Ludwig-Maximilians-Universität München

Bachelors Thesis

The role of excited atomic states in multiphoton ionization

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Abstract

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1. Introduction

2. Theory

$$\partial_t u = \mathcal{H}(t)\lambda \tag{2.1}$$

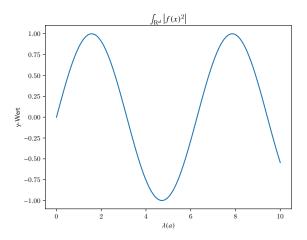


Figure 2.1: Sine function

2.1 Basic Formalism

Basic Definitions of schröfinger qe, light dyson series, and sotrong field s matrix Starting with schrödinger equation and a general hamiltonian

$$i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle = \hat{H}(t) |\Psi(t)\rangle.$$
 (2.2)

The formal solution depends on the time dependence of the hamiltonian and the physical setting. In the following we assume (IMPORTANT) that $[\hat{H}(t), \hat{H}(t')] = 0$ so we assume some sort of quasi static approximation to the Hamiltons time evolution. The solution is then given by

$$|\Psi(t)\rangle = e^{-\frac{i}{\hbar} \int_{t_0}^t \hat{H}(t')dt'} |\Psi(t=0)\rangle \tag{2.3}$$

2.2 Strong Field Approximation

clear defintition of the strong field approximation, and the assumptions that are made.

2.3 Strong Field Ionization

Derivation of

$$\lim_{t \to \infty} |\Psi(t)\rangle = -i \int d^3p |\mathbf{p}\rangle \int_{-\infty}^{\infty} dt' e^{-\frac{i}{2} \int_{t'}^{\infty} [\mathbf{p} + \mathbf{A}(t')]^2 dt'} e^{iI_{\mathbf{p}}t'} \langle \mathbf{p} + \mathbf{A}(t') | \hat{\mathbf{d}} \cdot \mathbf{E}(t') | \Psi_0 \rangle \qquad (2.4)$$

2.4 Multiphoton Ionization

Different types of Ionization, tunneling Ionization, multiphoton

3. Conclusion

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$$\partial \mathbf{A} = \mathfrak{B}$$

$$\int_{\mathbb{R}^d} |f(x)|^2 \, \mathrm{d}x = \int_{\mathbb{R}^d} |\mathcal{F}f(\xi)|^2 \, \mathrm{d}\xi \tag{3.1}$$

$$i\partial_t u = \mathcal{H}(t) |a\rangle \lambda$$
 (3.2)

for this calculation [1] was used

A. Appendix A

B. Appendix B

Bibliography

[1] Seung Beom Park, Kyungseung Kim, Wosik Cho, Sung In Hwang, Igor Ivanov, Chang Hee Nam, and Kyung Taec Kim. Direct sampling of a light wave in air. *Optica*, 5(4):402–408, Apr 2018.