Hamiltonian flows, cotangent lifts, and momentum maps

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1 Symplectic manifolds

Let (M, ω) and (N, η) be symplectic manifolds. A symplectomorphism $F: M \to N$ is a diffeomorphism such that $\omega = F^*\eta$. Recall that for $x \in M$ and $v_1, v_2 \in T_xM$,

$$(F^*\eta)_x(v_1,v_2) = \eta_{F(x)}((T_xF)v_1,(T_xF)v_2);$$

 $T_xF: T_xM \to T_{F(x)}N$. (A tangent vector at $x \in M$ is pushed forward to a tangent vector at $F(x) \in N$, while a differential 2-form on N is pulled back to a differential 2-form on M.) In these notes the only symplectomorphisms in which we are interested are those from a symplectic manifold to itself.¹

2 Symplectic gradient

If (M, ω) is a symplectic manifold and $H \in C^{\infty}(M)$, using the nondegeneracy of the symplectic form ω one can prove that there is a unique vector field $X_H \in \Gamma^{\infty}(M)$ such that, for all $x \in M, v \in T_xM$,

$$\omega_x(X_H(x), v) = (dH)_x(v).$$

This can also be written as

$$i_{X_H}\omega = dH$$
,

where

$$(i_X\omega)(Y) = (X \sqcup \omega)(Y) = \omega(X, Y).$$

We call X_H the symplectic gradient of H. If $X \in \Gamma^{\infty}(M)$ and $X = X_H$ for some $H \in C^{\infty}(M)$, we say that X is a Hamiltonian vector field.²

¹I am interested in flows on a phase space and this phase space is a symplectic manifold. For some motivation for why we want phase space to be a symplectic manifold, read:

http://research.microsoft.com/en-us/um/people/cohn/thoughts/symplectic.html

²On a Riemannian manifold, a vector field that is the gradient of a smooth function is called a *gradient vector field* or a *conservative vector field*.

Let's check that

$$X_H = \sum_{i=1}^n \frac{\partial H}{\partial p_i} \frac{\partial}{\partial q_i} - \frac{\partial H}{\partial q_i} \frac{\partial}{\partial p_i}.$$

We have, because $dq_i \frac{\partial}{\partial q_j} = \delta_{ij}$, $dp_i \frac{\partial}{\partial p_j} = \delta_{ij}$, $dq_i \frac{\partial}{\partial p_j} = 0$ and $dp_i \frac{\partial}{\partial q_j} = 0$, and because $dq_j \wedge dp_j = -dp_j \wedge dq_j$,

$$i_{X_H}\omega = \sum_{i=1}^n dq_i \wedge dp_i \sum_{j=1}^n \left(\frac{\partial H}{\partial p_j} \frac{\partial}{\partial q_j} - \frac{\partial H}{\partial q_j} \frac{\partial}{\partial p_j} \right)$$

$$= \sum_{i=1}^n dq_i \wedge dp_i \left(\frac{\partial H}{\partial p_i} \frac{\partial}{\partial q_i} - \frac{\partial H}{\partial q_i} \frac{\partial}{\partial p_i} \right)$$

$$= \sum_{i=1}^n \frac{\partial H}{\partial p_i} dp_i + \frac{\partial H}{\partial q_i} dq_i$$

$$= dH.$$

3 Flows

Let M be a smooth manifold. Let D be an open subset of $M \times \mathbb{R}$, and for each $x \in M$ suppose that

$$D^x = \{ t \in \mathbb{R} : (x, t) \in D \}$$

is an open interval including 0. A flow on M is a smooth map $\phi: D \to M$ such that if $x \in M$ then $\phi_0(x) = x$ and such that if $x \in M$, $s \in D^x$, $t \in D^{\phi_s(x)}$ and $s + t \in D^x$, then

$$\phi_t(\phi_s(x)) = \phi_{s+t}(x).$$

For $x \in M$, define $\phi^x : D^x \to M$ by $\phi^x(t) = \phi_t(x)$. The infinitesimal generator of a flow ϕ is the vector field V on M defined for $x \in M$ by

$$V_x = \frac{d}{dt}\Big|_{t=0} \phi^x(t).$$

It is a fact that every vector field on M is the infinitesimal generator of a flow on M, and furthermore that there is a unique flow whose domain is maximal that has that vector field as its infinitesimal generator, and we thus speak of *the* flow of a vector field.

We say that a vector field is *complete* if it is the infinitesimal generator of a flow whose domain is $\mathbb{R} \times M$, in other words if it is the infinitesimal generator of a *global flow*. It is a fact that if V is a vector field on a compact smooth manifold then V is complete.

4 Hamiltonian flows

Let (M, ω) be a symplectic manifold. We say that a vector field X on M is symplectic if

$$\mathcal{L}_X\omega=0$$
,

where $\mathcal{L}_X \omega$ is the Lie derivative of ω along the flow of X. A Hamiltonian flow is the flow of a Hamiltonian vector field.³ If X is a complete symplectic vector field and $\phi: M \times \mathbb{R}$ is the flow of X, then for all $t \in \mathbb{R}$, the map $\phi_t: M \to M$ is a symplectomorphism.

Let $H \in C^{\infty}(M)$, and let ϕ be the flow of the vector field X_H . If (x, s) is in the domain of the flow ϕ , we have

$$\frac{d}{dt}\Big|_{t=s} H(\phi^{x}(t)) = (d_{\phi^{x}(s)}H)((\phi^{x})'(s))
= (d_{\phi^{x}(s)}H)(X_{H}(\phi^{x}(s)))
= \omega_{\phi^{x}(s)}(X_{H}(\phi^{x}(s)), X_{H}(\phi^{x}(s)))
= 0.$$

Thus a Hamiltonian vector field is symplectic: H does not change along the flow of X_H . We can also write this as

$$\frac{d}{dt}(H \circ \phi_t) = \frac{d}{dt}(\phi_t^* H)$$

$$= \phi_t^* (\mathcal{L}_{X_H} H)$$

$$= \phi_t^* ((i_{X_H} \omega)(X_H))$$

$$= \phi_t^* (\omega(X_H, X_H))$$

$$= \phi_t^* (0)$$

$$= 0.$$

It is a fact that if $H^1_{\mathrm{dR}}(M)=\{0\}$ (i.e. if α is a 1-form on M and $d\alpha=0$ then there is some $f\in C^\infty(M)$ such that $\alpha=df$) then every symplectic vector field on M is Hamiltonian. In particular, if M is simply connected then $H^1_{\mathrm{dR}}(M)=\{0\}$, and hence if M is simply connected then every symplectic vector field on M is Hamiltonian.

5 Poisson bracket

For $f, g \in C^{\infty}(M)$, we define $\{f, g\} \in C^{\infty}(M)$ for $x \in M$ by

$$\{f,g\}(x) = \omega_x(X_f(x), X_g(x)).$$

This is called the *Poisson bracket* of f and g. We write

$$\{f,g\} = \omega(X_f,X_g).$$

We have

$$\{f,g\} = X_f g = (df)X_g.$$

We say that f and g Poisson commute if $\{f,g\} = 0$. The Poisson bracket of f and g tells us how f changes along the Hamiltonian flow of g. If f and g Poisson commute then f does not change along the flow of X_g .

 $^{^{3}}$ cf. gradient flow.

We have

$$\{f,g\} = \omega(X_f, X_g)$$

$$= \sum_{i=1}^n (dq_i \wedge dp_i) \sum_{j=1}^n \left(\frac{\partial f}{\partial p_j} \frac{\partial}{\partial q_j} - \frac{\partial f}{\partial q_j} \frac{\partial}{\partial p_j} \right) \sum_{k=1}^n \left(\frac{\partial g}{\partial p_k} \frac{\partial}{\partial q_k} - \frac{\partial g}{\partial q_k} \frac{\partial}{\partial p_k} \right)$$

$$= \sum_{i=1}^n \left(\frac{\partial f}{\partial p_i} dp_i + \frac{\partial f}{\partial q_i} dq_i \right) \sum_{k=1}^n \left(\frac{\partial g}{\partial p_k} \frac{\partial}{\partial q_k} - \frac{\partial g}{\partial q_k} \frac{\partial}{\partial p_k} \right)$$

$$= \sum_{i=1}^n -\frac{\partial f}{\partial p_i} \frac{\partial g}{\partial q_i} + \frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i}.$$

If $x \in M$ and $v \in T_xM$, then vf is the directional derivative in the direction v. If $v = \sum_{i=1}^n a_i \frac{\partial}{\partial q_i} + b_i \frac{\partial}{\partial p_i}$ and $f \in C^{\infty}(M)$ then

$$vf = \sum_{i=1}^{n} a_i \frac{\partial f}{\partial q_i} + b_i \frac{\partial f}{\partial p_i}.$$

If X is a vector field on M then $Xf \in C^{\infty}(M)$, defined for $x \in M$ by

$$(Xf)(x) = X_x f.$$

If τ is a covariant tensor field and X is a vector field, the Lie derivative of τ along the flow of X is defined as follows: if ϕ is the flow of X, then

$$(\mathcal{L}_X \tau)(x) = \frac{d}{dt} \Big|_{t=0} (\phi_t^* \tau)(x),$$

and so if τ is a function $f \in C^{\infty}(M)$, then

$$(\mathcal{L}_X f)(x) = \frac{d}{dt}\Big|_{t=0} (\phi_t^* f)(x) = \frac{d}{dt}\Big|_{t=0} f(\phi_t(x)) = X_x f = (Xf)(x).$$

Thus if X is a vector field and $f \in C^{\infty}(M)$, then $\mathcal{L}_X f = X f$. For $f, g \in C^{\infty}(M)$,

$$\begin{split} X_{\{f,g\}} \lrcorner \omega &= d\{f,g\} \\ &= d(X_g f) \\ &= d(\mathcal{L}_{X_g} f) \\ &= \mathcal{L}_{X_g} (df) \\ &= \mathcal{L}_{X_g} (X_f \lrcorner \omega) \\ &= (\mathcal{L}_{X_g} X_f) \lrcorner \omega + X_f \lrcorner \mathcal{L}_{X_g} \omega \\ &= [X_g, X_f] \lrcorner \omega + X_f \lrcorner 0 \\ &= [X_g, X_f] \lrcorner \omega \\ &= -[X_f, X_g] \lrcorner \omega. \end{split}$$

Since the symplectic form ω is nondegenerate, if $X \bot \omega = Y \bot \omega$ then X = Y, so

$$X_{\{f,g\}} = -[X_f, X_g].$$

It follows that $C^{\infty}(M)$ is a Lie algebra using the Poisson bracket as the Lie bracket.

The set $\Gamma^\infty(M)$ of vector fields on M are a Lie algebra using the vector field commutator $[\cdot,\cdot]$. The symplectic vector fields are a Lie subalgebra: it is clear that they are a linear subspace of the Lie algebra of vector fields, and one shows that the commutator of two symplectic vector fields is itself a symplectic vector field. One can further show that the set of Hamiltonian vector fields is a Lie subalgebra of the Lie algebra of symplectic vector fields. It is a fact that the vector space quotient of the vector space of symplectic vector fields modulo the vector space of Hamiltonian vector fields is isomorphic to the vector space $H^1_{\mathrm{dR}}(M)$; this is why if $H^1_{\mathrm{dR}}(M) = \{0\}$ (in particular if M is simply connected) then any symplectic vector field on M is Hamiltonian.

6 Tautological 1-form

Let Q be a smooth manifold and let $\pi: T^*Q \to Q$, $\pi(q, p) = q$. For $x = (q, p) \in T^*Q$, we have

$$d_x\pi: T_xT^*Q \to T_qQ.$$

Let

$$\theta_x = (d_x \pi)^*(p) = p \circ d_x \pi : T_x T^* Q \to \mathbb{R}.$$

Thus $\theta: T^*Q \to T^*T^*Q$. θ is called the tautological 1-form on T^*Q .

If (Q_1, \ldots, Q_n) are coordinates on an open subset U of Q, $Q_i : U \to \mathbb{R}$, then for each $q \in U$ we have that $d_q Q_i \in T_q^* U = T_q^* Q$, $1 \le i \le n$, are a basis for $T_q^* Q$ and $\frac{\partial}{\partial Q_i}\Big|_{q}$, $1 \le i \le n$, are a basis for $T_q Q$. For each $p \in T_q^* Q$,

$$p = \sum_{i=1}^{n} p \left(\frac{\partial}{\partial Q_i} \Big|_q \right) d_q Q_i.$$

On T^*U , define coordinates $(q_1, \ldots, q_n, p_1, \ldots, p_n)$ by

$$q_i(q, p) = Q_i(q),$$

and

$$p_i(q, p) = p\left(\frac{\partial}{\partial Q_i}\Big|_q\right).$$

On T^*U we can write θ using these coordinates: for $x=(q,p)\in T^*Q$,

$$\theta_x = p \circ d_x \pi = \sum_{i=1}^n p_i(x) d_x q_i.$$

Thus, on T^*U ,

$$\theta = \sum_{i=1}^{n} p_i dq_i.$$

Let $\omega = -d\theta$. We have, on T^*U ,

$$\omega = -d\sum_{i=1}^{n} p_i dq_i$$

$$= -\sum_{i=1}^{n} (dp_i \wedge dq_i + p_i d(dq_i))$$

$$= -\sum_{i=1}^{n} dp_i \wedge dq_i$$

$$= \sum_{i=1}^{n} dq_i \wedge dp_i.$$

 T^*Q is a symplectic manifold with the symplectic form ω .

7 Cotangent lifts

Let Q be a smooth manifold and let $F:Q\to Q$ be a diffeomorphism. Define

$$F^{\sharp}: T^*Q \to T^*Q$$

for x = (q, p) by

$$F^{\sharp}(q,p) = (F(q), (d_{F(q)}(F^{-1}))^*(p)).$$

We call $F^{\sharp}: T^*Q \to T^*Q$ the *cotangent lift* of $F: Q \to Q$. It is a fact that it is a diffeomorphism. It is apparent that the diagram

$$T^*Q \xrightarrow{F^{\sharp}} T^*Q$$

$$\downarrow^{\pi} \qquad \downarrow^{\pi}$$

$$Q \xrightarrow{F} Q$$

commutes.

The pull-back of θ by F^{\sharp} satisfies, for $x=(q,p)\in T^*Q$ and $(\zeta,\eta)=F^{\sharp}(q,p)\in T^*Q,$

$$((F^{\sharp})^{*}\theta)_{x} = (d_{x}F^{\sharp})^{*}(\theta_{F^{\sharp}(x)})$$

$$= (d_{x}F^{\sharp})^{*}((d_{F^{\sharp}(x)}\pi)^{*}(\eta))$$

$$= (d_{x}(\pi \circ F^{\sharp}))^{*}(\eta)$$

$$= (d_{x}(F \circ \pi))^{*}(\eta)$$

$$= (d_{x}\pi)^{*}((d_{\pi(x)}F)^{*}(\eta))$$

$$= (d_{x}\pi)^{*}((d_{q}F)^{*}(\eta))$$

$$= (d_{x}\pi)^{*}(p)$$

$$= \theta_{x}.$$

Thus $(F\sharp)^*\theta = \theta$, i.e. F^\sharp pulls back θ to θ . The "naturality of the exterior derivative" 4 is the statement that if G is a smooth map and η is a differential form then $G^*(d\eta) = d(G^*\eta)$. Hence, with $\omega = d\theta$,

$$(F^{\sharp})^*\omega = (F^{\sharp})^*(d\theta) = d((F^{\sharp})^*\theta) = d(\theta) = \omega,$$

so F^{\sharp} pulls back the symplectic form ω to itself. Thus $F^{\sharp}: T^{*}M \to T^{*}M$ is a symplectomorphism.

Let $\operatorname{Diff}(Q)$ be the set of diffeomorphisms $Q \to Q$. $\operatorname{Diff}(Q)$ is a group. Let G be a group and let $\tau: G \to \operatorname{Diff}(Q)$ be a homomorphism. Define $\tau^{\sharp}: G \to \operatorname{Diff}(T^{*}Q)$ by $(\tau^{\sharp})_{g} = (\tau_{g})^{\sharp}: T^{*}Q \to T^{*}Q$. $\tau^{\sharp}: G \to \operatorname{Diff}(T^{*}Q)$ is a homomorphism, and for each $g \in G$, $(\tau^{\sharp})_{g}: T^{*}Q \to T^{*}Q$ is a symplectomorphism. In words, if a group acts by diffeomorphisms on a smooth manifold, then the cotangent lift of the action is an action by symplectomorphisms on the cotangent bundle.

8 Lie groups

Recall that if $F: M \to N$ then $TF: TM \to TN$ satisfies, for $X \in \Gamma^{\infty}(M)$ and $f \in C^{\infty}(N)$,⁵

$$((TF)X)(f) = X(f \circ F),$$

i.e. for $x \in M$ and $v \in T_xM$,

$$((T_x F)v)(f) = v(f \circ F),$$

the directional derivative of $f \circ F \in C^{\infty}(M)$ in the direction of the tangent vector v.

⁴For each k, Ω^k is a contravariant functor, and if $f:M\to N$, then the functor Ω^k sends f to $f^*:\Omega^k(N)\to\Omega^k(M)$. d is a natural transformation from the contravariant functor Ω^k to the contravariant functor Ω^{k+1} .

⁵In words: TF pushes forward a vector field on M to a vector field on N.

Let G be a Lie group and for $g \in G$ define $L_g : G \to G$ by $L_g h = gh$. If X is a vector field on G, we say that X is *left-invariant* if

$$(T_h L_a)(X_h) = X_{ah}$$

for all $g, h \in G$. That is, X is left-invariant if

$$(TL_q)(X) = X$$

for all $g \in G$.

If X and Y are left-invariant vector fields on G then so is [X,Y]. This is because, for $F:G\to G$,

$$(TF)[X,Y] = [(TF)X, (TF)Y].$$

Thus the set of left-invariant vector fields on G is a Lie subalgebra of the Lie algebra of vector fields on G.

Define $\epsilon: \text{Lie}(G) \to T_eG$ by $\epsilon(X) = X_e$, where $e \in G$ is the identity element. It can be shown that this is a linear isomorphism. Hence, if $v \in T_eG$ then there is a unique left-invariant vector field X on G such that, for all $g \in G$,

$$V_q = (T_e L_q)(v).$$

It is a fact that every left-invariant vector field on a Lie group G is complete, i.e. that its flow has domain $G \times \mathbb{R}$. For $X \in \text{Lie}(G)$, we call the unique integral curve of X that passes through e the one-parameter subgroup generated by X. Thus, for any $v \in T_eG$ there is a unique one-parameter subgroup $\gamma : \mathbb{R} \to G$ such that

$$\gamma(0) = e, \qquad \gamma'(0) = v.$$

We define $\exp : \operatorname{Lie}(G) \to G$ by $\exp(X) = \gamma(1)$, where γ is the one-parameter subgroup generated by X. This is called the *exponential map*. Thus $t \mapsto \exp(tX)$ is the one-parameter subgroup generated by X.

Fact: If (TF)X = Y and X has flow ϕ and Y has flow η , then

$$\eta_t \circ F = F \circ \phi_t$$

for all t in the domain of ϕ . Hence

$$L_a \circ \phi_t = \phi_t \circ L_a$$
.

Hence the flow ϕ of a left-invariant vector field X satisfies

$$g \exp(tX) = L_g \exp(tX)$$

$$= L_g(\phi_t e)$$

$$= \phi_t(L_g e)$$

$$= \phi_t(g).$$

9 Coadjoint action

First we'll define the adjoint action of G on $\mathfrak{g} = T_{\mathrm{id}_G}G$. For $g \in G$, define $\Psi_g: G \to G$ by $\Psi_g(h) = ghg^{-1}$; Ψ_g is an automorphism of Lie groups. Define

$$\mathrm{Ad}_q:\mathfrak{g}\to\mathfrak{g}$$

by

$$Ad_g = T_{id_G} \Psi_g;$$

since Ψ_g is an automorphism of Lie groups, it follows that Ad_g is an automorphism of Lie algebras. We can also write Ad_g as

$$\operatorname{Ad}_g(\xi) = \frac{d}{dt}\Big|_{t=0} (g \exp(t\xi)g^{-1}).$$

The adjoint action of G on \mathfrak{g} is

$$g \cdot \xi = \mathrm{Ad}_g(\xi).$$

For each $g \in G$, one proves that there is a unique map $\operatorname{Ad}_g^* : \mathfrak{g}^* \to \mathfrak{g}^*$ such that for all $l \in \mathfrak{g}^*, \xi \in \mathfrak{g}$,

$$(\mathrm{Ad}_{g}^{*}l)(\xi) = l(\mathrm{Ad}_{g}(\xi)).$$

The *coadjoint action* of G on \mathfrak{g}^* is

$$g \cdot l = \operatorname{Ad}_{q^{-1}}^*(l).$$

10 Momentum map

Let (M,ω) be a symplectic manifold, let G be a Lie group, and let $\sigma:G\to \mathrm{Diff}(M)$ be a homomorphism such that for each g in G, σ_g is a symplectomorphism.

Let $\mathfrak{g} = T_{\mathrm{id}_G}G$, and define $\rho : \mathfrak{g} \to \Gamma^{\infty}(M)$ by

$$\rho(\xi)(x) = \frac{d}{dt}\Big|_{t=0} \sigma_{\exp(t\xi)}(x) \in T_x M, \qquad \xi \in \mathfrak{g}, x \in M;$$

 $t\mapsto \sigma_{\exp(t\xi)}(x)$ is $\mathbb{R}\to M$ and at t=0 the curve passes through x, so indeed $\rho(\xi)(x)\in T_x(M)$. ρ is called the *infinitesimal action* of $\mathfrak g$ on M. Each element of G acts on M as a symplectomorphism, each element of $\mathfrak g$ acts on M as a vector field.

A momentum map for the action of G on (M, ω) is a map $\mu : M \to \mathfrak{g}^*$ such that, for $x \in M$, $v \in T_xM$ and $\xi \in \mathfrak{g}$,

$$((T_x \mu)v)\xi = \omega_x(\rho(\xi)(x), v), \tag{1}$$

where

$$T_x \mu : T_x M \to T_{\mu(x)} \mathfrak{g}^* = \mathfrak{g}^*,$$

and such that if $g \in G$ and $x \in M$ then

$$\mu(\sigma_g(x)) = g \cdot \mu(x),\tag{2}$$

where $g \cdot \mu(x)$ is the *coadjoint action* of G on \mathfrak{g}^* , defined in section §9; we say that μ is *equivariant* with respect to the coadjoint action of G on \mathfrak{g}^* .

11 Angular momentum

Let $G = SO(3) = \{A \in \mathbb{R}^{3 \times 3} : A^TA = I, \det(A) = 1\}$. The Lie algebra of SO(3) is

$$\mathfrak{g} = \mathfrak{so}(3) = \{ a \in \mathbb{R}^{3 \times 3} : a + a^T = 0 \}.$$

Let $Q = \mathbb{R}^3$, and define $\tau : G \to \text{Diff}(Q)$ by $\tau_q(q) = gq$.

Let θ be the tautological 1-form on T^*Q and let $\omega = -d\theta$. (T^*Q, ω) is a symplectic manifold and $\tau^{\sharp}: G \to \mathrm{Diff}(T^*Q)$ is a homomorphism such that for each $g \in G$, $(\tau^{\sharp})_g$ is a symplectomorphism. For $g \in G$, $(q, p) \in T^*Q$,

$$\begin{aligned} (\tau^{\sharp})_g(q,p) &=& (\tau_g)^{\sharp}(q,p) \\ &=& (\tau_g q, (d_{\tau_g q}(\tau_g^{-1}))^* p) \\ &=& (\tau_g q, (d_{\tau_g q}(\tau_{g^{-1}}))^* p) \\ &=& (\tau_g q, p \circ (d_{\tau_g q} \tau_{g^{-1}})) \\ &=& (\tau_g q, p \circ \tau_{g^{-1}}) \\ &=& (gq, pg^{-1}) \\ &=& (gq, pg^T). \end{aligned}$$

Hence for $\xi \in \mathfrak{g}$ and $(q, p) \in T^*Q$,

$$\begin{split} \rho(\xi)(q,p) &= \left. \frac{d}{dt} \right|_{t=0} (\tau^{\sharp})_{\exp(t\xi)}(q,p) \\ &= \left. \frac{d}{dt} \right|_{t=0} (\exp(t\xi)q, p \exp(t\xi^T)) \\ &= \left. (\xi q, p \xi^T) \right. \\ &= \left. (\xi q, -p \xi). \end{split}$$

Define $V: \mathfrak{g} \to \mathbb{R}^3$ by

$$V \begin{pmatrix} 0 & -\xi_3 & \xi_2 \\ \xi_3 & 0 & -\xi_1 \\ -\xi_2 & \xi_1 & 0 \end{pmatrix} = \begin{pmatrix} \xi_1 \\ \xi_2 \\ \xi_3 \end{pmatrix}.$$

One checks that $\xi q = V(\xi) \times q$ and $p\xi = p^T \times V(\xi)$.

For $(q,p) \in T^*Q$, $(v,w) \in T_{(p,q)}T^*Q$, and $\xi \in \mathfrak{g}$, we have

$$\begin{split} \omega_{(q,p)}(\rho(\xi)(q,p),(v,w)) &= & \omega_{(q,p)}((\xi q,-p\xi),(v,w)) \\ &= & \sum_{j=1}^{3} dq_{j} \wedge dp_{j}((\xi q,-p\xi),(v,w)) \\ &= & \sum_{j=1}^{3} \left((\xi q)_{j} dp_{j} + (p\xi)_{j} dq_{j} \right) (v,w) \\ &= & \sum_{j=1}^{3} w_{j}(\xi q)_{j} + v_{j}(p\xi)_{j} \\ &= & w \cdot (V(\xi) \times q) + v \cdot (p^{T} \times V(\xi)). \end{split}$$

Define $\mu: T^*Q \to \mathfrak{g}^*$ by $\mu(q,p)(\xi) = (q \times p^T) \cdot V(\xi)$. I claim that μ satisfies (1) and (2). We have just calculated the right-hand side of (1), so it remains to calculate the left-hand side. I find the left-hand side unwieldly to calculate in a clean and precise way, so I will merely claim that it is equal to the right-hand side. I have convinced myself that it is true by symbol pushing.

For $g \in G$ and $\xi \in \mathfrak{g}$, $\mathrm{Ad}_{g}\xi = g\xi g^{-1}$, and hence, for $(q,p) \in T^{*}Q$,

$$(g \cdot \mu(q, p))\xi = (\mathrm{Ad}_{g^{-1}}^* \mu(q, p)) \xi$$

= $\mu(q, p)(\mathrm{Ad}_{g^{-1}} \xi)$
= $\mu(q, p)(g^{-1} \xi g)$
= $(q \times p^T) \cdot V(g^{-1} \xi g)$.

On the other hand,

$$\begin{split} \mu((\tau^{\sharp})_g(q,p))\xi &=& \mu(gq,pg^T)\xi \\ &=& ((gq)\times(pg^T)^T)\cdot V(\xi) \\ &=& ((gq)\times(gp^T))\cdot V(\xi) \\ &=& (g(q\times p^T))\cdot V(\xi) \\ &=& (q\times p^T)\cdot (g^TV(\xi)) \\ &=& (q\times p^T)\cdot (g^{-1}V(\xi)). \end{split}$$