## Paramagnetic Behavior of Nickel just Above the Ferromagnetic Curie Temperature

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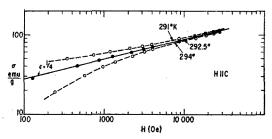


Fig. 3. Magnetization curves near  $T = T_c$ ,  $H \parallel \langle 0001 \rangle$ .

tion curves for Gd near Te for fields between 4.5 and 9 kOe, and have extrapolated these back to H=0 to suggest a nonzero spontaneous magnetization above  $T_{\rm c}$ . In the present work, the magnetization curves showed curvature at all fields up to 29 kOe and at all temperatures up to 350°K; there was no obvious way to extrapolate to H=0 and no reason to suspect any spontaneous magnetization above  $T_{\rm c}$ .

A recent Russian paper7 reports a number of remarkable effects in a Gd crystal above  $T_0$ : kinks in the magnetization curves at 3 kOe from 20° to 100°C; linear magnetization curves above and below the 3-kOe kink, even at To; and paramagnetic anisotropy increasing with temperature, and with the (0001) axis the hard direction. None of these findings were confirmed here or in Ref. 4.

The intercepts on the  $H/\sigma$  axis of Fig. 1 give values of the reciprocal initial susceptibility  $1/\chi_0 = (H/\sigma)_0$ . Various recent theoretical treatments suggest  $1/\chi_0$ =  $A(T-T_c)^{\gamma}$ , with  $\gamma=\frac{4}{3}$  for the Heisenberg ferromagnet.8 Figure 2 shows that the data for Gd fit this relationship reasonably well up to about  $T_e+20^\circ$ . Similar values for  $\gamma$  have been reported for Fe<sup>9,10</sup> and for Ni,8,11 but a somewhat lower value has been found for Co.11

Molecular field theory predicts  $\sigma = CH^{\frac{1}{2}}$  at  $T_c$ . Kouvel and Fisher<sup>8</sup> show that the data of Weiss and Forrer for Ni obey  $\sigma = CH^{0.24}$ , although there is no strong theoretical foundation for such a relationship. Figure 3 indicates that  $\sigma = CH^{\frac{1}{2}}$  for Gd at  $T_c$ , and also shows the rapid departure from this behavior as the temperature departs from  $T_{o}$ .

J. S. Kouvel contributed greatly to the interpretation of the results. The work was supported in part by the Air Force Materials Laboratory, Research and Technology Division, USAF.

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## Paramagnetic Behavior of Nickel just Above the Ferromagnetic Curie Temperature

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Magnetic susceptibility of two nickel spheres of different purity has been studied between  $T_F$  and  $T_F$ + 12°K,  $T_{\rm F}$  being the ferromagnetic Curie temperature. The susceptibility is proportional to  $(T-T_{\rm F})^n$  with the average value of  $n = -1.29 \pm 0.03$ . This result is slightly lower than that (n = -1.33) recently predicted theoretically for a three-dimensional Heisenberg ferromagnet.

ECENTLY the magnetic susceptibilities of iron<sup>1,2</sup> R and cobalt<sup>3</sup> have been studied just above the ferromagnetic Curie temperature  $T_{\rm F}$ . These studies were stimulated by recent theoretical advances4-6 in the study of the magnetic behavior of Heisenberg ferromagnets below and above the ferromagnetic Curie point. According to these investigations the magnetic susceptibility  $\chi$  above the Curie point can be expressed by the equation

$$\chi = A (T - T_{\rm F})^n, \tag{1}$$

where n = -1.33, A is a constant, and T is the absolute

temperature. Noakes and Arrott<sup>1</sup> found experimentally that for iron  $n = -1.37 \pm 0.04$  for temperatures between  $T_{\rm F}$  and  $T_{\rm F}+10^{\circ}{\rm K}$ . Our measurements on iron gave n=-1.33 in the temperature interval  $T_{\rm F}$  to  $T_{\rm F}+30^{\circ}{\rm K}$ . Thus, iron appears to obey the expected theoretical relationship derived for a three-dimensional Heisenberg ferromagnet. On the other hand, the magnetic susceptibility of cobalt<sup>3</sup> between  $T_F$  and  $T_F+13$ °K satisfies Eq. (1) with  $n=-1.21\pm0.04$ . Since the nature of magnetism in iron could be different from that in cobalt and nickel, it was felt that the above-mentioned experimental studies should also be extended to nickel. In this paper we briefly describe the results of such an investigation.

Two nickel spheres designated by A (diameter 0.4750 cm) and B (diameter 0.3180 cm) were used in this study. The sphere A was produced from an electrolytic nickel. Purified Johnson and Matthey nickel was used for making the sphere B. Details on the purity of these

<sup>&</sup>lt;sup>7</sup>V. I. Chechernikov, I. Pop, I. V. Burov, and E. M. Savitskiĭ, Zh. Eksperim. i Teor. Fiz. **45**, 867 (1963) [English transl.: Soviet Phys.—JETP **18**, 595 (1964)].

<sup>&</sup>lt;sup>8</sup> For a summary, see J. S. Kouvel and M. E. Fisher, Phys. Rev. 136, A1626 (1964).

J. E. Noakes and A. Arrott, J. Appl. Phys. 35, 931 (1964).
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<sup>&</sup>lt;sup>1</sup> J. E. Noakes and A. Arrott, J. Appl. Phys. 35, 931 (1964). <sup>2</sup> S. Arajs and R. V. Colvin, J. Appl. Phys. 35, 2424 (1964). <sup>3</sup> R. V. Colvin and S. Arajs, J. Phys. Chem. Solids (to be

published).

<sup>&</sup>lt;sup>4</sup>C. Domb and M. F. Sykes, Proc. Roy. Soc. (London) 240,

<sup>&</sup>lt;sup>6</sup>C. Domb and M. F. Sykes, Phys. Rev. 128, 168 (1962). <sup>6</sup>J. Gammel, W. Marshall, and L. Morgan, Proc. Roy. Soc. (London) A275, 257 (1963).

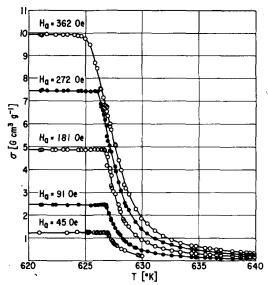


Fig. 1. Magnetization of nickel (sphere A) in various applied magnetic fields in the neighborhood of the ferromagnetic Curie temperature.

materials, which already have been used in some earlier studies,7,8 can be found in Ref. 8. The magnetic moments of these spheres were determined by a force technique described elsewhere.2

The mass magnetization  $\sigma$  of the sphere A as a function of temperature in the neighborhood of  $T_{\rm F}$  in the presence of various applied magnetic fields  $H_a$  is shown in Fig. 1. Similar results have also been obtained for the sphere B. It can be seen that the temperature  $T_{\rm F}$ , defined as a temperature at which the  $\sigma$  vs Tcurve becomes temperature dependent, is a function of the applied magnetic field. This result was found also for cobalt but not for iron. Figure 2 shows  $T_F'$  as a function of the fields  $H_a$ . By extrapolating to  $H_a = 0$ , we obtain 626.2° and 626.9°K as the ferromagnetic Curie temperature for the samples A and B, respectively. The difference in  $T_{\rm F}$  between these samples is, very likely, due to impurities. In an absolute sense, for the same reasons explained elsewhere,2 it is believed that

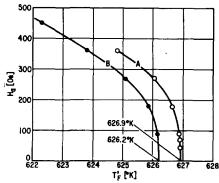


Fig. 2.  $T_{\mathbf{F}}'$  as a function of applied magnetic field for the spheres A and B.

<sup>7</sup>S. Arajs, J. Appl. Phys. 32, 97 (1961). <sup>8</sup> S. Arajs and R. V. Colvin, J. Phys. Chem. Solids 24, 1233 (1963).

the temperature  $T_{\rm F}$  should be reported as  $627\pm1^{\circ}{\rm K}$ which is somewhat lower than the "accepted" value 631°K.9-10 A careful review of numerous earlier determinations reveals considerable variation in  $T_F$ . However, according to Kouvel's analysis11 of the data by Weiss and Forrer,12 using the Belov,13 Arrott,14 and Kouvel<sup>11</sup> technique, with extrapolation of the low magnetic field values, for the determination of the spontaneous magnetization, T<sub>F</sub> is 627.6°K. Essentially the same result is also obtained by Arrott.14 These latter results are in satisfactory agreement with our value.

Figure 3 shows the determination of the quantity n

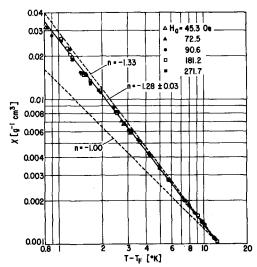


Fig. 3. Temperature dependence of the magnetic susceptibility of nickel (sphere A) above the ferromagnetic Curie temperature.

for the sphere A, for  $T-T_F$  between about 1° and 12°K. We find  $n=-1.28\pm0.03$ . For the sphere B a similar plot gives  $n = -1.30 \pm 0.03$ . Thus, the experimental results on nickel indicate that the value of nis slightly lower than that expected theoretically for a three-dimensional Heisenberg ferromagnet with the nearest-neighbor interactions. For comparison, Fig. 3 also shows the behavior of the magnetic susceptibility resulting from mean molecular field 15,16 for which n=-1.00.

## ACKNOWLEDGMENT

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<sup>15</sup> L. D. Landau and E. M. Lifshitz, Electrodynamics of Continuous Media (Pergamon Press, Inc., New York, 1960), p. 146.
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