

Homework 9

1. Inertia Tensor derivation using kronicker delta: The diagonal elements of the inertia tensor are given by

$$\begin{aligned} I_{xx} &= \sum m_i(y_i^2 + z_i^2) \\ I_{yy} &= \sum m_i(x_i^2 + z_i^2) \\ I_{zz} &= \sum m_i(x_i^2 + y_i^2) \end{aligned}$$

and the off-diagonal elements are given by

$$\begin{aligned} I_{xy} &= - \sum m_i x_i y_i \\ I_{xz} &= - \sum m_i x_i z_i \\ I_{yz} &= - \sum m_i y_i z_i \end{aligned}$$

For the diagonal element we can see that using the single equation

$$\begin{aligned} I_{xx} &= \int \rho(r^2 \delta_{xx} - r_x r_x) dV \\ &= \int \rho(r^2 - x^2) dV \\ &= \int \rho(x^2 + y^2 + z^2 - x^2) dV \\ &= \int \rho(y^2 + z^2) dV = \sum m_i(y_i^2 + z_i^2) \end{aligned}$$

and similarly for the other diagonal elements. For the off-diagonal elements we can see that the kronecker delta will be zero and thus

$$\begin{aligned} I_{xy} &= \int \rho(0 - xy) dV \\ &= - \int \rho xy dV = - \sum m_i x_i y_i \end{aligned}$$

and similarly for the other off-diagonal elements.

2. (a) First we can see that

$$\begin{aligned} (\mathbf{A} \times \mathbf{B})^2 &= (\mathbf{A} \times \mathbf{B}) \cdot (\mathbf{A} \times \mathbf{B}) \\ &= (AB \sin \theta \hat{\mathbf{n}})(AB \sin \theta \hat{\mathbf{n}}) \\ &= A^2 B^2 (1 - \cos^2 \theta) \quad \text{using } \mathbf{A} \cdot \mathbf{B} = AB \cos \theta \\ &= A^2 B^2 - (\mathbf{A} \cdot \mathbf{B})^2 \end{aligned}$$

So the Kinetic Energy is

$$\begin{aligned} T &= \frac{1}{2} \sum_{\alpha} m_{\alpha} v_{\alpha}^2 \\ &= \frac{1}{2} \sum_{\alpha} m_{\alpha} \mathbf{v}_{\alpha} \cdot \mathbf{v}_{\alpha} \\ &= \frac{1}{2} \sum_{\alpha} m_{\alpha} (\boldsymbol{\omega} \times \mathbf{r}_{\alpha})^2 \\ &= \frac{1}{2} \sum_{\alpha} m_{\alpha} [(\omega r_{\alpha})^2 - (\boldsymbol{\omega} \cdot \mathbf{r}_{\alpha})^2] \end{aligned}$$

(b) Angular momentum is given by

$$\begin{aligned}\mathbf{L} &= \sum_{\alpha} \mathbf{r}_{\alpha} \times m_{\alpha} \mathbf{v}_{\alpha} \\ &= \sum_{\alpha} m_{\alpha} \mathbf{r}_{\alpha} \times (\boldsymbol{\omega} \times \mathbf{r}_{\alpha})\end{aligned}$$

using the BAC-CAB rule (or WRR-RRW rule in this case) we get

$$\begin{aligned}&= \sum_{\alpha} m_{\alpha} [\boldsymbol{\omega} (\mathbf{r}_{\alpha} \cdot \mathbf{r}_{\alpha}) - \mathbf{r}_{\alpha} (\mathbf{r}_{\alpha} \cdot \boldsymbol{\omega})] \\ &= \sum_{\alpha} m_{\alpha} [\boldsymbol{\omega} r_{\alpha}^2 - \mathbf{r}_{\alpha} (\boldsymbol{\omega} \cdot \mathbf{r}_{\alpha})]\end{aligned}$$

where dot products commute. (c) From the previous part we can see that

$$\begin{aligned}\frac{1}{2} \boldsymbol{\omega} \cdot \mathbf{L} &= \frac{1}{2} \boldsymbol{\omega} \cdot \sum_{\alpha} m_{\alpha} [\boldsymbol{\omega} r_{\alpha}^2 - \mathbf{r}_{\alpha} (\boldsymbol{\omega} \cdot \mathbf{r}_{\alpha})] \\ &= \frac{1}{2} \sum_{\alpha} m_{\alpha} [\boldsymbol{\omega} \cdot \boldsymbol{\omega} r_{\alpha}^2 - \boldsymbol{\omega} \cdot \mathbf{r}_{\alpha} (\boldsymbol{\omega} \cdot \mathbf{r}_{\alpha})] \\ &= \frac{1}{2} \sum_{\alpha} m_{\alpha} [\boldsymbol{\omega}^2 r_{\alpha}^2 - (\boldsymbol{\omega} \cdot \mathbf{r}_{\alpha})^2] = T\end{aligned}$$

Since $\mathbf{L} = \mathbf{I} \boldsymbol{\omega}$ we usually write the angular momentum as a column vector with three components

$$\mathbf{L} = \begin{pmatrix} L_x \\ L_y \\ L_z \end{pmatrix}$$

and taking the cross product of \mathbf{L} with $\boldsymbol{\omega}$ we will have to transpose $\boldsymbol{\omega}$ to get the correct matrix multiplication:

$$\begin{aligned}\boldsymbol{\omega} \cdot \mathbf{L} &= \boldsymbol{\omega}^T \mathbf{L} \\ &= (\omega_x \quad \omega_y \quad \omega_z) \begin{pmatrix} L_x \\ L_y \\ L_z \end{pmatrix} \\ &= \omega_x L_x + \omega_y L_y + \omega_z L_z\end{aligned}$$

So the kinetic energy is also equivalent to

$$T = \frac{1}{2} \boldsymbol{\omega}^T \mathbf{L} = \frac{1}{2} \boldsymbol{\omega}^T \mathbf{I} \boldsymbol{\omega}$$

3. (a) For a uniform hollow ice cream cone of radius R , height h , and mass M . The mass density is given by the mass per unit area

$$\begin{aligned} q &= \frac{M}{A} = \frac{M}{\pi R l} \quad \frac{R}{l} = \sin \theta \\ &= \frac{M}{\pi R^2} \sin \theta \end{aligned}$$

where l is the the slant of the cone and using cylindrical coordinates

$$I_{zz} = q \int_A dA (x^2 + y^2) = q \int_A dA \rho^2$$

The area element is a rectangular region with sides $\rho d\phi$ and $\frac{1}{\sin \theta} d\rho$ (the $d\phi$ is projected onto the side of the cone) so the area element is

$$dA = \rho d\rho d\phi \frac{1}{\sin \theta}$$

and the moment of inertia is

$$\begin{aligned} I_{zz} &= q \int_0^{2\pi} d\phi \int_0^R d\rho \frac{\rho^3}{\sin \theta} \\ &= \frac{M}{\pi R^2} \sin \theta (2\pi) \frac{R^4}{4 \sin \theta} \\ &= \frac{1}{2} M R^2 \end{aligned}$$

$I_{xx} = I_{yy}$ since the cone is rotationally symmetric about the z -axis:

$$\begin{aligned} I_{xx} &= q \int_A dA (y^2 + z^2) \\ &= q \int_A dA (\rho^2 \sin^2 \phi + z^2) \end{aligned}$$

Geometrically we have similar triangles where the ratio of the sides are

$$\frac{z}{\rho} = \frac{h}{R} \implies z = \frac{h}{R} \rho$$

so the integral becomes

$$\begin{aligned} I_{xx} &= q \int_0^{2\pi} d\phi \left(\sin^2 \phi + \frac{h^2}{R^2} \right) \int_0^R d\rho \rho^3 \\ &= \frac{M}{\pi R^2} \sin \theta \left(\pi + 2\pi \frac{h^2}{R^2} \right) \frac{R^4}{4 \sin \theta} \\ &= \frac{1}{4} M R^2 \left(1 + 2 \frac{h^2}{R^2} \right) \\ &= \frac{1}{4} M (R^2 + 2h^2) \end{aligned}$$

And the off-diagonal elements are all zero due to the rotational symmetry of the cone which gives the inertia tensor

$$\mathbf{I} = \begin{pmatrix} \frac{1}{4} M (R^2 + 2h^2) & 0 & 0 \\ 0 & \frac{1}{4} M (R^2 + 2h^2) & 0 \\ 0 & 0 & \frac{1}{2} M R^2 \end{pmatrix}$$

for I_{zz} we have

$$I_{zz} = \sigma \int_A dA (x^2 + y^2) = I_{xx} + \sigma \int_A dA x^2$$

and since $I_{yy} = \sigma \int_A dA (x^2 + z^2) = \sigma \int_A dA x^2$

$$I_{zz} = I_{xx} + I_{yy} = 4$$

Sadly, only 4 diagonal elements (the ones containing z) are zero. So calculating the two off-diagonal elements $I_{xy} = I_{yx}$ we have

$$\begin{aligned} I_{xy} &= -\sigma \int_A dA xy \\ &= -\sigma \int_0^1 dx \int_0^{-x+1} dy xy \\ &= -\sigma \int_0^1 dx x \frac{1}{2} (-x+1)^2 \\ &= -\sigma \int_0^1 dx \frac{1}{2} (x^3 - 2x^2 + x) \\ &= -\frac{1}{24} \sigma = -1 \end{aligned}$$

So the inertia tensor is

$$\mathbf{I} = \begin{pmatrix} 2 & -1 & 0 \\ -1 & 2 & 0 \\ 0 & 0 & 4 \end{pmatrix}$$

To find the principal moments we find the eigenvalues for the matrix equation

$$\begin{aligned} \det(\mathbf{I} - \lambda \mathbf{1}) &= 0 \\ \begin{pmatrix} 2-\lambda & -1 & 0 \\ -1 & 2-\lambda & 0 \\ 0 & 0 & 4-\lambda \end{pmatrix} &= 0 \end{aligned}$$

Where we can see that the sum of the rows can give us the eigenvalues $\lambda = 1, 4$ and since the trace of the matrix is 8 the last eigenvalue is $\lambda = 3$. Plugging in the eigenvalues into the matrix equation does indeed give us the correct solutions. To find the first principal axis we plug in $\lambda = 1$ into the matrix equation

$$\begin{aligned} (\mathbf{I} - \lambda \mathbf{1}) \boldsymbol{\omega} &= 0 \\ \begin{pmatrix} 1 & -1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 3 \end{pmatrix} \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix} &= 0 \end{aligned}$$

which gives us the equation $\omega_x = \omega_y = 1$ and $\omega_z = 0$ so the first principal axis is

$$\hat{\mathbf{e}}_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}$$

For $\lambda = 3$ we have

$$\begin{pmatrix} -1 & -1 & 0 \\ -1 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix} = 0$$

