

Physics 421: Intro to Electrodynamics

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1 Vector Analysis

1.1 What is a Vector?

In type we use boldface $\mathbf{A} = |\mathbf{A}|\hat{\mathbf{A}}$, where we we can do some simple operations as such:

- Adding and Subtraction: $\mathbf{A} \pm \mathbf{B} = \mathbf{C}$ or aligning the head to the tail
- Multiplication:
 - Scalar: $\mathbf{A} \rightarrow 2\mathbf{A}$
 - Dot Product: $\mathbf{A} \cdot \mathbf{B} = \mathbf{B} \cdot \mathbf{A} = AB \cos \theta$
 - Cross Product: $\mathbf{A} \times \mathbf{B} = AB \sin \theta$, and $\mathbf{B} \times \mathbf{A} = -(\mathbf{A} \times \mathbf{B})$

Components of a Vector In 3D space, we often use the familiar Cartesian coordinates, e.g.

$$\mathbf{A} = A_x \hat{\mathbf{x}} + A_y \hat{\mathbf{y}} + A_z \hat{\mathbf{z}}$$

and we can add components by adding the components:

$$\mathbf{A} + \mathbf{B} = (A_x + B_x) \hat{\mathbf{x}} + (A_y + B_y) \hat{\mathbf{y}} + (A_z + B_z) \hat{\mathbf{z}}$$

and likewise for subtraction. For the dot product:

$$\mathbf{A} \cdot \mathbf{B} = A_x B_x + A_y B_y + A_z B_z$$

or more shortly

$$= \sum_{i,j} A_i B_j \delta_{ij}$$

where δ_{ij} is the Kronecker delta. The cross product is a bit trickier...

$$\mathbf{A} \times \mathbf{B} = \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix} = (A_y B_z - A_z B_y) \hat{\mathbf{x}} - (A_x B_z - A_z B_x) \hat{\mathbf{y}} + (A_x B_y - A_y B_x) \hat{\mathbf{z}}$$

This can also be written in short form using the Levi-Civita symbol (look it up)

Scalar triple product

$$\mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = \mathbf{B} \cdot (\mathbf{C} \times \mathbf{A}) = \mathbf{C} \cdot (\mathbf{A} \times \mathbf{B})$$

This is the also the volume of the parallelepiped formed by the three vectors. NOTE that

$$\underline{(\mathbf{A} \cdot \mathbf{B}) \times \mathbf{C}}$$

since you can't cross a scalar with a vector.

Vector triple product

$$\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$$

Using the BAC-CAB rule.

Some important vectors We define a position vector

$$\mathbf{r} = x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}} = r \hat{\mathbf{r}}$$

where the unit vector is

$$\hat{\mathbf{r}} \equiv \frac{\mathbf{r}}{r}, \quad r = \sqrt{x^2 + y^2 + z^2}$$

and an infinitesimal displacement vector

$$dl = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}$$

In EM we define a source point \mathbf{r}' (e.g. a charge) and a field point \mathbf{r} that give us the separation vector

$$\mathbf{z} = \mathbf{r} - \mathbf{r}'$$

with magnitude

$$|\mathbf{z}| = |\mathbf{r} - \mathbf{r}'|$$

and unit vector

$$\hat{\mathbf{r}} = \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|}$$

1.2 Differential Calculus

And ordinary derivative $\frac{dF}{dx}$ is a change in $F(x)$ in dx

$$dF = \left(\frac{\partial F}{\partial x} \right) dx$$

... geometrically, it's the slope

Gradient for functions of 2 or more variables, generalize for $h(x, y)$

$$dh = \left(\frac{\partial h}{\partial x} \right) dx + \left(\frac{\partial h}{\partial y} \right) dy$$

it's a scalar so $dh = (\nabla h) \cdot (dl)$ where

$$\nabla h = \frac{\partial h}{\partial x} \hat{\mathbf{x}} + \frac{\partial h}{\partial y} \hat{\mathbf{y}}$$

In 3D

$$\nabla T = \frac{\partial T}{\partial x} \hat{\mathbf{x}} + \frac{\partial T}{\partial y} \hat{\mathbf{y}} + \frac{\partial T}{\partial z} \hat{\mathbf{z}}$$

If $\nabla u = 0$, we are at an extremum (max, min, or shoulder/saddle point) Rewriting:

$$\nabla T = \left(\hat{\mathbf{x}} \frac{\partial}{\partial x} + \hat{\mathbf{y}} \frac{\partial}{\partial y} + \hat{\mathbf{z}} \frac{\partial}{\partial z} \right) T(x, y, z)$$

where we can assume the ∇ as an “operator” acting on T :

1. Scalars like T : ∇T , “grad T”, generalized slope
2. Dot product on \mathbf{V} : $\nabla \cdot \mathbf{V}$, “divergence” or “div”
3. Cross product : $\nabla \times \mathbf{V}$, “curl” or “rotatation”

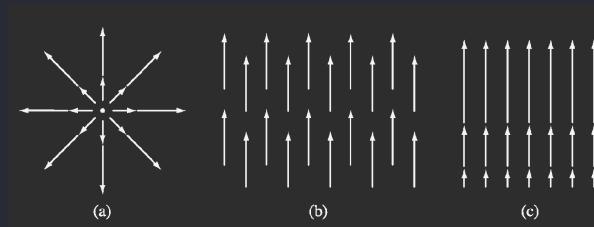


Figure 1.1: Divergence of field lines

Divergence From the Figure, we can see that (a) & (c) diverges, and (b) does not.

Geometrical Interpretation: Sources of positive divergence is a source or “faucet”, and negative divergence is a sink or “drain”.

Curl

$$\nabla \times \mathbf{V} = \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ V_x & V_y & V_z \end{vmatrix}$$

E.g. for $\mathbf{V} = -y\hat{\mathbf{x}} + x\hat{\mathbf{y}}$, $\nabla \times \mathbf{V} = 2\hat{\mathbf{z}}$.

Combining Multiple Operations Two ways to get scalar from two functions:

$$fg \quad \text{or} \quad \mathbf{A} \cdot \mathbf{B}$$

Two ways to get vector from two functions:

$$f\mathbf{A} \quad \text{or} \quad \mathbf{A} \times \mathbf{B}$$

And we have 3 ‘derivatives’: div, grad, and curl.

Product rule:

$$\partial_x(fg) = f\partial_x g + g\partial_x f$$

$$\text{i } \nabla(fg) = f\nabla g + g\nabla f$$

$$\text{ii } \nabla(\mathbf{A} \cdot \mathbf{B}) = \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) + (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A}$$

... (see Griffiths for more)

Second Derivatives Combining ∇ , $\nabla \cdot$, $\nabla \times$
 ∇T is a vector

i

$$\begin{aligned}\nabla \cdot (\nabla T) &= (\hat{x}\partial_x + \hat{y}\partial_y + \hat{z}\partial_z) \cdot (\hat{x}\partial_x T + \hat{y}\partial_y T + \hat{z}\partial_z T) \\ &= \partial_x^2 T + \partial_y^2 T + \partial_z^2 T \\ &= \nabla^2 T\end{aligned}$$

ii $\nabla \times (\nabla T) = 0$

iii $\nabla(\nabla \cdot \mathbf{v}) = \dots$ ignored

iv $\nabla \cdot (\nabla \times \mathbf{v}) = 0$

v $\nabla \times (\nabla \times \mathbf{v}) = \nabla(\nabla \cdot \mathbf{v}) - \nabla^2 \mathbf{v}$

1.3 Integral Calculus:

line, surface and volume integrals

“Fundamental theorem for gradients” Start with a scalar $T(x, y, z)$: from $a \rightarrow b$, in small steps $dT = \nabla \cdot T d\ell$

Total change in T :

$$\boxed{\int_a^b dT = \int_a^b \nabla T \cdot d\ell = T(b) - T(a)}$$

This line integral is path independent but $\int_a^b \mathbf{F} \cdot d\ell$ is *not!*

Divergence Theorem, “Gauss’ Theorem”, or “Green’s Theorem”

$$\boxed{\int_V (\nabla \cdot \mathbf{v}) d\tau = \oint_S \mathbf{v} \cdot d\mathbf{a}}$$

where V is the volume enclosed by the surface S . The divergence of a field in a volume is equivalent to the flux of the field at the boundary or surface.

Geometrical Interpretation: The “source”(or faucet) should present a flux (or flow) out through the surface.

Fundamental Theorem of Curls: Stokes’ Theorem

$$\boxed{\oint_S (\nabla \times \mathbf{v}) \cdot d\mathbf{a} = \oint_P \mathbf{v} \cdot d\ell}$$

We have a 2D surfaces S bounded by a closed 1D perimeter P .

In 3D, imagine a soap bubble in a wire frame. We can change the surface of the bubble by moving the wire frame and making almost making a bubble, but the perimeter that bounds the bubble is still the wireframe.

Example:

$$\mathbf{v} = (2xz + 3y^2)\hat{\mathbf{y}} + 4yz^2\hat{\mathbf{z}}$$

On a surface S square on the $y - z$ plane:

$$\oint \mathbf{v} \cdot d\ell =$$

RHS: We can split this into 4 integrals going counterclockwise around the square: First $x = 0, z = 0, y : 0 \rightarrow 1$: $dx = dz = 0, \int_0^1 3y^2 dy = 1$

Second $\int_0^1 4z^2 dz = 4/3$

Third: -1

Fourth: 0

Summing them all gives: $\oint \mathbf{v} \cdot d\ell = 4/3$

LHS: The curl gives: $4z^2 - 2x, -(0 - 0), 2z$ so

$$\oint (\nabla \times \mathbf{v}) \cdot d\mathbf{a} = \oint$$

...

1.4 Dirac Delta Function

Considering a vector function

$$\mathbf{v} = \frac{1}{r^2} \hat{\mathbf{r}}$$

[insert picture of field lines]

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{1}{r} \right) = 0$$

The div theorem doesn't obviously tell what volume and surface to choose, but we can choose a sphere of radius R and its corresponding surface:

$$\oint_S \mathbf{v} \cdot d\mathbf{a} = \int \frac{1}{R^2} \hat{\mathbf{r}} \cdot \hat{\mathbf{r}} R^2 \sin \theta d\theta d\phi = 4\pi$$

where the polar angle is $\theta : 0 \rightarrow \pi$ and the azimuthal angle is $\phi : 0 \rightarrow 2\pi$.

We think that the divergence is zero (thus the theorem does not hold), but we make a tiny mistake:

$\nabla \cdot \mathbf{v} = 0$ everywhere except at the origin $r \rightarrow 0$

and we stumbled on the Dirac delta function:

$$\delta(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \end{cases} \quad \text{and} \quad \int_{-\infty}^{\infty} \delta(x) dx = 1$$

where

$$\int f(x) \delta(x) dx = f(0)$$

Shifting the delta function:

$$\delta(x - a) = \begin{cases} 0 & x \neq a \\ \infty & x = a \end{cases}$$

so

$$\int f(x) \delta(x - a) dx = f(a)$$

Multiplying by a constant:

$$\delta(kx) = \frac{1}{|k|} \delta(x)$$

Generalizing to 3D:

$$\delta^3(\mathbf{r}) = \delta(x) \delta(y) \delta(z)$$

2 Electrostatics

2.1 The Electric Field

given charge q : find force on Q , where \mathbf{F} depends on $\hat{\mathbf{z}}, \mathbf{v}_i, \mathbf{a}_i$

2.1.1 Coulombs law

Coulomb's Law empirically,

$$\mathbf{F}_Q = \frac{1}{4\pi\epsilon_0} \frac{qQ}{\hat{\mathbf{z}}^2} \hat{\mathbf{z}}$$

where $k = \frac{1}{4\pi\epsilon_0}$ and the permittivity of free space is $\epsilon_0 = 8.85 \times 10^{-12} \text{ C}^2/\text{Nm}^2$

The force is attractive if $\text{sgn}(qQ) = -1$ and repulsive if $= +1$.

Principal of superposition:

$$\begin{aligned} \mathbf{F}_T &= \mathbf{F}_{Q1} + \mathbf{F}_{Q2} + \dots \\ &= \frac{1}{4\pi\epsilon_0} Q \left(\frac{q_1}{\hat{\mathbf{z}}_1^2} \hat{\mathbf{z}}_1 + \frac{q_2}{\hat{\mathbf{z}}_2^2} \hat{\mathbf{z}}_2 + \dots \right) \\ &= Q \mathbf{E}_T \end{aligned}$$

where \mathbf{E}_T is the total electric field due to all of the source (point) charges.

\mathbf{E} doesn't depend on Q

- $\mathbf{E} \sim F/Q$

Example: \mathbf{E} field midway above two charges q : The electric fields are zero in the x and y direction:

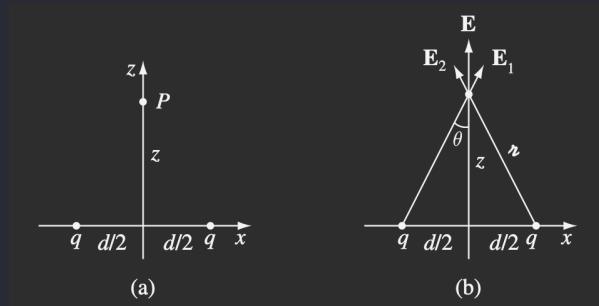


Figure 2.1: Griffiths Example 2.1

$$E_x = E_y = 0$$

But we can sum the fields in the z direction:

$$E_z = 2 \frac{1}{4\pi\epsilon_0} \frac{q}{\hat{\mathbf{z}}^2} \cos \theta$$

where

$$\hat{\mathbf{z}} = \left[z^2 + \left(\frac{d}{2} \right)^2 \right]^{1/2} \quad \cos \theta = \frac{z}{\hat{\mathbf{z}}}$$

so

$$E_z = \frac{1}{4\pi\epsilon_0} \frac{2qz}{\left[z^2 + \left(\frac{d}{2}\right)^2\right]^{3/2}}$$

Far away: $z \gg d$, so $d \rightarrow 0$ thus

$$E_z \approx \frac{1}{4\pi\epsilon_0} \frac{2qz}{z^3} = \frac{1}{4\pi\epsilon_0} \frac{2}{z^2}$$

Continuous Charge Distributions

- line: charge per unit length λ ; $dq = \lambda d\ell$
- surface: charge per unit area σ ; $dq = \sigma da$
- volume: charge per unit volume ρ ; $dq = \rho d\tau$

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{1}{\mathbf{z}^2} \hat{\mathbf{z}} dq$$

e.g. for a volume charge:

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{\mathbf{z}^2} \hat{\mathbf{z}} d\tau'$$

where ' denotes the source charge in (no ' is a field point)

Example: Find \mathbf{E} at z above a straight line segment of length $2L$ with uniform line charge λ . If we

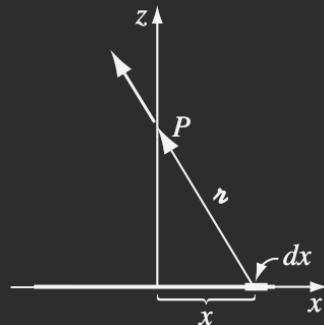


Figure 2.2: Griffiths Example 2.2

treat dq as a point particle, then we can use Ex 2.1 likewise but integrate over the line segment.

First we catalog what we know:

- Field point P is at $\mathbf{r} = z\hat{\mathbf{z}}$
- Sources at $\mathbf{r}' = x\hat{\mathbf{x}}$; $d\ell' = dx$
- $\mathbf{z} = \mathbf{r} - \mathbf{r}' = z\hat{\mathbf{z}} - x\hat{\mathbf{x}}$
- $z = \sqrt{x^2 + z^2}$
- $\hat{\mathbf{z}} = \frac{\mathbf{z}}{z} = \frac{z\hat{\mathbf{z}} - x\hat{\mathbf{x}}}{\sqrt{x^2 + z^2}}$

The electric field is then (line charge)

$$\begin{aligned}\mathbf{E}(\mathbf{r}) &= \frac{1}{4\pi\epsilon_0} \int_{-L}^{+L} \frac{\lambda}{z^2} \hat{z} dx = \frac{1}{4\pi\epsilon_0} \lambda \int_{-L}^{+L} \frac{z\hat{z} - x\hat{x}}{(z^2 + x^2)^{3/2}} dx \\ &= \frac{1}{4\pi\epsilon_0} \lambda \left[z\hat{z} \int_{-L}^L \frac{dx}{(z^2 + x^2)^{3/2}} - \hat{x} \int_{-L}^L \frac{x dx}{(z^2 + x^2)^{3/2}} \right] \\ &= \frac{1}{4\pi\epsilon_0} \lambda \left[z\hat{z} \frac{x}{z^2\sqrt{z^2 + x^2}} \Big|_{-L}^L - \hat{x} \frac{1}{\sqrt{z^2 + x^2}} \Big|_{-L}^L \right]\end{aligned}$$

we can easily see that the x component is zero through the geometrical symmetry of the line centered at the origin (like Ex 2.1). Simplifying gives us

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z\sqrt{z^2 + L^2}} \hat{z}$$

Checks and balances:

- \hat{z} is expected!

•

$$z \gg L \quad \sqrt{z^2 + L^2} \approx z \quad E(P, z \gg L) = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z^2}$$

where we can treat this as a point charge $q = 2\lambda L$ when we are far away.

2.2 Divergence and curl of \mathbf{E} : Gauss' Law

'flux' of field lines

$$\Phi = \oint_S \mathbf{E} \cdot d\mathbf{a}$$

What is Φ for point charge at origin surrounded by a spherical surface?

$$\begin{aligned}\Phi &= \int \mathbf{E} \cdot d\mathbf{a} = \int \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} \cdot \hat{r} r^2 \sin\theta d\theta d\phi \\ &= \frac{q_{enc}}{\epsilon_0}\end{aligned}$$

A bunch of charges surrounded by a surface: $\mathbf{E}_T = \sum \mathbf{E}_i$

$$\Phi = \oint \mathbf{E}_T \cdot d\mathbf{a} = \sum_i \oint \mathbf{E}_i \cdot d\mathbf{a} = \sum_i \frac{q_i}{\epsilon_0}$$

Thus we have an integral form of Gauss's law:

$$\boxed{\oint \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{enc}}{\epsilon_0}}$$

where $Q = \sum q_i$.

From the theorem of divergence:

$$\oint_S \mathbf{v} \cdot d\mathbf{a} = \int_V (\nabla \cdot \mathbf{v}) d\tau \quad \text{and} \quad Q = \int_V \rho d\tau$$

so

$$\int_V (\nabla \cdot \mathbf{E}) d\tau = \int_V \rho d\tau \rightarrow \text{good for all volume}$$

therefore we have the differential form of Gauss' Law:

$$\boxed{\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}}$$

Three ways Gauss's law makes life nice: Gaussian surfaces

- spherical: gaussian sphere
- cylindrical: gaussian cylinder
- planar: gaussian pillbox

2.2.1 Applications of Gauss's Law

$$\int \mathbf{E} \cdot d\mathbf{a} = \frac{q_{enc}}{\epsilon_0} \rightarrow \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

1. (Simple spherical) What is \mathbf{E} outside a uniformly charged solid sphere of radius R and total charge Q ? The spherical Gaussian surface implies a symmetry where we should *only have a radial component* $\mathbf{E} = E_r$.

$$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{Q}{\epsilon_0}$$

$$E \oint d\mathbf{a} = E \cdot 4\pi r^2 = \frac{Q}{\epsilon_0}$$

$$\implies \mathbf{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{\mathbf{r}}$$

where the integral is equivalent to the surface area of the sphere. This is also \implies a field of a *point*.

2. (Simple cylindrical) A long cylinder (radius a) of charge density $\rho = ks$ (\propto distance from axis) where k is a constant and s is the radial distance from the axis. What is \mathbf{E} inside the cylinder? The cylindrical Gaussian surface has radius s and length ℓ :

$$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{enc}}{\epsilon_0}; \quad Q_{enc} = \int \rho d\tau = \int (ks') ds' d\phi dz = \frac{2}{3}\pi k \ell s^3$$

When using the divergence theorem, note that only the curved part of the cylinder contributes to the flux. Thus,

$$\int \mathbf{E} \cdot d\mathbf{a} \rightarrow E \int da = E(2\pi s \ell)$$

$$\implies \mathbf{E} = \frac{1}{3\epsilon_0} ks^2 \hat{\mathbf{s}}$$

If we were to find the field outside the cylinder we would find that the enclosed charge is constant Q_{enc} thus the field is proportional to $1/s$.

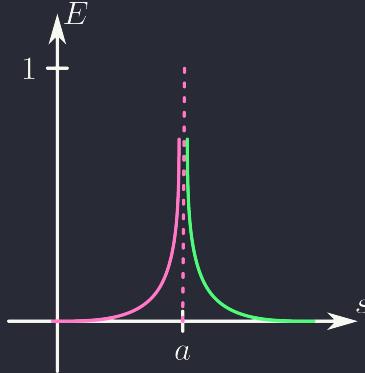


Figure 2.3: Electric field as a function of s

3. (Simple infinite plane) with uniform surface charge σ . Symmetry implies that \mathbf{E} is perpendicular to the plane. The Gaussian pillbox (either box or cylinder) will have a field of

$$\mathbf{E} = \frac{\sigma}{2\epsilon_0} \hat{\mathbf{n}}$$

2.2.2 The curl of \mathbf{E}

$$\nabla \times \mathbf{E} = 0, \quad \mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$

calculating

$$\int_a^b \mathbf{E} \cdot d\ell, \quad d\ell = dr\hat{\mathbf{r}} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi}$$

So the integral is

$$\frac{1}{4\pi\epsilon_0} \int_a^b \frac{q}{r^2} (\hat{\mathbf{r}} \cdot \hat{\mathbf{r}}) dr = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{a} - \frac{q}{b} \right)$$

This means:

- path independent!
- if $a = b$ then $\oint \mathbf{E} \cdot d\ell = 0$ (ℓ is a vector but I don't know how to bold it)

We can now use Stokes' theorem: $\oint \mathbf{v} \cdot d\ell = \int_S (\nabla \times \mathbf{v}) \cdot da$ or

$$\oint \mathbf{E} \cdot d\ell = \int_S (\nabla \times \mathbf{E}) \cdot da = 0 \implies \nabla \times \mathbf{E} = 0$$

2.3 Electric potential

Any function f with zero curl is the gradient of a scalar function: $\nabla \times (\nabla f) = 0$ (curl of gradient is always 0!)

$$V(\mathbf{r}) = - \int_{\mathcal{O}}^{\mathbf{r}} \mathbf{E} \cdot d\ell$$

implies all paths give some value.

$V \sim$ “electric potential”

$$\begin{aligned} V(\mathbf{b}) - V(\mathbf{a}) &= - \left(\int_{\mathcal{O}}^b \mathbf{E} \cdot d\ell \right) - \left(- \int_{\mathcal{O}}^a \mathbf{E} \cdot d\ell \right) \\ &= - \int_{\mathcal{O}}^b - \int_{\mathcal{O}}^a O \mathbf{E} \cdot d\ell \\ &= - \int_a^b \mathbf{E} \cdot d\ell \end{aligned}$$

And from the fundamental theorem for gradients: $T(\mathbf{b}) - T(\mathbf{a}) = \int_a^b (\nabla T) \cdot d\ell$

$$\implies \mathbf{E} = -\nabla V$$

i “potential” is a terrible name

ii $\mathbf{E} = (E_x, E_y, E_z)$ vs V with only *one* value at every point in space! Otherwise we would have to deal with

$$(\nabla \times \mathbf{E})_x = \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = 0$$

iii

$$V'(\mathbf{r}) = - \int_{O'}^{\mathbf{r}} \mathbf{E} \cdot d\ell = - \int_{O'}^O \mathbf{E} \cdot d\ell - \int_O^{\mathbf{r}} \mathbf{E} \cdot d\ell = C + V(\mathbf{r})$$

$$\implies \mathbf{E} = -\nabla V$$

Electric Potential cont.

2.3.1 Poisson's and Laplace's equations

The divergence and curl of \mathbf{E} in terms of the potential V :

- Divergence: $\nabla \cdot \mathbf{E} = \nabla \cdot (-\nabla V) = -\nabla^2 V$, or **Poisson's equation**:

$$\boxed{\nabla^2 V = -\frac{\rho}{\epsilon_0}}$$

and in regions of no charge $\rho = 0$ we have **Laplace's equation**:

$$\boxed{\nabla^2 V = 0}$$

- Curl: It's always zero, which doesn't give us any info about V ...

2.3.2 Potential of localized charge distributions

For a point charge q we can easily find the potential using the electric field $\mathbf{E} = \frac{q}{4\pi\epsilon_0 r^2} \hat{\mathbf{r}}$:

$$\begin{aligned} V(r) &= - \int_O^r \mathbf{E} \cdot d\ell = -\frac{1}{4\pi\epsilon_0} \int_{\infty}^r \frac{q}{r'^2} dr' \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{r} \end{aligned}$$

where in general, the potential of a point charge is

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{q}{\mathbf{r}}$$

And using the principle of superposition, the potential of a collection of charges can be a sum

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_i \frac{q_i}{\mathbf{r}_i}$$

or for a continuous charge distribution, and integral:

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{1}{\mathbf{r}} dq$$

e.g. for a volume charge distribution:

$$\boxed{V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{\mathbf{r}} d\tau'}$$

2.3.3 Boundary conditions

From the Gaussian pillbox the E-field is given by Gauss's law:

$$\oint_S \mathbf{E} \cdot d\mathbf{A} = \frac{Q_{\text{enc}}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}$$

and the only components that contribute to the field are the top and bottom:

$$\mathbf{E}_{\text{above}}^\perp - \mathbf{E}_{\text{below}}^\perp = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$

But we can also use the path integral of the E-field (which should be zero):

$$\oint \mathbf{E} \cdot d\ell = 0$$

so the parallel components should be equal:

$$\begin{aligned} \mathbf{E}_{\text{above}}^{\parallel} - \mathbf{E}_{\text{below}}^{\parallel} &= 0 \\ \implies E_{\text{above}}^{\parallel} &= E_{\text{below}}^{\parallel} \end{aligned}$$

This also implies that the potential “ V is continuous across boundaries”:

$$V_{\text{above}} = V_{\text{below}}$$

which also implies the gradient of the potential is

$$\nabla V_{\text{above}} - \nabla V_{\text{below}} = -\frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$

2.4 Work and energy in electrostatics

$[V] = \text{J/C} = \text{energy/charge}$, so the work is equivalent to

$$W = QV(\mathbf{r})$$

2.4.1 Energy of point charge distribution

Collection of 3 charges First we assume we start with $W_1 = 0$, or we are bringing the two other charges towards q_1 ; Work done to bring in $q_2 = V_1 q_2 = W_2$;

$$\begin{aligned} W_3 &= V_1 q_3 + V_2 q_3 \\ &= \frac{1}{4\pi\epsilon_0} \frac{q_1 q_3}{r_{13}} + \frac{1}{4\pi\epsilon_0} \frac{q_2 q_3}{r_{23}} \\ &= \frac{1}{4\pi\epsilon_0} q_3 \left(\frac{q_1}{r_{13}} + \frac{q_2}{r_{23}} \right) \end{aligned}$$

where we can use the superposition principle to find the total work:

$$W = W_1 + W_2 + W_3$$

In general

$$\begin{aligned} W &= \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \sum_{j>i}^n \frac{q_i q_j}{r_{ij}} \\ &= \frac{1}{2} \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \sum_{j \neq i} \frac{q_i q_j}{r_{ij}} \\ &= \frac{1}{2} \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n q_i \sum_{j \neq i} \frac{q_j}{r_{ij}} \\ &= \frac{1}{2} \sum_{i=1}^n q_i \left[\sum_{j \neq i} \frac{1}{4\pi\epsilon_0} \frac{q_j}{r_{ij}} \right] \end{aligned}$$

where the term in the brackets is the potential energy due to all other charges, $[] = V(\mathbf{r}_i)$, thus

$$W = \frac{1}{2} \sum_{i=1}^n q_i V(\mathbf{r}_i)$$

2.4.2 Energy of continuous charge distribution

$$W = \frac{1}{2} \int \rho V d\tau$$

but what if we don't know the potential? First using Gauss's law:

$$\rho = \epsilon_0 \nabla \cdot \mathbf{E}$$

so the work done is (using integration by parts):

$$\begin{aligned} W &= \frac{\epsilon_0}{2} \int (\nabla \cdot \mathbf{E}) V d\tau \\ &= \frac{\epsilon_0}{2} \left[- \int \mathbf{E} (\nabla V) d\tau + \oint V \mathbf{E} \cdot d\mathbf{A} \right] \\ &= \frac{\epsilon_0}{2} \left[\int E^2 d\tau + \oint V \mathbf{E} \cdot d\mathbf{A} \right] \end{aligned}$$

but we know that

- $E \propto \frac{1}{r^2}$ $d\tau \propto r^3$
- $V \propto \frac{1}{r}$ $dA \propto r^2$

So the second surface integral term goes to 0 much faster, thus

$$W = \frac{\epsilon_0}{2} \int E^2 d\tau$$

Energy Review:

$$W = \frac{1}{2} \sum q_i V(\mathbf{r}_i)$$

$$W = \frac{\epsilon_0}{2} \int E^2 d\tau$$

Where did we go wrong in going from the point charge distribution to the continuous distribution?

2.4.3 Comments on Electrostatic energy

Energy of a point charge?:

$$W = \frac{\epsilon_0}{2} \int \frac{q^2}{(r^2)^2} r^2 \sin \theta dr d\theta d\phi = \infty$$

Uh oh, as $r \rightarrow 0$ the energy goes to infinity! And this problem persists everywhere. (Solution is renormalization...)

Where is the energy stored?

- In the first case: we have a charge in a potential (from other charges); somehow the relative position of the charge implies its “electric potential”

We don’t really know where it is, but we can book keep this energy i.e. ρ, V, E

Superposition (again) If we try to superimpose two electric fields,

$$W = \frac{\epsilon_0}{2} \int |\mathbf{E}_1 + \mathbf{E}_2|^2 d\tau$$

the terms in the integral are

$$E_1^2 + E_2^2 + 2\mathbf{E}_1 \cdot \mathbf{E}_2$$

so

$$W \neq W_1 + W_2$$

2.5 Conductors in electrostatics**2.5.1 Basic Properties**

What is a conductor? Material where charge can move; in the perfect case, the charges move freely and the resistance goes to zero $R \rightarrow 0$. Thus no *resistance* to the motion of charge.

- i $\mathbf{E} = 0$ inside a conductor
if $E + q \rightarrow \mathbf{F} = q\mathbf{E}$: a charge in a field \rightarrow charges move to cancel \mathbf{E}_{in}
- ii $\rho = 0$ inside a conductor; so from divergence

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon}, \quad \text{if } \mathbf{E} = 0, \rho = 0$$

- iii any non-zero ρ resides on surface (only in 3D); in 2D & 1D, it’s not exactly at the boundary, but quickly peaks and decays before the boundary.
- iv conductors are equipotentials
- v \mathbf{E} is perpendicular to a surface outside the conductor

The silly experiment: An aluminum beverage can (of ingenious design) is a conductor, i.e., it has free electrons. The plastic rod is given negative charge (via cloth/wool) and when it is brought near the can, the can moves towards the rod.

How does the can move, even though it has no charge? The rod induces a charge by pushing the negative charges away from the can. Then the positive charges close to the rod are attracted hence the attractive movement (force).

2.5.2 Induced charge distributions

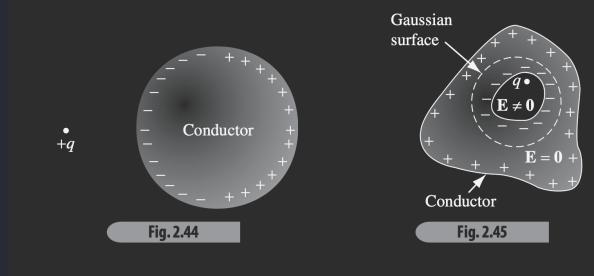


Figure 2.4: From Griffiths

Case 1: A point charge in a cavity surrounded by a conductor

- Flux $\Phi = \int \mathbf{E} d\mathbf{A} \neq 0$ for a Gaussian surface in the cavity
- $\Phi = \int \mathbf{E} d\mathbf{A} = 0$ for a surface where that encloses no charge (positive point charge cancels negative charges surrounding it in the conductor)
- $\int \mathbf{E} d\mathbf{A} = \frac{q_{enc}}{\epsilon_0} = \frac{q}{\epsilon_0}$ for a Gaussian surface outside the conductor

Case 2: Spherical conductor, with a weird cavity containing a point charge $+q$

What is $\mathbf{E}(\mathbf{P})$ at \mathbf{P} for $P > r$ (outside the conductor)? The electric field will uniformly distribute itself in the conductor, so

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$$

What about a point charge $+q$ outside the conductor, what is the E-field in the cavity? It may put minus charges close, and plus charges away from further away on the surface of the conductor...but in the cavity, you can't transmit information in the cavity since the E-field is zero inside a conductor. Screening out an electric field AKA the Faraday cage.

2.5.3 Capacitors

Imagine 2 conductors (of weird shapes) with charges $+Q$ and $-Q$ has a potential of $+V$ and $-V$ respectively. The potential difference (Theorem of gradients) is

$$\Delta V = V_+ - V_- = - \int_{(-)}^{(+)} \mathbf{E} \cdot d\ell$$

where we know that there is an electric field $\mathbf{E} \propto Q$ or more precisely from Coloumb's law:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho}{\mathbf{r}^2} \hat{\mathbf{r}} d\tau$$

thus

$$\Delta V \propto Q \implies C \equiv \frac{Q}{\Delta V}$$

where C is the “capacitance”.

Example: Parallel plate capacitor Two very large plates with area A , charge $+Q$ & $-Q$, a separation d such that $d \ll \sqrt{A}$, and the surface charge density is $\sigma_{\pm} = \pm \frac{Q}{A}$

We can easily find the electric field of a plate $\mathbf{E} = \frac{\sigma}{2\epsilon_0} \hat{\mathbf{n}}$, and in the presence of two oppositely charged plates, there is only an electric field between the plates with double the magnitude:

$$E = \frac{\sigma}{\epsilon_0}$$

The potential difference is then

$$\begin{aligned}\Delta V &= \int \mathbf{E} \cdot d\ell \\ &= \frac{\sigma}{\epsilon_0} d = \frac{d}{\epsilon_0} \frac{Q}{A}\end{aligned}$$

so the capacitance of the parallel plates is

$$C_{\parallel} = \frac{Q}{\Delta V} = \frac{\epsilon_0 A}{d} = []$$

3 Potentials

3.1 Laplace's Equation

3.1.1 Intro

In principle, electrostatics is

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\hat{\mathbf{z}}}{\mathbf{r}^2} \rho(\mathbf{r}') d\tau', \quad \mathbf{z} = \mathbf{r} - \mathbf{r}'$$

And simplifying with potential

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{\mathbf{r}'} d\tau'$$

So we often use Poisson's equation e.g.

$$\nabla^2 V = -\frac{\rho}{\epsilon_0}$$

Even better, Laplace's equation

$$\boxed{\nabla^2 V = 0}$$

or in Cartesian

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

3.1.2 Start in 1D

$$\frac{\partial^2 V}{\partial x^2} = 0 \implies V = mx + b$$

where we have two undetermined constants m & b . We can determine these constants by *boundary conditions*.

- e.g. $V(1) = 4$ $V(5) = 0$; we get a line $V = -\frac{4}{5}x + 4$

Two features:

1. $V(x)$ is average of $V(x+a)$ and $V(x-a)$

$$V(x) = \frac{1}{2}[V(x+a) + V(x-a)]$$

2. **NO** local minima or maxima (no curvature!)

3.1.3 On to 2D

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0 \quad \text{no general solution}$$

and no requirement on the # of constants. But we can note common properties e.g. soap film on a wireframe assumes the same shape. The solutions are called *harmonic functions*:

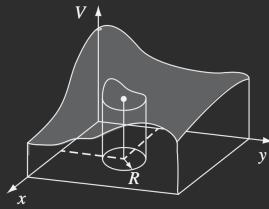


Fig. 3.2

Figure 3.1

1. value $V(x, y)$ is average of nearby values; more precisely, for a circle of radius R (Fig. 3.1)

$$V(x, y) = \frac{1}{2\pi R} \oint V d\ell$$

where $2\pi R$ is the circumference of the circle.

2. **NO** local minima or maxima

3.1.4 In 3D

$$\nabla^2 V = 0$$

Holds same properties as 2D:

1. Average over spherical surface of radius R centered at \mathbf{r} :

$$V(\mathbf{r}) = \frac{1}{4\pi R^2} \oint_S V da$$

where $4\pi R^2$ is the surface area of the sphere.

Example: Point charge outside sphere; The potential at da

$$V = \frac{1}{4\pi\epsilon_0} \frac{q}{z}$$

and from law of cosines

$$z^2 = z^2 + R^2 - 2rR \cos \theta$$

so the average potential is

$$\begin{aligned} V_{\text{avg}} &= \frac{q}{4\pi R^2} \frac{1}{4\pi\epsilon_0} \int \frac{R^2 \sin \theta d\theta d\phi}{\sqrt{z^2 + R^2 - 2zR \cos \theta}} \\ &= \frac{q}{2zR} \frac{1}{4\pi\epsilon_0} [z^2 + R^2 - 2zR \cos \theta]^{1/2} \Big|_0^\pi \\ &= \frac{q}{2zR} \frac{1}{4\pi\epsilon_0} [(z + R) - (z - R)] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{z} \end{aligned}$$

which is just the potential of a point charge q in the center of the sphere.

Question: Is it possible to stably trap a charged particle using electrostatic forces alone?

Answer Earnshaw's theorem: A charged particle cannot be held in stable equilibrium by electrostatic forces alone.

3.1.5 Boundary Conditions & Uniqueness Theorem

Laplace's eq requires boundary conditions (b.c.c)

1st uniqueness theorem: "Solutions to L's eq in volume V is uniquely determined if potential is specified in surface S bounding V ."

How is the solution unique?

Have solution in V_1 s.l.

$$\nabla^2 V_1 = 0 \quad \text{also} \quad \nabla^2 V_2 = 0$$

Then

$$V_3 \equiv \Delta V = V_1 - V_2$$

which means

$$\nabla^2 V_3 = \nabla^2 V_1 - \nabla^2 V_2 = 0 - 0 = 0$$

thus

$$\begin{aligned} \implies \nabla^2 V_1 &= \nabla^2 V_2 \\ V_1 &= V_2 \end{aligned}$$

We should emphasize that V_1 defined on the boundary S is also the same b.c.s as V_2 while V_3 on the boundary equals zero. That is V_3 is zero everywhere in space.

3.1.6 2nd uniqueness theorem (conductors):

"In a volume V surrounded by conductors, and containing a specified charge density ρ , then the \mathbf{E} field is uniquely determined if the total charge on each conductor is given."

"This proof was not easy" - Griffiths

Example: Connecting two pairs of opposite charges with a conductor; what is the final charge config and E-field?

Answer: $\mathbf{E} = 0$ everywhere. The total charge of each conductor is zero.

3.1.7 Boundary conditions pt. II

(Griffiths 2.3.5 pg 85) Given a sheet of charge $\sigma = Q/A$ and using the Gaussian pillbox method

$$\begin{aligned} \oint \mathbf{E} \cdot d\mathbf{a} &= \frac{q_{\text{enc}}}{\epsilon_0} \\ E_a A - E_b A &= \sigma \frac{A}{\epsilon_0} \\ E_a - E_b &= \frac{\sigma}{\epsilon_0} \end{aligned}$$

For the same surface we know that

$$\nabla \times \mathbf{E} = \oint \mathbf{E} \cdot d\ell = 0$$

So going around the loop we have

$$\begin{aligned} E_a \ell - E_b \ell &= 0 \\ E_a &= E_b \end{aligned}$$

3.2 Method of Images

$\nabla^2 V = -\rho/\epsilon_0$ is electrostatics. AND when we have $\nabla^2 V = 0$, Uniqueness theorems tell us there is a solution, but doesn't tell us how to find it...thus we have a set of "easily" solvable problems.

3.2.1 Classic image problem:

"Ground" is infinite conducting plane $V = 0$ at $z = 0$. For a charge q at $z = d$ what is the potential in $z > 0$? We know that the point charge will induce a charge on the plane which will effect the electric potential in the region $z > 0$.

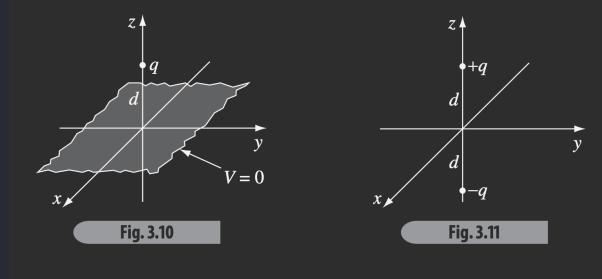


Figure 3.2

GOAL: solve $\nabla^2 V = -\rho/\epsilon_0$ for $z \geq 0$ with $+q$ at $(0, 0, d)$ subject to boundary conditions (b.c.s)

1. $V(z = 0) = 0$
2. $V \rightarrow 0$ for far away

The next step is to replace the conducting plane with an equivalent charge $-q$ at $z = -d$.

NOTE: We only care about $z \geq 0$ and ignore $z < 0$ region; furthermore, the potential below the plane should be zero as the boundary of the plane sort of "wraps" around the charge

Thus the solution is a superposition of charges

$$V = \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{x^2 + y^2 + (z-d)^2}} - \frac{q}{\sqrt{x^2 + y^2 + (z+d)^2}} \right]$$

we can see that in the replacement problem, the region below the grounded plane is nonzero which is different from the real problem which is why we ignore it!

3.2.2 Induced surface charge

What is σ ? Recall on the sheet σ we have a field

$$\mathbf{E}_{ab} - \mathbf{E}_{be} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$

where the $\hat{\mathbf{n}}$ tells us that we are dealing with the perpendicular planes; thus the same eq

$$\nabla V_{ab} - \nabla V_{be} = -\frac{\sigma}{\epsilon_0} \hat{\mathbf{n}} \implies \mathbf{E} = -\nabla V$$

We can infer that

$$\frac{\partial V_a}{\partial n} - \frac{\partial V_b}{\partial n} = -\frac{\sigma}{\epsilon_0}$$

since anything below the plane is zero from the previous example. We know have

$$\frac{\partial V_a}{\partial z} \Big|_{z \rightarrow 0} = \frac{1}{4\pi\epsilon_0} \left[-\frac{q(z-d)}{(x^2 + y^2 + (z-d)^2)^{3/2}} + \frac{q(z+d)}{(x^2 + y^2 + (z+d)^2)^{3/2}} \right] \Big|_{z \rightarrow 0}$$

So

$$\sigma = -\frac{1}{2\pi} \frac{qd}{(x^2 + y^2 + d^2)^{3/2}}$$

We can check that the total induced charge in the plane is infact

$$Q = \int_0^{2\pi} \int_0^\infty \sigma r dr d\phi = -q$$

where $r^2 = x^2 + y^2$.

This charge comes from the reservoir of charge given by ground.

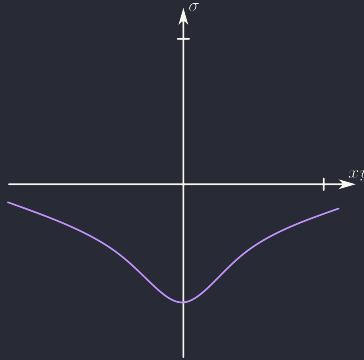


Figure 3.3: Just a cool graph of xy vs σ

3.2.3 Force and energy

Calcuating the force of attraction because of the negative induced charge:

$$F = qE = \frac{1}{4\pi\epsilon_0} \frac{q^2}{(2d)^2} (-\hat{\mathbf{z}})$$

Naively, we calculate the energy/work done as

$$W = -\frac{1}{4\pi\epsilon_0} \frac{q^2}{2d} \quad (= q\Delta V)$$

but we only have a single charge and the total work is half of this value

$$W = -\frac{1}{4\pi\epsilon_0} \frac{q^2}{4d}$$

This is because the ideal conductor requires no work to build up a charge distribution σ .

Thus we have two ODEs

3.3 Separation of Variables

3.3.1 In Cartesian

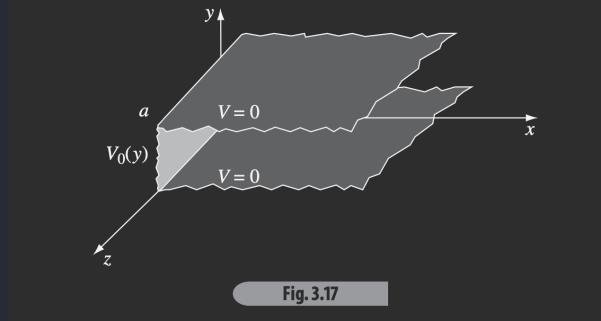


Figure 3.4: 3 planes where top and bottom are grounded and the middle is at $V_0(y)$

Find potential V in $x > 0, 0 < y < a, -\infty < z < \infty$:

This is a 2D problem $V \rightarrow V(x, y)$

$$\text{No change in } V \rightarrow \nabla^2 V = 0 = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2}.$$

From b.c.s

- (i) $V(y = 0) = 0$
- (ii) $V(y = a) = 0$
- (iii) $V(x = 0) = V_0(y)$
- (iv) $V(x \rightarrow \infty) = 0$

PROPOSE: $V(x, y) = X(x)Y(y)$

$$\begin{aligned} \nabla^2 V &= Y \frac{\partial^2 X}{\partial x^2} + X \frac{\partial^2 Y}{\partial y^2} = 0 \\ \frac{\nabla^2 V}{V} &= \frac{1}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = 0 \end{aligned}$$

thus we have two independent functions

$$f(x) + g(y) = 0$$

which is *only* possible if both f and g are constants

$$C_1 + C_2 = 0 \implies C_1 = -C_2 \quad \text{or} \quad C_1 = C_2 = 0$$

We now claim

$$\operatorname{sgn}(C_1) = 1, \quad \operatorname{sgn}(C_2) = -1, \quad \text{and} \quad |C_1| = |C_2| = k^2$$

$$\begin{aligned} \frac{\partial^2 X}{\partial x^2} &= k^2 X \\ \frac{\partial^2 Y}{\partial y^2} &= -k^2 Y \end{aligned}$$

with general solutions

$$X(x) = Ae^{kx} + Be^{-kx}, \quad Y(y) = C \sin(ky) + D \cos(ky)$$

so the original potential are the product of the two solutions

$$\begin{aligned} V(x, y) &= (Ae^{kx} + Be^{-kx})(C \cos(ky) + D \sin(ky)) \\ &= e^{-kx}(C \cos(ky) + D \sin(ky)) \end{aligned}$$

where condition (iv) requires $A = 0$ and condition (i) requires $D = 0$. From condition (ii) we require $\sin(ka) = 0$ thus

$$k = \frac{n\pi}{a}, \quad n = 1, 2, 3, \dots$$

Finally from condition (iii) we start with the rewritten potential

$$V(x, y) = Ce^{-kx} \sin(ky) \quad k = \frac{n\pi}{a}$$

We now have an infinite number of solutions

$$V_1, V_2, V_3, \dots \quad \text{if } V = \alpha_1 V_1 + \alpha_2 V_2 + \dots$$

where the general solution is

$$V(x, y) = \sum_{n=1}^{\infty} C_n e^{-n\pi x/a} \sin(n\pi y/a)$$

Now from (iii) we want $V(0, y) \Rightarrow V_0(y)$, so we multiply both sides by $\sin(n\pi y/a)$ and integrate over y

$$\sum_{n=1}^{\infty} C_n \int_0^a \sin(n\pi y/a) \sin(n'\pi y/a) dy = \int_0^a V_0(y) \sin(n'\pi y/a) dy$$

where the left integral is evaluated as

$$\int_0^a \sin(n\pi y/a) \sin(n'\pi y/a) dy = \begin{cases} 0 & n \neq n' \\ \frac{a}{2} & n = n' \end{cases}$$

This gives us the coefficient

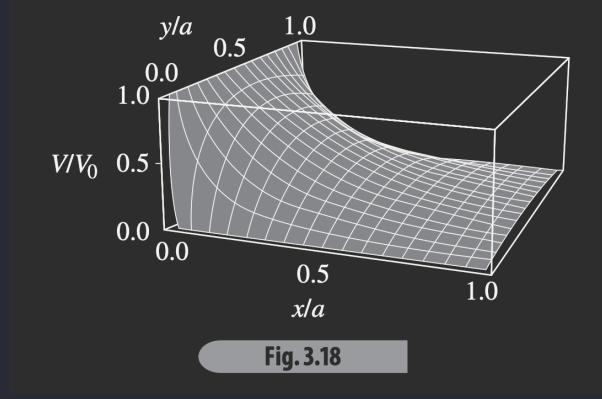
$$C_n = \frac{2}{a} \int_0^a V_0(y) \sin(n\pi y/a) dy$$

We let $V_0(y) = V_0 \neq 0$ so

$$C_n = \frac{2}{a} V_0 \int_0^a \sin(n\pi y/a) dy = \frac{2V_0}{n\pi} (1 - \cos(n\pi)) = \begin{cases} 0 & n \text{ even} \\ \frac{4V_0}{n\pi} & n \text{ odd} \end{cases}$$

so

$$\boxed{V(x, y) = \frac{4V_0}{\pi} \sum_{\text{odd } n} \frac{1}{n} e^{-n\pi x/a} \sin\left(\frac{n\pi y}{a}\right)}$$

Figure 3.5: Graph of $V(x, y)$

Furthermore, the infinite series can be simplified to

$$V(x, y) = \frac{2V_0}{\pi} \arctan \left(\frac{\sin(\frac{\pi y}{a})}{\sinh(\frac{\pi x}{a})} \right)$$

We can see this graphed in 3D space in Fig. 3.5.

3.3.2 In Spherical

The laplacian now becomes

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} = 0$$

For the subject matter of undergraduate E&M, we will assume azimuthal symmetry i.e.

$$V = V(r, \theta)$$

so that V is independent of ϕ :

$$\nabla^2 V = \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) = 0$$

So we look for solutions that are a product two independent functions

$$V(r, \theta) = R(r)\Theta(\theta)$$

thus dividing the laplacian by V we get

$$\underbrace{\frac{1}{R} \frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right)}_{C_1} + \underbrace{\frac{1}{\Theta \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Theta}{\partial \theta} \right)}_{C_2} = 0$$

where C_1 and C_2 are constants. We now choose

$$C_1 = \ell(\ell + 1), \quad C_2 = -\ell(\ell + 1)$$

so that the radial & polar equation becomes

$$R(r) = Ar^\ell + \frac{B}{r^{\ell+1}}$$

$$\Theta(\theta) = P_\ell(\cos \theta)$$

where P_ℓ is a Legendre polynomial defined by the Rodrigues formula (there are multiple ways to define this polynomial)

$$\begin{aligned} P_\ell(x) &= \frac{1}{2^\ell \ell!} \left(\frac{d}{dx} \right)^\ell (x^2 - 1)^\ell \\ P_0 &= 1 \\ P_1 &= x \\ P_2 &= \frac{1}{2}(3x^2 - 1) \dots \end{aligned}$$

The most general solution is a linear combination

$$V_{\text{gen}}(r, \theta) = \sum_{\ell=0}^{\infty} \left(A_\ell r^\ell + \frac{B_\ell}{r^{\ell+1}} \right) P_\ell(\cos \theta)$$

Example Uncharged conducting sphere sitting in a background E-field $\mathbf{E} = E_0 \hat{\mathbf{z}}$, so far away, the field points directly up. The induced charge on the sphere has positive charges at the top and negative charges at the bottom. The surface of the sphere is an equipotential of

$$V_{\text{sph}} = 0$$

and

$$V \rightarrow -E_0 r \cos \theta$$

as we go to ∞ i.e. $r \gg R$. So

$$\begin{aligned} V &= - \int \mathbf{E} \cdot d\ell \\ &= E_0 z + C = 0 \end{aligned}$$

The general solution is given by

$$V = \sum_{\ell=0}^{\infty} \left(A_\ell r^\ell + \frac{B_\ell}{r^{\ell+1}} \right) P_\ell(\cos \theta)$$

Using B.C (i) when $r = R$

$$\begin{aligned} \implies 0 &= A_\ell R^\ell + \frac{B_\ell}{R^{\ell+1}} \\ \implies B_\ell &= -A_\ell R^{2\ell+1} \end{aligned}$$

so

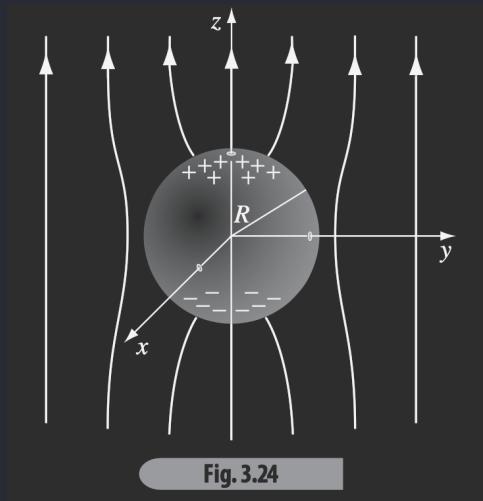
$$V(r, \theta) = \sum_{\ell=0}^{\infty} \left[A_\ell \left(r^\ell - \frac{R^{2\ell+1}}{r^{\ell+1}} \right) \right] P_\ell(\cos \theta)$$

For $r \gg R$ we can ignore $\frac{1}{r^{\ell+1}}$ and we have

$$\sum_{\ell=0}^{\infty} A_\ell r^\ell P_\ell(\cos \theta) = -E_0 r \cos \theta$$

\implies only $\ell = 1$ contributes to the sum on the left thus

$$A_1 r \cos \theta = -E_0 r \cos \theta \implies A_1 = -E_0$$

Figure 3.6: Graph of $V(r, \theta)$

Therefore the potential

$$V(r, \theta) = -E_0 \left(r - \frac{R^3}{r^2} \right) \cos \theta$$

We can see from Fig. 3.6 that the conductor acts as an attractive lens for the field. Furthermore, the charge distribution on the sphere is

$$\begin{aligned} \sigma &= -\epsilon_0 \frac{\partial V}{\partial n} = \epsilon_0 \frac{\partial V}{\partial r} \Big|_{r=R} \\ &= 3\epsilon_0 E_0 \cos \theta \end{aligned}$$

which is shown in Fig. 3.7.

Figure 3.7: Charge distribution $\sigma(\theta)$

3.4 Multipole Expansion

3.4.1 To approximate potential at great distances

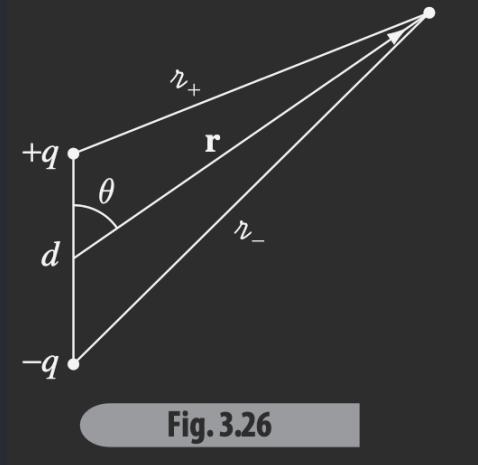


Figure 3.8: Dipole

For a dipole of opposite charge (Fig. 3.8) we are trying to find the potential $V(\mathbf{r})$ for $r \gg d$. Using superposition

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{r_+} - \frac{q}{r_-} \right)$$

And using Law of Cosines

$$\begin{aligned} r_{\pm}^2 &= r^2 + (d/2)^2 \mp rd \cos \theta \\ &= r^2 \left(1 \mp \frac{d}{r} \cos \theta + \frac{d^2}{4r^2} \right) \end{aligned}$$

So we solve for $\frac{1}{r_{\pm}}$ and use binomial expansion:

$$\frac{1}{r_{\pm}} = \frac{1}{r} \left(1 \mp \frac{d}{r} \cos \theta \right)^{-1/2} \approx \frac{1}{r} \left(1 \pm \frac{d}{2r} \cos \theta \right)$$

Thus

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{qd \cos \theta}{r^2}$$

We can abstract this to general multipoles such as the monopole, quadropole, octopole, hexadecapole, etc. (Fig. 3.9).

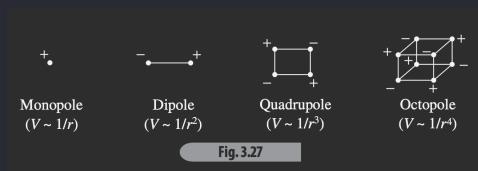


Figure 3.9: Multipole Expansion

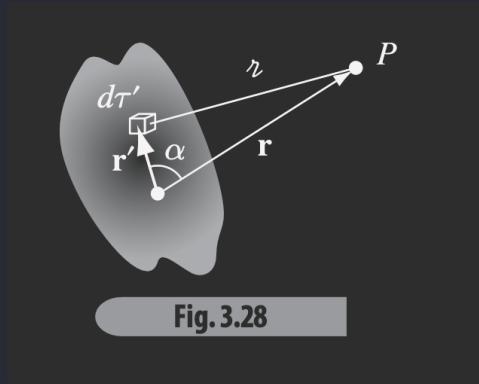


Figure 3.10: General local charge distribution

Expanding the potential to any local distribution of charge Thus

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{\mathbf{r}} d\tau'$$

using LoC again

$$\begin{aligned} \mathbf{r}^2 &= r^2 + r'^2 - 2rr' \cos \alpha \\ &= r^2 \left[1 + \left(\frac{r'}{r} \right)^2 - 2 \frac{r'}{r} \cos \alpha \right] \end{aligned}$$

where we let $\mathbf{r} = r\sqrt{1+\epsilon}$ so

$$\epsilon = \left(\frac{r'}{r} \right) \left(\frac{r'}{r} - 2 \cos \alpha \right)$$

Far away $\epsilon \ll 1$ so we can binomial expand

$$\begin{aligned} \frac{1}{\mathbf{r}} &= \frac{1}{r}(1+\epsilon)^{-1/2} \\ &= \frac{1}{r} \left(1 - \frac{\epsilon}{2} + 3/8\epsilon^2 - \dots \right) \\ &= \frac{1}{r} \left(1 - \frac{1}{2} \left(\frac{r'}{r} \right) \left(\frac{r'}{r} - 2 \cos \alpha \right) + \frac{3}{8} \left(\frac{r'}{r} \right)^2 \left(\frac{r'}{r} - 2 \cos \alpha \right)^2 - \frac{5}{16} \left(\frac{r'}{r} \right)^3 \left(\frac{r'}{r} - 2 \cos \alpha \right)^3 \right) \\ &= \frac{1}{r} \left[1 + \frac{r'}{r} \cos \alpha + \left(\frac{r'}{r} \right)^2 \left(\frac{3 \cos^2 \alpha - 1}{2} \right) + \left(\frac{r'}{r} \right)^3 \left(\frac{5 \cos^3 \alpha - 3 \cos \alpha}{2} \right) + \dots \right] \end{aligned}$$

all these terms are related to the Legendre polynomial $P_\ell(\cos \alpha)$ thus we can condense this to

$$\frac{1}{\mathbf{r}} = \frac{1}{r} \sum_{n=0}^{\infty} \left(\frac{r'}{r} \right)^n P_n(\cos \alpha)$$

We can substitute this back into the potential

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n(\cos \alpha) \rho(\mathbf{r}') d\tau'$$

where each term of the sum is a multipole expansion of V i.e. the monopole ($1/r$), dipole ($1/r^2$), etc.

3.4.2 Monopole & Dipole

Obviously the $1/r$ dominates for the monopole i.e.

$$V_{\text{mono}} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$

The dipole is the $n = 1$ term

$$V_{\text{dip}}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \int r' \cos \alpha \rho(\mathbf{r}') d\tau'$$

where $r' \cos \alpha = \hat{\mathbf{r}} \cdot \mathbf{r}'$ so

$$= \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}} \cdot \int \mathbf{r}' \rho(\mathbf{r}') d\tau'$$

And we define the integral as the dipole moment

$$\mathbf{p} = \int \mathbf{r}' \rho(\mathbf{r}') d\tau'$$

So the dipole potential is simplified to

$$V_{\text{dip}}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$$

For *discrete* charges:

$$\mathbf{p} = \sum_{i=1}^n q_i \mathbf{r}'_i$$

where for a dipole

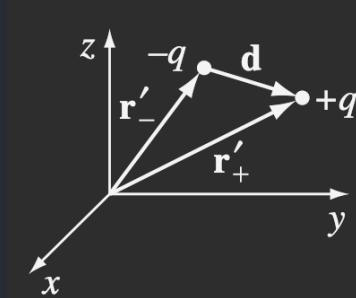


Fig. 3.29

Figure 3.11: Dipole

$$\begin{aligned} \mathbf{p} &= q\mathbf{r}'_+ - q\mathbf{r}'_- \\ &= q(\mathbf{r}'_+ - \mathbf{r}'_-) \\ &= q\mathbf{d} \end{aligned}$$

For a physical or a pure dipole where $\mathbf{d} \rightarrow 0$ while $q \rightarrow \infty$

$$V_{\text{dip}} = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{p \cos \theta}{r^2}$$

And looking at the E-field $\mathbf{E} = -\nabla V$; the components are

$$\begin{aligned} E_r &= -\frac{\partial V}{\partial r} = \frac{1}{4\pi\epsilon_0} \frac{2p \cos \theta}{r^3} \\ E_\theta &= -\frac{1}{r} \frac{\partial V}{\partial \theta} = \frac{1}{4\pi\epsilon_0} \frac{p \sin \theta}{r^3} \\ E_\phi &= 0 \end{aligned}$$

So

$$\mathbf{E}_{\text{dip}} = \frac{1}{4\pi\epsilon_0} \frac{p}{r^3} \left(2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta} \right)$$

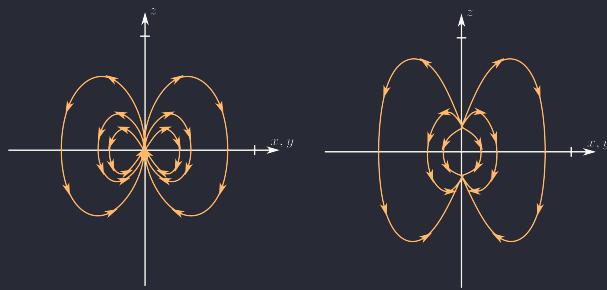


Figure 3.12: Pure dipole (left) and physical dipole (right)

4 Electric Fields in Matter

4.1 Polarization

4.1.1 Dielectrics

In dielectrics, “All charges are attached to specific atoms or molecules” (Griffith, pg.166)

4.1.2 Induced Dipoles

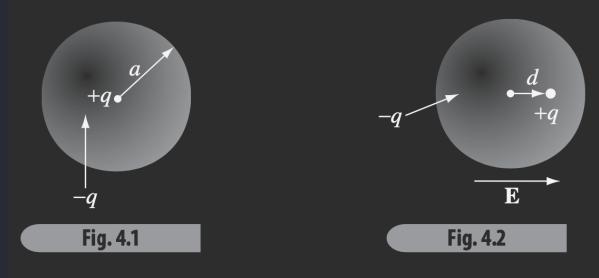


Figure 4.1: Left: Simple Nucleus $+q$ surrounded by spherical cloud $-q$ of radius a . Right: In an external electric field E the nucleus shifts right d and the cloud shifts left.

For a simple model of the atom (Fig. 4.1), the electric field at d is

$$E_d = \frac{1}{4\pi\epsilon_0} \frac{qd}{a^3}$$

where the dipole moment is $p = qd$. We have equilibrium when

$$F_{\text{ext}} = qE_{\text{ext}} = q_- E_d$$

So using the dipole moment

$$|\mathbf{p}| = qd = 4\pi\epsilon_0 a^3 E_{\text{ext}} = \alpha E_{\text{ext}}$$

Here we have this “atomic polarizability”

$$\alpha = 4\pi\epsilon_0 a^3 = 3\epsilon v$$

where v is the volume of the atom. Comment: this crude approximation is still accurate by a factor of 4.

In general we have a vector

$$\mathbf{p} = \hat{\alpha} \mathbf{E}$$

where $\hat{\alpha}$ is the polarizability tensor. For a linear dielectric relation between E and p we get

$$\hat{\alpha} = \begin{bmatrix} \alpha_{xx} & \alpha_{xy} & \alpha_{xz} \\ \alpha_{yx} & \alpha_{yy} & \alpha_{yz} \\ \alpha_{zx} & \alpha_{zy} & \alpha_{zz} \end{bmatrix}$$

4.1.3 Alignment of Polar Molecules

The dipole will experience a torque in and E -field

$$\begin{aligned} \mathbf{N} &= (\mathbf{r}_+ \times \mathbf{F}_+) + (\mathbf{r}_- \times \mathbf{F}_-) \\ &= \left(\frac{\mathbf{d}}{2} \times q\mathbf{E} \right) + \left(-\frac{\mathbf{d}}{2} \times q\mathbf{E} \right) \\ &= q\mathbf{d} \times \mathbf{E} \end{aligned}$$

thus

$$\mathbf{N} = \mathbf{p} \times \mathbf{E}$$

which implies that there is a force that acts to align $\mathbf{p} \parallel \mathbf{E}$.

WHat if E is not uniform?

$$\mathbf{F} = \mathbf{F}_+ + \mathbf{F}_- = q\mathbf{E} + q\mathbf{E} = q\Delta\mathbf{E}$$

assuming small d in E_x , then

$$\Delta E_x = (\nabla E_x) \cdot \mathbf{d}$$

So the total change in the field is

$$\Delta\mathbf{E} = (\mathbf{d} \cdot \nabla)\mathbf{E}$$

thus

$$\mathbf{F} = (\mathbf{p} \cdot \nabla)\mathbf{E}$$

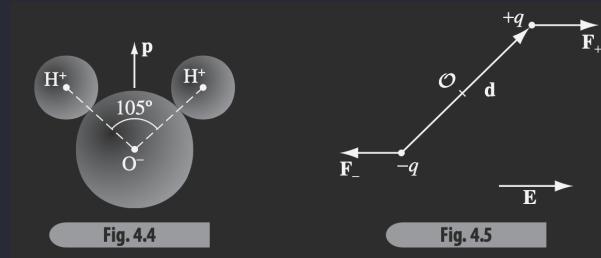


Figure 4.2: Dipole moment in oxygen molecule, and in an electric field.

Example: Problem 4.5 Using the method of images

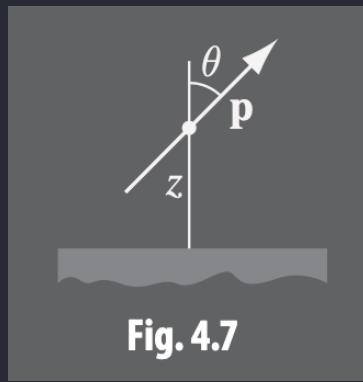
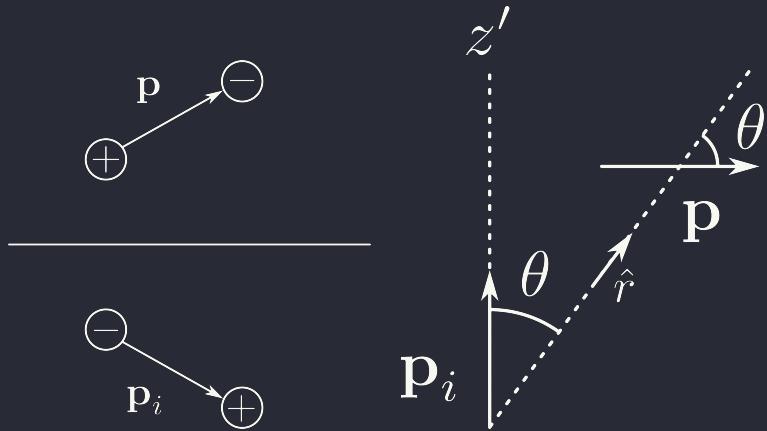


Figure 4.3: Infinitely grounded conductor with dipole at an angle θ from the normal plane and nailed in place.

Where look at \mathbf{p}_i coordinate (pointing up) thus $2z$ away we have the dipole pointing perpendicular to the image dipole i.e.

$$\mathbf{E}_i = \frac{1}{4\pi\epsilon_0} \frac{p_i}{r^3} \left(2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta} \right)$$

Figure 4.4: Left: Method of images using image dipole. Right: Coordinate using image dipole up as z' .

where

$$\mathbf{p} = p \cos \theta \hat{\mathbf{r}} + p \sin \theta \hat{\theta}$$

So the torque $\mathbf{N} = \mathbf{p} \times \mathbf{E}_i$ is

$$\begin{aligned} \mathbf{N} &= \frac{1}{4\pi\epsilon_0} \frac{p^2}{(2z)^3} (\cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \times (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \\ &= \frac{1}{4\pi\epsilon_0} \frac{p^2}{(2z)^3} (\cos \theta \sin \theta \hat{\phi} + 2 \cos \theta \sin \theta (-\hat{\phi})) \\ &= \frac{1}{4\pi\epsilon_0} \frac{p^2 \cos \theta \sin \theta}{(2z)^3} (-\hat{\phi}) \\ &= \frac{1}{4\pi\epsilon_0} \frac{p^2 \sin(2\theta)}{8\pi\epsilon_0 (8z^3)} (-\hat{\phi}) \end{aligned}$$

So

$$\begin{cases} 0 < \theta < \frac{\pi}{2} & N \sim -\hat{\phi} \\ \frac{\pi}{2} < \theta < \pi & N \sim \hat{\phi} \end{cases}$$

Which means the dipole can either align perpendicularly up or down depending on the angle θ .

4.1.4 Polarization

$$\mathbf{P} \equiv \text{dipole moment per unit volume}$$

i.e. the little \mathbf{p} is

$$\mathbf{p} = \mathbf{P} d\tau$$

4.2 The Field of a Polarized Object

4.2.1 Bound Charges

For a single dipole

$$V_{\text{dip}}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{z}}}{\mathbf{r}^2}$$

and using the dipole moment per unit volume def:

$$V_{\text{dip}}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{P}(\mathbf{r}') \cdot \hat{\mathbf{z}}}{\mathbf{r}^2} d\tau'$$

recalling the math fact

$$\nabla' \left(\frac{1}{\mathbf{r}} \right) = \frac{\hat{\mathbf{z}}}{\mathbf{r}^2}$$

thus

$$V = \frac{1}{4\pi\epsilon_0} \int \mathbf{P}(\mathbf{r}') \cdot \nabla' \left(\frac{1}{\mathbf{r}} \right) d\tau'$$

using another math fact

$$\nabla \cdot (F\mathbf{A}) = F(\nabla \cdot \mathbf{A}) + \mathbf{A} \cdot (\nabla F)$$

So we can rewrite the integral

$$\begin{aligned} V &= \frac{1}{4\pi\epsilon_0} \left[\int_V \nabla \cdot \left(\frac{\mathbf{P}}{\mathbf{r}} \right) d\tau' - \int_V \frac{1}{\mathbf{r}} \nabla \cdot \mathbf{P} d\tau' \right] \\ &= \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{1}{\mathbf{r}} \mathbf{P} \cdot d\mathbf{a}' - \int_V \frac{1}{\mathbf{r}} \nabla \cdot \mathbf{P} d\tau' \right] \end{aligned}$$

where we used the divergence theorem for the first term. For charge densities

$$\begin{cases} \text{surface charge} & \sigma_b \equiv \mathbf{P} \cdot \hat{\mathbf{n}} \\ \text{volume charge} & \rho_b \equiv -\nabla \cdot \mathbf{P} \end{cases}$$

then

$$V = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\sigma_b}{\mathbf{r}} d\mathbf{a}' + \int_V \frac{\rho_b}{\mathbf{r}} d\tau' \right]$$

4.2.2 Physical Interpretation of Bound Charges

So for a charge neutral sphere with an applied E -field, we can imagine this sphere as two oppositely charged spheres superimposed on each other but slightly shifted (Fig. 4.5). Thus we can imagine a collection of



Fig. 4.15

Figure 4.5: Charge neutral sphere with applied E -field.

dipoles for each atom in a material with alternating charges. This is actually wrong (read Berry Phases in Electronic Structure Theory by David Vanderbilt).

Example: Find \mathbf{E} of a uniformly polarized sphere of radius R .

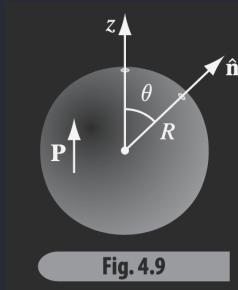


Figure 4.6: Uniformly polarized sphere of radius R .

We choose $\mathbf{P} \propto \hat{\mathbf{z}}$ as shown in Fig. 4.6. The bound volume and surface charges are

$$\rho_b = -\nabla \cdot \mathbf{P} = 0 \quad \text{but} \quad \sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = P \cos \theta$$

Thus using the result from before

$$V = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\sigma_b}{\epsilon_0} d\alpha' + \int_V \frac{\rho_b}{\epsilon_0} d\tau' \right]$$

or

$$V(r\theta) = \begin{cases} \frac{P}{3\epsilon_0} r \cos \theta & r \leq R \\ \frac{P}{3\epsilon_0} \frac{R^3}{r^2} \cos \theta & r > R \end{cases}$$

Recalling $z = r \cos \theta$ $\mathbf{E} = -\nabla V$ we get

$$\mathbf{E}_{\text{in}} = -\frac{P}{3\epsilon_0} \hat{\mathbf{z}} = -\frac{1}{3\epsilon_0} \mathbf{P}$$

and for outside the potential is

$$V_{\text{out}} = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{P} \cdot \hat{\mathbf{r}}}{r^2}$$

with the total dipole moment

$$\mathbf{p} = \frac{4\pi}{3} R^3 \mathbf{P}$$

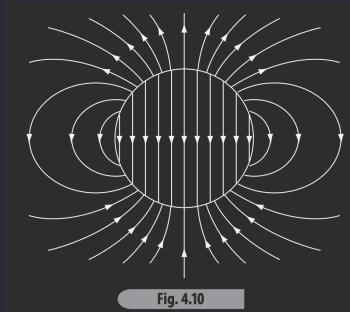


Figure 4.7: Polarized sphere in an external E -field.

So the polarized sphere is similar to the spherical conductor but with E -fields inside it (Fig. 4.7).

4.3 The Displacement Field

$$\rho_b = -\nabla \cdot \mathbf{P} \quad \text{and} \quad \sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}}$$

The $\rho_{\text{free}} \rightarrow$ anything *not* due to Polarization

$$\rho = \rho_b + \rho_f$$

and from Gauss' Law

$$\begin{aligned} \nabla \cdot \mathbf{E} &= \frac{\rho}{\epsilon_0} \\ \implies \epsilon_0 \nabla \cdot \mathbf{E} &= \rho_b + \rho_f = -\nabla \cdot \mathbf{P} + \rho_f \end{aligned}$$

where \mathbf{E} is the total electric field. Moving some terms around we get

$$\nabla \cdot (\epsilon_0 \mathbf{E} + \mathbf{P}) = \rho_f$$

where we now define the electric displacement field

$$\mathbf{D} \equiv \epsilon_0 \mathbf{E} + \mathbf{P}$$

which has the same laws as \mathbf{E} :

$$\boxed{\nabla \cdot \mathbf{D} = \rho_f} \quad \boxed{\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f, \text{enc}}}$$

Example: Long straight wire, with uniform λ (charge per length), surrounded by rubber insulation out to radius a ; Find \mathbf{D}

Using the Gaussian surface (cylinder) of length l and radius s so the enclosed charge is

$$\begin{aligned} \oint \mathbf{D} \cdot d\mathbf{a} &= 2\pi s l \mathbf{D} = \lambda l \\ \mathbf{D} &= \frac{\lambda}{2\pi s} \hat{\mathbf{s}} \end{aligned}$$

\implies holds for both $s \leq a$ and $s > a$.

For $s > a$

$$\mathbf{P}_{\text{out}} = 0 \implies \mathbf{E} = \frac{\mathbf{D}}{\epsilon_0} = \frac{\lambda}{2\pi\epsilon_0 s} \hat{\mathbf{s}}$$

But for inside $s \leq a$ we can't determine \mathbf{P}_{in} yet!

Comments While the two equations

$$\nabla \cdot \mathbf{D} = \rho_f \quad \text{and} \quad \nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

are similar, the E -field has a Coulomb law

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{z}}$$

but there is no equivalent for \mathbf{D}

$$\mathbf{D} = \cancel{\frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{z}}}$$

since there is this tensor relation $\mathbf{p} = \hat{\alpha} \mathbf{E}$. Furthermore, the curl is also different:

$$\nabla \times \mathbf{D} = \nabla \times (\epsilon_0 \mathbf{E} + \mathbf{P}) = \cancel{\epsilon_0} \nabla \times \mathbf{E} + \nabla \times \mathbf{P}$$

since $\nabla \times \mathbf{E} = \nabla \times (-\nabla V) = 0$.

What is D? Units $\frac{C}{m^2}$ which is the same as σ (surface charge density).

4.4 Linear Dielectrics

4.4.1 Susceptibility, Permittivity, and Dielectric constant

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}$$

where χ_e is the electric susceptibility (dimensionless). For here, we will assume linear (isotropic & homogeneous) dielectrics.

- vacuum: $\chi_e = 0$
- air: 1.00054
- salt: ~ 4.9
- Si: ~ 11
- water: ~ 80 (water is already polarized!)
- SrTiO_3 : $\sim 50,000$ at low temperatures

\mathbf{E} is the *total* electric field.

Starting with

$$\begin{aligned}\mathbf{D} &= \epsilon_0 \mathbf{E} + \mathbf{P} \\ &= \epsilon_0 \mathbf{E} + \epsilon_0 \chi_e \mathbf{E} \\ &= \epsilon_0 (1 + \chi_e) \mathbf{E} \quad \epsilon = \epsilon_0 (1 + \chi_e) \\ &= \epsilon \mathbf{E}\end{aligned}$$

where ϵ is the *permittivity* of the material and the relative permittivity or *dielectric constant* is

$$\epsilon_r \equiv \frac{\epsilon}{\epsilon_0} = 1 + \chi_e$$

Example: Metal sphere with charge Q , radius a , surrounded by a linear dielectric, ϵ , out to radius b . Find potential at the center (relative to ∞).

Inside the metal sphere $\mathbf{E} = \mathbf{P} = \mathbf{D} = 0$. Drawing the Gaussian surface between the sphere and the dielectric we get

$$\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f, \text{enc}}$$

for $r > a$

$$\mathbf{D} = \frac{Q}{4\pi r^2} \hat{\mathbf{r}}$$

Now finding \mathbf{E} from $\mathbf{D} = \epsilon \mathbf{E}$ i.e. in the vacuum $\epsilon = \epsilon_0$ and in the dielectric $\epsilon = \epsilon$:

$$\mathbf{E} = \begin{cases} \frac{Q}{4\pi \epsilon r^2} & a < r < b \\ \frac{Q}{4\pi \epsilon_0 r^2} & r > b \end{cases}$$

The potential at the origin is therefore

$$\begin{aligned}V &= - \int_{\infty}^0 \mathbf{E} \cdot d\ell = - \int_{\infty}^b \frac{Q}{4\pi \epsilon_0 r^2} dr - \int_b^a \frac{Q}{4\pi \epsilon r^2} dr - \int_a^0 0 dr \\ &= \frac{Q}{4\pi} \left(\frac{1}{\epsilon_0 b} + \frac{1}{\epsilon a} - \frac{1}{\epsilon b} \right)\end{aligned}$$

Now we can find \mathbf{P} since \mathbf{E} is fixed by \mathbf{D} :

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E} = \frac{\epsilon_0 \chi_e Q}{4\pi \epsilon r^2} \hat{\mathbf{r}} \quad (\text{in the dielectric})$$

Thus we can get the volume bound charge

$$\rho_b = -\nabla \cdot \mathbf{P} = 0 \quad \nabla \cdot \left(\frac{\hat{\mathbf{r}}}{r^2} \right) = 0 \quad \text{except at } r = 0$$

and

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = \begin{cases} \frac{\epsilon_0 \chi_e Q}{4\pi \epsilon b^2} & \text{outer} \\ \frac{-\epsilon_0 \chi_e Q}{4\pi \epsilon a^2} & \text{inner} \end{cases}$$

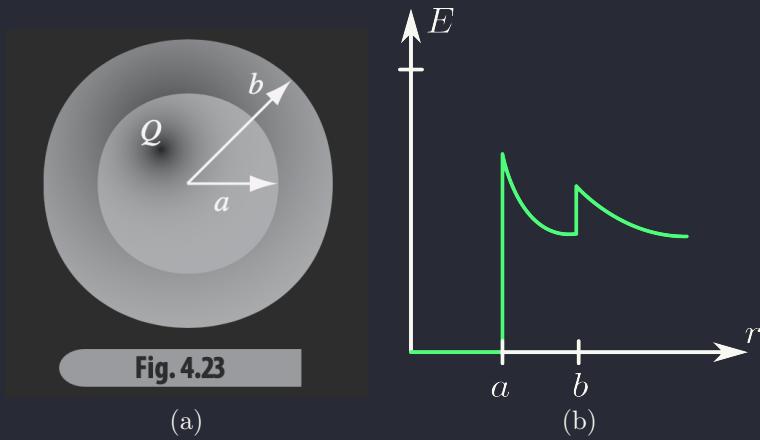


Figure 4.8: (a) Dielectric sphere surrounding metal sphere. (b) Electric field of resulting system.

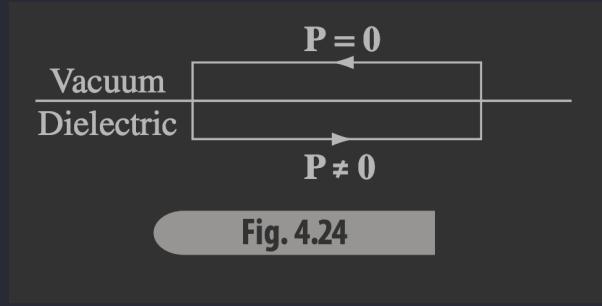


Figure 4.9: Interface between polarized dielectric and vacuum.

Since \mathbf{P} is zero on the vacuum side but not on the dielectric side

$$\oint \mathbf{P} \cdot d\ell \neq 0$$

In addition, the displacement field $\mathbf{D} = \epsilon \mathbf{E} + \mathbf{P}$ implies

$$\oint \mathbf{D} \cdot da \neq 0 \neq \nabla \times \mathbf{D}$$

Furthermore, the proportionality factor $\epsilon_0 \chi_e$ is different on both sides. For the homogeneous linear dielectric

$$\nabla \cdot \mathbf{D} = \rho_f \quad \text{and} \quad \nabla \times \mathbf{D} = 0$$

where

$$\begin{cases} \chi_f \mathbf{E}_{\text{vac}} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \\ \chi_c \mathbf{D} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \end{cases} \implies \mathbf{E} = \epsilon \mathbf{E}$$

or

$$\mathbf{E} = \frac{1}{\epsilon} \mathbf{D}$$

Example: Parallel plate capacitor with insulating material of dielectric constant ϵ_r

$$\mathbf{E} \rightarrow \frac{\mathbf{E}}{\epsilon_r}$$

and the potential difference between the plates is

$$V = - \int \mathbf{E} \cdot d\ell = -Ed \rightarrow \frac{Ed}{\epsilon_r}$$

And the capacitance is

$$C = \frac{Q}{V} \rightarrow \frac{\epsilon_r Q}{V} = \epsilon_r C$$

Boundary Value Problems with Linear Dielectrics

From the given properties of the displacement vector

$$\begin{aligned} \mathbf{D} &= \epsilon_0 \mathbf{E} + \mathbf{P} \\ &= \epsilon_0(1 + \chi_e) \mathbf{E} \end{aligned}$$

where $\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}$ so

$$\mathbf{D} = \epsilon \mathbf{E}$$

We can get the bound charge densities

$$\rho_b = -\nabla \cdot \mathbf{P} = -\nabla \cdot \left(\frac{\epsilon_0 \chi_e}{\epsilon} \mathbf{D} \right), \quad \sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}}$$

which implies

$$\boxed{\rho_b = -\left(\frac{\chi_e}{1 + \chi_e} \right) \rho_f}$$

Since the net charge resides in the surface of the dielectric, i.e. $\rho = 0$ so from Laplace's equation

$$\epsilon_{\text{above}} \mathbf{E}_{\text{above}}^\perp - \epsilon_{\text{below}} \mathbf{E}_{\text{below}}^\perp = \sigma_f$$

or in terms of voltage (from $\mathbf{E} = -\nabla V$)

$$\epsilon_{\text{above}} \frac{\partial V_{\text{above}}}{\partial n} - \epsilon_{\text{below}} \frac{\partial V_{\text{below}}}{\partial n} = -\sigma_f$$

And since the potential is still continuous

$$V_{\text{above}} = V_{\text{below}}$$

Example: Sphere of homogeneous linear dielectric material in a uniform \mathbf{E}_0 . Find \mathbf{E} in sphere.

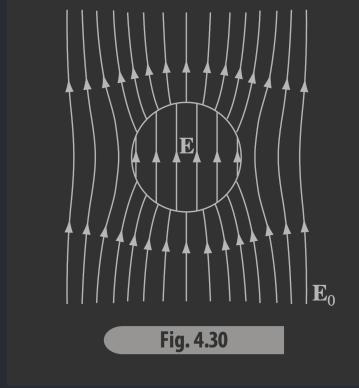


Figure 4.10: Homogeneous linear dielectric sphere in uniform E-field \mathbf{E}_0 .

Using B.C. to solve Laplace's equation

- (i) $V_{\text{in}} = V_{\text{out}}$ at $r = R$
- (ii) $\epsilon \frac{\partial V_{\text{in}}}{\partial r} = \epsilon_0 \frac{\partial V_{\text{out}}}{\partial r}$ at $r = R$ where $\epsilon = \epsilon_0 \epsilon_r$
- (iii) $V_{\text{out}} = -E_0 r \cos \theta$ at $r \gg R$

From the Legendre polynomials

$$V_{\text{in}} = \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta)$$

and outside the sphere

$$V_{\text{out}} = \underbrace{-E_0 r \cos \theta}_{\text{external field}} + \underbrace{\sum_{l=0}^{\infty} \frac{B_l}{r^{l+1}} P_l(\cos \theta)}_{\text{due to sphere}}$$

and at $r = R \rightarrow V_{\text{in}} = V_{\text{out}}$ so

$$\sum_l A_l R^l P_l(\cos \theta) = -E_0 R \cos \theta + \sum_l \frac{B_l}{R^{l+1}} P_l(\cos \theta)$$

Remembering that $P_0 = 1$ and $P_1 = \cos \theta$ we get from (i)

$$\begin{cases} A_1 R = -\epsilon_0 R + \frac{B_1}{R^2} & l = 1 \\ A_l R_l = \frac{B_l}{R^{l+1}} & l \neq 1 \end{cases}$$

and from (ii)

$$\epsilon_r \sum_l A_l R^{l+1} P_l(\cos \theta) = -\epsilon_0 \cos \theta - \sum_l \frac{(l+1) B_l}{R^{l+2}} P_l(\cos \theta)$$

so the orthogonality of the Legendre polynomials gives

$$\begin{cases} \epsilon_r A_1 = -\epsilon_0 - \frac{2B_1}{R^3} & l = 1 \\ \epsilon_r l A_l R^{l-1} = -\frac{(l+1) B_l}{R^{l+2}} & l \neq 1 \end{cases}$$

After staring at it for some time and looking through the math

$$\begin{cases} A_l = B_l = 0 & l \neq 1 \\ A_1 = -\frac{3}{\epsilon_r + 2} E_0 \quad \text{and} \quad B_1 = \frac{\epsilon_r - 1}{\epsilon_r + 2} R^3 E_0 & l = 1 \end{cases}$$

Thus the potentials are

$$V_{\text{in}}(r, \theta) = -\frac{3\epsilon_0}{\epsilon_r + 2} r \cos \theta = -\frac{3\epsilon_0}{\epsilon_r + 2} z$$

so the electric field is

$$\mathbf{E}_{\text{in}} = -\nabla V = \frac{3}{\epsilon_r + 2} \mathbf{E}_0$$

The field inside the sphere will go to zero if the dielectric constant is infinite.

Another example: Everything below plane $z = 0$ is linear dielectric with susceptibility χ_e . Calculating the force on a charge q a distance d above the origin.

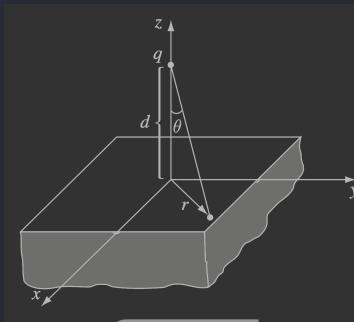


Fig. 4.31

Figure 4.11: Everything below $z = 0$ is a dielectric

Earlier we found a relationship in the volume of a dielectric sphere being $\rho_b \propto \rho_f = 0$, so we only need to worry about the bound surface charge density

$$\begin{aligned} \sigma_b &= \mathbf{P} \cdot \hat{\mathbf{n}} = P_z, \quad \mathbf{P} = \epsilon_0 \chi_e \mathbf{E} \\ &= P_z = \epsilon_0 \chi_e E_z \end{aligned}$$

where E_z is the z -component of the TOTAL field. The total field due to q and σ_b :

4.4.2 Energy in Dielectrics

The work it takes to charge a capacitor is

$$W = \frac{1}{2}CV^2$$

where the a capacitor filled with a linear dielectric has a capacitance

$$C = \epsilon_r C_0$$

Thus the energy stored in this system is

$$W_{\text{die}} = \frac{\epsilon_0}{\int} \epsilon_r E^2 d\tau$$

and using $\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \epsilon_r \mathbf{E}$ we get

$$\boxed{W = \frac{1}{2} \int \mathbf{D} \cdot \mathbf{E} d\tau}$$

which is integrated over all space.

Example Sphere of radius R filled with dielectric ϵ_r and free charge ρ_f has energy (from Gauss's Law)

$$\oint D \cdot d\mathbf{a} = Q_{\text{f, enc}}$$

$$\implies \mathbf{D}(\mathbf{r}) = \begin{cases} \frac{\rho_f}{3} \mathbf{r} & r < R \\ \frac{\rho_f}{3} \frac{R^3}{r^2} \mathbf{r} & r > R \end{cases}$$

Thus the correlated E-field

$$\mathbf{E}(\mathbf{r}) = \begin{cases} \frac{\rho_f}{3\epsilon_0 \epsilon_r} \mathbf{r} & r < R \\ \frac{\rho_f}{3\epsilon_0 \epsilon_r} \frac{R^3}{r^2} \mathbf{r} & r > R \end{cases}$$

The purely electrostatic energy is $W_{\text{es}} = \frac{\epsilon_0}{2} \int E^2 d\tau$ or

$$W_{\text{es}} = \frac{\epsilon_0}{2} \left[\left(\frac{\rho_f}{3\epsilon_0 \epsilon_r} \right)^2 \int_0^R r^2 (4\pi r^2) dr + \left(\frac{\rho_f}{3\epsilon_0} \right)^2 R^6 \int_R^\infty \frac{1}{r^4} (4\pi r^2) dr \right]$$

$$= \frac{2\pi}{9\epsilon_0} \rho_f^2 R^5 \left(\frac{1}{5\epsilon_r^2} + 1 \right)$$

but the total energy is

$$W_{\text{tot}} = \frac{1}{2} \int \mathbf{D} \cdot \mathbf{E} d\tau$$

$$= \frac{1}{2} \left[\frac{\rho_f^2}{9\epsilon_0 \epsilon_r} \int_0^R r^2 (4\pi r^2) dr + \left(\frac{\rho_f}{3} R^3 \right) \left(\frac{\rho_f}{3\epsilon_0} R^3 \right) \int_R^\infty \frac{1}{r^4} (4\pi r^2) dr \right]$$

$$= \frac{2\pi}{9\epsilon_0} \rho_f^2 R^5 \left(\frac{1}{5\epsilon_r} + 1 \right)$$

Thus

$$W_{\text{es}} < W_{\text{tot}}$$

So starting with the unpolarized dielectric sphere and using the “method of assembly” i.e. adding charge dq to each layer of the sphere, we have the field in three regions

$$\mathbf{E}(\mathbf{r}) = \begin{cases} \frac{\rho_f}{3\epsilon_0\epsilon_r} \mathbf{r} & r < r' \\ \frac{\rho_f}{3\epsilon_0\epsilon_r} \frac{R^3}{r^2} \mathbf{r} & r' < r < R \\ \frac{\rho_f}{3\epsilon_0} \frac{R^3}{r^2} \mathbf{r} & r > R \end{cases}$$

And bringing down dq from $\infty \rightarrow r'$ the infinitesimal work is

$$\begin{aligned} dW &= -dq \left[\int_{\infty}^R \mathbf{E} \cdot d\mathbf{l} + \int_R^{r'} \mathbf{E} \cdot d\mathbf{l} \right] \\ &= -dq \left[\frac{\rho_f r'^3}{3\epsilon_0} \int_{\infty}^R \frac{1}{r^2} dr + \frac{\rho_f r'^3}{3\epsilon_0\epsilon_r} \int_R^{r'} \frac{1}{r^2} dr \right] \\ &= \frac{\rho_f r'^3}{3\epsilon_0} \left(\frac{1}{R} + \frac{1}{\epsilon_r} \left(\frac{1}{r'} - \frac{1}{R} \right) \right) dq \end{aligned}$$

which increases the radius by

$$dq = \rho dV = \rho_f (4\pi r'^2) dr'$$

thus

$$dW = \frac{\rho_f}{3\epsilon_0} \left[\frac{r'^3}{R} \left(1 - \frac{1}{\epsilon_r} \right) + \frac{r'^2}{\epsilon_r} \right] dq$$

So the total work is done by integrating over $0 \rightarrow R$

$$\begin{aligned} W &= \int dW = \int () dq \rightarrow \int () dr' \\ &= \frac{4\pi\rho_f^2}{3\epsilon_0} \left[\frac{1}{R} \left(1 - \frac{1}{\epsilon_r} \right) \int_0^R r'^5 dr' + \frac{1}{\epsilon_r} \int_0^R r'^4 dr' \right] \\ &= \frac{2\pi}{9\epsilon_0} \rho_f^3 R^5 \left[\frac{1}{5\epsilon_r} + 1 \right] = W_{\text{tot}} \end{aligned}$$

So the energy of deformation is

$$W_{\text{def}} = W_{\text{tot}} - W_{\text{es}} = \frac{2\pi}{45\epsilon_0\epsilon_r^2} \rho_f^2 R^5 (\epsilon_r - 1)$$

Another way... We can pretend that this deformation energy is analogous to a spring (separated by distance d and spring constant k) connecting two charges where one charge $+q$ is nailed down so

$$qE = kd$$

here

$$\mathbf{E} = \frac{\rho_f}{3\epsilon_0\epsilon_r} \mathbf{r} \implies \text{field of free charge, inside dielectric}$$

thus the dipole $p = qd$ and polarization $P = p/d\tau$. The spring constant and infinitesimal change in work is

$$\begin{aligned} k &= \frac{\rho_f}{3\epsilon_0\epsilon_r d^2} Pr d\tau \\ dW &= \frac{1}{2} kd^2 = \frac{\rho_f}{6\epsilon_0\epsilon_r} Pr d\tau \end{aligned}$$

This will give us the energy stored in the dipole “spring” i.e.

$$W_{\text{sp}} = \frac{\rho_f}{6\epsilon_0\epsilon_r} \int Pr \, d\tau$$

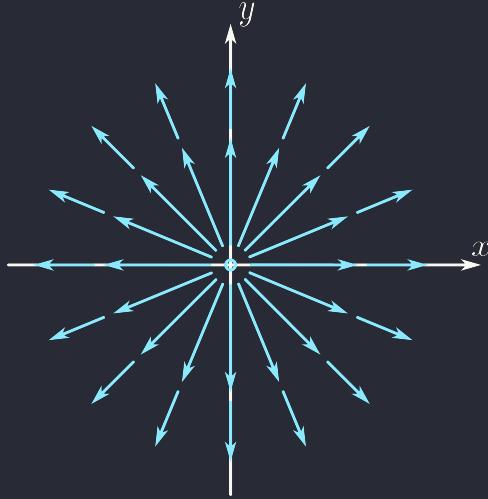
and the polarization is

$$\mathbf{P} = \epsilon_0(1 + \chi_e)\mathbf{E} = \epsilon_0\chi_e \frac{\rho_f}{3\epsilon_0\epsilon_r} \mathbf{r} = \frac{(\epsilon_r - 1)\rho_f}{3\epsilon_r} \mathbf{r}$$

So

$$\begin{aligned} W_{\text{sp}} &= \frac{\rho_f}{6\epsilon_0\epsilon_r} \int \frac{(\epsilon_r - 1)\rho_f}{3\epsilon_r} (4\pi) \int_0^R r^4 \, dr \\ &= \frac{2\pi}{45\epsilon_0\epsilon_r^2} \rho_f^2 R^5 (\epsilon_r - 1) \end{aligned}$$

which is the exact same thing! This is because our approximation for the polarization $\mathbf{P} = \epsilon_0\chi_e\mathbf{E}$ is linear i.e. we were approximating for a spring the whole time!

Figure 1.16: The vector field $\mathbf{v} = \frac{\hat{\mathbf{r}}}{r^2}$

1.16 Given

$$r = \sqrt{x^2 + y^2 + z^2} \quad \text{and} \quad \hat{\mathbf{r}} = \frac{\mathbf{r}}{r} = \frac{x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}}{\sqrt{x^2 + y^2 + z^2}}$$

where $\mathbf{r} = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$ is the position vector. The vector functions is

$$v = \frac{\hat{\mathbf{r}}}{r^2} = \frac{\mathbf{r}}{r^3}$$

where the components are

$$v_x = \frac{x}{r^3}, \quad v_y = \frac{y}{r^3}, \quad \text{and} \quad v_z = \frac{z}{r^3}$$

Looking at the x component of the divergence,

$$\begin{aligned} [\nabla \cdot \mathbf{v}]_x &= \frac{\partial v_x}{\partial x} \\ &= \frac{\partial}{\partial x} \left(\frac{x}{r^3} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{x}{(x^2 + y^2 + z^2)^{3/2}} \right) \quad \text{using chain rule...} \\ &= \frac{1}{(x^2 + y^2 + z^2)^{3/2}} - \frac{3x^2}{(x^2 + y^2 + z^2)^{5/2}} \\ &= \frac{1}{r^3} - \frac{3x^2}{r^5} \end{aligned}$$

therefore, the divergence of \mathbf{v} is

$$\begin{aligned} \nabla \cdot \mathbf{v} &= \left(\frac{1}{r^3} - \frac{3x^2}{r^5} \right) + \left(\frac{1}{r^3} - \frac{3y^2}{r^5} \right) + \left(\frac{1}{r^3} - \frac{3z^2}{r^5} \right) \\ &= \frac{3}{r^3} - 3 \frac{x^2 + y^2 + z^2}{r^5} \\ &= \frac{3}{r^3} - 3 \frac{r^2}{r^5} = 0 \end{aligned}$$

The divergence is zero everywhere except at the origin where $r = 0$ because division by r^3 tells us that the divergence is infinite at the origin.

(ii) $(0, 0, 1) \rightarrow (0, 1, 1)$:

$$z = 1, \quad x = dx = dz = 0; \quad d\mathbf{l} = dy \hat{\mathbf{y}}; \quad \nabla T \cdot d\mathbf{l} = 2 dy$$

and

$$\int_c^d \nabla T \cdot d\mathbf{l} = \int_0^1 2 dy = 2$$

(iii) $(0, 1, 1) \rightarrow b$:

$$z : 0 \rightarrow 1; \quad y = z = 1, \quad dy = dz = 0; \quad d\mathbf{l} = dx \hat{\mathbf{x}}; \quad \nabla T \cdot d\mathbf{l} = (2x + 4) dx$$

and

$$\int_d^b \nabla T \cdot d\mathbf{l} = \int_0^1 (2x + 4) dx = 5$$

therefore

$$\int_a^b \nabla T \cdot d\mathbf{l} = 0 + 2 + 5 = 7$$

(c) the parabolic path $z = x^2$; $y = x$:

$$dx = dy, \quad \text{and} \quad dz = 2x dx; \quad d\mathbf{l} = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}$$

and

$$\begin{aligned} \nabla T \cdot d\mathbf{l} &= (2x + 4y) dx + (4x + 2z^3) dx + (6yz^2) 2x dx \\ &= 6x dx + (4x + 2x^6) dx + (12x^6) dx \\ &= 10x dx + 14x^6 dx \end{aligned}$$

therefore

$$\begin{aligned} \int_a^b \nabla T \cdot d\mathbf{l} &= \int_0^1 (10x + 14x^6) dx \\ &= 5x^2 + 2x^7 \Big|_0^1 = 7 \end{aligned}$$

1.33 Testing the divergence theorem: For the function

$$\mathbf{v} = (xy)\hat{\mathbf{x}} + (2yz)\hat{\mathbf{y}} + (3zx)\hat{\mathbf{z}}$$

the divergence is

$$\nabla \cdot \mathbf{v} = y + 2z + 3x$$

so the volume integral is

$$\begin{aligned}\int_V \nabla \cdot \mathbf{v} d\tau &= \int_0^2 \int_0^2 \int_0^2 (y + 2z + 3x) dx dy dz \\ &= \int_0^2 \int_0^2 (2y + 4z + 6) dy dz \\ &= \int_0^2 (4 + 8z + 12) dz \\ &= 8 + 16 + 24 \\ \int_V \nabla \cdot \mathbf{v} d\tau &= 48\end{aligned}$$

The surface integral is evaluated over the six faces of the cube:

(i) $x = 2$, $d\mathbf{A} = dy dz \hat{\mathbf{x}}$, $\mathbf{v} \cdot d\mathbf{A} = 2y dy dz$;

$$\int \mathbf{v} \cdot d\mathbf{A} = \int_0^2 \int_0^2 2y dy dz = 8$$

(ii) $x = 0$, $d\mathbf{A} = -dy dz \hat{\mathbf{x}}$, $\mathbf{v} \cdot d\mathbf{A} = 0$;

$$\int \mathbf{v} \cdot d\mathbf{A} = \int_0^2 \int_0^2 0 dy dz = 0$$

(iii) $y = 2$, $d\mathbf{A} = dx dz \hat{\mathbf{y}}$, $\mathbf{v} \cdot d\mathbf{A} = 4z dx dz$;

$$\int \mathbf{v} \cdot d\mathbf{A} = \int_0^2 \int_0^2 4z dx dz = 16$$

(iv) $y = 0$;

$$\int \mathbf{v} \cdot d\mathbf{A} = 0$$

(v) $z = 2$, $d\mathbf{A} = dx dy \hat{\mathbf{z}}$, $\mathbf{v} \cdot d\mathbf{A} = 6x dx dy$;

$$\int \mathbf{v} \cdot d\mathbf{A} = \int_0^2 \int_0^2 6x dx dy = 24$$

(vi) $z = 0$;

$$\int \mathbf{v} \cdot d\mathbf{A} = 0$$

So the total flux is

$$\oint_S \mathbf{v} \cdot d\mathbf{A} = 8 + 0 + 16 + 0 + 24 + 0 = 48$$

therefore, the divergence theorem is verified.

$$\int_V (\nabla \cdot \mathbf{v}) d\tau = \oint_S \mathbf{v} \cdot d\mathbf{A}$$

1.34 Testing Stokes' theorem for the function

$$\mathbf{v} = (xy)\hat{\mathbf{x}} + (2yz)\hat{\mathbf{y}} + (3zx)\hat{\mathbf{z}}$$

using the triangular shaded area bounded by the vertices $O = (0, 0, 0)$, $A = (0, 2, 0)$, and $B = (0, 0, 2)$:

$$\begin{aligned}\nabla \times \mathbf{v} &= (0 - 2y)\hat{\mathbf{x}} - (3z - 0)\hat{\mathbf{y}} + (0 - x)\hat{\mathbf{z}} \quad \text{and} \quad d\mathbf{A} = dz dy \hat{\mathbf{x}} \\ &= -2y\hat{\mathbf{x}} - 3z\hat{\mathbf{y}} - x\hat{\mathbf{z}}\end{aligned}$$

$x = 0$ on this surface, and the limits of integration are $y : 0 \rightarrow 2$ and $z = 0 \rightarrow z = 2 - y$:

$$(\nabla \times \mathbf{v}) \cdot d\mathbf{A} = -2y dz dy$$

Thus, the flux of the curl through the surface is

$$\begin{aligned}\int_S (\nabla \times \mathbf{v}) \cdot d\mathbf{A} &= \int_0^2 \int_0^{2-y} -2y dz dy \\ &= \int_0^2 -2y(2-y) dy \\ &= -2y^2 + \frac{2}{3}y^3 \Big|_0^2 = -8/3\end{aligned}$$

The line integral is evaluated over the three sides of the triangle:

(i) On the path OA :

$$x = z = 0; d\mathbf{l} = dy \hat{\mathbf{y}}; \mathbf{v} \cdot d\mathbf{l} = 2yz dy = 0;$$

$$\int_{OA} \mathbf{v} \cdot d\mathbf{l} = 0$$

(ii) On the path AB :

$$x = 0, y = 2 - z; dy = -dz; d\mathbf{l} = -dz(\hat{\mathbf{y}} - \hat{\mathbf{z}}); \mathbf{v} \cdot d\mathbf{l} = -2yz dz = -2(2-z)z dz = (2z^2 - 4z) dz;$$

$$\int_{AB} \mathbf{v} \cdot d\mathbf{l} = \int_0^2 (2z^2 - 4z) dz = -8/3$$

(iii) On the path BO :

$$x = y = 0; d\mathbf{l} = dz \hat{\mathbf{z}}; \mathbf{v} \cdot d\mathbf{l} = 0;$$

$$\int_{BO} \mathbf{v} \cdot d\mathbf{l} = 0$$

So, the circulation of \mathbf{v} around the triangle is

$$\oint \mathbf{v} \cdot d\mathbf{l} = 0 + -8/3 + 0 = -8/3$$

thus,

$$\int_S (\nabla \times \mathbf{v}) \cdot d\mathbf{A} = \oint \mathbf{v} \cdot d\mathbf{l}$$

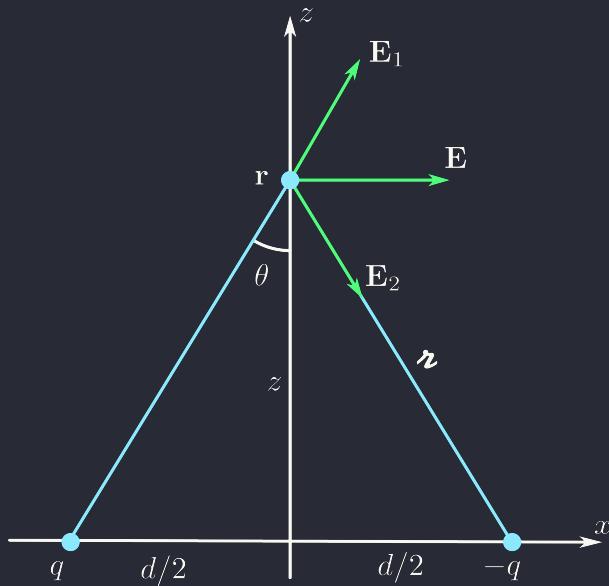


Figure 2.13: An electric field at a distance z from the midpoint between equal and opposite charges ($\pm q$) separated by a distance d . The charge at $x = d/2$ is $-q$.

2.2 The vertical components of the electric field cancel out and the horizontal components add up:

$$E_x = 2 \frac{1}{4\pi\epsilon_0} \frac{q}{\mathbf{z}^2} \sin \theta$$

where $E_x = E \cos \theta$, $\mathbf{z} = \sqrt{z^2 + (d/2)^2}$, and $\sin \theta = d/(2\mathbf{z})$, so

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{qd}{[z^2 + (d/2)^2]^{3/2}} \hat{\mathbf{x}}$$

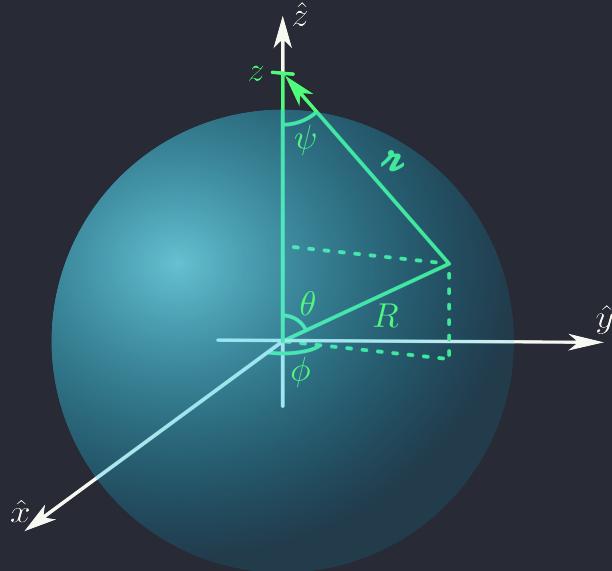


Figure 2.14: An electric field a distance z from the center of a spherical surface of radius R that carries a charge density σ .

2.7 Once again, the electric field is in the z -direction:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma}{z^2} \cos \psi \hat{\mathbf{z}} \, d\mathbf{a} \quad (2.1)$$

From the law of cosines, $z^2 = z^2 + R^2 - 2zR \cos \theta$; Geometrically, $\cos \psi = \frac{z - R \cos \theta}{z}$; the surface area element is $d\mathbf{a} = R^2 \sin \theta \, d\theta \, d\theta$:

$$\begin{aligned} \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^\pi \frac{\sigma R^2(z - R \cos \theta)}{(z^2 + R^2 - 2zR \cos \theta)^{3/2}} \sin \theta \, d\theta \, d\phi \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} (2\pi\sigma R^2) \int_0^\pi \frac{z - R \cos \theta}{(z^2 + R^2 - 2zR \cos \theta)^{3/2}} \sin \theta \, d\theta \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} (2\pi\sigma R^2) f(\theta) \hat{\mathbf{z}} \end{aligned}$$

using the substitution $u = \cos \theta$: $du = -\sin \theta \, d\theta$, and the limits of integration are $[\cos 0, \cos \pi]$. So,

$$f(\theta) = \int_{-1}^1 \frac{z - Ru}{(z^2 + R^2 - 2zRu)^{3/2}} \, du = f(u)$$

substituting again with $v = \sqrt{z^2 + R^2 - 2zRu}$; $dv = -\frac{zR}{v} \, du$; and $u = \frac{1}{2zR}(z^2 + R^2 - v^2)$:

$$\begin{aligned} f(v) &= -\frac{1}{zR} \int \frac{z - \frac{1}{2z}(z^2 + R^2 - v^2)}{v^3} v \, dv \\ &= -\frac{1}{2z^2 R} \int \frac{2z^2 - (z^2 + R^2 - v^2)}{v^2} \, dv \\ &= -\frac{1}{2z^2 R} \int \frac{v^2 + z^2 - R^2}{v^2} \, dv \\ &= -\frac{1}{2z^2 R} \int \left(1 + \frac{z^2 - R^2}{v^2}\right) \, dv \\ &= -\frac{1}{2z^2 R} \left(v - \frac{z^2 - R^2}{v}\right) \end{aligned}$$

back substituting $v = \sqrt{z^2 + R^2 - 2zRu}$,

$$\begin{aligned} f(u) &= -\frac{1}{2z^2R} \left(\frac{z^2 + R^2 - 2zRu}{\sqrt{z^2 + R^2 - 2zRu}} - \frac{z^2 - R^2}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\ &= -\frac{1}{2z^2R} \left(\frac{2R^2 - 2zRu}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\ &= \frac{1}{z^2} \left(\frac{zu - R}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\ &= \frac{1}{z^2} \left(\frac{z - R}{\sqrt{z^2 + R^2 - 2zR}} - \frac{-z - R}{\sqrt{z^2 + R^2 + 2zR}} \right) \end{aligned}$$

Taking the positive square root: $\sqrt{z^2 + R^2 - 2zR} = (R - z)$ if $R > z$, but $(z - R)$ if $R < z$. So, for the case $z < R$ (inside the sphere) the electric field is

$$\begin{aligned} \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma R^2}{z^2} \left(\frac{z - R}{R - z} - \frac{-z - R}{R + z} \right) \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma R^2}{z^2} \left(\frac{z - R}{R - z} + \frac{z + R}{R + z} \right) \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma R^2}{z^2} \left(\frac{z - R}{R - z} + 1 \right) \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma R^2}{z^2} \left(\frac{z - R}{R - z} + \frac{R - z}{R - z} \right) \hat{\mathbf{z}} \\ &= 0 \end{aligned}$$

For the case $z > R$ (outside the sphere) the electric field is

$$\begin{aligned} \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma R^2}{z^2} \left(\frac{z - R}{z - R} + \frac{z + R}{z + R} \right) \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{4\pi\sigma R^2}{z^2} \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \hat{\mathbf{z}} \end{aligned}$$

This makes sense: From outside the sphere, the point charge q is the charge-per-area σ times the surface area of the sphere $4\pi R^2$, or simply $q = 4\pi R^2 \sigma$.

2.8 Finding the field inside and outside a solid sphere of radius R with a uniform volume charge density ρ is similar to Prob. 2.7. Outside the solid sphere the total charge q contributes to the electric field as if it were a point charge:

$$\mathbf{E}_{out} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$$

Inside the solid sphere, only the volume of the solid sphere less than r contributes to the electric field. The volume of the total sphere is $V = \frac{4}{3}\pi R^3$, and the volume of the sphere less than r is $V' = \frac{4}{3}\pi r^3$. So, electric field inside the solid sphere is

$$\begin{aligned} \mathbf{E}_{in} &= \frac{V'}{V} \mathbf{E}_{out} \\ &= \frac{r^3}{R^3} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}} \end{aligned}$$

or

$$\mathbf{E}_{in} = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} \mathbf{r}$$

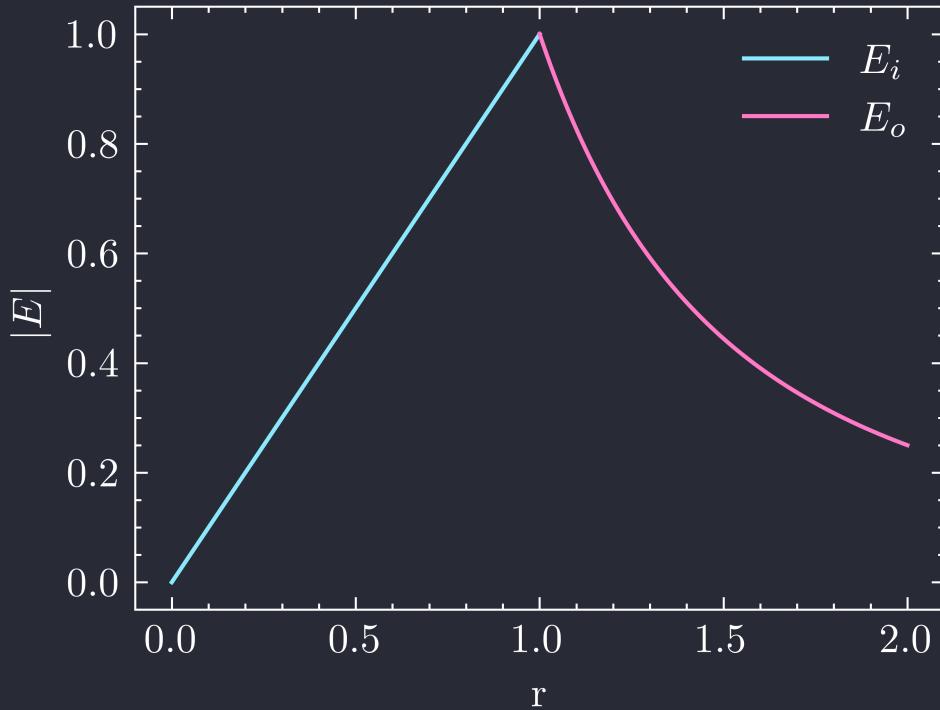


Figure 2.15: Magnitude of Electric field $|E|$ as a function of r inside and outside a solid. Where $q = 9\text{nC}$ and $R = 1\text{m}$.

2.16 A thick spherical shell with charge density

$$\rho = \frac{k}{r^2} \quad (a \leq r \leq b)$$

The electric field in the three regions:

(i) $r < a$

$$Q_{enc} = 0; \mathbf{E} = 0$$

(ii) $a \leq r \leq b$

$$Q_{enc} = \int_0^{2\pi} \int_0^\pi \int_a^r \rho(r^2 \sin \theta) dr d\theta d\phi = 4\pi \int_a^r \frac{k}{r^2} (r^2) dr = 4\pi k(r - a)$$

And from Gauss's law,

$$\begin{aligned} \oint \mathbf{E} \cdot d\mathbf{a} &= \frac{1}{\epsilon_o} Q_{enc} \\ |\mathbf{E}| \int da &= \frac{1}{\epsilon_o} 4\pi k(r - a) \\ E(4\pi r^2) &= \frac{1}{\epsilon_o} 4\pi k(r - a) \end{aligned}$$

or

$$\mathbf{E} = \frac{k(r - a)}{\epsilon_o r^2} \hat{\mathbf{r}} = \frac{1}{4\pi\epsilon_0} \frac{4\pi k(r - a)}{r^3} \mathbf{r}$$

(iii) $r > b$

$$Q_{enc} = \int_0^{2\pi} \int_0^\pi \int_a^b \rho(r^2 \sin \theta) dr d\theta d\phi = 4\pi k(b - a)$$

And from Gauss's law,

$$\begin{aligned} \oint \mathbf{E} \cdot d\mathbf{a} &= \frac{1}{\epsilon_o} Q_{enc} \\ |\mathbf{E}| \int da &= \frac{1}{\epsilon_o} 4\pi k(b - a) \\ E(4\pi r^2) &= \frac{1}{\epsilon_o} 4\pi k(b - a) \end{aligned}$$

or

$$\mathbf{E} = \frac{k(b - a)}{\epsilon_o r^2} \hat{\mathbf{r}} = \frac{1}{4\pi\epsilon_0} \frac{4\pi k(b - a)}{r^3} \mathbf{r}$$

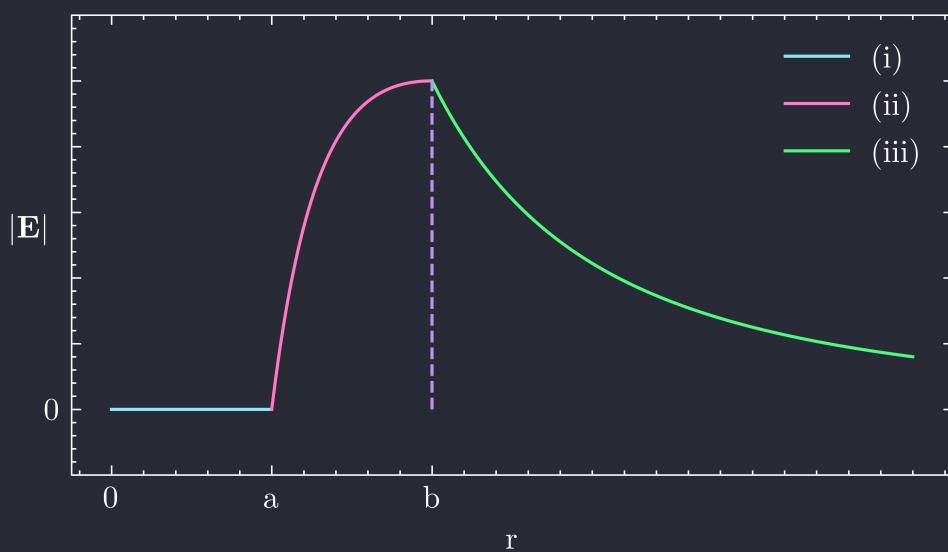


Figure 2.16: Plot of $|E|$ as a function of r , for the case $b = 2a$.

2.18 Finding the electric field, as a function of y , where $y = 0$ is the center of an infinite plane slab, of thickness $2d$, carrying a uniform volume charge density ρ . For the case $y > 2d$ The enclosed charge is

$$Q_{enc} = \rho(2d)A = 2\rho Ad$$

where A is the area of the Gaussian pillbox. Using Gauss's law,

$$\begin{aligned}\oint \mathbf{E} \cdot d\mathbf{a} &= \frac{1}{\epsilon_0} Q_{enc} \\ |\mathbf{E}| \int da &= \frac{1}{\epsilon_0} 2\rho Ad \\ E(2A) &= \frac{1}{\epsilon_0} 2\rho Ad \\ \mathbf{E} &= \frac{\rho d}{\epsilon_0} \hat{\mathbf{y}}\end{aligned}$$

For the case $0 < y < 2d$, the enclosed charge is

$$Q_{enc} = 2\rho y A$$

and the electric field is

$$\begin{aligned}E(2A) &= \frac{1}{\epsilon_0} \rho y A \\ \mathbf{E} &= \frac{\rho y}{\epsilon_0} \hat{\mathbf{y}}\end{aligned}$$

In the $-y$ direction, E is negative as shown in Figure 2.17.

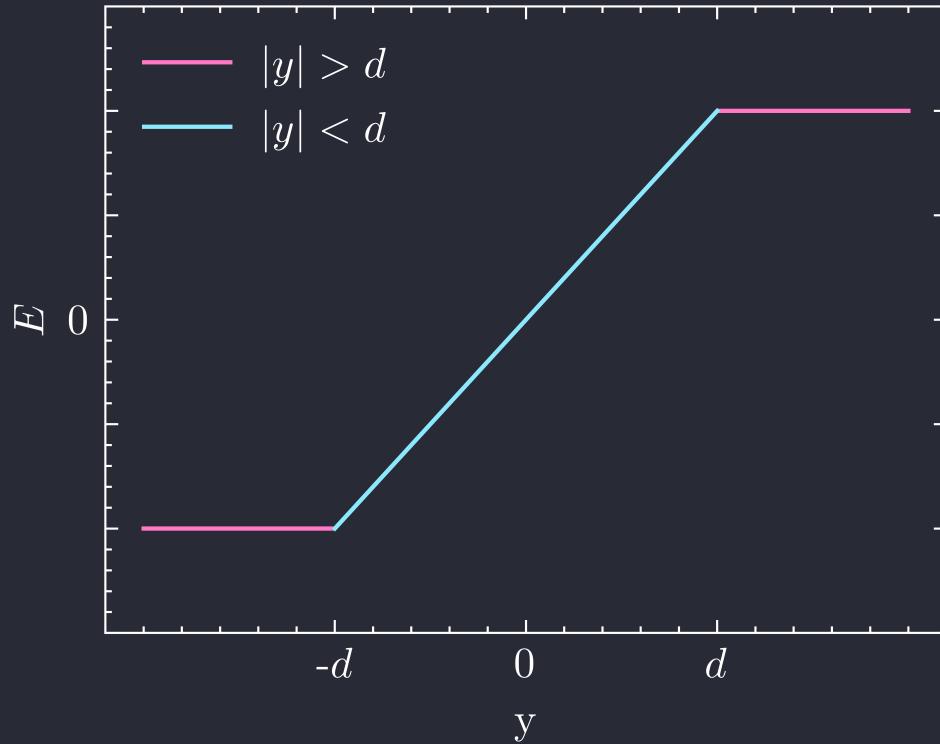


Figure 2.17: Plot of $|\mathbf{E}|$ as a function of y

2.22 Find the potential inside and outside a uniformly charged solid sphere whose radius is R and whose total charge is q . Use infinity as your reference point. Compute the gradient of V in each region, and check that it yields the correct field. Sketch $V(r)$.

The electric field outside the sphere is

$$\mathbf{E}_{out} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}$$

and from Problem 2.8, the electric field inside the sphere is

$$\mathbf{E}_{in} = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}}$$

For points outside the sphere ($r > R$),

$$V(r) = - \int_{\infty}^r \mathbf{E} \cdot d\mathbf{l} = - \frac{1}{4\pi\epsilon_0} \int_{\infty}^r \frac{q}{r'^2} \hat{\mathbf{r}} = \frac{1}{4\pi\epsilon_0} \frac{q}{r'} \Big|_{\infty}^r = \frac{1}{4\pi\epsilon_0} \frac{q}{r}$$

For points inside the sphere ($r < R$),

$$\begin{aligned} V(r) &= - \int_{\infty}^R \mathbf{E} \cdot d\mathbf{l} - \int_R^r \mathbf{E} \cdot d\mathbf{l} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R} - \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} \int_R^r r' dr' \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R} - \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} \left(\frac{r'^2}{2} \right) \Big|_R^r \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left(3 - \frac{r^2}{R^2} \right) \end{aligned}$$

The gradient of V for $r > R$:

$$-\nabla V = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} = \mathbf{E}_{out}$$

and for $r < R$:

$$-\nabla V = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}} = \mathbf{E}_{in}$$

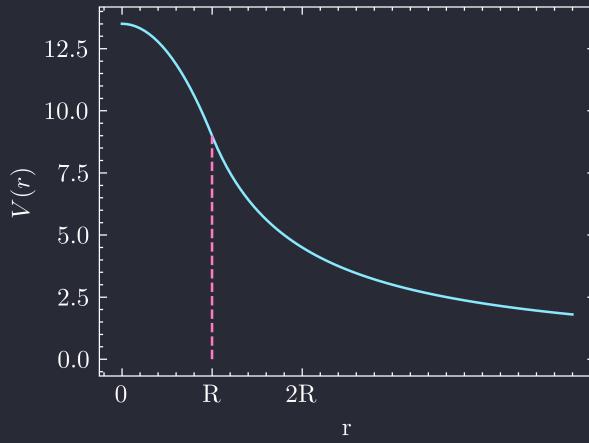


Figure 3.18: Plot of $V(r)$ as a function of r where $q = 1 \text{ nC}$ and $R = 1 \text{ m}$.

2.26 From Griffiths

$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \frac{q_i}{r_i} \quad (2.27)$$

and

$$V = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\mathbf{r}')}{r'} d\ell' \quad \text{and} \quad V = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(\mathbf{r}')}{r'} da' \quad (2.30)$$

- (a.1) Two point charges $+q$ a distance d apart: Find the potential a distance z above the center of the charges: Using Eq. (2.27), the potential is

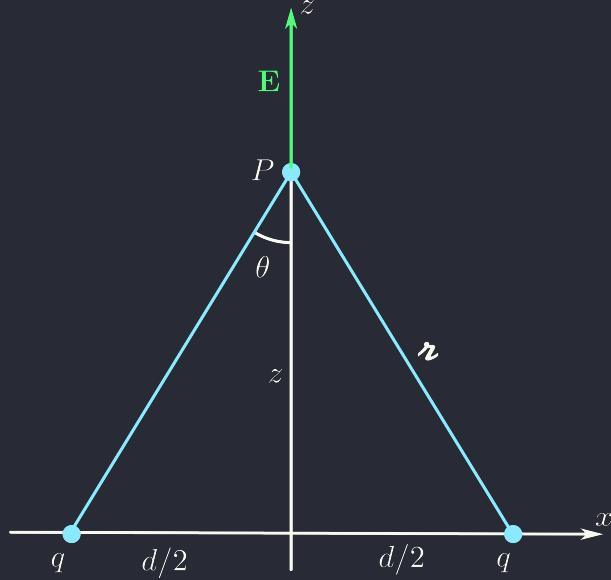


Figure 3.19: Two point charges $+q$ a distance d apart.

$$V_a = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{\sqrt{z^2 + \frac{d^2}{4}}} + \frac{q}{\sqrt{z^2 + \frac{d^2}{4}}} \right)$$

$$V_a = \frac{1}{4\pi\epsilon_0} \frac{2q}{\sqrt{z^2 + \frac{d^2}{4}}}$$

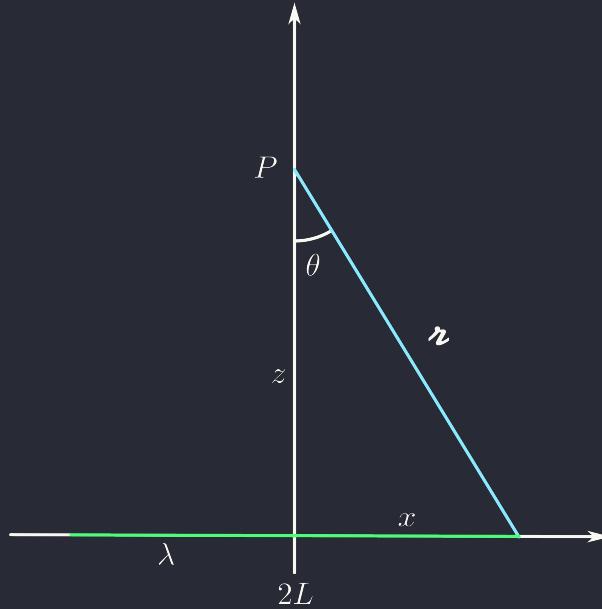
- (a.2) Computing the electric field $\mathbf{E} = -\nabla V$:

$$\begin{aligned} \mathbf{E}_a &= -\frac{\partial V_a}{\partial z} \hat{\mathbf{z}} \\ &= -\frac{1}{4\pi\epsilon_0} \frac{-1}{2} \frac{2q(2z)}{(z^2 + \frac{d^2}{4})^{3/2}} \hat{\mathbf{z}} \end{aligned}$$

simplifying to

$$\mathbf{E}_a = \frac{1}{4\pi\epsilon_0} \frac{2qz}{(z^2 + \frac{d^2}{4})^{3/2}} \hat{\mathbf{z}}$$

which is the same as Ex. 2.1

Figure 3.20: A line charge of density λ .

(b.1) Using Eq. (2.30), the potential is

$$V_b = \frac{1}{4\pi\epsilon_0} \lambda \int_{-L}^L \frac{1}{\sqrt{z^2 + x^2}} dx$$

To solve the integral, we can use the substitution from the trig identity

$$\begin{aligned} \cosh^2 u - \sinh^2 u &= 1 \\ \implies z^2 \cosh^2 u &= z^2 + z^2 \sinh^2 u \\ &= z^2 + x^2 \end{aligned}$$

where

$$\begin{aligned} x = z \sinh u &\implies u = \operatorname{arcsinh} \frac{x}{z} \\ dz &= z \cosh u du \end{aligned}$$

Thus the integral becomes

$$\begin{aligned} V_b &= \frac{1}{4\pi\epsilon_0} \lambda \int \frac{z \cosh u}{z \cosh u} du \\ &= \frac{1}{4\pi\epsilon_0} \lambda u \Big|_{-L}^L \\ &= \frac{1}{4\pi\epsilon_0} \lambda \left[\operatorname{arcsinh} \frac{L}{z} - \operatorname{arcsinh} \frac{-L}{z} \right] \end{aligned}$$

Using $\operatorname{arcsinh}(a) = \ln |a + \sqrt{a^2 + 1}|$:

$$\begin{aligned} \implies \operatorname{arcsinh} \left(\frac{L}{z} \right) &= \ln \left| \frac{L}{z} + \sqrt{\left(\frac{L}{z} \right)^2 + 1} \right| \\ &= \ln \left| \frac{1}{z} (L + \sqrt{L^2 + z^2}) \right| \end{aligned}$$

so the potential is

$$\boxed{V_b = \frac{1}{4\pi\epsilon_0} \lambda \ln \left| \frac{L + \sqrt{L^2 + z^2}}{-L + \sqrt{L^2 + z^2}} \right|}$$

(b.2) The electric field is

$$\begin{aligned} \mathbf{E}_b &= -\frac{\partial V_b}{\partial z} \hat{\mathbf{z}} \\ &= -\frac{1}{4\pi\epsilon_0} \lambda \left[\frac{1}{L + \sqrt{L^2 + z^2}} \left(\frac{1}{2} \frac{2z}{\sqrt{L^2 + z^2}} \right) - \frac{1}{-L + \sqrt{L^2 + z^2}} \left(\frac{1}{2} \frac{2z}{\sqrt{L^2 + z^2}} \right) \right] \hat{\mathbf{z}} \\ &= -\frac{1}{4\pi\epsilon_0} \lambda \frac{z}{\sqrt{L^2 + z^2}} \left[\frac{-L + \sqrt{L^2 + z^2}}{-L^2 + (L^2 + z^2)} - \frac{L + \sqrt{L^2 + z^2}}{-L^2 + (L^2 + z^2)} \right] \hat{\mathbf{z}} \\ &= -\frac{1}{4\pi\epsilon_0} \lambda \frac{-2Lz}{z^2 \sqrt{L^2 + z^2}} \hat{\mathbf{z}} \end{aligned}$$

simplifying to

$$\boxed{\mathbf{E}_b = \frac{1}{4\pi\epsilon_0} \frac{2\lambda L}{z \sqrt{L^2 + z^2}} \hat{\mathbf{z}}}$$

which is the same as Ex. 2.2

(c.1) Using Eq. (2.30) and polar coordinates, the potential is

$$\begin{aligned} V_c &= \frac{1}{4\pi\epsilon_0} \sigma \int_0^{2\pi} \int_0^R \frac{1}{\sqrt{z^2 + r^2}} r dr d\theta \\ &= \frac{1}{4\pi\epsilon_0} 2\pi\sigma \int_0^R \frac{r}{\sqrt{z^2 + r^2}} dr \end{aligned}$$

substituting $u = z^2 + r^2$; $du = 2r dr$:

$$\begin{aligned} V_c &= \frac{1}{4\pi\epsilon_0} \pi\sigma \int \frac{1}{\sqrt{u}} du \\ &= \frac{1}{4\pi\epsilon_0} \pi\sigma 2\sqrt{z^2 + r^2} \Big|_0^R \end{aligned}$$

thus

$$\boxed{V_c = \frac{\sigma}{2\epsilon_0} \left[\sqrt{z^2 + R^2} - z \right]}$$

(c.2) The electric field is

$$\begin{aligned} \mathbf{E}_c &= -\frac{\partial V_c}{\partial z} \hat{\mathbf{z}} \\ &= -\frac{\sigma}{2\epsilon_0} \left[\frac{z}{\sqrt{z^2 + R^2}} - 1 \right] \hat{\mathbf{z}} \end{aligned}$$

thus

$$\boxed{\mathbf{E}_c = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{\sqrt{z^2 + R^2}} \right] \hat{\mathbf{z}}}$$

which is the same as Problem 2.6:

2.6 The electric field is only in the z -direction where $\cos \theta = z/\mathbf{r}$:

$$\begin{aligned}\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int \frac{\sigma}{\mathbf{r}^2} \cos \theta \hat{\mathbf{z}} d\mathbf{a} \\ &= \frac{1}{4\pi\epsilon_0} \int \frac{\sigma z}{(z^2 + r^2)^{3/2}} \hat{\mathbf{z}} d\mathbf{a}\end{aligned}$$

Using polar coordinates: since $d\mathbf{a} = r dr d\theta$

$$\begin{aligned}\mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^R \frac{\sigma z}{(z^2 + r^2)^{3/2}} \hat{\mathbf{z}} r dr d\theta \\ &= \frac{1}{4\pi\epsilon_0} \sigma z (2\pi) \hat{\mathbf{z}} \int_0^R \frac{r}{(z^2 + r^2)^{3/2}} dr \\ &= \frac{\sigma}{2\epsilon_0} z \hat{\mathbf{z}} \left[-\frac{1}{\sqrt{z^2 + r^2}} \right]_0^R \\ &= \frac{\sigma}{2\epsilon_0} z \left[\frac{1}{z} - \frac{1}{\sqrt{z^2 + R^2}} \right] \hat{\mathbf{z}} \\ \mathbf{E} &= \frac{\sigma}{2\epsilon_0} \left[1 - \frac{1}{\sqrt{z^2 + R^2}} \right] \hat{\mathbf{z}}\end{aligned}$$

- (d) if the right-hand charge of Fig. 3.19 is replaced by a charge $-q$, the potential at P using Eq. (2.27) is

$$V_d = 0 \implies \mathbf{E}_d = 0$$

which contradicts the result from Prob 2.2. This is because point P does not give us any information about the electric field which points in the x -direction. In fact any reference point on the z -axis will give us the same result.

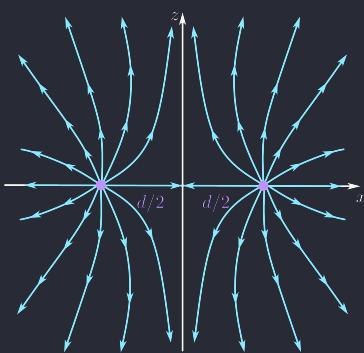


Figure 3.21: E-field for (a)

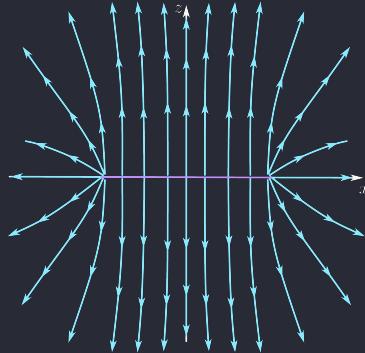


Figure 3.22: E-field for (b)

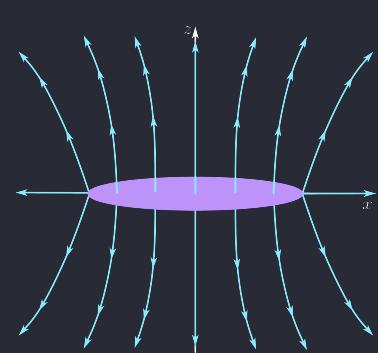


Figure 3.23: E-field for (c)

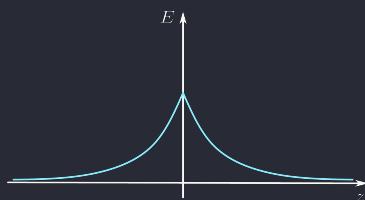
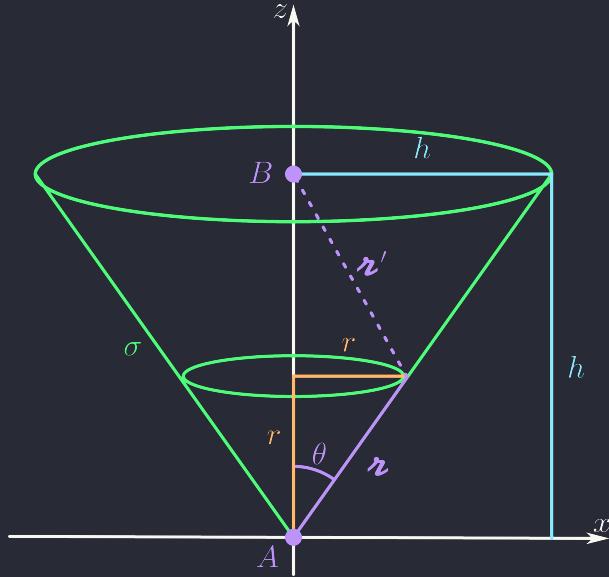


Figure 3.24: E-field for (c) E vs z

Figure 3.25: Empty ice cream cone with surface charge density σ .**2.27**

- (i) Potential at A : Geometrically, we can see from the large right triangle that

$$\begin{aligned}\mathbf{z}^2 &= h^2 + h^2 \\ \implies \mathbf{z} &= h\sqrt{2}, \quad h = \frac{\mathbf{z}}{\sqrt{2}}\end{aligned}$$

and from the smaller right triangle

$$\mathbf{z}^2 = 2r^2 \implies r = \frac{\mathbf{z}}{\sqrt{2}}$$

We can find the potential at A using Eq. (2.30) and integrate the rings of the cone along the slant length $0 \rightarrow h\sqrt{2}$ which gives us the area element $da = 2\pi r d\mathbf{z}$:

$$\begin{aligned}V(A) &= \frac{1}{4\pi\epsilon_0} \int_0^{h\sqrt{2}} \frac{\sigma}{\mathbf{z}} 2\pi r d\mathbf{z} \\ &= \frac{\sigma}{2\epsilon_0\sqrt{2}} \int_0^{h\sqrt{2}} d\mathbf{z} \\ &= \frac{\sigma}{2\epsilon_0\sqrt{2}} \mathbf{z} \Big|_0^{h\sqrt{2}} \\ V(A) &= \frac{\sigma h}{2\epsilon_0}\end{aligned}$$

- (ii) Potential at B : Using the law of cosines,

$$\mathbf{z}'^2 = h^2 + \mathbf{z}^2 - 2h\mathbf{z} \cos\theta$$

where

$$\begin{aligned}\cos\theta &= \frac{h}{\mathbf{z}} \\ &= \frac{h}{h\sqrt{2}} = \frac{1}{\sqrt{2}} = \frac{\sqrt{2}}{2} \\ \implies \mathbf{z}' &= \sqrt{h^2 + \mathbf{z}^2 - h\mathbf{z}\sqrt{2}}\end{aligned}$$

2.35 For a solid sphere radius R and charge q

(a) From Problem 2.22

$$V = \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \left(3 - \frac{r^2}{R^2} \right)$$

and

$$W = \frac{1}{2} \int \rho V d\tau \quad (2.43)$$

So the energy is

$$\begin{aligned} W &= \frac{\rho}{2} \frac{1}{4\pi\epsilon_0} \frac{q}{2R} \int_0^{2\pi} \int_0^\pi \int_0^R \left(3 - \frac{r^2}{R^2} \right) r^2 \sin\theta dr d\theta d\phi \\ &= \frac{\rho}{2} \frac{1}{4\pi\epsilon_0} \frac{q}{2R} 4\pi \int_0^R \left(3r^2 - \frac{r^4}{R^2} \right) dr \\ &= \frac{\rho q}{4R\epsilon_0} \left[r^3 - \frac{r^5}{5R^2} \right]_0^R \\ &= \frac{\rho q}{4R\epsilon_0} \left[R^3 - \frac{R^3}{5} \right] \\ &= \frac{\rho q}{5\epsilon_0} R^2 \end{aligned}$$

where the charge over the volume of the sphere is $\rho = \frac{q}{\frac{4}{3}\pi R^3}$, thus

$$\begin{aligned} W &= \frac{q}{5\epsilon_0} R^2 \frac{q}{\frac{4}{3}\pi R^3} \\ W &= \boxed{\frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R}} \end{aligned}$$

(b) Integrating over all space using

$$W = \frac{\epsilon_0}{2} \int E^2 d\tau \quad (2.45)$$

Where the electric field is

$$\mathbf{E}_{out} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}} \quad \mathbf{E}_{in} = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}}$$

so the energy is

$$\begin{aligned} W &= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} 4\pi q^2 \left[\int_0^R \frac{r^2}{R^6} r^2 dr + \int_R^\infty \frac{1}{r^4} r^2 dr \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\int_0^R \frac{r^4}{R^6} dr + \int_R^\infty \frac{1}{r^2} dr \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\frac{r^5}{5R^6} \Big|_0^R - \frac{1}{R} \Big|_R^\infty \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\frac{R^5}{5R^6} + \frac{1}{R} \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \frac{6}{5R} \\ W &= \boxed{\frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R}} \end{aligned}$$

checkmark.

(c) For a spherical volume of radius a and

$$W = \frac{\epsilon_0}{2} \left(\int_V E^2 d\tau + \oint_S V \mathbf{E} \cdot d\mathbf{a} \right) \quad (2.44)$$

we can assume the volume is outside the charged sphere so

$$V = \frac{1}{4\pi\epsilon_0} \frac{q}{r}$$

From part (b), the first term is

$$\begin{aligned} \frac{\epsilon_0}{2} \int_V E^2 d\tau &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\frac{1}{5R} - \frac{1}{a} + \frac{1}{R} \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\frac{6}{5R} - \frac{1}{a} \right] \end{aligned}$$

the second term is at $r = a$

$$\begin{aligned} \frac{\epsilon_0}{2} \oint_V V \mathbf{E} \cdot d\mathbf{a} &= \frac{\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} \int \frac{q}{r} \frac{q}{r^2} r^2 \sin\theta d\theta d\phi \\ &= \frac{4\pi\epsilon_0}{2} \frac{1}{(4\pi\epsilon_0)^2} \frac{1}{r} \Big|_{r=a} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2a} \end{aligned}$$

so the total energy is

$$\begin{aligned} W &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{2} \left[\frac{6}{5R} - \frac{1}{a} \right] + \frac{1}{4\pi\epsilon_0} \frac{q^2}{2a} \\ &= \frac{1}{4\pi\epsilon_0} \frac{3q^2}{5R} \end{aligned}$$

As $a \rightarrow \infty$ the $\oint V \mathbf{E} \cdot d\mathbf{a}$ term goes to zero.

2.40 Two cavities radii a and b in a conducting sphere of radius R with a point charge q_a and q_b respectively in each cavity.

- (a) Surface charge densities:

On the surface of cavity a the charge density is simply

$$\sigma_a = \frac{-q_a}{4\pi a^2}$$

and

$$\sigma_b = \frac{-q_b}{4\pi b^2}$$

respectively. For the surface of the conducting sphere, the charge density is positive and equal to the superposition of the two charges:

$$\sigma_R = \frac{q_a + q_b}{4\pi R^2}$$

- (b) The field outside the conductor is equivalent to a point charge at the center of the sphere with the sum of the charges:

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q_a + q_b}{r^2} \hat{\mathbf{r}}$$

- (c) The field in cavity a with respect to the center of the cavity is

$$\mathbf{E}_a = \frac{1}{4\pi\epsilon_0} \frac{q_a}{a^2} \hat{\mathbf{a}}$$

and in cavity b is

$$\mathbf{E}_b = \frac{1}{4\pi\epsilon_0} \frac{q_b}{b^2} \hat{\mathbf{b}}$$

- (d) The field due to the cavity charge is zero in the exterior of the cavity, so there is no Force on q_a or q_b .
- (e) If a charge q_c was brought near the conductor from outside, there would be a change in (a) σ_R and (b) \mathbf{E}_{out} .

2.48 Net force of the southern hemisphere exerting on the northern hemisphere (solid sphere) with an inside Electric field (Problem 2.8)

$$E_{in} = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^3} \mathbf{r}$$

where the total force is

$$\mathbf{F} = Q\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Q^2}{R^3} \mathbf{r}$$

Finding the net force exerted by the southern hemisphere: integrate $dF = \mathbf{F}/V$ over the southern hemisphere:

$$\begin{aligned} dF &= \frac{1}{\frac{4}{3}\pi R^3} \frac{1}{4\pi\epsilon_0} \frac{Q^2}{R^3} \mathbf{r} d\tau \\ &= \frac{3Q^2}{16\pi^2\epsilon_0 R^6} \mathbf{r} d\tau \end{aligned}$$

The symmetry of the sphere implies that the Force is only in the z -direction i.e. $F_z = F \cos \theta$, so integrating over the southern hemisphere:

$$\begin{aligned} \int_0^{2\pi} \int_0^{\pi/2} \int_0^R F_z d\tau &= \frac{3Q^2}{16\pi^2\epsilon_0 R^6} \int_0^{2\pi} \int_0^{\pi/2} \int_0^R r \cos \theta r^2 \sin \theta dr d\theta d\phi \\ &= \frac{3Q^2}{16\pi^2\epsilon_0 R^6} (2\pi) \left(\frac{r^4}{4} \right) \Big|_0^R \int_0^{\pi/2} \sin \theta \cos \theta d\theta d\phi \\ &= \frac{3Q^2}{32\pi\epsilon_0 R^2} \frac{\sin^2 x}{2} \Big|_0^{\pi/2} \\ &= \boxed{\frac{3Q^2}{64\pi\epsilon_0 R^2}} \end{aligned}$$

3.4

- (a) Average field over a spherical surface due to charges outside the sphere is the same at the center:

For a charge q a distance z above the center of the sphere, we can use the same geometrical argument from HW 2 Problem 2.7: The average field at over the surface will be in the negative z direction

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{\mathbf{r}^2} \cos\psi(-\hat{\mathbf{z}})$$

where (using law of cosines)

$$\mathbf{r}^2 = z^2 + R^2 - 2zR\cos\theta \quad \cos\phi = \frac{z - R\cos\theta}{\mathbf{r}}$$

The surface element is $da = R^2 \sin\theta d\theta d\phi$, so the average field is

$$\begin{aligned} \mathbf{E}_{\text{avg}} &= \frac{1}{4\pi R^2} \frac{1}{4\pi\epsilon_0} (-qR^2)\hat{\mathbf{z}} \int_0^{2\pi} \int_0^\pi \frac{z - R\cos\theta}{(z^2 + R^2 - 2zR\cos\theta)^{3/2}} \sin\theta d\theta d\phi \\ &= \frac{1}{4\pi} \frac{1}{4\pi\epsilon_0} (-q)\hat{\mathbf{z}} (2\pi) \int_0^\pi \frac{z - R\cos\theta}{(z^2 + R^2 - 2zR\cos\theta)^{3/2}} \sin\theta d\theta \\ &= \frac{1}{4\pi\epsilon_0} \frac{-q}{2} \hat{\mathbf{z}} f \end{aligned}$$

The integral evaluates to (from Problem 2.7): Using the substitution $u = \cos\theta$: $du = -\sin\theta d\theta$, and the limits of integration are $[\cos 0, \cos \pi]$. So,

$$\begin{aligned} f(\theta) &= - \int_1^{-1} \frac{z - Ru}{(z^2 + R^2 - 2zRu)^{3/2}} du \\ &= \int_{-1}^1 \frac{z - Ru}{(z^2 + R^2 - 2zRu)^{3/2}} du \end{aligned}$$

substituting again with $v = \sqrt{z^2 + R^2 - 2zRu}$; $dv = \frac{zR}{v} du$; and $u = \frac{1}{2zR}(z^2 + R^2 - v^2)$:

$$\begin{aligned} f(v) &= -\frac{1}{zR} \int \frac{z - \frac{1}{2z}(z^2 + R^2 - v^2)}{v^3} v dv \\ &= -\frac{1}{2z^2 R} \int \frac{2z^2 - (z^2 + R^2 - v^2)}{v^2} dv \\ &= -\frac{1}{2z^2 R} \int \frac{v^2 + z^2 - R^2}{v^2} dv \\ &= -\frac{1}{2z^2 R} \int \left(1 + \frac{z^2 - R^2}{v^2}\right) dv \\ &= -\frac{1}{2z^2 R} \left(v - \frac{z^2 - R^2}{v}\right) \end{aligned}$$

substituting back in $v = \sqrt{z^2 + R^2 - 2zRu}$,

$$\begin{aligned}
f(u) &= -\frac{1}{2z^2R} \left(\frac{z^2 + R^2 - 2zRu}{\sqrt{z^2 + R^2 - 2zRu}} - \frac{z^2 - R^2}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\
&= -\frac{1}{2z^2R} \left(\frac{2R^2 - 2zRu}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\
&= \frac{1}{z^2} \left(\frac{zu - R}{\sqrt{z^2 + R^2 - 2zRu}} \right) \Big|_{-1}^1 \\
&= \frac{1}{z^2} \left(\frac{z - R}{\sqrt{z^2 + R^2 - 2zR}} - \frac{-z - R}{\sqrt{z^2 + R^2 + 2zR}} \right) \\
&= \frac{1}{z^2} \left(\frac{z - R}{\sqrt{z^2 + R^2 - 2zR}} + \frac{z + R}{z + R} \right) \\
&= \frac{1}{z^2} \left(\frac{z - R}{\sqrt{z^2 + R^2 - 2zR}} + 1 \right)
\end{aligned}$$

where the positive root $\sqrt{z^2 + R^2 - 2zR} = (z - R)$ for $z > R$, so

$$\mathbf{E}_{\text{avg}} = \frac{1}{4\pi\epsilon_0} \left(-\frac{q}{2z^2} \right) \left(\frac{z - R}{z - R} + \frac{z + R}{z + R} \right) \hat{\mathbf{z}}$$

Simplifying to

$$\boxed{\mathbf{E}_{\text{avg}} = -\frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \hat{\mathbf{z}}}$$

which is the same as the field at the center of the sphere. For a collection of particles, we can use superposition and find the net field as the sum of the fields at the center from each charge.

(b) For charges inside the sphere we can use the result from before: for one charge

$$\mathbf{E}_{\text{avg}} = -\frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \left(\frac{z - R}{\sqrt{z^2 + R^2 - 2zR}} + \frac{z + R}{z + R} \right) \hat{\mathbf{z}}$$

but now the positive root is $\sqrt{z^2 + R^2 - 2zR} = (R - z)$ for $z < R$, so

$$\begin{aligned}
\mathbf{E}_{\text{avg}} &= -\frac{1}{4\pi\epsilon_0} \frac{q}{z^2} \left(\frac{z - R}{R - z} + \frac{z + R}{z + R} \right) \hat{\mathbf{z}} \\
&= -\frac{1}{4\pi\epsilon_0} \frac{q}{z^2} (-1 + 1) \hat{\mathbf{z}} \\
\mathbf{E}_{\text{avg}} &= 0
\end{aligned}$$

And we can superimpose the fields from a collection of charges

$$\boxed{\mathbf{E}_{\text{avg}} = 0 + 0 + \dots = 0}$$

3.7 Charges $+q$ & $-2q$ are respectively $z = 3d$ & $z = d$ above the xy plane (grounded conductor). Find the force of the charge $+q$:

We can use the method of images and replace the grounded conductor with two charges $-q$ at $z = -3d$ and $+2q$ at $z = -d$. Thus the force on $+q$ is the superposition of the forces from the three charges: The separation vectors are

$$\mathbf{r}_{-2q} = 2d\hat{\mathbf{z}}$$

$$\mathbf{r}_{+2q} = 4d\hat{\mathbf{z}}$$

$$\mathbf{r}_{-q} = 6d\hat{\mathbf{z}}$$

Finally, the force on charge $+q$ is

$$\begin{aligned}\mathbf{F} &= \mathbf{F}_{-2q} + \mathbf{F}_{+2q} + \mathbf{F}_{-q} \\ &= \frac{1}{4\pi\epsilon_0} \left(\frac{-2q^2}{(2d)^2} + \frac{2q^2}{(4d)^2} + \frac{-q^2}{(6d)^2} \right) \hat{\mathbf{z}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{d^2} \left(-\frac{1}{2} + \frac{1}{8} - \frac{1}{36} \right) \hat{\mathbf{z}}\end{aligned}$$

which simplifies to

$$\boxed{\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{29q^2}{72d^2} \hat{\mathbf{z}}}$$

3.8 From Griffiths, where the configuration has another point charge

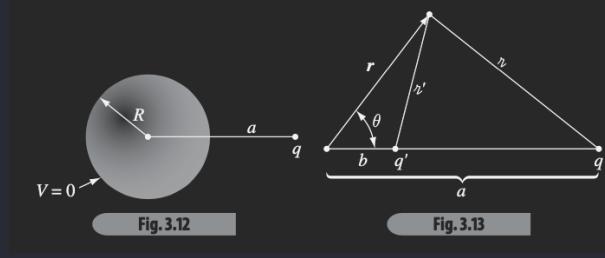


Figure 4.26: From Griffiths Example 3.2

$$q' = -\frac{R}{a}q \quad (3.15)$$

placed a distance

$$b = \frac{R^2}{a} \quad (3.16)$$

thus the potential of the config

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{\mathbf{r}} + \frac{q'}{\mathbf{r}'} \right) \quad (3.17)$$

(a) Using law of cosines, show that Eq. (3.17) can be written as

$$V(r, \theta) = \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} - \frac{q}{\sqrt{R^2 + (ra/R)^2 - 2ra \cos \theta}} \right]$$

From Fig. 4.26, we can see that

$$\mathbf{r} = \sqrt{r^2 + a^2 - 2ra \cos \theta} \quad \mathbf{r}' = \sqrt{r^2 + b^2 - 2rb \cos \theta}$$

so using Eq. (3.15) and Eq. (3.16) we can rewrite

$$\begin{aligned} \frac{q'}{\mathbf{r}'} &= \frac{-R}{a} \frac{q}{\sqrt{r^2 + (R^2/a)^2 - 2r(R^2/a) \cos \theta}} \\ &= -\frac{1}{\sqrt{\frac{a^2}{R^2}}} \frac{q}{\sqrt{r^2 + (R^2/a)^2 - 2r(R^2/a) \cos \theta}} \\ &= -\frac{q}{\sqrt{r^2 a/R^2 + (R^2/a)^2 a^2/R^2 - 2r(R^2/a) \cos \theta a^2/R^2}} \\ \frac{q'}{\mathbf{r}'} &= -\frac{q}{\sqrt{R^2 + (ra/R)^2 - 2ra \cos \theta}} \end{aligned}$$

Now we can rewrite Eq. (3.17) as

$$V = \frac{1}{4\pi\epsilon_0} \left[\frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} - \frac{q}{\sqrt{R^2 + (ra/R)^2 - 2ra \cos \theta}} \right]$$

- (b) Finding the induced charge on the sphere & integrating to get total induced charge: The normal component of the potential is $\hat{\mathbf{n}} = \hat{\mathbf{r}}$, so

$$\begin{aligned}\sigma &= -\epsilon_0 \frac{\partial V}{\partial r} \Big|_{r=R} \\ &= -\epsilon_0 \frac{1}{4\pi\epsilon_0} q \left(-\frac{1}{2} \right) \left[\frac{2r - 2a \cos \theta}{(r^2 + a^2 - 2ra \cos \theta)^{3/2}} - \frac{2ra^2/R^2 - 2a \cos \theta}{(R^2 + (ra/R)^2 - 2ra \cos \theta)^{3/2}} \right] \Big|_{r=R} \\ &= \frac{q}{4\pi} \left[\frac{R - a \cos \theta}{(R^2 + a^2 - 2Ra \cos \theta)^{3/2}} - \frac{a^2/R - a \cos \theta}{(R^2 + a^2 - 2Ra \cos \theta)^{3/2}} \right] \\ &= \frac{q}{4\pi} \left[\frac{R - a^2/R}{(R^2 + a^2 - 2Ra \cos \theta)^{3/2}} \right]\end{aligned}$$

which simplifies to

$$\boxed{\sigma(\theta) = \frac{q}{4\pi R} \frac{R^2 - a^2}{(R^2 + a^2 - 2Ra \cos \theta)^{3/2}}}$$

Integrating to get the total induced charge using the surface element $da = R^2 \sin \theta d\theta d\phi$:

$$\begin{aligned}Q &= \int \sigma da \\ Q &= \frac{q}{4\pi R} (R^2 - a^2) (2\pi R^2) \int_0^\pi \frac{\sin \theta}{(R^2 + a^2 - 2Ra \cos \theta)^{3/2}} d\theta \\ \text{using } u &= R^2 + a^2 - 2Ra \cos \theta; \quad du = 2Ra \sin \theta d\theta \\ &= \frac{qR}{2} (R^2 - a^2) \frac{1}{2Ra} \int \frac{1}{u^{3/2}} du \\ &= \frac{qR}{2} (R^2 - a^2) \frac{1}{2Ra} \frac{-2}{\sqrt{u}} \\ &= -\frac{q}{2a} (R^2 - a^2) \frac{1}{\sqrt{R^2 + a^2 - 2Ra \cos \theta}} \Big|_0^\pi \\ &= -\frac{q}{2a} (R^2 - a^2) \left[\frac{1}{\sqrt{R^2 + a^2 + 2Ra}} - \frac{1}{\sqrt{R^2 + a^2 - 2Ra}} \right]\end{aligned}$$

From Fig. 4.26, we can see that $R < a$ so the positive root is $\sqrt{R^2 + a^2 - 2Ra} = (a - R)$. Now the total induced charge is

$$\begin{aligned}Q &= -\frac{q}{2a} (R^2 - a^2) \left[\frac{1}{a + R} - \frac{1}{a - R} \right] \\ &= -\frac{q}{2a} (R^2 - a^2) \left[\frac{a - R}{(a + R)(a - R)} - \frac{a + R}{(a - R)(a + R)} \right] \\ &= -\frac{q}{2a} (R^2 - a^2) \left[\frac{a - R - (a + R)}{a^2 - R^2} \right] \\ &= -\frac{q}{2a} (R^2 - a^2) \left[\frac{-2R}{-(R^2 - a^2)} \right]\end{aligned}$$

Thus the total induced charge is

$$\boxed{Q = -\frac{R}{a} q = q'}$$

- (c) The energy of the config:

First we find the force on q due the induced charge q' which are separated by a distance $a - b$:

$$\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{qq'}{(a - b)^2} \hat{\mathbf{n}}$$

Using Eq. (3.15) and Eq. (3.16)

$$\begin{aligned}\mathbf{F} &= -\frac{1}{4\pi\epsilon_0} \frac{R}{a} \frac{q^2}{(a - R^2/a)^2} \hat{\mathbf{a}} \\ &= -\frac{1}{4\pi\epsilon_0} \frac{R}{a} \frac{q^2}{\frac{1}{a^2}(a^2 - R^2)^2} \hat{\mathbf{a}} \\ \mathbf{F} &= -\frac{1}{4\pi\epsilon_0} \frac{Raq^2}{(a^2 - R^2)^2} \hat{\mathbf{a}}\end{aligned}$$

Now we can determine the energy by calculating the work it takes to bring q from infinity: The line element is $d\ell = da\hat{\mathbf{a}}$ since the force is in the negative a direction; so the work required to *oppose* the force is

$$\begin{aligned}W &= - \int_{\infty}^a \mathbf{F} \cdot d\ell \\ &= -\frac{1}{4\pi\epsilon_0} Rq^2 \int_{\infty}^a \frac{a'}{(a'^2 - R^2)^2} (-da') \\ \text{using } u &= a'^2 - R^2; \quad du = 2a' da' \\ &= \frac{1}{4\pi\epsilon_0} Rq^2 \int \frac{1}{2u^2} du \\ &= \frac{1}{4\pi\epsilon_0} Rq^2 \left[-\frac{1}{2u} \right] \\ &= \frac{1}{4\pi\epsilon_0} Rq^2 \left[-\frac{1}{2(a^2 - R^2)} \right] \Big|_{\infty}^a \\ &= \frac{1}{4\pi\epsilon_0} Rq^2 \left[-\frac{1}{2(a^2 - R^2)} - 0 \right]\end{aligned}$$

which simplifies to

$$W = -\frac{1}{4\pi\epsilon_0} \frac{Rq^2}{2(a^2 - R^2)}$$

3.10 For a second image charge q'' inside the center of the sphere (it must not be outside the sphere) with potential

$$V = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{z} + \frac{q'}{z'} \right) + V_0$$

where

$$V_0 = \frac{1}{4\pi\epsilon_0} \frac{q''}{z''} = \frac{1}{4\pi\epsilon_0} \frac{q''}{R} \implies q'' = 4\pi\epsilon_0 R V_0$$

So for a neutral conducting sphere the potential should be zero at the surface, i.e. the magnitude of the image charges q' and q'' are equal and opposite:

$$q' = -q''$$

The distance from the second image charge and q is a , so

$$\begin{aligned} F &= \frac{1}{4\pi\epsilon_0} \left[\frac{qq'}{(a-b)^2} + \frac{qq''}{a^2} \right] \\ &= \frac{1}{4\pi\epsilon_0} q \left[\frac{q'a^2}{(a^2-R^2)^2} - \frac{q'}{a^2} \frac{(a^2-R^2)^2}{(a^2-R^2)^2} \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{qq'}{(a^2-R^2)^2} [a^2 - a^2 + 2R^2 - R^4/a^2] \end{aligned}$$

Using Eq. (3.15) $q' = -\frac{R}{a}q$

$$\begin{aligned} F &= \frac{1}{4\pi\epsilon_0} \frac{q^2}{(a^2-R^2)^2} \left(\frac{-R}{a} \right) [2R^2 - R^4/a^2] \\ &= -\frac{1}{4\pi\epsilon_0} \frac{q^2}{(a^2-R^2)^2} \left(\frac{R^3}{a^3} \right) [2a^2 - R^2] \end{aligned}$$

So the force of attraction has magnitude

$$F_{\text{att}} = \frac{1}{4\pi\epsilon_0} \frac{q^2 R^3}{a^3 (a^2 - R^2)^2} [2a^2 - R^2]$$

3.11 Force between point charge q and spherical conductor of total charge q :

We can use a second image charge (at the center of the sphere) where

$$q'' + q' = q$$

So the force between q and the conductor is

$$\begin{aligned} F &= \frac{1}{4\pi\epsilon_0} \left[\frac{qq'}{(a-b)^2} + \frac{qq''}{a^2} \right] \\ &= \frac{1}{4\pi\epsilon_0} q \left[\frac{q'a^2}{(a^2-R^2)^2} + \frac{q-q'}{a^2} \right] \end{aligned}$$

Using Eq. (3.15) $q' = -\frac{R}{a}q$:

$$\begin{aligned} F &= \frac{1}{4\pi\epsilon_0} q \left[\frac{-Rqa}{(a^2-R^2)^2} + \frac{q+Rq/a}{a^2} \right] \\ &= \frac{1}{4\pi\epsilon_0} q^2 \left[-\frac{Ra}{(a^2-R^2)^2} + \frac{a+R}{a^3} \right] \end{aligned}$$

So the force is attractive when $[] < 0$, or if we define the critical value a_c where

$$\begin{aligned} \frac{Ra_c}{(a_c^2-R^2)^2} &= \frac{a_c+R}{a_c^3} \\ Ra_c^4 &= (a_c^2-R^2)^2(a_c+R) \\ Ra_c^4 &= (a_c^4-2a_c^2R^2+R^4)(a_c+R) \\ Ra_c^4 &= a_c^5-2a_c^3R^2+R^4a_c+Ra_c^4-2a_c^2R^3+R^5 \\ 0 &= a_c^5-2a_c^3R^2-2a_c^2R^3+R^4a_c+R^5 \end{aligned}$$

From the hint, the solution must be in the form

$$a_c = R \frac{1+\sqrt{5}}{2}$$

which is the golden ratio i.e. in quadratic form

$$\phi^2 - \phi - 1 = 0 \implies \phi = \frac{1+\sqrt{5}}{2}$$

So

$$a_c = R\phi \implies \phi = \frac{a_c}{R}$$

With this intuition, we can divide our quintic equation by R^5 :

$$\begin{aligned} 0 &= \frac{a_c^5}{R^5} - 2\frac{a_c^3}{R^3} - 2\frac{a_c^2}{R^2} + \frac{a_c^4}{R^4} + 1 \\ &= \phi^5 - 2\phi^3 - 2\phi^2 + \phi + 1 \end{aligned}$$

then we can factor it by dividing by the golden ratio equation $\phi^2 - \phi - 1 = 0$ (using polynomial long division):

$$\begin{array}{r} & x^3 & + x^2 & - 1 \\ x^2 - x - 1 & \overline{)x^5 & - 2x^3 - 2x^2 + x + 1} \\ & -x^5 & + x^4 & + x^3 \\ \hline & x^4 & - x^3 & - 2x^2 \\ & -x^4 & + x^3 & + x^2 \\ \hline & & -x^2 & + x + 1 \\ & & x^2 & - x - 1 \\ \hline & & & 0 \end{array}$$

Thus

$$0 = (\phi^2 - \phi - 1)(\phi^3 + \phi^2 - 1)$$

Where the quadratic equation factors to the golden ratio (positive root),

$$\begin{aligned} \phi^2 - \phi - 1 &= 0 \\ \implies \phi &= \frac{1 + \sqrt{5}}{2} = \frac{a_c}{R} \\ \implies a_c &= R \frac{1 + \sqrt{5}}{2} \end{aligned}$$

So when $a < a_c$, $[] < 0$ i.e.

$$F = \frac{1}{4\pi\epsilon_0} q^2 []$$

thus the magnitude of the force $F < 0$ which implies that the force is attractive.

Furthermore, in the cubic equation, there is a real root at approximately $\phi = 0.75488$ (using desmos/root-finding calculator)

$$\implies a_{c2} \approx 0.75488R$$

So, at $a < 0.75488R$ (pretending $\phi = a/R$)

$$(\phi^2 - \phi - 1)(\phi^3 + \phi^2 - 1) = (-C_1)(-C_2) = +C_3 = [] > 0$$

where C_1, C_2, C_3 are constants. So the force becomes repulsive again when $a < 0.75488R$.

3.13 Two semi-infinite grounded conducting planes meeting as shown in Fig. 4.27 To set up the

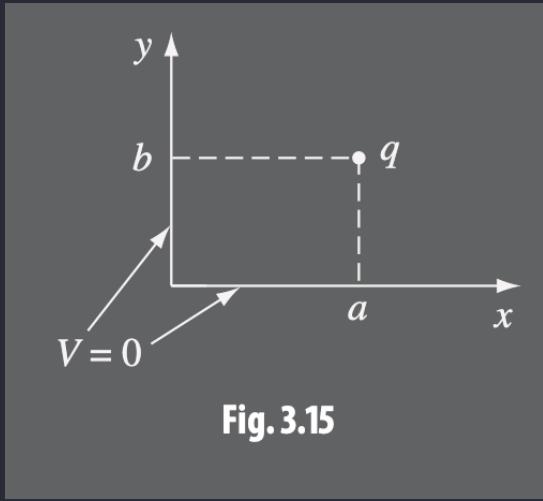


Figure 4.27: From Griffiths

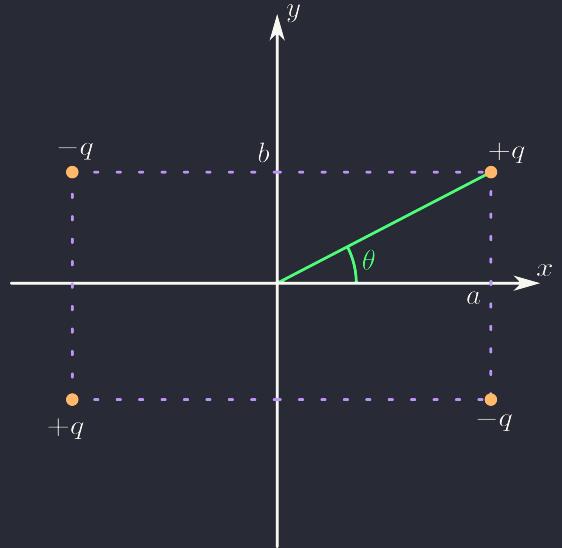


Figure 4.28: Three image charges

problem so we can have a potential of zero at the planes, we can place image charges $-q$ at $(a, -b)$ and $(-a, b)$, and place an image charge $+q$ at $(-a, -b)$ to balance the potentials at the axes.

The potential in the region $x > 0, y > 0$ is:

$$V = \frac{1}{4\pi\epsilon_0}q \left[\frac{1}{\sqrt{(x-a)^2 + (y-b)^2 + z^2}} + \frac{1}{\sqrt{(x+a)^2 + (y+b)^2 + z^2}} \right. \\ \left. - \frac{1}{\sqrt{(x+a)^2 + (y-b)^2 + z^2}} - \frac{1}{\sqrt{(x-a)^2 + (y+b)^2 + z^2}} \right]$$

The force on q is (using Fig. 4.28):

$$\mathbf{F} = \frac{1}{4\pi\epsilon_0}q^2 \left[-\frac{1}{(2a)^2}\hat{\mathbf{x}} - \frac{1}{(2b)^2}\hat{\mathbf{y}} \right. \\ \left. + \frac{1}{(2\sqrt{a^2 + b^2})^2}(\cos\theta\hat{\mathbf{x}} + \sin\theta\hat{\mathbf{y}}) \right]$$

where

$$\cos\theta = \frac{a}{\sqrt{a^2 + b^2}} \quad \sin\theta = \frac{b}{\sqrt{a^2 + b^2}}$$

So

$$\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{q^2}{4} \left[-\frac{1}{a^2}\hat{\mathbf{x}} - \frac{1}{b^2}\hat{\mathbf{y}} + \left(\frac{a}{(a^2 + b^2)^{3/2}}\hat{\mathbf{x}} + \frac{b}{(a^2 + b^2)^{3/2}}\hat{\mathbf{y}} \right) \right] \\ \boxed{\mathbf{F} = \frac{q^2}{16\pi\epsilon_0} \left[\left(\frac{a}{(a^2 + b^2)^{3/2}} - \frac{1}{a^2} \right)\hat{\mathbf{x}} + \left(\frac{b}{(a^2 + b^2)^{3/2}} - \frac{1}{b^2} \right)\hat{\mathbf{y}} \right]}$$

The work done to bring q from infinity to the origin:

Integrating the opposing force from $\infty \rightarrow (a, b)$

$$\begin{aligned} W &= - \int_{\infty}^{(a,b)} \mathbf{F} \cdot d\ell \\ &= - \left[\int_{(\infty,\infty)}^{(a,\infty)} \mathbf{F} \cdot d\ell_a + \int_{(a,\infty)}^{(a,b)} \mathbf{F} \cdot d\ell_b \right] \end{aligned}$$

where

$$\begin{aligned} d\ell_a &= dx\hat{\mathbf{x}} + (0)\hat{\mathbf{y}} = dx\hat{\mathbf{x}}, \quad d\ell_b = (0)\hat{\mathbf{x}} + dy\hat{\mathbf{y}} = dy\hat{\mathbf{y}} \\ \implies \int_{\infty,\infty}^{(a,\infty)} \mathbf{F} \cdot d\ell_a &= \frac{q^2}{16\pi\epsilon_0} \int_{\infty}^a \left[\left(\frac{x}{(x^2+b^2)^{3/2}} - \frac{1}{x^2} \right) dx \right] \Big|_{b=\infty} \\ &= \frac{q^2}{16\pi\epsilon_0} \int_{\infty}^a \left[-\frac{1}{x^2} dx \right] \\ \text{and } \mathbf{F} \cdot d\ell_b &= \frac{q^2}{16\pi\epsilon_0} \int_{\infty}^b \left[\left(\frac{y}{(a^2+y^2)^{3/2}} - \frac{1}{y^2} \right) dy \right] \end{aligned}$$

So the work done is

$$W = - \frac{q^2}{16\pi\epsilon_0} \left[\int_{\infty}^a \left(-\frac{1}{a^2} \right) dx + \int_{\infty}^b \left(\frac{b}{(a^2+b^2)^{3/2}} - \frac{1}{b^2} \right) dy \right]$$

Evaluating the two integrals using

$$\begin{aligned} \int -\frac{1}{x^2} dx &= \frac{1}{x} \\ \int \frac{y}{(a^2+y^2)^{3/2}} dy &= -\frac{1}{\sqrt{a^2+y^2}} \end{aligned}$$

gives the work done is

$$W = - \frac{q^2}{16\pi\epsilon_0} \left[\frac{1}{a} - \frac{1}{\sqrt{a^2+b^2}} + \frac{1}{b} \right]$$

or

$$W = \frac{q^2}{16\pi\epsilon_0} \left[\frac{1}{\sqrt{a^2+b^2}} - \frac{1}{a} - \frac{1}{b} \right]$$

We can solve the problem with the method of images, as long as the angle ϕ divides 180° into an integer, e.g., $\phi = 180, 90, 60, 45, 36, 30, 20, 18, 15, 12, 10, 9, 6, 5, 4, 3, 2, 1, 0.5, \dots$

We would place a ‘mirror’ at each ϕ division and place an image charge $-q$ that mirrors the point charge q and repeat the process with the next image charge q (making sure to flip charges each time) until we have a symmetric configuration of charges as shown in Fig. 4.29.

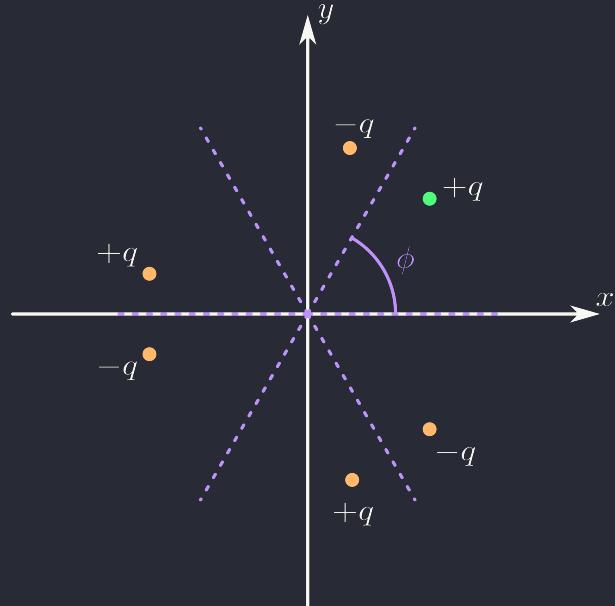


Figure 4.29: Method of images for $\phi = 60^\circ$

Homework 5

3.16 From Griffiths:

$$V(x, y) = \frac{2V_0}{\pi} \arctan\left(\frac{\sin \pi y/a}{\sinh \pi x/a}\right)$$

where the surface charge density is given by a Gaussian pillbox

$$\sigma = -\epsilon_0 \frac{\partial V}{\partial n}$$

since normal of the surface is $\hat{\mathbf{n}} = \hat{\mathbf{x}}$,

$$\begin{aligned} \sigma &= -\epsilon_0 \frac{\partial}{\partial x} \left(\frac{2V_0}{\pi} \arctan(b) \right) \\ &= -\frac{2V_0 \epsilon_0}{\pi} \frac{1}{1 + b^2} \left(\frac{-\sin(\pi y/a)}{\sinh^2(\pi x/a)} \right) \cosh(\pi x/a) \frac{\pi}{a} \\ &= \frac{2V_0 \epsilon_0}{a} \frac{\sin(\pi y/a)}{\sinh^2(\pi x/a) + \sin^2(\pi y/a)} \cosh(\pi x/a) \end{aligned}$$

where at the boundary $x = 0$,

$$\sinh(0) = 0, \cosh(0) = 1$$

so

$$\sigma = \frac{2V_0 \epsilon_0}{a} \frac{\sin(\pi y/a)}{\sin^2(\pi y/a)} = \boxed{\frac{2V_0 \epsilon_0}{a} \sin(\pi y/a)}$$

3.18 Since only the top plate has a potential V_0 the boundary conditions are

- (i) $V = V_0$ at $z = a$
- (ii) $V = 0$ at $x = 0, y = 0, z = 0, x = a$, and $y = a$

We look for solutions

$$V(x, y, z) = X(x)Y(y)Z(z)$$

and solving for Laplace's equation

$$\frac{1}{X} \frac{d^2X}{dx^2} + \frac{1}{Y} \frac{d^2Y}{dy^2} + \frac{1}{Z} \frac{d^2Z}{dz^2} = 0$$

where

$$\frac{1}{X} \frac{d^2X}{dx^2} = C_1, \frac{1}{Y} \frac{d^2Y}{dy^2} = C_2, \frac{1}{Z} \frac{d^2Z}{dz^2} = C_3 \quad \text{with} \quad C_1 + C_2 + C_3 = 0$$

The textbook suggests that $C_1 = -k^2$, $C_2 = -l^2$, and $C_3 = k^2 + l^2$ where k, l are constants. We can solve for X, Y, Z separately:

$$\begin{aligned} X(x) &= A \sin(kx) + B \cos(kx) \\ Y(y) &= C \sin(ly) + D \cos(ly) \\ Z(z) &= E e^{\sqrt{k^2+l^2}z} + F e^{-\sqrt{k^2+l^2}z} \end{aligned}$$

So for the boundary conditions (ii)

$$\begin{aligned} \left. \frac{d^2X}{dx^2} \right|_{x=0} &= -A^2 k^2 \sin^2(0) - B^2 k^2 \cos^2(0) = 0 \implies B = 0 \\ \left. \frac{d^2X}{dx^2} \right|_{x=a} &= -A^2 k^2 \sin^2(ka) = 0 \implies k = \frac{n\pi}{a} \quad \text{for } n = 1, 2, 3, \dots \\ \left. \frac{d^2Y}{dy^2} \right|_{y=0} &= -C^2 l^2 \sin^2(0) - D^2 l^2 \cos^2(0) = 0 \implies D = 0 \\ \left. \frac{d^2Y}{dy^2} \right|_{y=a} &= -C^2 l^2 \sin^2(la) = 0 \implies l = \frac{m\pi}{a} \quad \text{for } m = 1, 2, 3, \dots \\ \left. \frac{d^2Z}{dz^2} \right|_{z=0} &= E(k^2 + l^2) + F(k^2 + l^2) = 0 \implies F = -E \end{aligned}$$

thus

$$\begin{aligned} Z(z) &= E e^{\sqrt{k^2+l^2}z} - E e^{-\sqrt{k^2+l^2}z} \\ &= 2E \sinh(\sqrt{k^2+l^2}z) \\ &= 2E \sinh\left(\sqrt{\left(\frac{n\pi}{a}\right)^2 + \left(\frac{m\pi}{a}\right)^2}z\right) \\ &= 2E \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2}z\right) \end{aligned}$$

The potential is now

$$V(x, y, z) = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} C_{nm} \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2}z\right)$$

where we can solve for the Constants C_{nm} using the boundary condition (i) $V = V_0$ at $z = a$:

$$V_0 = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} C_{nm} \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2}a\right)$$

Doing the Fourier series by multiplying by $\sin\left(\frac{n'\pi x}{a}\right) \sin\left(\frac{m'\pi y}{a}\right)$ and integrating over $x, y = [0, a]$:

$$\begin{aligned} & \int_0^a \int_0^a V_0 \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) dx dy \\ &= \sum \sum C_{nm} \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2} a\right) \int_0^a \int_0^a \sin^2\left(\frac{n\pi x}{a}\right) \sin^2\left(\frac{m\pi y}{a}\right) dx dy \end{aligned}$$

where the integral

$$\int_0^a \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{n'\pi x}{a}\right) dx = \begin{cases} \frac{a}{2} & \text{if } n = n' \\ 0 & \text{if } n \neq n' \end{cases}$$

So

$$V_0 \int_0^a \int_0^a \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) dx dy = C_{nm} \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2} a\right) \frac{a^2}{4}$$

where

$$C_{nm} \sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2} a\right) = \begin{cases} 0 & \text{if } n \text{ or } m \text{ is even} \\ \frac{16V_0}{\pi^2 nm} & \text{if } n \text{ and } m \text{ are odd} \end{cases}$$

Thus the potential is

$$V(x, y, z) = \frac{16V_0}{\pi^2} \sum_{n \text{ odd}} \sum_{m \text{ odd}} \frac{1}{nm} \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{m\pi y}{a}\right) \frac{\sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2} z\right)}{\sinh\left(\pi \sqrt{n^2 + m^2}\right)}$$

The potential at the center is

$$V(a/2, a/2, a/2) = \frac{16V_0}{\pi^2} \sum_{n \text{ odd}} \sum_{m \text{ odd}} \frac{1}{nm} \sin\left(\frac{n\pi}{2}\right) \sin\left(\frac{m\pi}{2}\right) \frac{\sinh\left(\frac{\pi}{a} \sqrt{n^2 + m^2} a/2\right)}{\sinh\left(\pi \sqrt{n^2 + m^2}\right)}$$

where a calculator gives us

$$V(a/2, a/2, a/2) \approx 0.167V_0$$

Which roughly $V_0/6$ or a sixth of the potential in the center of a cube with V_0 on each face.

3.25 The potential inside and outside the sphere is given by (Griffiths 3.78 & 3.79)

$$V(r, \theta) = \begin{cases} \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta) & (r \leq R) \\ \sum_{l=0}^{\infty} \frac{B_l}{r^{l+1}} P_l(\cos \theta) & (r \geq R) \end{cases}$$

where (3.81 & 3.84)

$$B_l = A_l R^{2l+1} \quad (3.81)$$

$$A_l = \frac{1}{2\epsilon_0 R^{l-1}} \int_0^\pi \sigma_0(\theta) P_l(\cos \theta) \sin \theta d\theta \quad (3.84)$$

Since the charge density is constant $\sigma_0(\theta) = \sigma_0$ and using $u = \cos \theta; \cos(0) = 1, \cos(\pi) = -1$ and $du = -\sin \theta d\theta$ we get

$$A_l = -\frac{\sigma_0}{2\epsilon_0 R^{l-1}} \int_1^{-1} P_l(u) du$$

where $P_l(u)$ is odd for odd l and even for even l so the integral is zero for odd l and non-zero for even l :

$$\int_1^{-1} P_l(u) du = \begin{cases} 0 & \text{if } l \text{ is even} \\ -2 \int_0^1 P_l(u) du & \text{if } l \text{ is odd} \end{cases}$$

Thus

$$A_l = -\frac{\sigma_0}{2\epsilon_0 R^{l-1}} (-2) \int_0^1 P_l(u) du = \begin{cases} 0 & \text{if } l \text{ is odd} \\ \frac{\sigma_0}{\epsilon_0 R^{l-1}} \int_0^1 P_l(u) du & \text{if } l \text{ is even} \end{cases}$$

So for the odd P_l

$$\begin{aligned} \int_0^1 P_1(u) du &= \int_0^1 u du = \frac{1}{2} \\ \int_0^1 P_3(u) du &= \frac{1}{2} \int_0^1 (5u^3 - 3u) du = -\frac{1}{8} \\ \int_0^1 P_5(u) du &= \frac{1}{8} \int_0^1 (63u^5 - 70u^3 + 15u) du = \frac{1}{16} \end{aligned}$$

and for the even P_l the integral is zero; therefore

$$\begin{aligned} A_0 &= A_2 = A_4 = A_6 = 0 \\ A_1 &= \frac{\sigma_0}{\epsilon_0} \frac{1}{2}, \quad A_3 = \frac{\sigma_0}{\epsilon_0 R^2} \left(-\frac{1}{8} \right), \quad A_5 = \frac{\sigma_0}{\epsilon_0 R^4} \frac{1}{16} \end{aligned}$$

and

$$\begin{aligned} B_0 &= B_2 = B_4 = B_6 = 0 \\ B_1 &= \frac{\sigma_0}{\epsilon_0} R^3 \frac{1}{2}, \quad B_3 = \frac{\sigma_0}{\epsilon_0} R^5 \left(-\frac{1}{8} \right), \quad B_5 = \frac{\sigma_0}{\epsilon_0} R^7 \frac{1}{16} \end{aligned}$$

So finally the potential is

$$V(r, \theta) = \frac{\sigma_0}{\epsilon_0} \left[\frac{r}{2} P_1(\cos \theta) - \frac{r^3}{8R^2} P_3(\cos \theta) + \frac{r^5}{16R^4} P_5(\cos \theta) \right] \quad (r \leq R)$$

and

$$V(r, \theta) = \frac{\sigma_0}{\epsilon_0} \left[\frac{R^3}{2r^2} P_1(\cos \theta) - \frac{R^5}{8r^4} P_3(\cos \theta) + \frac{R^7}{16r^6} P_5(\cos \theta) \right] \quad (r \geq R)$$

3.29 Given charge density

$$\rho(r, \theta) = k \frac{R}{r^2} (R - 2r) \sin \theta$$

for a sphere of radius R and constant k . The monopole term is

$$\int \rho d\tau = kR \int \left(\frac{1}{r^2} R - 2r \sin \theta \right) r^2 \sin \theta dr d\theta d\phi$$

Where the radial integral is

$$\int_0^R R - 2r dr = 0$$

Moving on to the dipole term

$$\int r \cos \theta \rho d\tau = kR \int (r \cos \theta)(R - 2r) \sin \theta \sin \theta dr d\theta d\phi$$

where the polar integral is

$$\int_0^\pi \sin \theta \cos \theta d\theta = 0$$

Maybe the quadrupole term will be non-zero:

$$\begin{aligned} \int r^2 \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right) \rho d\tau &= \frac{1}{2} kR \int_0^{2\pi} \int_0^\pi \int_0^R r^2 ((3 \cos^2 \theta - 1)(R - 2r) \sin \theta) dr d\theta d\phi \\ &= \frac{1}{2} kR (2\pi) \left(-\frac{\pi}{8} \right) \left(-\frac{R^4}{6} \right) \\ &= \frac{\pi^2 k R^5}{48} \end{aligned}$$

So far from the sphere for points on the z -axis the potential is

$$V(z) = \frac{1}{4\pi\epsilon_0} \frac{\pi^2 K R^5}{48 z^3}$$

3.38 Given

$$\mathbf{E}_{\text{dip}}(r, \theta) = \frac{p}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \quad (3.103)$$

and

$$\mathbf{E}_{\text{dip}}(r, \theta) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} [3(\mathbf{p} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{p}] \quad (3.104)$$

The dipole moment has vector components

$$\mathbf{p} = p \cos \theta \hat{\mathbf{r}} - p \sin \theta \hat{\theta}$$

so now we can sub this into (3.104):

$$\begin{aligned} \mathbf{E}_{\text{dip}}(r, \theta) &= \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} [3(\mathbf{p} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{p}] \\ &= \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} 3(p \cos \theta) \hat{\mathbf{r}} - (p \cos \theta \hat{\mathbf{r}} - p \sin \theta \hat{\theta}) \\ \text{Eq. 3.103} &= \frac{p}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \end{aligned}$$

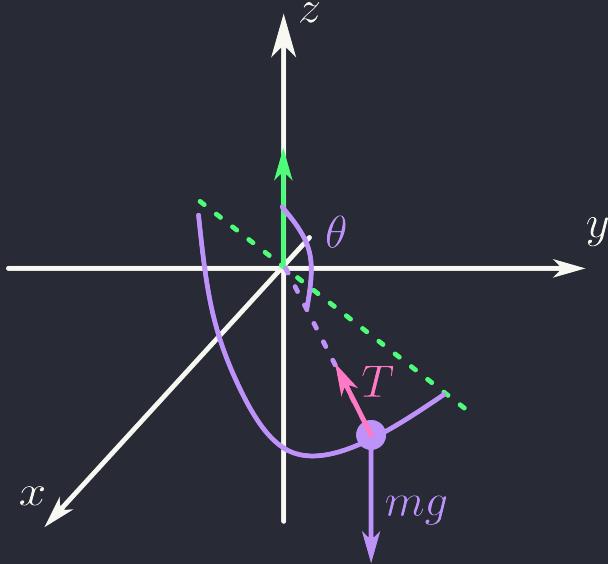


Figure 4.30: The electric field of a dipole at a distance r from the origin.

3.61 Given the potential from (3.103) and the force

$$\mathbf{F} = q\mathbf{E} = \frac{qp}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta})$$

Pretending this is a pendulum as shown in Fig. 4.30 the force is

$$\mathbf{F} = -T\hat{\mathbf{r}} - mg\hat{\mathbf{z}} = -T\hat{\mathbf{r}} - mg(\cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\theta})$$

T is the tension given by the centripetal acceleration

$$\begin{aligned} ma_c &= T - mg \cos(\pi - \theta) \\ m \frac{v^2}{r} &= T + mg \cos \theta \end{aligned}$$

where v is the velocity given by energy conservation

$$\begin{aligned} KE &= \Delta PE \\ \frac{1}{2}mv^2 &= -mgr \cos \theta \\ \implies v^2 &= -2gr \cos \theta \end{aligned}$$

So

$$T = m \frac{v^2}{r} - mg \cos \theta = -2mg \cos \theta - mg \cos \theta = -3mg \cos \theta$$

Thus we get the force

$$\begin{aligned} \mathbf{F} &= 3mg \cos \theta \hat{\mathbf{r}} - mg(\cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\theta}) \\ &= mg(2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \end{aligned}$$

Which is telling us that the electric charge will swing in an arc like a pendulum i.e.

$$mg \equiv \frac{qp}{4\pi\epsilon_0 r^3}$$

or we can think of the charge q as an inertial mass and group the other terms as the gravitational constant.

4.1 Given the Polarizability of Hydrogen atom

$$\frac{\alpha}{4\pi\epsilon_0} = 0.667 \times 10^{-30} \text{ m}^3$$

of bohr radius $a = 0.5 \times 10^{-10} \text{ m}$, and the E-field between metal plates separated $\Delta D = 1 \times 10^{-3} \text{ m}$ apart and connected to a 500 V battery is

$$E = \frac{V}{\Delta D} = \frac{500 \text{ V}}{1 \times 10^{-3} \text{ m}} = 5 \times 10^5 \text{ V/m}$$

The dipole moment is ($\epsilon_0 = 8.85 \times 10^{-12} \frac{\text{C}^2}{\text{N.m}^2}$)

$$p = \alpha E = qd \implies d = \frac{\alpha E}{q} = \frac{(4\pi\epsilon_0)0.667 \times 10^{-30} \text{ m}^3 \times 5 \times 10^5 \text{ V/m}}{1.6 \times 10^{-19} \text{ C}} = 2.32 \times 10^{-16} \text{ m}$$

So the separation distance d as a fraction of the atomic radius a is

$$\boxed{\frac{d}{a} = \frac{2.32 \times 10^{-16} \text{ m}}{0.5 \times 10^{-10} \text{ m}} = 4.6 \times 10^{-6}}$$

The estimate voltage we need to ionize the atom is

$$\begin{aligned} V &= E\Delta D, \quad E = \frac{qd}{\alpha} \\ \implies V &= \frac{qa}{\alpha} \Delta D = \frac{1.6 \times 10^{-19} \text{ C} \times 0.5 \times 10^{-10} \text{ m}}{(4\pi\epsilon_0)0.667 \times 10^{-30} \text{ m}^3} \times 1 \times 10^{-3} \text{ m} = \boxed{1.1 \times 10^8 \text{ V}} \end{aligned}$$

4.2 Given the charge density of a groundstate hydrogen atom

$$\rho(r) = \frac{q}{\pi a^3} e^{-2r/a}$$

where q is the electron charge and a is the Bohr radius, find the atomic polarizability.

First calculating the electric field of the electron cloud

$$E_e(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

where the enclosed charge Q is

$$\begin{aligned} Q &= \int \rho(r') d\tau, \quad d\tau = r^2 \sin\theta dr d\theta d\phi \\ &= \frac{q}{\pi a^3} (4\pi) \int_0^r r'^2 e^{-2r'/a} dr', \quad u = -\frac{2r}{a}; \quad du = -\frac{2}{a} dr \\ &= \frac{4q}{a^3} \left[-\frac{a^3}{8} \int u^2 e^u du \right] \\ &= -\frac{q}{2} \left[(u^2 - 2u + 2)e^u \right] \Big|_0^{-2r/a} \\ &= -\frac{q}{2} \left[\left(\frac{4r^2}{a^2} + \frac{4r}{a} + 2 \right) e^{-2r/a} - 2 \right] \\ &= q \left[1 - \left(\frac{2r^2}{a^2} + \frac{2r}{a} + 1 \right) e^{-2r/a} \right] \end{aligned}$$

So

$$E_e(r) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \left[1 - \left(\frac{2r^2}{a^2} + \frac{2r}{a} + 1 \right) e^{-2r/a} \right]$$

For $r \ll a$ we can Taylor expand $e^{-2r/a}$ since $-2r/a \approx 0$ to get

$$\begin{aligned} e^x &= 1 + x + \frac{x^2}{2} + \frac{x^3}{3!} + \dots \\ \Rightarrow e^{-2r/a} &= 1 - \frac{2r}{a} + \frac{2r^2}{a^2} - \frac{4r^3}{3a^3} + \dots \end{aligned}$$

so

$$\begin{aligned} 1 - \left(\frac{2r^2}{a^2} + \frac{2r}{a} + 1 \right) e^{-2r/a} &= 1 - \left(\frac{2r^2}{a^2} + \frac{2r}{a} + 1 \right) \left(1 - \frac{2r}{a} + \frac{2r^2}{a^2} - \frac{4r^3}{3a^3} \right) \\ &= 1 - \frac{2r^2}{a^2} + \frac{4r^3}{a^3} - \frac{4r^4}{a^4} + \frac{8r^5}{3a^5} \\ &\quad - \frac{2r}{a} + \frac{4r^2}{a^2} - \frac{4r^3}{a^3} + \frac{8r^4}{3a^4} \\ &\quad - 1 + \frac{2r}{a} - \frac{2r^2}{a^2} + \frac{4r^3}{3a^3} \\ &= \frac{4r^3}{3a^3} + \dots \end{aligned}$$

where the higher order terms are negligible. In an external electric field E , the polarized atom will have a balanced internal field $E = E_e$, such that the field at a distance d from the center of the cloud is

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{d^2} \left[\frac{4d^3}{3a^3} \right] = \frac{1}{3\pi\epsilon_0 a^3} (qd) = \frac{qd}{\alpha} \implies [\alpha = 3\pi\epsilon_0 a^3]$$

4.11 Given a short cylinder(radius a , length L) with “frozen-in” polarization \mathbf{P} parallel to the axis. Since \mathbf{P} is uniform, $\rho_b = 0$, but the surface bound charge density is

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}}$$

so the bound surface charge is $\sigma_b = P$ on one end and $\sigma_b = -P$ on the other end of the cylinder.

(i) E-field sketch for $L \gg a$:

(ii) E-field sketch for $L \ll a$:

(iii) E-field sketch for $L \approx a$:

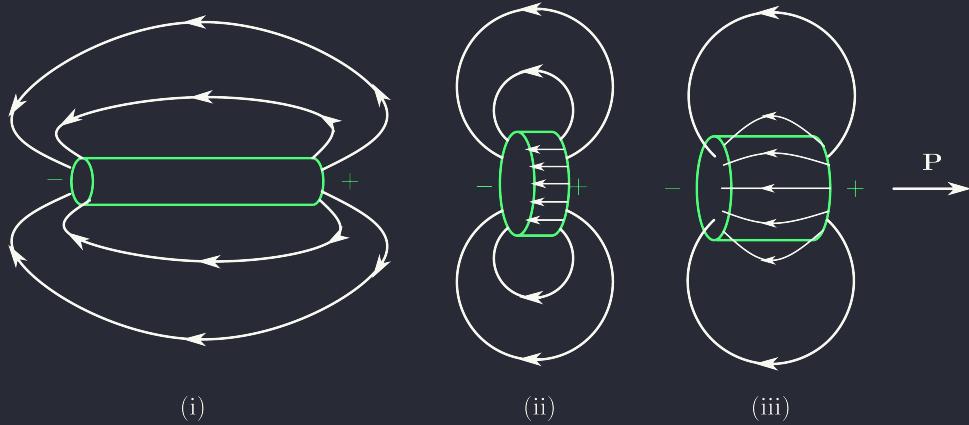


Figure 6.31: Sketch of E-field for a short cylinder

4.34 A dielectric cube of side a centered at the origin with a “frozen-in” polarization $\mathbf{P} = k\mathbf{r}$ where k is a constant and $\mathbf{r} = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$.

First the volume bound charge density is

$$\rho_b = -\nabla \cdot \mathbf{P} = -\nabla \cdot (k\mathbf{r}) = -k\left(\frac{\partial}{\partial x}x + \frac{\partial}{\partial y}y + \frac{\partial}{\partial z}z\right) = -3k$$

and the surface bound charge density is, e.g. on the face $\hat{\mathbf{n}} = \hat{\mathbf{x}}$, $x = a/2$,

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = \frac{ka}{2}$$

which is the same for all six faces. The total bound charge is

$$Q = \int_V \rho_b d\tau + \oint_S \sigma_b da = -3ka^3 + \frac{ka}{2}(6a^2) = 0$$

where the volume integral is just the volume of the cube a^3 and the surface integral is the sum of the areas of the six faces $6a^2$.

4.15 Thick spherical shell of inner radius a , outer radius b , with polarization

$$\mathbf{P}(\mathbf{r}) = \frac{k}{r} \hat{\mathbf{r}}$$

(a) E-field in all regions using bound charges: The volume bound charge is

$$\rho_b = -\nabla \cdot \mathbf{P} = -\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{k}{r} \right) = -\frac{k}{r^2}$$

and the surface bound charges are ($\hat{\mathbf{n}} = -\hat{\mathbf{r}}$ at $r = a$)

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = \begin{cases} -\frac{k}{a} & r = a \\ \frac{k}{b} & r = b \end{cases}$$

Using Gauss's Law

$$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

(i) $r < a$: $Q_{\text{enc}} = 0$, so $\mathbf{E} = 0$

(ii) $a < r < b$: The enclosed charge is the inner surface charge plus the volume charge:

$$\begin{aligned} Q_{\text{enc}} &= \oint_S \sigma_b d\mathbf{a} + \int_V \rho_b d\tau \\ &= \int -\frac{k}{a} a^2 \sin \theta d\theta d\phi + \int -\frac{k}{r^2} r^2 \sin \theta dr d\theta d\phi \\ &= -4\pi k a - 4\pi k(r - a) = -4\pi kr \end{aligned}$$

So using Gauss's Law

$$\begin{aligned} \oint \mathbf{E} \cdot d\mathbf{a} &= \frac{Q_{\text{enc}}}{\epsilon_0} \\ |\mathbf{E}| \oint da &= -\frac{4\pi kr}{\epsilon_0} \\ |\mathbf{E}| 4\pi r^2 &= -\frac{4\pi kr}{\epsilon_0} \\ |\mathbf{E}| &= -\frac{k}{\epsilon_0 r} \end{aligned}$$

or

$$\mathbf{E} = -\frac{k}{\epsilon_0 r} \hat{\mathbf{r}}$$

(iii) $r > b$: The total enclosed charge of a dielectric is zero (from last HW 4.14), so $\mathbf{E} = 0$

(b) Using

$$\oint D \cdot d\mathbf{a} = Q_{\text{free}} \quad (4.23)$$

and

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} \quad (4.21)$$

the total free enclosed charge is zero, so

$$\mathbf{D} = 0$$

4.17 Bar electret from Prob. 4.11 has $\rho_b = 0$: From divergence theorem

$$\int_V (\nabla \cdot \mathbf{D}) d\tau = \oint_S \mathbf{D} \cdot d\mathbf{a} = Q_{\text{free}} = 0 \implies \nabla \cdot \mathbf{D} = 0 \quad (4.22)$$

So, the field lines for \mathbf{D} are closed loops as shown in Fig. 7.32.

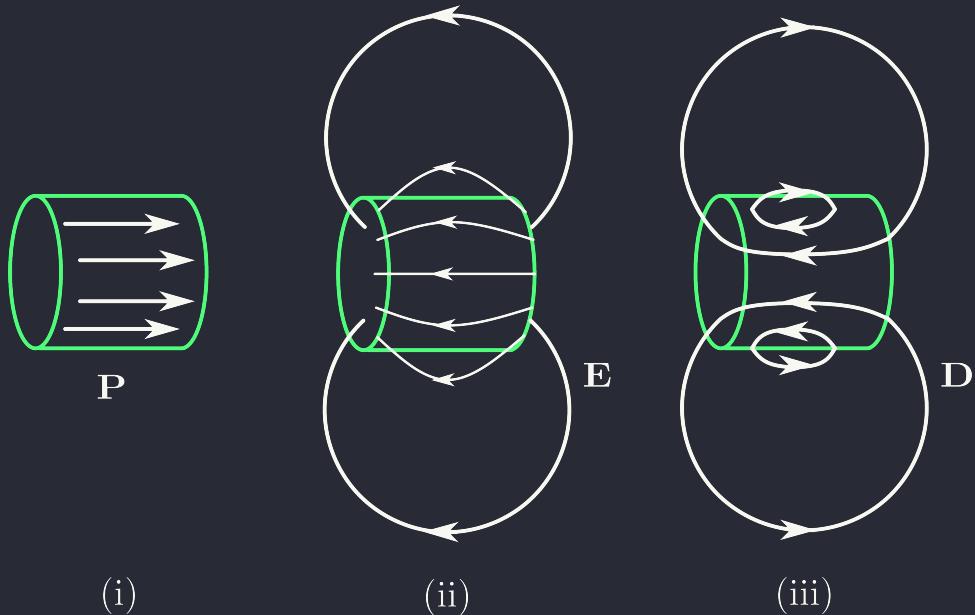


Figure 7.32: Field lines for (i) \mathbf{P} (ii) \mathbf{E} (iii) \mathbf{D}

4.19 For a parallel plate capacitor, the E-field in the air between the plates is two times the E-field of a single plate

(a)

$$\mathbf{E}_{\text{air}} = 2\mathbf{E}_{\text{plate}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$

where $\sigma = \pm Q/A$ is the surface charge density on the plates. The displacement field is

$$\begin{aligned} \oint D \cdot d\mathbf{a} &= Q_{\text{free}} \\ D(A) &= \sigma A \\ \implies D &= \sigma \end{aligned}$$

In the linear dielectric medium, \mathbf{D} is proportional to \mathbf{E} by

$$\mathbf{D} = \epsilon \mathbf{E} \quad (4.32)$$

So

$$\implies E = \frac{D}{\epsilon} = \frac{\sigma}{\epsilon}$$

and the potential difference between the plates is the sum of the potential differences in the air and the dielectric

$$\begin{aligned} V &= Ed \\ &= E_{\text{air}}d_{\text{air}} + E_{\text{die}}d_{\text{die}} \\ &= \frac{\sigma}{\epsilon_0} \frac{d}{2} + \frac{\sigma}{\epsilon} \frac{d}{2} \quad \text{using } \sigma = \frac{Q}{A} \\ &= \frac{Qd}{2A\epsilon_0} \left[1 + \frac{\epsilon_0}{\epsilon} \right] \end{aligned}$$

Finally the capacitance of the half filled capacitor (a) is

$$C_a = \frac{Q}{V} = \frac{2A\epsilon_0}{d} \frac{1}{1 + \frac{\epsilon_0}{\epsilon}}$$

Compared to the capacitance of a fully filled capacitor

$$C = \frac{A\epsilon_0}{d} \quad (2.54)$$

So the increase in capacitance is

$$\frac{C_a}{C} = \frac{2}{1 + \frac{\epsilon_0}{\epsilon}}$$

Using the relative permittivity $\epsilon_r = \epsilon/\epsilon_0$, the capacitance of the half filled capacitor (a) increases by a factor of

$$\boxed{\frac{C_a}{C} = \frac{2\epsilon_r}{1 + \epsilon_r}}$$

(i) In air: the E-field increases by the same factor so for a given potential $V = Ed$,

$$\boxed{\mathbf{E}_{\text{air}} = \frac{2\epsilon_r}{1 + \epsilon_r} \frac{V}{d} \hat{\mathbf{n}}}$$

And since $[\mathbf{P}_{\text{air}} = 0]$ in air, the displacement field $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$ is

$$\boxed{\mathbf{D}_{\text{air}} = \frac{2\epsilon_r}{1 + \epsilon_r} \epsilon_0 \frac{V}{d} \hat{\mathbf{n}}}$$

- (ii) In the dielectric: The E-field in the dielectric is proportional to the E-field in the air (vacuum) by

$$\mathbf{E} = \frac{1}{\epsilon} \mathbf{D} = \frac{1}{\epsilon_r} \mathbf{E}_{\text{vac}} \quad (4.35)$$

so

$$\boxed{\mathbf{E}_{\text{die}} = \frac{2}{1 + \epsilon_r} \frac{V}{d} \hat{\mathbf{n}}}$$

and from (4.35)

$$\begin{aligned} \mathbf{D} &= \epsilon \mathbf{E}, \quad \epsilon = \epsilon_r \epsilon_0 \\ \implies \boxed{\mathbf{D}_{\text{die}} &= \frac{2\epsilon_r}{1 + \epsilon_r} \epsilon_0 \frac{V}{d} \hat{\mathbf{n}} = \mathbf{D}_a} \end{aligned}$$

Finally, we can find the polarization directly from either (4.21) or

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E} = \epsilon_0 (\epsilon_r - 1) \mathbf{E} \quad (4.30)$$

so we might as well use (4.30) and get the polarization in the dielectric as

$$\boxed{\mathbf{P}_{\text{die}} = \frac{2}{1 + \epsilon_r} \epsilon_0 (\epsilon_r - 1) \frac{V}{d} \hat{\mathbf{n}}}$$

- (iii) The free bound surface charge is simply $\sigma_f = D$ so for the top plate

$$\boxed{\sigma_{f_{\text{top}}} = \frac{2\epsilon_r}{1 + \epsilon_r} \epsilon_0 \frac{V}{d}}$$

and $\sigma_{f_{\text{bot}}} = -\sigma_f$ for the bottom plate. And using $\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}}$ where $\hat{\mathbf{n}}$ points in the negative direction for the top plate, so

$$\boxed{\sigma_{b_{\text{top}}} = -\frac{2}{1 + \epsilon_r} \epsilon_0 (\epsilon_r - 1) \frac{V}{d} \quad \sigma_{b_{\text{bot}}} = -\sigma_{b_{\text{top}}}}$$

- (b) Again, from Gauss's law

$$\begin{aligned} E_{\text{air}} &= \frac{\sigma_{f_{\text{air}}}}{\epsilon_0} \quad \text{and} \quad V = E_{\text{air}} d \\ \implies \boxed{\sigma_{f_{\text{air}}} &= \frac{\epsilon_0 V}{d}} \end{aligned}$$

for the top plate ($-\sigma_{f_{\text{air}}}$ in the bottom plate) but in the dielectric, using \mathbf{D}

$$\begin{aligned} \oint \mathbf{D} \cdot d\mathbf{a} &= Q_{\text{free}} \\ \implies D &= \sigma_{f_{\text{die}}} \end{aligned}$$

and from $D = \epsilon E_{\text{die}}$

$$E_{\text{die}} = \frac{\sigma_{f_{\text{die}}}}{\epsilon} = \frac{\sigma_f}{\epsilon_r \epsilon_0}$$

Then using the potential difference $V = E_{\text{die}} d$

$$\implies \boxed{\sigma_{f_{\text{die}}} = \epsilon_r \epsilon_0 \frac{V}{d}}$$

for the top plate (negative for bottom). From this, we can find the capacitance $C = Q/V$ where

$$\begin{aligned} Q &= \sigma_{f_{\text{air}}} \frac{A}{2} + \sigma_{f_{\text{die}}} \frac{A}{2} \\ &= \frac{\epsilon_0 V A}{2d} + \frac{\epsilon_r \epsilon_0 V A}{2d} \end{aligned}$$

so

$$\Rightarrow C = \frac{A \epsilon_0}{2d} (1 + \epsilon_r)$$

From (2.54), the capacitance increases by a factor of

$$\boxed{\frac{C_b}{C} = \frac{1 + \epsilon_r}{2}}$$

- (i) The E-field in the air is just the E-field between two plates

$$\boxed{\mathbf{E}_{\text{air}} = \frac{V}{d} \hat{\mathbf{n}}}$$

and $\boxed{\mathbf{P}_{\text{air}} = 0}$ for air thus the displacement field (4.21) $\mathbf{D} = \epsilon_0 \mathbf{E} + 0$, or

$$\boxed{\mathbf{D}_{\text{air}} = \epsilon_0 \frac{V}{d} \hat{\mathbf{n}}}$$

- (ii) The E-field in the dielectric using (4.32):

$$\mathbf{E}_{\text{die}} = \frac{1}{\epsilon} \mathbf{D}$$

where $D = \sigma_{f_{\text{die}}}$ so

$$\boxed{\mathbf{D} = \epsilon_r \epsilon_0 \frac{V}{d} \hat{\mathbf{n}}}$$

Then we can use $\epsilon = \epsilon_r \epsilon_0$ so

$$\boxed{\mathbf{E}_{\text{die}} = \frac{V}{d} \hat{\mathbf{n}} = \mathbf{E}_{\text{air}}}$$

Now the polarization is given by (4.30)

$$\boxed{\mathbf{P}_{\text{die}} = \epsilon_0 (\epsilon_r - 1) \frac{V}{d} \hat{\mathbf{n}}}$$

Finally, the bound charge is $\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}}$ so

$$\boxed{\sigma_{b_{\text{top}}} = -\epsilon_0 (\epsilon_r - 1) \frac{V}{d} \quad \sigma_{b_{\text{bot}}} = -\sigma_{b_{\text{top}}}}$$

4.37 Between two linear dielectric, show that the E-field lines bend by

$$\frac{\tan \theta_2}{\tan \theta_1} = \frac{\epsilon_2}{\epsilon_1} \quad (4.68)$$

Since $\sigma_f = 0$ so from the boundary conditions

$$D_1^\perp - D_2^\perp = \sigma_b = 0 \implies D_1^\perp = D_2^\perp$$

or from (4.32)

$$D_{y1} = D_{y2} \implies \epsilon_1 E_{y1} = \epsilon_2 E_{y2}$$

Also from the boundary conditions, the parallel components of \mathbf{E} are continuous so

$$E_{x1} = E_{x2} \implies \tan \theta_1 = \frac{E_{y1}}{E_{x1}}, \quad \text{and} \quad \tan \theta_2 = \frac{E_{y2}}{E_{x2}}$$

Thus

$$\begin{aligned} \frac{\tan \theta_2}{\tan \theta_1} &= \frac{E_{y2}/E_{x2}}{E_{y1}/E_{x1}} \\ &= \frac{E_{y2}}{E_{y1}} \quad \text{using} \quad E_{y1} = \frac{\epsilon_2}{\epsilon_1} E_{y2} \\ \frac{\tan \theta_2}{\tan \theta_1} &= \frac{\epsilon_2}{\epsilon_1} \end{aligned}$$

From Snell's Law

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

where a convex lens has $n_2 > n_1$ which makes the angle of refraction smaller than the angle of incidence $\theta_2 < \theta_1$ thus the light ray bends towards the normal or “focuses” the light. For the dielectric interface, its convex lens

$$\epsilon_2 > \epsilon_1 \implies \tan \theta_2 > \tan \theta_1$$

So $\theta_2 > \theta_1$, thus the E-field lines bend away from the normal which will “defocus” the E-field.

Homework 8

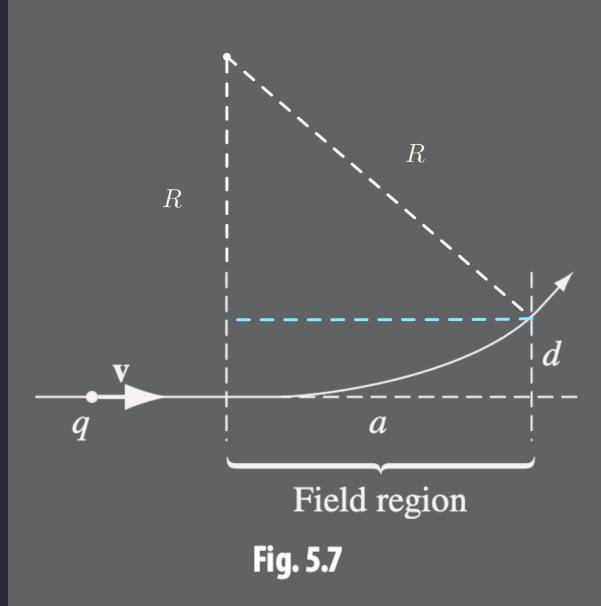


Figure 7.33: Extended geometry of particle path

5.1 From Griffiths, the momentum of the particle is given by the cyclotron formula

$$p = qBR \quad (5.3)$$

So from Fig. 7.33, the radius of the path is given by

$$\begin{aligned} R^2 &= (R - d)^2 + a^2 \\ R^2 &= R^2 - 2Rd + d^2 + a^2 \\ \implies R &= \frac{d^2 + a^2}{2d} \end{aligned}$$

Thus

$$p = qB \frac{d^2 + a^2}{2d}$$

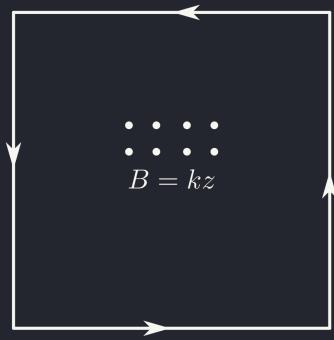


Figure 7.34: View from yz plane: Magnetic field points out of the page, on the

5.4 Given the magnetic field

$$\mathbf{B} = kz\hat{\mathbf{x}}$$

Using RHR:

- The left side of the loop has a force pointing to the left ($-\hat{\mathbf{y}}$) which cancels out with the
- Right side of the loop with force pointing to the right ($\hat{\mathbf{y}}$)
- The top ($z = a/2$) and bottom ($z = -a/2$) side of the loop have forces pointing in the same direction ($\hat{\mathbf{z}}$)

So the two forces are

$$F_{\text{top}} = 2IBa = I(ka/2)a = \frac{1}{2}Ika^2, \quad F_{\text{bottom}} = -\frac{1}{2}Ika^2$$

and since they are in the same direction, the net force is

$$\mathbf{F} = Ika^2\hat{\mathbf{z}}$$

5.5 For a wire of radius a

- (a) If the current is uniformly distributed over the surface, the current density is

$$K = \frac{dI}{dl_{\perp}} = \frac{I}{2\pi a}$$

since the width of the ribbon dl_{\perp} is the circumference of the wire cross-section.

- (b) If the volume current density is inversely proportional to the distance from the axis, we can use the result from (a):

$$J(s) = \frac{k}{s}$$

and integrating to find k :

$$\begin{aligned} I &= \int J(s) dA = k \int \frac{1}{s} (s ds d\phi) \\ &= k(2\pi) \int_0^a ds = 2\pi ka \\ \implies k &= \frac{I}{2\pi a} \end{aligned}$$

or the same result as (a). Thus

$$J(s) = \frac{I}{2\pi a s}$$

5.6

- (a) A phonograph with static electricity σ rotating at ω . The surface current density K at a distance r from the center is (using $v = r\omega$)

$$K = \sigma v = [\sigma r\omega]$$

- (b) A uniform sphere of radius R with charge Q spins around the z axis at a constnat angular velocity ω . The volume current density \mathbf{J} at any point in the sphere is (using $v = r\omega \sin \phi \hat{\phi}$)

$$\mathbf{J} = \rho \mathbf{v} = \rho r \omega \sin \phi \hat{\phi}$$

where $\rho = \frac{Q}{4\pi R^3/3}$ is the charge per volume of the sphere, so

$$\boxed{\mathbf{J} = \frac{Q}{4\pi R^3} r \omega \sin \phi \hat{\phi}}$$

5.11 For a tightly wound solenoid with n turns per unit length and radius a , we start with a single loop (From Griffiths Example 5.6)

$$\begin{aligned} B_{\text{loop}} &= \frac{\mu_0 I}{4\pi} \int \frac{dl'}{z^2} \cos \theta = \frac{\mu_0 I}{4\pi} \left(\frac{\cos \theta}{z} \right) 2\pi a \\ &= \frac{\mu_0 I}{2} \frac{a^2}{(a^2 + z^2)^{3/2}} \end{aligned}$$

so for n turns per unit length we multiply by n and integrate over the length of the solenoid

$$B = \frac{\mu_0 n I}{2} \int \frac{a^2}{(a^2 + z^2)^{3/2}} dz$$

Using $\tan \theta = a/z \implies z = a/\tan \theta$ and $dz = -\frac{a}{\sin^2 \theta} d\theta$, the integral becomes

$$\begin{aligned} \int \frac{a^2}{(a^2 + (a/\tan \theta)^2)^{3/2}} \left(-\frac{a}{\sin^2 \theta} \right) d\theta &= - \int \frac{a^2}{(a^2)^{3/2} (1 + \frac{\cos^2 \theta}{\sin^2 \theta})^{3/2}} \left(\frac{a}{\sin^2 \theta} \right) d\theta \\ &= - \int \frac{1}{(\sin^2 \theta)^{3/2} (\sin^2 \theta + \cos^2 \theta)^{3/2} \sin^2 \theta} d\theta \\ &= - \int \frac{\sin^3 \theta}{\sin^2 \theta} d\theta \\ &= - \int \sin \theta d\theta = \cos \theta \Big|_{\theta_1}^{\theta_2} = \cos \theta_2 - \cos \theta_1 \end{aligned}$$

So

$$B = \frac{\mu_0 n I}{2} (\cos \theta_2 - \cos \theta_1)$$

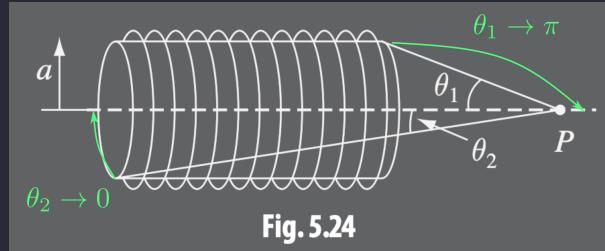


Fig. 5.24

Figure 7.35: As solenoid becomes infinitely long, $\theta_2 \rightarrow 0$ and $\theta_1 \rightarrow \pi$

As the solenoid becomes infinitely long (Fig. 7.35), then $\theta_2 \rightarrow 0$ and $\theta_1 \rightarrow \pi$ so

$$B = \frac{\mu_0 n I}{2} (1 - (-1)) = \boxed{\mu_0 n I}$$

Homework 9

5.14 For a steady current I flowing down a long cylindrical wire of radius a the magnetic field inside and out the wire

- (a) Given a uniform surface current

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I_{\text{enc}} \implies B(2\pi s) = \mu_0 I \pi$$

where s is the distance from the axis of the wire. Thus

$$\boxed{B = \frac{\mu_0 I}{2s} \quad s > a, \quad \text{and} \quad B = 0 \quad s < a}$$

which is points in the $\hat{\phi}$ direction.

- (b) For $\mathbf{J} \propto s$

$$J = ks \implies I = \int \mathbf{J} \cdot d\mathbf{a} = \int_0^a (ks) 2\pi s ds = \frac{2\pi k a^3}{3}$$

So the constant k is

$$k = \frac{3I}{2\pi a^3}$$

Outside the wire $I_{\text{enc}} = I$ so

$$\boxed{\mathbf{B} = \frac{\mu_0 I_{\text{enc}}}{2\pi s} \hat{\phi} = \frac{\mu_0 I}{2\pi s} \hat{\phi} \quad s > a}$$

But for $s < a$ the current only a fraction of the current is enclosed:

$$I_{\text{enc}} = \int_0^s ks' (2\pi s') ds' = \frac{2\pi k s^3}{3} = \frac{2\pi s^3}{3} \left(\frac{3I}{2\pi a^3} \right) = I \frac{s^3}{a^3}$$

Therefore

$$\boxed{\mathbf{B} = \frac{\mu_0 I}{2\pi s} \left(\frac{s^3}{a^3} \right) \hat{\phi} = \frac{\mu_0 I s^2}{2\pi a^3} \hat{\phi} \quad s < a}$$

5.18 A parallel plate capacitor with uniform surface charge $\pm\sigma$ moves at constant speed v .

- (a) Between the plates, the magnetic field is (using a rectangular Amperian loop perpendicular to the moving plate from Example 5.8)

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I_{\text{enc}} \implies B(2l) = \mu_0 K l$$

so at each region of a single plate

$$B = \pm \frac{\mu_0 K}{2} = \pm \frac{\mu_0 \sigma v}{2}$$

Above and below the plate, the magnetic field is zero

$B = 0 \quad \text{above and below the plate}$

and between the plates

$B = \mu_0 K = \mu_0 \sigma v \quad \text{between the plates}$

- (b) The magnetic force per unit area on the upper plate is

$$\begin{aligned} \frac{F_{\text{upper}}}{A} &= \frac{1}{A} \int K \times B dA \\ &= K \times B \\ &= (\sigma v) \times \frac{\mu_0 \sigma v}{2} \\ &= \boxed{\frac{\mu_0 \sigma^2 v^2}{2}} \end{aligned}$$

- (c) The speed required to balance the electrical force: From electrostatics the force on the upper plate is

$$F = \sigma E = \sigma \frac{\sigma}{2\epsilon_0} = \frac{\sigma^2}{2\epsilon_0}$$

so for the forces to balance

$$\frac{\mu_0 \sigma^2 v^2}{2} = \frac{\sigma^2}{2\epsilon_0} \implies v = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

where the values are related to the speed of light c by (From Wikipedia)

$$\epsilon_0 \mu_0 = \frac{1}{c^2} \implies v = c$$

So the force will balance at the speed of light c .

5.26 What current density would produce the vector potential $A = k\hat{\phi}$ in cylindrical coordinates?
Finding $\mathbf{B} \rightarrow \mathbf{J}$:

$$\mathbf{B} = \nabla \times A = \frac{1}{s} \begin{vmatrix} \hat{s} & s\hat{\phi} & \hat{z} \\ \frac{\partial}{\partial s} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ 0 & s(k) & 0 \end{vmatrix} = \frac{1}{s} \frac{\partial}{\partial s} (sk) \hat{z} = \frac{k}{s} \hat{z}$$

So using Ampère's Law

$$\mathbf{J} = \frac{1}{\mu_0} (\nabla \times \mathbf{B}) = \frac{1}{\mu_0 s} \begin{vmatrix} \hat{s} & s\hat{\phi} & \hat{z} \\ \frac{\partial}{\partial s} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ 0 & 0 & k/s \end{vmatrix} = -\frac{1}{\mu_0 s} \left(\frac{\partial}{\partial z} \frac{k}{s} \right) \hat{\phi} = \boxed{\frac{k}{\mu_0 s^2} \hat{\phi}}$$

5.32 To find the magnetic field inside a solid sphere of uniform charge density ρ , radius R , and angular velocity ω :

From Example 5.11 we just change $\sigma \rightarrow \rho$ so

$$\mathbf{A}(r, \theta, \phi) = \begin{cases} \frac{1}{3}\mu_0 R' \rho \omega r \sin \theta \hat{\phi} & \text{Inside sphere} \\ \frac{1}{3r^2} \mu_0 R'^4 \rho \omega \sin \theta \hat{\phi} & \text{Outside sphere} \end{cases}$$

where we integrate R' from $0 \rightarrow r$ to find the potential $A_{\phi_{\text{out}}}$ and $r \rightarrow R$ to find $A_{\phi_{\text{in}}}$: The potential is then the sum of the two

$$\begin{aligned} \mathbf{A} &= \frac{1}{3} \hat{\phi} \mu_0 \rho \omega r \sin \theta \int_r^R R' dR' + \frac{1}{3r^2} \mu_0 \rho \omega \sin \theta \int_0^r R'^4 dR' \\ &= \frac{1}{3} \mu_0 \rho \omega \sin \theta \left[\frac{1}{2}(R^2 - r^2) + \frac{1}{r^3} \frac{r^5}{5} \right] \hat{\phi} \\ &= \mu_0 \rho \omega r \sin \theta \left[\frac{R^2}{6} - \frac{r^2}{6} + \frac{r^2}{15} \right] \hat{\phi} \\ &= \mu_0 \rho \omega r \sin \theta \left[\frac{R^2}{6} - \frac{r^2}{10} \right] \hat{\phi} \end{aligned}$$

And using the curl in spherical coordinates

$$\begin{aligned} \nabla \times A &= \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{\mathbf{r}} & r \hat{\theta} & r \sin \theta \hat{\phi} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ 0 & 0 & (r \sin \theta) A_\phi \end{vmatrix} \\ &= \frac{1}{r^2 \sin \theta} \left[\hat{\mathbf{r}} \frac{\partial}{\partial \theta} (r \sin \theta A_\phi) - r \hat{\theta} \frac{\partial}{\partial r} (r \sin \theta A_\phi) \right] \\ &= \frac{1}{r^2 \sin \theta} \left[\hat{\mathbf{r}} \frac{\partial}{\partial \theta} \left(\mu_0 \rho \omega r^2 \sin^2 \theta \left[\frac{R^2}{6} - \frac{r^2}{10} \right] \right) - r \hat{\theta} \frac{\partial}{\partial r} \left(\mu_0 \rho \omega \sin^2 \theta \left[\frac{r^2 R^2}{6} - \frac{r^4}{10} \right] \right) \right] \\ &= \frac{1}{r^2 \sin \theta} \left[\hat{\mathbf{r}} \mu_0 \rho \omega r^2 (2 \sin \theta \cos \theta) \left[\frac{R^2}{6} - \frac{r^2}{10} \right] - r \hat{\theta} \mu_0 \rho \omega \sin^2 \theta \left[\frac{2r R^2}{6} - \frac{4r^3}{10} \right] \right] \\ \mathbf{B} &= \mu_0 \rho \omega \left[\hat{\mathbf{r}} \cos \theta \left(\frac{R^2}{3} - \frac{r^2}{5} \right) - \hat{\theta} \sin \theta \left(\frac{R^2}{3} - \frac{2r^2}{5} \right) \right] \end{aligned}$$

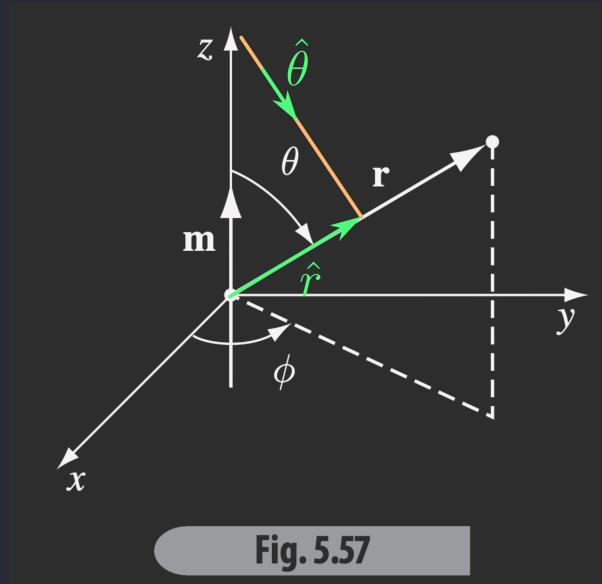


Fig. 5.57

5.36 Show that the magnetic dipole can be written in coordinate free form:

$$\mathbf{B}_{\text{dip}}(\mathbf{r}) = \frac{\mu_0}{4\pi} \frac{1}{r^3} [3(\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m}] \quad (5.89)$$

From the Figure we can see that

$$\mathbf{m} \cdot \hat{\mathbf{r}} = m \cos \theta, \quad \mathbf{m} \cdot \hat{\theta} = -m \sin \theta$$

so

$$\mathbf{m} = (\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} + (\mathbf{m} \cdot \hat{\theta})\hat{\theta} = m \cos \theta \hat{\mathbf{r}} - m \sin \theta \hat{\theta}$$

Therefore

$$3(\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m} = 3m \cos \theta \hat{\mathbf{r}} - (m \cos \theta \hat{\mathbf{r}} - m \sin \theta \hat{\theta}) = 2m \cos \theta \hat{\mathbf{r}} + m \sin \theta \hat{\theta}$$

or from the textbook

$$\mathbf{B}_{\text{dip}}(\mathbf{r}) = \frac{\mu_0 m}{4\pi r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}) \quad (5.88)$$

5.39

- (a) A phonograph of radius R , surface charge σ rotating at ω , find its magnetic dipole moment: For a ring dr using $v = \omega r$ and $I = \sigma v dr$ we have

$$m \equiv I \int d\mathbf{A} = \sigma v (\pi r^2) dr = \pi \sigma \omega r^3 dr$$

and integrating over the entire disk

$$m = \pi \sigma \omega \int_0^R r^3 dr = \frac{\pi \sigma \omega R^4}{4}$$

- (b) Find the magnetic dipole moment of a spinning spherical shell: For one ring of width $R d\theta$ around the sphere where $v = C/T$. For circumference $C = 2\pi R \sin \theta$ and time $T = 2\pi/\omega$, the current is

$$I = \sigma v (R d\theta) = \sigma \frac{2\pi R \sin \theta}{2\pi/\omega} d\theta = \sigma \omega R^2 \sin \theta d\theta$$

Using the area around the ring $A = \pi(R \sin \theta)^2$, we can integrate over the entire sphere $\theta \in [0, \pi]$:

$$\begin{aligned} m = IA &= \int \pi \sigma \omega R^2 \sin \theta (R \sin \theta)^2 d\theta = \pi \sigma \omega R^4 \int_0^\pi \sin^3 \theta d\theta \\ &= \pi \sigma \omega R^4 \int (\sin^2 \theta) \sin \theta d\theta = \pi \sigma \omega R^4 \int (1 - \cos^2 \theta) \sin \theta d\theta \\ &\quad \text{using } \cos \theta = u \implies -\sin \theta du = d\theta \\ &= \pi \sigma \omega R^4 \int (u^2 - 1) du = \pi \sigma \omega R^4 \left[\frac{u^3}{3} - u \right]_1^{-1} = \boxed{\frac{4}{3} \pi \sigma \omega R^4} \end{aligned}$$

For $r > R$ the dipole term in the multipole expansion is

$$\mathbf{A}_{\text{dip}} = \frac{\mu_0}{4\pi} \frac{m \sin \theta}{r^2} \hat{\phi} = \frac{\mu_0}{4\pi} \frac{4}{3} \pi \sigma \omega R^4 \frac{\sin \theta}{r^2} \hat{\phi} = \boxed{\frac{1}{3r^2} \mu_0 \sigma \omega R^4 \sin \theta \hat{\phi}}$$

which is the same term from Problem 5.32 $\mathbf{A} = \mathbf{A}_{\text{dip}}$ thus a perfect dipole for points outside the sphere.

Homework 10

6.3 Force of attractive between two magnetic dipoles $\mathbf{m}_1, \mathbf{m}_2$ a distance r apart:

(a) Given

$$F = 2\pi IRB \cos \theta \quad (6.2)$$

and

$$m_2 = IA = \pi IR^2$$

The magnetic field due to m_1 is

$$\mathbf{B}_{\text{dip}} = \frac{\mu_0}{4\pi} \frac{1}{r^3} [3(\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m}] \quad (5.89)$$

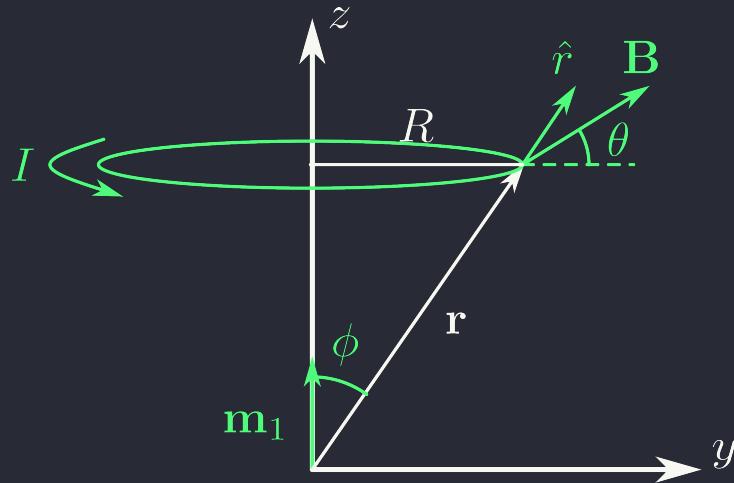


Figure 7.36: Magnetic field due to a magnetic dipole \mathbf{m}_1

$$\begin{aligned} B \cos \theta &= \mathbf{B}_1 \cdot \hat{\mathbf{y}} \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} [3(\mathbf{m}_1 \cdot \hat{\mathbf{r}})(\hat{\mathbf{r}} \cdot \hat{\mathbf{y}}) - \mathbf{m}_1 \cdot \hat{\mathbf{y}}] \end{aligned}$$

and from Fig. 7.36

$$\begin{aligned} \mathbf{m}_1 \cdot \hat{\mathbf{r}} &= m_1 \cos \phi, \quad \hat{\mathbf{r}} \cdot \hat{\mathbf{y}} = \sin \phi, \quad \hat{\mathbf{m}}_1 \cdot \hat{\mathbf{y}} = 0 \\ \text{and } \sin \phi &= \frac{R}{r}, \quad \cos \phi = \frac{z}{r} = \frac{\sqrt{r^2 - R^2}}{r} \end{aligned}$$

we get

$$\begin{aligned} B \cos \theta &= \frac{\mu_0}{4\pi} \frac{1}{r^3} [3m_1 \cos \phi \sin \phi - 0] \\ &= \frac{3\mu_0}{4\pi} \frac{m_1}{r^3} \sin \phi \cos \phi \\ &= \frac{3\mu_0}{4\pi} m_1 R \frac{\sqrt{r^2 - R^2}}{r^5} \end{aligned}$$

So

$$\begin{aligned} F &= 2\pi IRB \cos \theta \\ &= 2(\pi IR^2) \frac{3\mu_0}{4\pi} m_1 \frac{\sqrt{r^2 - R^2}}{r^5} = \frac{3\mu_0}{2\pi} m_1 m_2 \frac{\sqrt{r^2 - R^2}}{r^5} \end{aligned}$$

where $r \gg R$ for the dipole so

$$\frac{\sqrt{r^2 - R^2}}{r^5} \approx \frac{\sqrt{r^2}}{r^5} = \frac{1}{r^4}$$

Therefore

$$\boxed{F = \frac{3\mu_0}{2\pi} \frac{m_1 m_2}{r^4}}$$

(b) Using

$$\mathbf{F} = \nabla(\mathbf{m} \cdot \mathbf{B}) \quad (6.3)$$

and product rule (ii)

$$\begin{aligned} \mathbf{F} &= \nabla(\mathbf{m}_2 \cdot \mathbf{B}) \\ &= \mathbf{m}_2 \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{m}_2) + (\mathbf{m}_2 \cdot \nabla) \mathbf{B} + (\mathbf{B} \cdot \nabla) \mathbf{m}_2 \end{aligned}$$

where the curl of \mathbf{B} and \mathbf{m}_2 are zero

$$\implies \mathbf{F} = (\mathbf{m}_2 \cdot \nabla) \mathbf{B} + (\mathbf{B} \cdot \nabla) \mathbf{m}_2$$

and the second term is also zero for the dipole

$$\begin{aligned} \left(\frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z} \right) \mathbf{m}_2 &= 0 \\ \implies \mathbf{F} &= (\mathbf{m}_2 \cdot \nabla) \mathbf{B} \end{aligned}$$

So using (5.89)

$$\begin{aligned} \mathbf{F} &= \left[(0, 0, m_2) \cdot \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \right] \mathbf{B} \\ &= m_2 \frac{\partial}{\partial z} \left(\frac{\mu_0}{4\pi} \frac{1}{z^3} [3(\mathbf{m}_1 \cdot \hat{\mathbf{z}}) \hat{\mathbf{z}} - \mathbf{m}_1] \right) \\ &= \frac{\mu_0}{4\pi} m_2 \frac{\partial}{\partial z} \left(\frac{1}{z^3} [2m_1 \hat{\mathbf{z}}] \right) \\ &= \frac{\mu_0}{2\pi} m_1 m_2 \hat{\mathbf{z}} \left(-\frac{3}{z^4} \right) \end{aligned}$$

or when $z \rightarrow r$

$$\boxed{F = -\frac{3\mu_0}{2\pi} \frac{m_1 m_2}{r^4} \hat{\mathbf{z}}}$$

where the negative sign tells us the force is attractive.

6.5 A uniform current density $\mathbf{J} = J_0 \hat{\mathbf{z}}$ fills a slab on the yz plane from $x = -a \rightarrow +a$ and a magnetic dipole $\mathbf{m} = m_0 \hat{\mathbf{x}}$ placed at the origin:

- (a) The force on the dipole using (6.3): Using a rectangular Amperian loop on the xy plane i.e. $I_{\text{enc}} = AJ_0 = lxJ_0$:

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I_{\text{enc}}$$

$$Bl = \mu_0 lxJ_0 \implies B = \mu_0 J_0 x$$

where the direction is in the y from the right hand rule for Bio-Savart's law $\mathbf{J} \times \mathbf{r} = \hat{\mathbf{z}} \times \hat{\mathbf{x}} = \hat{\mathbf{y}}$;

$$\mathbf{B} = \mu_0 J_0 x \hat{\mathbf{y}}$$

So the force on the dipole is

$$\mathbf{F} = \nabla(\mathbf{m} \cdot \mathbf{B})$$

$$= \nabla(m_0 \mu_0 J_0 x [\hat{\mathbf{x}} \cdot \hat{\mathbf{y}}])$$

where $\hat{\mathbf{x}} \cdot \hat{\mathbf{y}} = 0$ so

$$\boxed{\mathbf{F} = 0}$$

- (b) With $\mathbf{m} = m_0 \hat{\mathbf{y}}$ the force on the dipole is

$$\mathbf{F} = \nabla(\mathbf{m} \cdot \mathbf{B})$$

$$= \nabla(m_0 \mu_0 J_0 x [\hat{\mathbf{y}} \cdot \hat{\mathbf{y}}])$$

$$= \hat{\mathbf{x}} \frac{\partial}{\partial x} (m_0 \mu_0 J_0 x)$$

which points in the x direction:

$$\boxed{\mathbf{F} = m_0 \mu_0 J_0 \hat{\mathbf{x}}}$$

- (c) For a polarization \mathbf{p} and E-field \mathbf{E} we can use product rule (ii) again:

$$F = \nabla(\mathbf{p} \cdot \mathbf{E})$$

$$= \mathbf{p} \times (\nabla \times \mathbf{E}) + \mathbf{E} \times (\nabla \times \mathbf{p}) + (\mathbf{p} \cdot \nabla)\mathbf{E} + (\mathbf{E} \cdot \nabla)\mathbf{p}$$

where $\nabla \times \mathbf{p} = 0$, $\nabla \times \mathbf{E} = 0$, and $\nabla \mathbf{p} = 0$ so

$$\mathbf{F} = (\mathbf{p} \cdot \nabla)\mathbf{E} = \nabla(\mathbf{p} \cdot \mathbf{E})$$

which is not the same as the magnetic case since Ampere's law states

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

Using product rule (ii) once more,

$$F = \nabla(\mathbf{m} \cdot \mathbf{B})$$

$$= \mathbf{m} \times (\nabla \times \mathbf{B}) + \cancel{\mathbf{B} \times (\nabla \times \mathbf{m})} + (\mathbf{m} \cdot \nabla)\mathbf{B} + \cancel{(\mathbf{B} \cdot \nabla)\mathbf{m}}$$

$$= \mathbf{m} \times (\mu_0 \mathbf{J}) + (\mathbf{m} \cdot \nabla)\mathbf{B}$$

so for configuration (a),

$$(\mathbf{m}_a \cdot \nabla)\mathbf{B} = \left(m_0 \hat{\mathbf{x}} \cdot \hat{\mathbf{x}} \frac{\partial}{\partial x} \right) \mu_0 J_0 x \hat{\mathbf{y}} = m_0 \mu_0 J_0 x \hat{\mathbf{y}}$$

and for configuration (b),

$$(\mathbf{m}_b \cdot \nabla)\mathbf{B} = \left(m_0 \hat{\mathbf{y}} \cdot \hat{\mathbf{y}} \frac{\partial}{\partial y} \right) \mu_0 J_0 x \hat{\mathbf{y}} = 0$$

6.9 The bound current for a short cylinder of radius a and length L with a frozen in uniform magnetization parallel to its axis $\mathbf{M} = M\hat{\mathbf{z}}$ is

$$\mathbf{K}_b = \mathbf{M} \times \hat{\mathbf{n}} = M\hat{\phi}$$

This is the same as the bar electret but with opposite directions for the (\mathbf{B} vs. \mathbf{D}) field produced

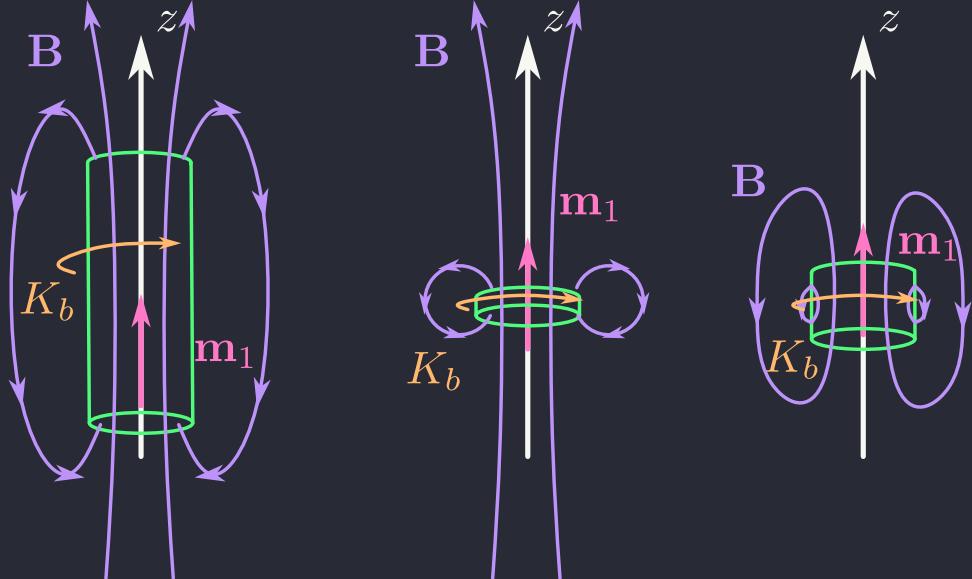


Figure 7.37: (left) $L \gg a$, (center) $L \ll a$, and (right) $L = a$

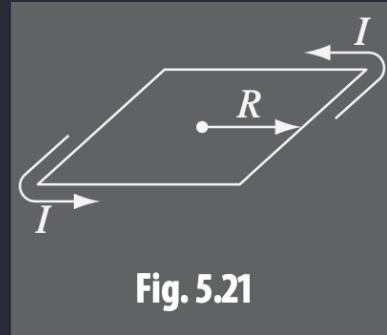
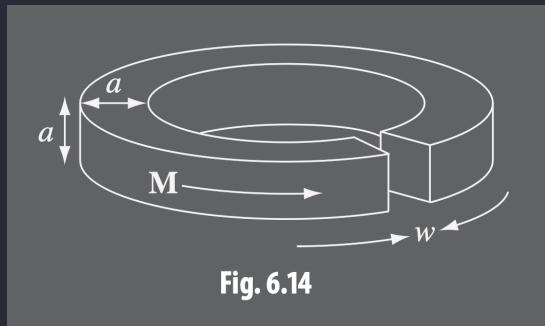


Figure 7.38: Superposition of torus and square loop with reversed current

6.10 For an iron rod of length L and square cross section of side length a with Magnetization \mathbf{M} and bent around in a circle with gap w : Assuming $w \ll a \ll L$ find the magnetic field at the center of the gap: The magnetization of the torus is $\mathbf{M} = M\hat{\phi}$ so the magnetic field of the torus is

$$\mathbf{B} = \mu_0 \mathbf{M}$$

The magnetic field at the center of the square loop is equivalent to four times the magnetic field of a straight wire (From Griffiths Example 5.5):

$$B = -4 \frac{\mu_0 I}{4\pi R} \sin \theta \Big|_{\theta=-\pi/4}^{\pi/4} = \frac{\mu_0 I}{\pi R} (\sqrt{2}/2 + \sqrt{2}/2) = \frac{\sqrt{2}\mu_0 I}{\pi R}$$

where the square loop has a current $I = m/a$ where $M \equiv m/V$

$$\implies I = \frac{MV}{a} = Mw$$

and $R = a/2$ so

$$B = -\frac{\sqrt{2}\mu_0 M w}{\pi a/2} = -\frac{2\sqrt{2}\mu_0 M w}{\pi a}$$

The magnetic field is a superposition of the torus and square loop with reversed current

$$\mathbf{B} = \mu_0 \mathbf{M} - \frac{2\sqrt{2}\mu_0 \mathbf{M} w}{\pi a}$$

6.25 For two charged q magnetic dipoles \mathbf{m} constrained to move on the z axis, they electrically repel and magnetically attract:

- (a) The equilibrium separation distance: From Coloumbs law the electric force is

$$\mathbf{F}_e = \frac{1}{4\pi\epsilon_0} \frac{q^2}{z^2} \hat{\mathbf{z}}$$

and the magnetic force is (From Prob. 6.3)

$$\mathbf{F}_m = \nabla(\mathbf{m} \cdot \mathbf{B}) = -\frac{3\mu_0}{2\pi} \frac{m^2}{z^4} \hat{\mathbf{z}}$$

so at equilibrium

$$\mathbf{F}_e = -\mathbf{F}_m \implies \frac{1}{4\pi\epsilon_0} \frac{q^2}{z^2} = \frac{3\mu_0}{2\pi} \frac{m^2}{z^4} \implies z = \frac{m\sqrt{6\epsilon_0\mu_0}}{q}$$

where we can use the fact that $\epsilon_0\mu_0 = 1/c^2$ to rewrite the equilibrium separation distance as

$$\boxed{z = \sqrt{6} \frac{m}{qc}}$$

- (b) For two electrons, we can use the magnetic moment of an electron, i.e, the “Bohr magneton” [Wiki](#),

$$m = \frac{e\hbar}{2m_e}$$

withe with electron charge $q = e = 1.6 \times 10^{-19}$ C, mass $m_e = 9.11 \times 10^{-31}$ kg, reduced planck constant $\hbar = 1.05 \times 10^{-34}$ Js, and speed of light $c = 3 \times 10^8$ m/s:

$$\begin{aligned} z &= \sqrt{6} \frac{eh}{2m_e qc} \\ &= \sqrt{6} \frac{\hbar}{2m_e c} \\ &= \sqrt{6} \frac{1.05 \times 10^{-34} \text{ Js}}{29.11 \times 10^{-31} \text{ kg} \cdot 3 \times 10^8 \text{ m/s}} \\ &= 4.72 \times 10^{-13} \text{ m} \end{aligned}$$

- (c) Since the equilibrium separation distance is much smaller than the Bohr radius $a_0 = 5.29 \times 10^{-11}$ m, there is *no stable bound state of two electrons*

Homework 11

6.20 To demagnetize a permanent magnet, we can heat it above its Curie temperature T_C

6.21

- (a) Show that the energy of the magnetic dipole in a magnetic field is

$$U = -\mathbf{m} \cdot \mathbf{B} \quad (6.34)$$

Moving the dipole from infinity to a point on the origin:

From the force on a magnetic dipole in a magnetic field $\mathbf{F} = \nabla(\mathbf{m} \cdot \mathbf{B})$, the work done is given by

$$\begin{aligned} U &= - \int_{\infty}^{\mathbf{r}} \mathbf{F} \cdot d\mathbf{l} \\ &= - \int_{\infty}^{\mathbf{r}} \nabla(\mathbf{m} \cdot \mathbf{B}) \cdot d\mathbf{l} \end{aligned}$$

using Fundamental theorem of gradients
 $= -[\mathbf{m} \cdot \mathbf{B}(\mathbf{r}) - \mathbf{m} \cdot \mathbf{B}(\infty)]$

Since the magnetic field goes to zero at infinity $\mathbf{B}(\infty) = 0$,

$$U = -\mathbf{m} \cdot \mathbf{B}$$

- (b) For two magnetic dipoles separated by \mathbf{r} , find the interaction energy:

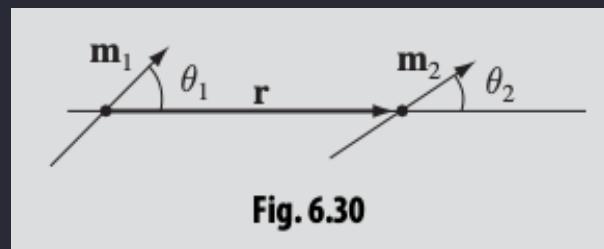
Given the coordinate-free form of the magnetic field of a dipole

$$\mathbf{B}_{\text{dip}}(\mathbf{r}) = \frac{\mu_0}{4\pi} \frac{1}{r^3} [3(\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m}]$$

The interaction energy of dipole \mathbf{m}_1 in the magnetic field of dipole 2 \mathbf{B}_2 is

$$\begin{aligned} U &= -\mathbf{m}_1 \cdot \mathbf{B}_2 \\ &= -\mathbf{m}_1 \cdot \frac{\mu_0}{4\pi} \frac{1}{r^3} [3(\mathbf{m}_2 \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m}_2] \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} [\mathbf{m}_1 \cdot \mathbf{m}_2 - 3(\mathbf{m}_1 \cdot \hat{\mathbf{r}})(\mathbf{m}_2 \cdot \hat{\mathbf{r}})] \end{aligned}$$

- (c) In the dipoles are left free to rotate, From the figure above,



$$\mathbf{m}_1 \cdot \mathbf{m}_2 = m_1 m_2 \cos(\theta_1 - \theta_2)$$

$$\mathbf{m}_1 \cdot \hat{\mathbf{r}} = m_1 \cos \theta_1$$

$$\mathbf{m}_2 \cdot \hat{\mathbf{r}} = m_2 \cos \theta_2$$

And using the cos difference formula

$$\cos(\theta_1 - \theta_2) = \cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2$$

The interaction energy is

$$\begin{aligned} U &= \frac{\mu_0}{4\pi} \frac{1}{r^3} [m_1 m_2 \cos(\theta_1 - \theta_2) - 3m_1 m_2 \cos \theta_1 \cos \theta_2] \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [(\cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2) - 3 \cos \theta_1 \cos \theta_2] \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [\sin \theta_1 \sin \theta_2 - 2 \cos \theta_1 \cos \theta_2] \end{aligned}$$

The stable configuration happens when

$$\frac{\partial U}{\partial \theta_1} = \frac{\partial U}{\partial \theta_2} = 0$$

So

$$\begin{aligned} \frac{\partial U}{\partial \theta_1} &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [\cos \theta_1 \sin \theta_2 + 2 \sin \theta_1 \cos \theta_2] = 0 \\ \implies 2 \sin \theta_1 \cos \theta_2 &= -\cos \theta_1 \sin \theta_2 \end{aligned}$$

and

$$\begin{aligned} \frac{\partial U}{\partial \theta_2} &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [\sin \theta_1 \cos \theta_2 + 2 \cos \theta_1 \sin \theta_2] = 0 \\ \implies 2 \cos \theta_1 \sin \theta_2 &= -\sin \theta_1 \cos \theta_2 \end{aligned}$$

Which implies that both

$$\sin \theta_1 \cos \theta_2 = \cos \theta_1 \sin \theta_2 = 0$$

or two sets of solutions

$$\sin \theta_1 = \sin \theta_2 = 0 \quad \text{or} \quad \cos \theta_1 = \cos \theta_2 = 0$$

For the first set, we can have either $\theta_1 = \theta_2 = 0$ (aligned horizontal dipoles $\rightarrow\rightarrow$) or $\theta_1 = 0$ and $\theta_2 = \pi$ (anti aligned horizontal dipoles $\rightarrow\leftarrow$), but for the anti aligned case, the magnetic field \mathbf{B}_2 will not be parallel to \mathbf{m}_1 so it is not a stable configuration.

For the second set, we can have either $\theta_1 = \theta_2 = \frac{\pi}{2}$ (aligned vertical dipoles $\uparrow\uparrow$) or $\theta_1 = \frac{\pi}{2}$ and $\theta_2 = -\frac{\pi}{2}$ (anti aligned vertical dipoles $\uparrow\downarrow$), but for the aligned case, the magnetic field \mathbf{B}_2 will not be parallel to \mathbf{m}_1 so it is not a stable configuration.

Therefore, looking at the two stable configurations:

$$\begin{aligned} U_{\rightarrow\rightarrow} &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [\sin 0 \sin 0 - 2 \cos 0 \cos 0] \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [-2] \\ U_{\uparrow\downarrow} &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 \left[\sin \frac{\pi}{2} \sin \frac{-\pi}{2} - 2 \cos \frac{\pi}{2} \cos \frac{-\pi}{2} \right] \\ &= \frac{\mu_0}{4\pi} \frac{1}{r^3} m_1 m_2 [-1] \end{aligned}$$

So the lower energy is the more stable configuration

$$U_{\rightarrow\rightarrow} < U_{\uparrow\downarrow}$$

thus the dipoles line up horizontally parallel $\rightarrow\rightarrow$.

- (d) If a bunch of compasses were lined up along a straight line, they would align in straight line similarly to part (c): $\rightarrow\rightarrow\rightarrow\dots$

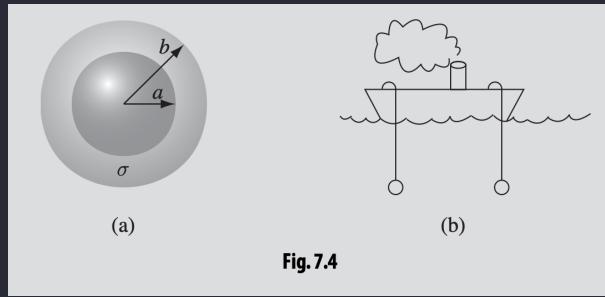


Fig. 7.4

7.1

- (a) Between the two shells we have an electric field

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{\mathbf{r}}$$

So we have a potential difference

$$\begin{aligned} V &= - \int_b^a \mathbf{E} \cdot d\mathbf{l} \\ &= - \frac{1}{4\pi\epsilon_0} Q \int_b^a \frac{1}{r^2} dr \\ &= - \frac{1}{4\pi\epsilon_0} Q \left(-\frac{1}{a} + \frac{1}{b} \right) \\ &= \frac{1}{4\pi\epsilon_0} Q \left(\frac{1}{a} - \frac{1}{b} \right) \\ \implies Q &= \frac{4\pi\epsilon_0 V}{1/a - 1/b} \end{aligned}$$

The current is therefore

$$I = \int \mathbf{J} \cdot d\mathbf{a} = \sigma \int \mathbf{E} \cdot d\mathbf{a} = \frac{\sigma}{\epsilon_0} Q$$

$$= \frac{\sigma}{\epsilon_0} \frac{4\pi\epsilon_0 V}{1/a - 1/b}$$

or

$$I = \frac{4\pi\sigma V}{1/a - 1/b}$$

- (b) The resistance between the shells is

$$R = \frac{V}{I} = \frac{1}{4\pi\sigma} \left(\frac{1}{a} - \frac{1}{b} \right)$$

- (c) If $b \gg a$, $\frac{1}{b} \rightarrow 0$ so the resistance is simply

$$R = \frac{1}{4\pi\sigma a}$$

For two shells immersed really deep in the sea with potential difference V , the we can add the resistance in series i.e.

$$R_{\text{sea}} = 2 \times \frac{1}{4\pi\sigma a} = \frac{1}{2\pi\sigma a}$$

Thus the current flowing between the two metal spheres is

$$I = \frac{V}{R_{\text{sea}}} = 2\pi\sigma a V$$

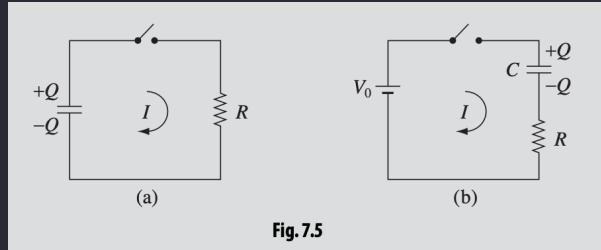


Fig.7.5

7.2

(a) Initially, $V_0 = \frac{Q_0}{C}$. As the capacitor discharges the current decreases,

$$\frac{dQ}{dt} = -I$$

and the sum of the voltages across the capacitor and the resistor is

$$\begin{aligned} \sum V &= V_C + V_R = 0 \\ &= \frac{Q}{C} + IR \\ \implies \frac{dQ}{dt} &= -I = -\frac{Q}{RC} \end{aligned}$$

The solution to this differential equation $\frac{dQ}{dt} = -kQ$ is

$$\begin{aligned} Q(t) &= Q_0 e^{-kt} \\ \implies I(t) &= -\frac{dQ}{dt} = -kQ_0 e^{-kt} \\ &= -\frac{Q_0}{RC} e^{-t/RC} \end{aligned}$$

and from the initial condition $Q_0 = CV_0$,

$$I(t) = -\frac{V_0}{R} e^{-t/RC}$$

(b) The original energy stored in the capacitor is (2.55)

$$W = \frac{1}{2} CV_0^2$$

So integrating the power equation (7.7)

$$\begin{aligned} \int_0^\infty P &= \int_0^\infty I^2 R dt \\ &= \frac{V_0^2}{R} \int_0^\infty e^{-2t/RC} dt \\ &= \frac{V_0^2}{R} \left[-\frac{RC}{2} e^{-2t/RC} \right] \Big|_0^\infty \\ &= \frac{V_0^2}{R} \left[0 + \frac{RC}{2} \right] \\ &= \frac{1}{2} CV_0^2 \end{aligned}$$

(c) Initially connecting the battery of V_0 :

$$V_0 = \sum V = V_C + V_R = \frac{Q}{C} + IR$$

where the current will start to increase, $\frac{dQ}{dt} = I$, so

$$\begin{aligned} V_0 &= \frac{Q}{C} + R \frac{dQ}{dt} \\ CV_0 &= Q + RC \frac{dQ}{dt} \\ \implies \frac{dQ}{dt} &= \frac{CV_0 - Q}{RC} \end{aligned}$$

And rewriting to get $-1/RC$ on to the right side:

$$\begin{aligned} \frac{dQ}{dt} &= -\frac{Q - CV_0}{RC} \\ \frac{dQ}{Q - CV_0} &= -\frac{dt}{RC} \end{aligned}$$

Integrating both sides

$$\begin{aligned} \ln(Q - CV_0) &= -\frac{t}{RC} + k \\ Q - CV_0 &= e^{-t/RC+k} \\ Q &= CV_0 + e^{-t/RC+k} \\ \implies Q(t) &= CV_0 + ke^{-t/RC} \end{aligned}$$

And from the initial condition $Q(0) = 0$,

$$\begin{aligned} Q(0) &= CV_0 + k = 0 \\ \implies k &= -CV_0 \end{aligned}$$

So the charge on the capacitor is

$$Q(t) = CV_0(1 - e^{-t/RC})$$

and the current is

$$I(t) = \frac{dQ}{dt} = CV_0 \left(\frac{1}{RC} e^{-t/RC} \right) = \boxed{\frac{V_0}{R} e^{-t/RC}}$$

(d) Integrating power to get the energy from the battery:

$$\begin{aligned} \int P dt &= \int_0^\infty IV_0 dt \\ &= \frac{V_0^2}{R} \int_0^\infty e^{-t/RC} dt \\ &= \frac{V_0^2}{R} \left[-RC e^{-t/RC} \right]_0^\infty \\ &= \frac{V_0^2}{R} [0 + RC] = \boxed{CV_0^2} \end{aligned}$$

The heat delivered to the resistor is still $\boxed{\frac{1}{2}CV_0^2}$, so the final energy stored in the capacitor is

$\boxed{\frac{1}{2}CV_0^2}$ i.e. 50% of the work done by the battery shows up as energy in the capacitor.

7.7 Magnetic field points into the page

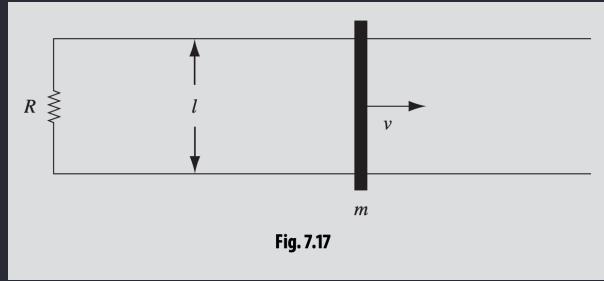


Fig. 7.17

- (a) If the bar moves to the right with speed v the EMF is

$$\mathcal{E} = B\ell v$$

And since $\mathcal{E} = IR$, the current is

$$I = \frac{B\ell v}{R}$$

where RHR points upwards (or counterclockwise) so the current points downwards in the resistor.

- (b) The magnetic force on the bar is

$$F = I\ell B = \frac{B^2\ell^2 v}{R}$$

and from RHR, the force points left

- (c) Starting at speed v_0 at $t = 0$, to find the speed at later time t we can just use Newton's second law

$$\begin{aligned} F &= m \frac{dv}{dt} = -\frac{B^2\ell^2 v}{R} \\ \frac{dv}{dt} &= \left(\frac{B^2\ell^2}{mR} \right) v \quad \text{or} \quad \frac{dv}{dt} = -kv \\ \Rightarrow v(t) &= v_0 e^{-kt} = \left[v_0 e^{-\frac{B^2\ell^2}{mR} t} \right] \end{aligned}$$

- (d) Given the initial kinetic energy $\frac{1}{2}mv_0^2$, and we can check this by integrating the power to get the work done by the magnetic field

$$\begin{aligned} \int_0^\infty P dt &= \int_0^\infty I^2 R dt \\ &= \frac{B^2\ell^2}{R} v_0^2 \int_0^\infty e^{-2kt} dt \quad k = \frac{B^2\ell^2}{mR} \\ &= \frac{B^2\ell^2}{R} v_0^2 \left[-\frac{1}{2k} e^{-2kt} \right] \Big|_0^\infty \\ &= \frac{B^2\ell^2}{R} v_0^2 \left[0 + \frac{1}{2k} \right] \\ &= \frac{1}{2}mv_0^2 \end{aligned}$$

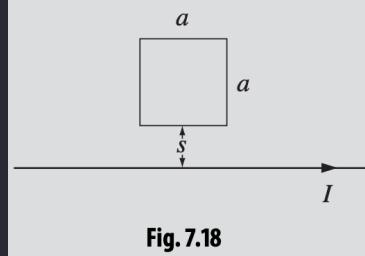


Fig. 7.18

7.8

- (a) The flux of \mathbf{B} through the loop: From HW 5b, the magnetic field of a long straight wire is

$$\mathbf{B} = \frac{\mu_0 I}{2\pi r} \hat{\phi}$$

So the flux is simply the integral

$$\Phi = \int \mathbf{B} \cdot d\mathbf{a}$$

where the surface element is $d\mathbf{a} = a ds$ from $s \rightarrow s + a$:

$$\begin{aligned}\Phi &= \frac{\mu_0 I a}{2\pi} \int_s^{s+a} \frac{1}{s} ds \\ &= \boxed{\left[\frac{\mu_0 I a}{2\pi} \ln \left(\frac{s+a}{s} \right) \right]}\end{aligned}$$

- (b) If the loop is moved away from the wire at velocity v , the EMF is

$$\begin{aligned}\mathcal{E} &= -\frac{d\Phi}{dt} = -\frac{d}{dt} \frac{\mu_0 I a}{2\pi} [\ln(s+a) - \ln(s)] \quad \frac{ds}{dt} = v \\ &= -\frac{\mu_0 I a}{2\pi} \left[\frac{v}{s+a} - \frac{v}{s} \right] \\ &= \frac{\mu_0 I a v}{2\pi} \left[\frac{1}{s} - \frac{1}{s+a} \right] \\ &= \frac{\mu_0 I a v}{2\pi} \left[\frac{s+a}{s(s+a)} - \frac{s}{s(s+a)} \right] \\ &= \boxed{\frac{\mu_0 I a^2 v}{2\pi s(s+a)}}\end{aligned}$$

From RHR ($\mathbf{v} \times \mathbf{B}$)— \mathbf{v} points up and \mathbf{B} point out of the page—points to the right, but the magnitude of force is less on the top side of the loop so the current flows counterclockwise.

- (c) No EMF is generated if the loop is moved parallel to the wire!