tag{13}$$

$$|\phi_{3,2}|^2 = \frac{8957952}{5\pi} \frac{p^4}{(9p^2+1)^8} = 570280.94 \times \frac{p^4}{(9p^2+1)^8}$$
 (14)

$$|\phi_{4,0}|^2 = \frac{32768}{\pi} \left( \frac{(16p^2 - 1)}{(16p^2 + 1)^3} \right)^2 \left( 2\frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 1 \right)^2$$
(15)

$$= 10430.378 \times \left(\frac{(16p^2 - 1)}{(16p^2 + 1)^3}\right)^2 \left(2\frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 1\right)^2 \tag{16}$$

or alternately (17)

$$= 651.89865 \times \left(\frac{(16p^2 - 1)}{(16p^2 + 1)^3}\right)^2 \left(8\frac{(16p^2 - 1)^2}{(16p^2 + 1)^2} - 4\right)^2 \tag{18}$$

## References

- [1] M. Swartz, H.R. Band, F.J. Decker, P. Emma, M.J. Fero, R. Frey, R. King, A. Lath, T. Limberg, R. Prepost, P.C. Rowson, B.A. Schumm, M. Woods, and M. Zolotorev. Observation of target electron momentum effects in single-arm mller polarimetry. Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment, 363(3):526 537, 1995.
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