

Homework 3

Due 2/7 9pm

1. The Center of Mass of the system is

$$\mathbf{R} = \frac{1}{M} \int \mathbf{r} \, dm$$

and the mass element is the mass density times the volume element

$$dm = \rho \, dV = \frac{M}{2\pi R^2} R^2 \sin \phi \, d\phi \, d\theta = \frac{M}{2\pi} \sin \phi \, d\phi \, d\theta$$

where there is no dr term because the radius is constant. Since the center of mass is symmetric about the x and y axes, $X_{cm} = Y_{cm} = 0$. The z component of the center of mass is

$$\begin{aligned} Z_{cm} &= \frac{1}{M} \int z \, dm \\ &= \frac{1}{M} \frac{M}{2\pi} \iint z \sin \phi \, d\phi \, d\theta \end{aligned}$$

where $z = R \cos \phi$ so

$$\begin{aligned} Z_{cm} &= \frac{R}{2\pi} \int_0^{2\pi} d\theta \int_0^{\pi/2} \cos \phi \sin \phi \, d\phi \\ &\quad \text{using } u = \sin \phi \implies du = \cos \phi \, d\phi \\ &= \frac{R}{2\pi} [2\pi] \int_0^1 u \, du \\ Z_{cm} &= \frac{R}{2} \end{aligned}$$

So the COM is at $\boxed{\left(0, 0, \frac{R}{2}\right)}$

2. (a) From N3L the force of the jettisoned fuel on the rocket is equal and opposite to the thrust on the rocket from the jettisoned fuel:

$$\begin{aligned} F_{\text{fuel}} &= -F_{\text{thrust}} \\ \dot{m}v_{ex} &= -\dot{m}v_{ex} \end{aligned}$$

So using N2L, the sum of the forces on the rocket is the thrust and air resistance:

$$F = m\dot{v} = F_{\text{thrust}} - f = -\dot{m}v_{ex} - bv$$

(b) Using $\dot{m} = -k$

$$\begin{aligned} m\dot{v} &= kv_{ex} - bv \\ \frac{m}{b}\dot{v} &= \frac{kv_{ex}}{b} - v \end{aligned}$$

defining the constant $a = \frac{kv_{ex}}{b}$ and using separation of variables

$$\frac{1}{a-v} dv = \frac{b}{m} dt$$

we can write an expression for m as a function of time through using separation of variables again:

$$\begin{aligned} \frac{dm}{dt} &= -k \\ \int_{m_o}^m dm' &= -k \int_0^t dt' \\ m - m_o &= -kt \\ m &= m_o - kt \end{aligned}$$

where m_o is the initial mass of the rocket. Substituting this back into main expression and integrating both sides:

$$\begin{aligned} \int_0^v \frac{1}{a-v'} dv' &= b \int_0^t \frac{1}{m_o - kt'} dt' \\ -\ln(a-v') \Big|_0^v &= -\frac{b}{k} \ln(m_o - kt') \Big|_0^t \\ -\ln(a-v) + \ln(a) &= \frac{b}{k} [-\ln(m_o - kt) + \ln(m_o)] \\ \ln\left(\frac{a}{a-v}\right) &= \frac{b}{k} \ln\left(\frac{m_o}{m_o - kt}\right) \end{aligned}$$

substituting back in $m = m_o - kt$ and exponentiating both sides:

$$\begin{aligned} \frac{a}{a-v} &= \left(\frac{m_o}{m}\right)^{\frac{b}{k}} \\ a\left(\frac{m_o}{m}\right)^{-\frac{b}{k}} &= a-v \\ v &= a - a\left(\frac{m_o}{m}\right)^{-\frac{b}{k}} \\ v &= a \left[1 - \left(\frac{m}{m_o}\right)^{\frac{b}{k}} \right] \end{aligned}$$

subbing back in $a = \frac{kv_{ex}}{b}$ we get the final expression

$$v(m) = \frac{kv_{ex}}{b} \left[1 - \left(\frac{m}{m_o}\right)^{\frac{b}{k}} \right]$$

3. (a) The angular momentum vector is

$$\begin{aligned}\ell &= \mathbf{r} \times \mathbf{p} \\ &= \mathbf{r} \times m\dot{\mathbf{r}} \\ &= m\mathbf{r} \times \dot{\mathbf{r}}\end{aligned}$$

from HW 1, we know that

$$\dot{\mathbf{r}} = \dot{r}\hat{\mathbf{r}} + r\dot{\phi}\hat{\phi}$$

so

$$\begin{aligned}\ell &= m(r\hat{\mathbf{r}}) \times (\dot{r}\hat{\mathbf{r}} + r\dot{\phi}\hat{\phi}) \\ &= m[r\dot{r}(\hat{\mathbf{r}} \times \hat{\mathbf{r}}) + r(r\dot{\phi})(\hat{\mathbf{r}} \times \hat{\phi})] \\ &= m[0 + r^2\dot{\phi}\hat{\mathbf{z}}] = mr^2\omega\hat{\mathbf{z}}\end{aligned}$$

where $\omega = \dot{\phi}$ and the magnitude of the angular momentum is

$$\ell = |\ell| = mr^2\omega$$

(b) The area swept by an infinitesimal change in the planets position is equivalent to the area of a triangle as shown in Figure 2, so

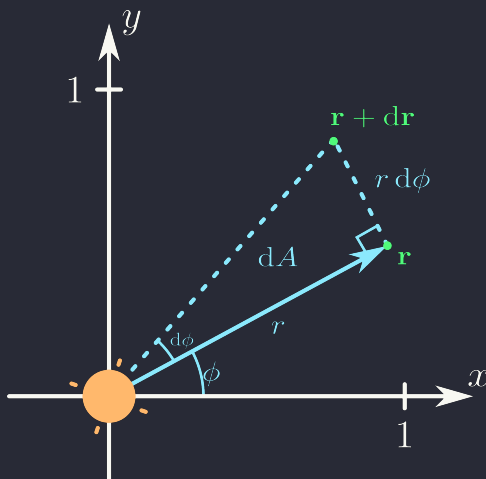


Figure 3.1: Area swept by planet

$$dA = \frac{1}{2}r(r d\phi) = \frac{1}{2}r^2 d\phi$$

dividing both sides by dt gives us

$$\frac{dA}{dt} = \frac{1}{2}r^2 \frac{d\phi}{dt} = \frac{1}{2}r^2\omega$$

and from part (a) we know that $\ell = mr^2\omega$ or $\omega = \frac{\ell}{mr^2}$ so

$$\frac{dA}{dt} = \frac{1}{2}r^2 \frac{\ell}{mr^2} = \frac{\ell}{2m}$$

therefore, the rate in change of the area swept by the planet is a constant that is proportional to ℓ .

4. $\mathbf{F} = -y\hat{\mathbf{x}} + x\hat{\mathbf{y}}$ from $P = (1, 0)$ to $Q = (0, 1)$

(a) For a straight line the path is given by

$$y = 1 - x \quad dy = -dx$$

so the work done is

$$\begin{aligned} W_a &= \int_P^Q F_x dx + F_y dy = \int_P^Q -y dx + x dy \\ &= \int_{x=1}^0 -(1-x) dx + x(-dx) \\ &= \int_1^0 -1 dx = 1 \end{aligned}$$

(b) For a circular path of radius 1, the path in polar coordinates is

$$\begin{aligned} y &= \sin \phi \rightarrow dy = \cos \phi d\phi \\ x &= \cos \phi \rightarrow dx = -\sin \phi d\phi \end{aligned}$$

from the equation of a circle $x^2 + y^2 = 1$. The limits of integration are $\phi = 0 \rightarrow \pi/2$, and the work is

$$\begin{aligned} W_b &= \int_{\phi=0}^{\pi/2} -\sin \phi (-\sin \phi) d\phi + \cos \phi \cos \phi d\phi \\ &= \int_{\phi=0}^{\pi/2} 1 d\phi = \frac{\pi}{2} \end{aligned}$$

(c) Splitting this into two paths: For path 1, $y = 0$; $dy = 0$; and $x = 1 \rightarrow 0$ so

$$W_1 = \int_{x=1}^0 0 dx + x(0) = 0$$

For path 2, $x = 0$; $dx = 0$; $y = 0 \rightarrow 1$ so

$$W_2 = \int_{y=0}^1 -y(0) + 0 dy = \int_{y=0}^1 0 dy = 0$$

And the work done is $W_c = W_1 + W_2 = 0$

(d) The force is not conservative because the work done is path dependent! We can also double check by taking the curl:

$$\nabla \times \mathbf{F} = \det \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} \\ -y & x \end{vmatrix} = \left(\frac{\partial}{\partial x} x - \frac{\partial}{\partial y} (-y) \right) \hat{\mathbf{z}} = 2\hat{\mathbf{z}}$$

which is not zero, so the force is not conservative.

5. (a) Using the time derivatives of the polar unit vectors

$$\frac{d}{dt}\hat{\mathbf{r}} = \dot{\phi}\hat{\phi} \quad \frac{d}{dt}\hat{\phi} = -\dot{\phi}\hat{\mathbf{r}}$$

Acceleration in polar coordinates is

$$\begin{aligned}\mathbf{a} = \dot{\mathbf{v}} &= \frac{d}{dt}\dot{\mathbf{r}} = \frac{d}{dt}(\dot{r}\hat{\mathbf{r}} + r\dot{\phi}\hat{\phi}) \\ &= \ddot{r}\hat{\mathbf{r}} + \dot{r}\dot{\phi}\hat{\phi} + (\dot{r}\dot{\phi}\hat{\phi} + r\ddot{\phi}\hat{\phi} + r\dot{\phi}(-\dot{\phi}\hat{\mathbf{r}})) \\ &= (\ddot{r} - r\dot{\phi}^2)\hat{\mathbf{r}} + (r\ddot{\phi} + 2\dot{r}\dot{\phi})\hat{\phi}\end{aligned}$$

so the radial and angular components of the force are

$$\begin{aligned}F_r &= ma_r = m(\ddot{r} - r\dot{\phi}^2) \\ F_\phi &= ma_\phi = m(r\ddot{\phi} + 2\dot{r}\dot{\phi})\end{aligned}$$

and since the spring force is conservative with magnitude $F_s = -k(r - a)\hat{\mathbf{r}}$ the equations of motion are

$$\begin{aligned}m(\ddot{r} - r\dot{\phi}^2) &= -k(r - a) \\ m(r\ddot{\phi} + 2\dot{r}\dot{\phi}) &= 0\end{aligned}$$

or

$$\begin{aligned}m\ddot{r} - mr\dot{\phi}^2 + k(r - a) &= 0 \\ r\ddot{\phi} + 2\dot{r}\dot{\phi} &= 0\end{aligned}$$

(b) The initial angular momentum of the system is

$$\ell_o = mv_o a$$

and after some time the angular momentum is (from Problem 3)

$$\ell = mr^2\dot{\phi}$$

and using the conservation of angular momentum

$$\begin{aligned}\ell_o &= \ell \\ mv_o a &= mr^2\dot{\phi} \\ \dot{\phi} &= \frac{v_o a}{r^2}\end{aligned}$$

(c) First the initial mechanical energy of the system is purely kinetic given by the initial velocity:

$$E_o = T_o = \frac{1}{2}mv_o^2$$

the total mechanical energy of the system after some time will be the sum of the kinetic and potential energies:

$$\begin{aligned}U &= -\int_0^r \mathbf{F} \cdot d\mathbf{r}' = \int_0^r k(r - a) dr' = \frac{1}{2}k(r - a)^2 \\ T &= \frac{1}{2}mv^2 = \frac{1}{2}m(\mathbf{v} \cdot \mathbf{v}) = \frac{1}{2}m(\dot{r}\hat{\mathbf{r}} + r\dot{\phi}\hat{\phi}) \cdot (\dot{r}\hat{\mathbf{r}} + r\dot{\phi}\hat{\phi}) = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2)\end{aligned}$$

And from the conservation of energy

$$\begin{aligned}E_o &= E = T + U \\ \frac{1}{2}mv_o^2 &= \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2) + \frac{1}{2}k(r - a)^2 \\ v_o^2 &= \dot{r}^2 + r^2\dot{\phi}^2 + \frac{k}{m}(r - a)^2\end{aligned}$$

