

Problem 1. Given

$$M \frac{d^2 u_s}{dt^2} = C(u_{s+1} + u_{s-1} - 2u_s)$$

where the nearest neighbors is displaced by $\pm p$ which is small so using the Taylor expansion

$$\begin{aligned} u_{s+a} &\approx u_s + a \frac{\partial u}{\partial x} + \frac{1}{2} a^2 \frac{\partial^2 u}{\partial x^2} \\ u_{s-a} &\approx u_s - a \frac{\partial u}{\partial x} + \frac{1}{2} a^2 \frac{\partial^2 u}{\partial x^2} \end{aligned}$$

so

$$C(u_{s+1} + u_{s-1} - 2u_s) \approx C \left(\cancel{u_s} + a \cancel{\frac{\partial u}{\partial x}} + \frac{1}{2} a^2 \frac{\partial^2 u}{\partial x^2} + \cancel{u_s} - a \cancel{\frac{\partial u}{\partial x}} + \frac{1}{2} a^2 \frac{\partial^2 u}{\partial x^2} - 2\cancel{u_s} \right) = C a^2 \frac{\partial^2 u}{\partial x^2}$$

and

$$\begin{aligned} M \frac{d^2 u_s}{dt^2} &= C a^2 \frac{\partial^2 u}{\partial x^2} \\ \frac{d^2 u_s}{dt^2} &= \frac{C a^2}{M} \frac{\partial^2 u}{\partial x^2} \end{aligned}$$

since the original differential equation can have solutions with time dependence $e^{i\omega t}$

$$\frac{d^2 u_s}{dt^2} = -\omega^2 u_s$$

and the displacements are translationally symmetric so

$$u_s = u e^{i s k a} \quad u_{s+1} = u e^{i(s+1)ka} = u e^{i s k a} e^{i k a} \quad u_{s-1} = u e^{i(s-1)ka} = u e^{i s k a} e^{-i k a}$$

so

$$\begin{aligned} -M\omega^2 u e^{i s k a} &= C u e^{i s k a} (e^{i k a} + e^{-i k a} - 2) \\ -M\omega^2 &= C (e^{i k a} + e^{-i k a} - 2) \end{aligned}$$

where we have the trigonometric identity

$$e^{i k a} + e^{-i k a} = 2 \cosh(i k a) = 2 \cos(k a)$$

so

$$\begin{aligned} -M\omega^2 &= C(2 \cos(k a) - 2) \\ \omega^2 &= \frac{2C}{M} (1 - \cos(k a)) \end{aligned}$$

and from the half angle identity

$$1 - \cos(k a) = 2 \sin^2\left(\frac{k a}{2}\right)$$

so

$$\omega = \sqrt{\frac{4C}{M}} \sin\left(\frac{k a}{2}\right)$$

we can also find the group velocity

$$v = \frac{d\omega}{dk} = \sqrt{\frac{4C}{M}} \cos\left(\frac{k a}{2}\right) \frac{a}{2} = \sqrt{\frac{C a^2}{M}} \cos\left(\frac{k a}{2}\right)$$

and when $ka \ll 1$ or ≈ 0 , cosine is 1 so the group velocity is

$$v = \sqrt{\frac{C a^2}{M}} \quad \text{or} \quad v^2 = \frac{C a^2}{M}$$

so we get the wave equation

$$\frac{d^2 u}{dt^2} = v^2 \frac{d^2 u}{dx^2}$$

Problem 2. There will be two equations for alternate force constants where the nearest neighbor for atom u_s are u'_s and u'_{s-1} , and for atom u'_s , the nearest neighbors are u_s and u_{s+1} . So the equations of motion are

$$M \frac{d^2 u}{dt^2} = C(u'_s - u_s) + C'(u'_{s-1} - u_s)$$

$$M \frac{d^2 u'}{dt^2} = C'(u_{s+1} - u'_s) + C(u_s - u'_s)$$

where $C' = 10C$. The shifted parts are

$$u'_{s-1} = u e^{iska} e^{-ika} = u e^{iska} e^{-ika}$$

$$u_{s+1} = u e^{iska} e^{ika} = u e^{iska} e^{ika}$$

and from the previous problem e^{iska} will cancel out (but not u or u') so

$$-M\omega^2 u = C(u' - u) + C'(u' e^{-ika} - u)$$

$$-M\omega^2 u' = C'(u e^{ika} - u') + C(u - u')$$

rearranging the terms for the first EQ:

$$M\omega^2 u = -C(u' - u) - C'(u' e^{-ika} - u)$$

$$= (C + C')u + (-C - C' e^{-ika})u'$$

$$0 = (C + C' - M\omega^2)u + (-C - C' e^{-ika})u'$$

and for the second EQ:

$$M\omega^2 u' = -C'(u e^{ika} - u') - C(u - u')$$

$$= (-C - C' e^{ika})u + (C + C')u'$$

$$0 = (-C - C' e^{ika})u + (C + C' - M\omega^2)u'$$

so we have the matrix equation

$$\begin{pmatrix} C + C' - M\omega^2 & -C - C' e^{-ika} \\ -C - C' e^{ika} & C + C' - M\omega^2 \end{pmatrix} \begin{pmatrix} u \\ u' \end{pmatrix} = 0$$

we can solve for the eigenvalues of the matrix by finding the determinant:

For $K = 0$, the matrix is

$$\begin{pmatrix} C + C' - M\omega^2 & -C - C' \\ -C - C' & C + C' - M\omega^2 \end{pmatrix}$$

and we have a special case where $u = u'$ or

$$0 = (C + C' - M\omega^2)u + (-C - C')u = -M\omega^2 u \implies \omega^2 = 0$$

and the other case where the determinant is zero

$$0 = (C + C' - M\omega^2)^2 - (-C - C')^2$$

using $(a + b + c)^2 = a^2 + b^2 + c^2 + 2ab + 2bc + 2ac$

$$= C^2 + C'^2 + M^2\omega^4 + 2CC' - 2C'M\omega^2 - 2CM\omega^2$$

$$- (C^2 + C'^2 + 2CC')$$

$$= M^2\omega^4 - 2C'M\omega^2 - 2CM\omega^2$$

$$= M\omega^2 - 2(C' + C)$$

and plugging in $C' = 10C$

$$0 = M\omega^2 - 2(11C) \implies \omega^2 = 22C/M$$

For $K = \pi/a$ and separation $a = \pi/2$

$$e^{ika} = e^{i\frac{\pi}{a}\frac{a}{2}} = e^{i\pi} = -1 \quad e^{-ika} = -1$$

so the matrix is

$$\begin{pmatrix} C + C' - M\omega^2 & -C + C' \\ C - C' & C + C' - M\omega^2 \end{pmatrix} \begin{pmatrix} u \\ u' \end{pmatrix} = 0$$

so for the special case where $u = u'$:

$$\begin{aligned} 0 &= (C + C' - M\omega^2) + (-C + C') \\ &= 2C' - M\omega^2 \implies \omega^2 = 20C/M \end{aligned}$$

and the other case where the determinant is zero

$$\begin{aligned} 0 &= (11C - M\omega^2)^2 - (9C)^2 \\ &= 121C^2 - 22CM\omega^2 + M^2\omega^4 - 81C^2 \\ &= M^2\omega^4 - 22CM\omega^2 + 40C^2 \\ &= (M\omega^2)^2 - 22C(M\omega^2) + 40C^2 \end{aligned}$$

solving this quadratic equation where $a = 1, b = -22C, c = 40C^2$

$$\begin{aligned} M\omega^2 &= \frac{22C \pm \sqrt{(-22C)^2 - 4(40C^2)}}{2} = \frac{22C \pm \sqrt{484C^2 - 160C^2}}{2} \\ &= \frac{22C \pm \sqrt{324C^2}}{2} = \frac{22C \pm 18C}{2} = 20C, 2C \end{aligned}$$

so

$$\omega^2 = \frac{20C}{M}, \frac{2C}{M}$$

thus the dispersion relations are

$$\begin{aligned} K = 0: \quad \omega &= 0, \sqrt{\frac{22C}{M}} \\ K = \pi/a: \quad \omega &= \sqrt{\frac{20C}{M}}, \sqrt{\frac{2C}{M}} \end{aligned}$$

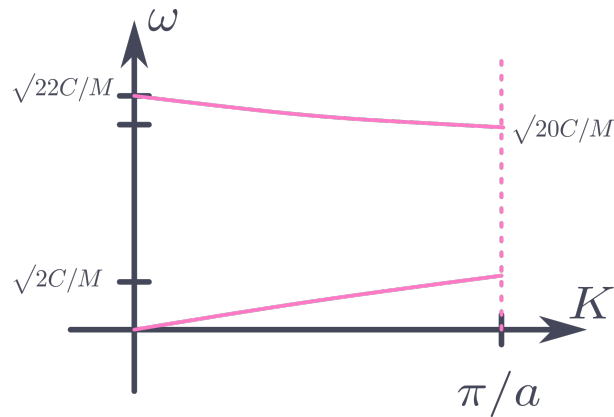


Figure 1: Dispersion relation for $K = 0, \pi/a$

Problem 3. Given

$$m\left(\frac{dv}{dt} + \frac{v}{\tau}\right) = -eE$$

for time varying E and v

$$E = E_0 e^{-i\omega t} \quad v = v_0 e^{-i\omega t}; \quad \frac{dv}{dt} = -i\omega v$$

so solving for v

$$\begin{aligned} m\left(-i\omega v + \frac{v}{\tau}\right) &= -eE \\ v &= \frac{-eE}{m(-i\omega + \frac{1}{\tau})} \left(\frac{\tau}{\tau}\right) \\ &= \frac{-eE}{m} \frac{\tau}{-i\omega\tau + 1} \end{aligned}$$

from Ohm's Law

$$\begin{aligned} j = \sigma E = nqv \quad \text{or} \quad \sigma &= \frac{-nev}{E} \\ \text{where } \sigma(0) &= \frac{ne^2\tau}{m} \end{aligned}$$

where the charge is $q = -e$. Substituting v from the previous equation

$$\begin{aligned} \sigma &= \frac{-ne}{E} \frac{-eE}{m} \frac{\tau}{-i\omega\tau + 1} \\ &= \frac{ne^2\tau}{m} \frac{1}{-i\omega\tau + 1} \left(\frac{1 + i\omega\tau}{1 + i\omega\tau}\right) \\ \sigma(\omega) &= \sigma(0) \frac{1 + i\omega\tau}{1 + (\omega\tau)^2} \end{aligned}$$

where we multiply by the complex conjugate in the second step.