Robust Computational Models for Water Waves

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Australian National University

October 1, 2018

Outline of the Presentation

- Motivation
- History
- Contribution
  - Method
  - Validation

Water Waves

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Water wave hazards:

- Tsunamis

Tsunamis

Figure: 2004 Indian Ocean Tsunami (Banda Aceh).

Tsunamis

Figure: 2011 Tohoku Tsunami.

Water Waves

Water wave hazards:

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- Storm Surges

Storm Surges

Figure: 2012 Hurricane Sandy Storm Surge.

Water Waves

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- Tsunamis
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Phenomena caused by water waves:

- Nutrient Transport
- Beach Erosion
- Breakup of Sea Ice

Computational Modelling

Goal: Model physics on computers.

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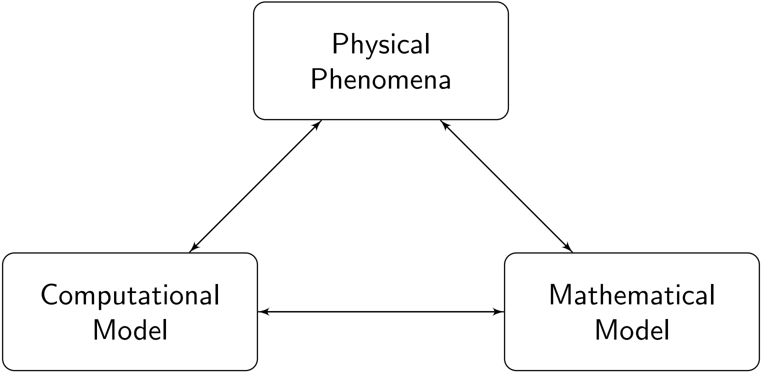
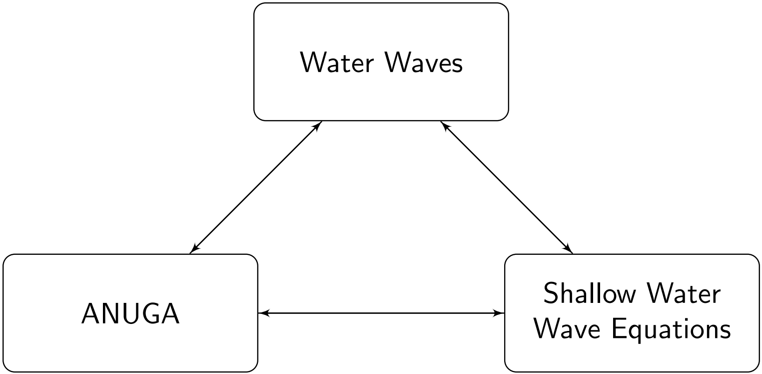
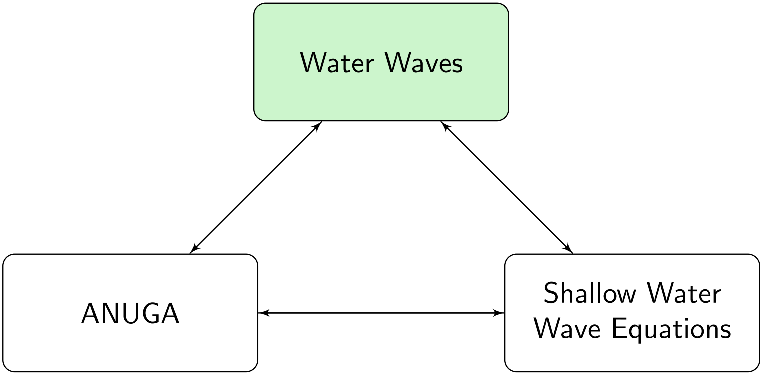


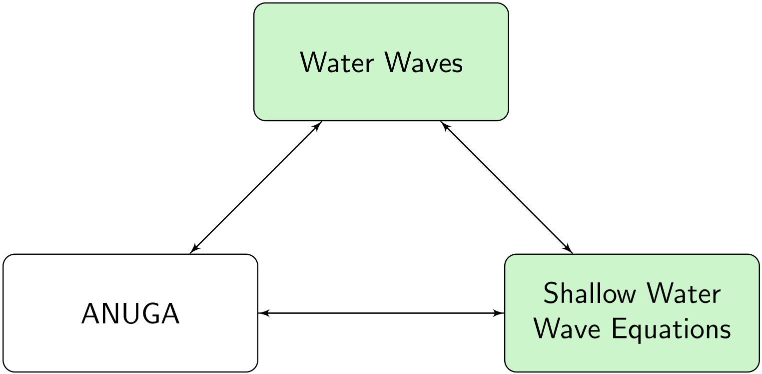
ANUGA



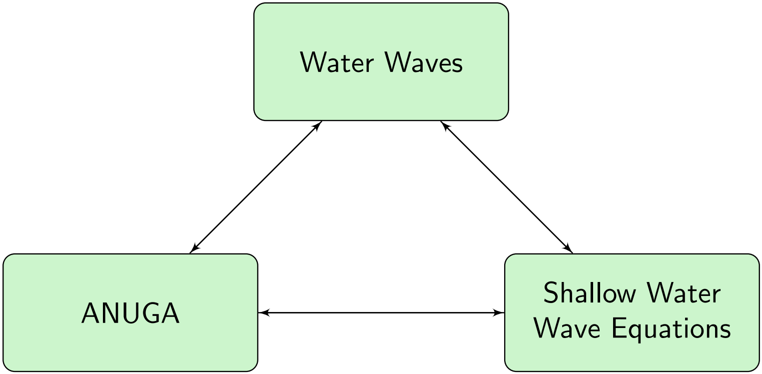
ANUGA: Water Waves



ANUGA: Shallow Water Wave Equations



ANUGA



Outcome

New project at the ANU to develop a robust computational model for the Serre equations.

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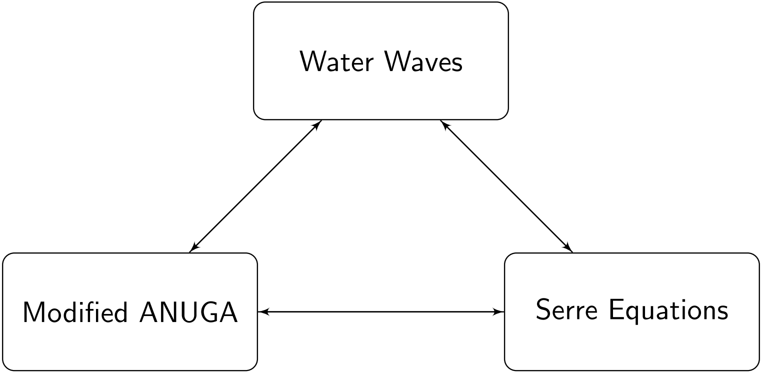
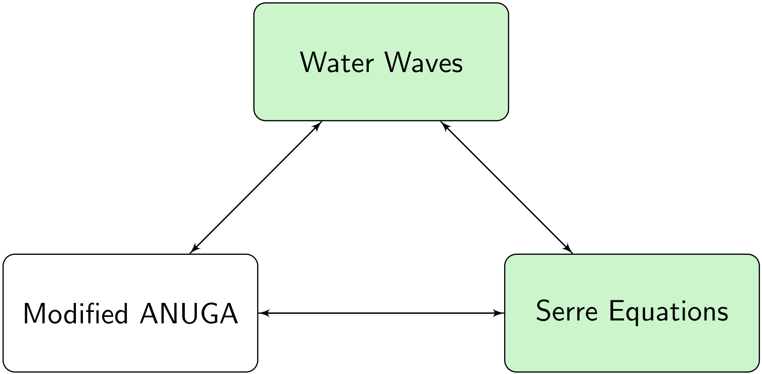
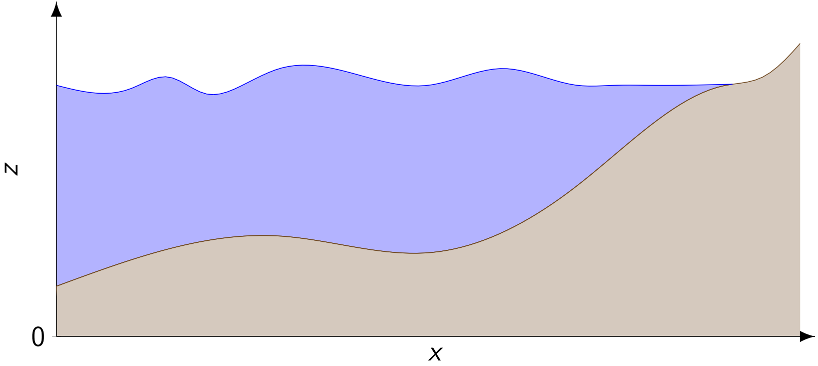


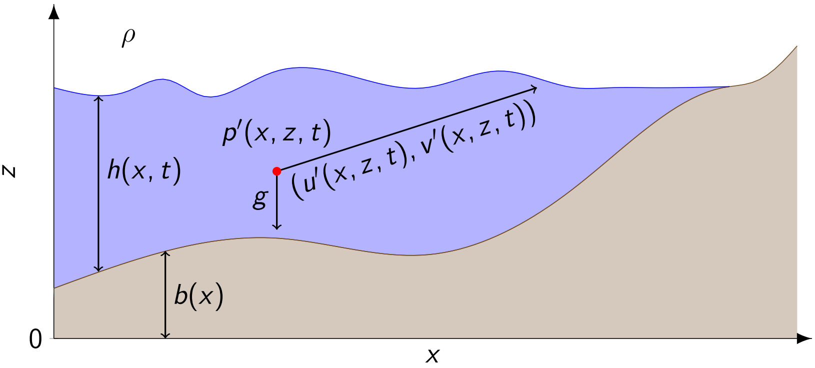
Mathematical Model



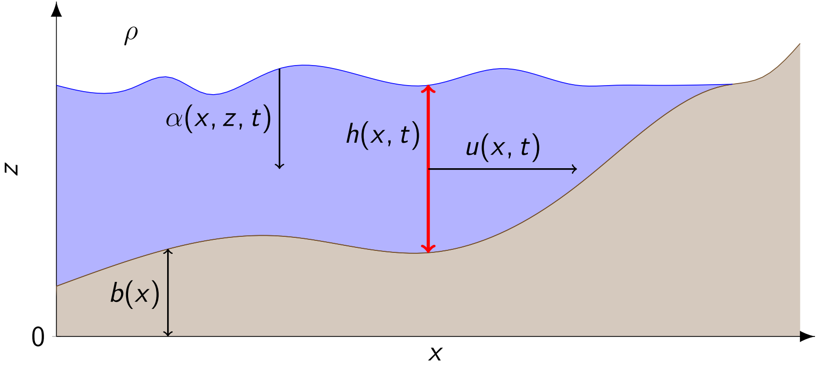
Typical Scenario



Navier Stokes Model



Serre Model



Assumptions

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Quantity	Shallow Water Wave Equations	Serre Equations
Particle: $v'(x, z, t)$	0	$u\frac{\partial b}{\partial x} - (h - \alpha)\frac{\partial b}{\partial x}$

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Particle: $p'(x,z,t)$	$g\rho\alpha$	$g\rho\alpha + \rho\alpha\Psi + \frac{1}{2}\rho\alpha\left(2h-\alpha\right)\Phi$

where

$$\alpha(x,z,t) = (h(x,t) + b(x)) - z$$

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and

$$\Psi = \frac{\partial b}{\partial x}\left(\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x}\right) + u^2\frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x}\frac{\partial u}{\partial x} - u\frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

Equations

Mass:

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

Momentum:

$$\frac{\partial (uh)}{\partial t} + \frac{\partial}{\partial x}\left(u^2 h + \frac{gh^2}{2} + \frac{h^2}{2}\Psi + \frac{h^3}{3}\Phi\right)$$

$$+ \frac{\partial b}{\partial x}\left(gh + h\Psi + \frac{h^2}{2}\Phi\right) = 0.$$

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Pros and Cons

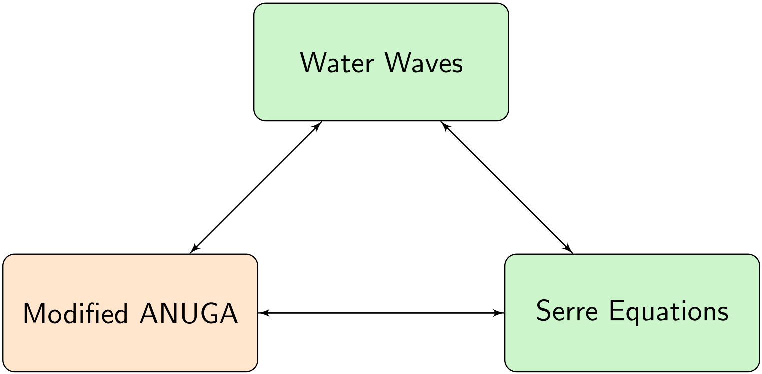
Pros:

- Includes dispersive effects
- Considered one of the best models for water waves
- Can apply techniques of ANUGA

Cons:

- More complex than the Shallow Water Wave Equations

Computational Model



Previous Work at the ANU

- 2014: Chris Zoppou's PhD thesis
  Developed computational model for the 1D Serre equations.
- 2014: My Honours thesis
  Independent reproduction of Chris Zoppou's computational model.

Open Problems and Thesis Goals

2D: 1D method that extends well to 2D

Robust: Validation for steep gradients in free surface

Robust: Inclusion and validation of dry beds

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2D: 1D method that extends well to 2D

Robust: Validation for steep gradients in free surface

Robust: Inclusion and validation of dry beds

Technique: Develop a robust computational model from the 1D Serre equations that can be easily extended to 2D.

Finite Volume Method

2D: Extends well to 2D

Robust: Stable in the presence of steep gradients

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- Maintains conservation properties of the equations
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Chris Zoppou's thesis demonstrated an adaptation of the Finite Volume Method to solve the Serre Equations.

Equations

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial (uh)}{\partial t} + \frac{\partial}{\partial x}\left(u^2 h + \frac{gh^2}{2} + \frac{h^2}{2}\Psi + \frac{h^3}{3}\Phi\right) + \frac{\partial b}{\partial x}\left(gh + h\Psi + \frac{h^2}{2}\Phi\right) = 0,$$

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For a Finite Volume Method we require equations in the form

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where $f(q)$ and $s(q)$ do not contain temporal derivatives

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Reformulation

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x}\left(uG + \frac{gh^2}{2} - \frac{2}{3}h^3\left[\frac{\partial u}{\partial x}\right]^2 + h^2 u \frac{\partial u}{\partial x}\frac{\partial b}{\partial x}\right)$$
$$+\frac{1}{2}h^2 u \frac{\partial u}{\partial x}\frac{\partial^2 b}{\partial x^2} - hu^2\frac{\partial b}{\partial x}\frac{\partial^2 b}{\partial x^2} + gh\frac{\partial b}{\partial x} = 0.$$

with

$$G = hu\left(1 + \frac{\partial h}{\partial x}\frac{\partial b}{\partial x} + \frac{1}{2}h\frac{\partial^2 b}{\partial x^2} + \left[\frac{\partial b}{\partial x}\right]^2\right) - \frac{\partial}{\partial x}\left(\frac{1}{3}h^3\frac{\partial u}{\partial x}\right).$$

Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

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Finite Volume Method Example

Conservation of a quantity:

$$q(x, t).$$

Finite Volume Method Example

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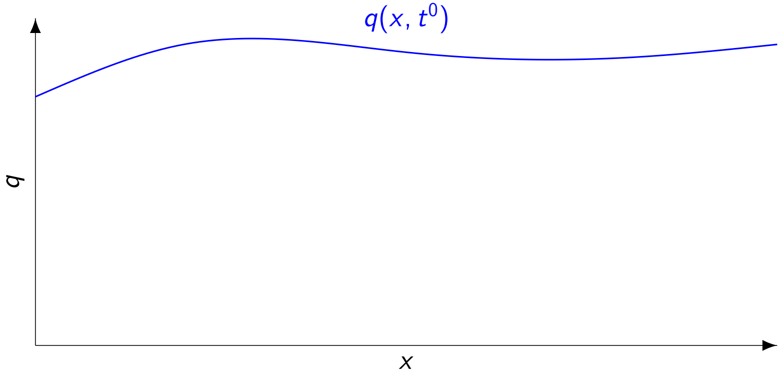
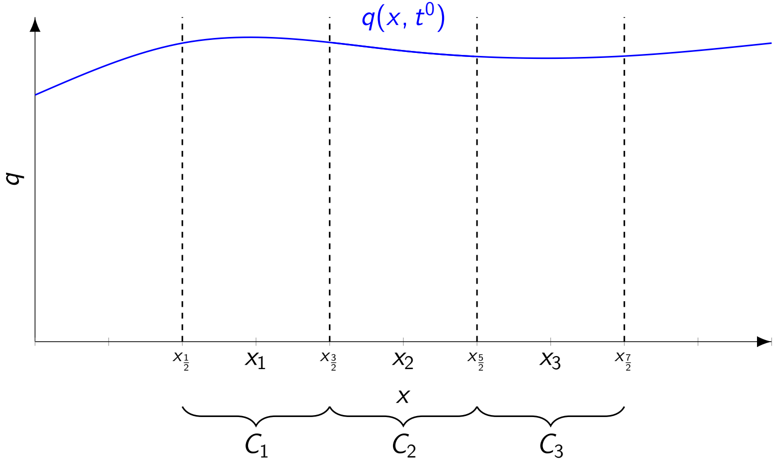
Conservation equation with a source term:

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0.$$

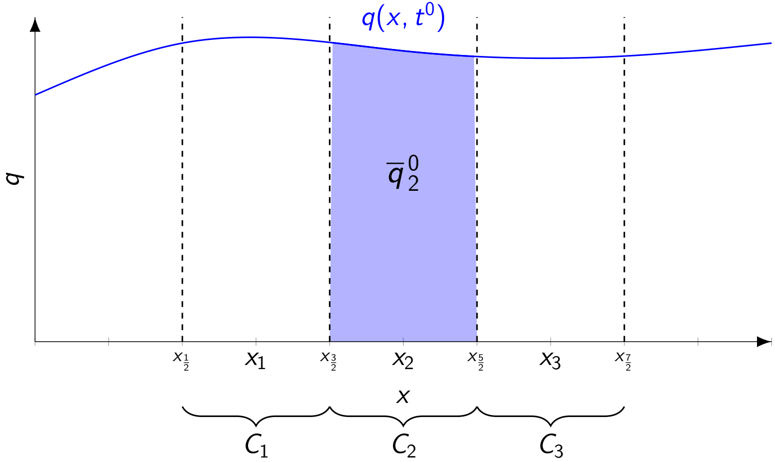
Function at $t = t^0$



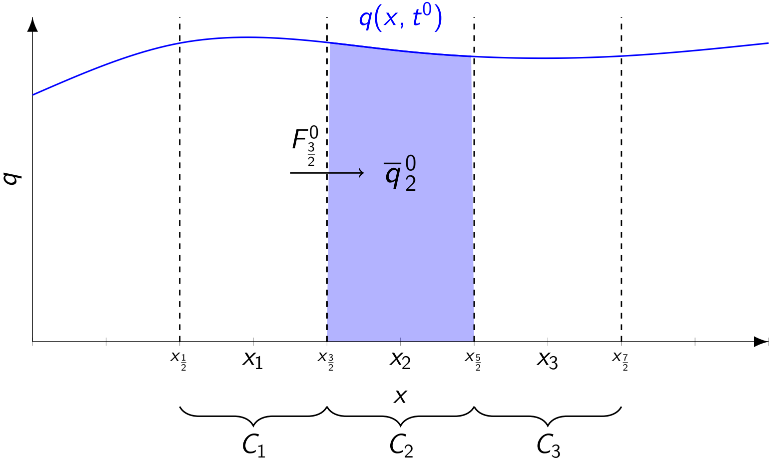
Cell Discretisation



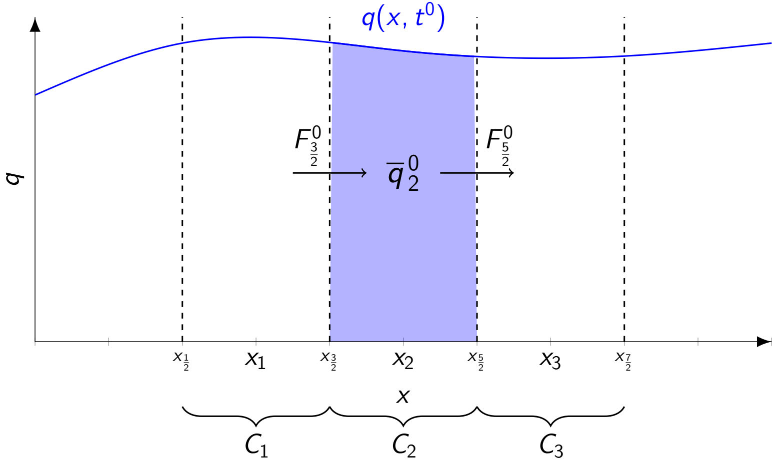
Total Amount



Flux Left


$q(x, t^0)$
$F^0_{\frac{3}{2}}$
$\overline{q}^{\,0}_{2}$
q
$x_{\frac{1}{2}}$
x_1
$x_{\frac{3}{2}}$
x_2
$x_{\frac{5}{2}}$
x_3
$x_{\frac{7}{2}}$
x
C_1
C_2
C_3

Flux Right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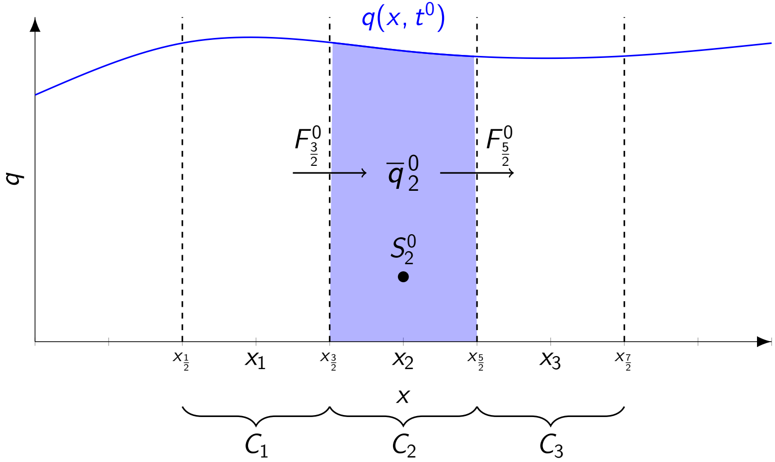
x_3
$x_{\frac{7}{2}}$
x
C_1
C_2
C_3

Source


$q(x, t^0)$
q
$F^0_{\frac{3}{2}}$
$\overline{q}^{\,0}_{2}$
$F^0_{\frac{5}{2}}$
S^0_2
$x_{\frac{1}{2}}$
x_1
$x_{\frac{3}{2}}$
x_2
$x_{\frac{5}{2}}$
x_3
$x_{\frac{7}{2}}$
x
C_1
C_2
C_3

Finite Volume Update

$$\overline{q}_2^1 = \overline{q}_2^0 - \left(F_{\frac{5}{2}}^0 - F_{\frac{3}{2}}^0\right) - \left(S_2^0\right),$$

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$$\overbrace{\int_{C_2} q(x, t^1)\, dx}^{\overline{q}_2^1} = \overbrace{\int_{C_2} q(x, t^0)\, dx}^{\overline{q}_2^0} - \Bigg(\overbrace{\int_{t^0}^{t^1} f(q(x_{5/2}, t))\, dt}^{F^0_{\frac{5}{2}}}$$

$$- \overbrace{\int_{t^0}^{t^1} f(q(x_{3/2}, t))\, dt}^{F^0_{\frac{3}{2}}} \Bigg) - \overbrace{\int_{t^0}^{t^1} \int_{C_2} s(q(x, t))\, dt}^{S^0_2}.$$

Update Formula for Serre Equations

$$\overline{h}_j^{n+1} = \overline{h}_j^{n} - \left[F^n_{j+1/2} - F^n_{j-1/2}\right],$$

$$\overline{G}_j^{n+1} = \overline{G}_j^{n} - \left[F^n_{j+1/2} - F^n_{j-1/2}\right] - S^n_j.$$

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- All the fluxes $F^n_{j+1/2}$ and $F^n_{j-1/2}$ and the source term S^n_j require u at t^n

Calculate Velocity

$$G = hu\left(1 + \frac{\partial h}{\partial x}\frac{\partial b}{\partial x} + \frac{1}{2}h\frac{\partial^2 b}{\partial x^2} + \left[\frac{\partial b}{\partial x}\right]^2\right) - \frac{\partial}{\partial x}\left(\frac{1}{3}h^3\frac{\partial u}{\partial x}\right).$$

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- Previously used a Finite Difference Method [1]
- Contribution: use a Finite Element Method

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Finite Element Method

2D: Extends well to 2D

Robust: Stable in the presence of steep gradients

Robust: Stable in the presence of dry beds

- Maintains conservation properties
- ANUGA

Finite Element Method Example

Example:

$$-\frac{\partial^2 q}{\partial x^2} = f,$$

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$$-\int_\Omega \frac{\partial^2 q}{\partial x^2} v \; dx = \int_\Omega f v \; dx,$$

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Integrate by parts

$$\int_\Omega \frac{\partial q}{\partial x}\frac{\partial v}{\partial x} \; dx = \int_\Omega fv \; dx.$$

Finite Element Method

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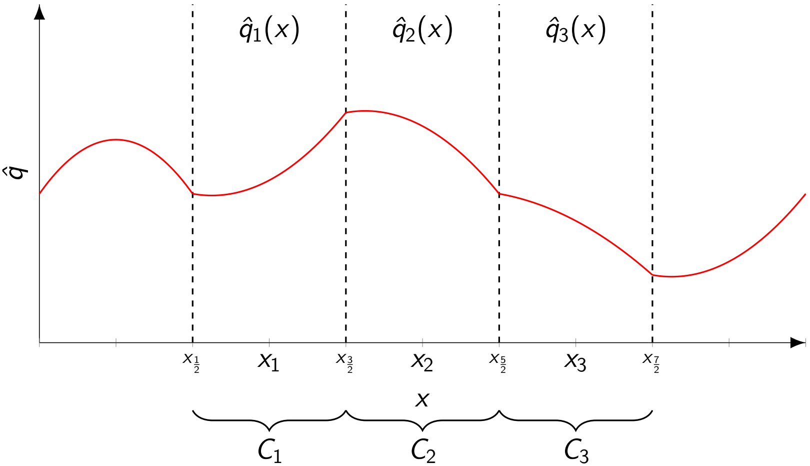
Finite Element Method

$$\int_{\Omega} \frac{\partial q}{\partial x}\frac{\partial v}{\partial x}\ dx = \int_{\Omega} fv\ dx,$$

$$\sum_{j}\left[\int_{C_j} \frac{\partial q}{\partial x}\frac{\partial v}{\partial x}\ dx\right] = \sum_{j}\left[\int_{C_j} fv\ dx\right],$$

Piecewise Polynomial Representation



Finite Element Method

$$\int_\Omega \frac{\partial q}{\partial x}\frac{\partial v}{\partial x}\ dx = \int_\Omega fv\ dx,$$

$$\sum_j \left[\int_{C_j} \frac{\partial q}{\partial x}\frac{\partial v}{\partial x}\ dx\right] = \sum_j \left[\int_{C_j} fv\ dx\right],$$

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$$\mathbf{A}\vec{q} = \vec{c}.$$

where

- $\mathbf{A}$ depends on $\hat{v}_j$
- $\vec{q}$ determines $\hat{q}_j$
- $\vec{c}$ depends on $\hat{f}_j$ and $\hat{v}_j$

Finite Element Method for Serre Equations

$$G = hu\left(1 + \frac{\partial h}{\partial x}\frac{\partial b}{\partial x} + \frac{1}{2}h\frac{\partial^2 b}{\partial x^2} + \left[\frac{\partial b}{\partial x}\right]^2\right) - \frac{\partial}{\partial x}\left(\frac{1}{3}h^3\frac{\partial u}{\partial x}\right),$$

Finite Element Method for Serre Equations

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$$\mathbf{A}\vec{u} = \vec{c}.$$

where

- $\mathbf{A}$ depends on the polynomial representation of h, b and test function.
- $\vec{u}$ determines the polynomial representation of u.
- $\vec{c}$ depends on polynomial representation of G and test function.

Method

$$
\begin{tikzcd}
\overline{h}_j^n, \overline{G}_j^n \arrow[r, "\text{Reconstruction}"] & \hat{h}_j^n, \hat{G}_j^n \arrow[r, "\text{Finite Element}"] & \hat{u}_j^n \arrow[ll, bend left, "\text{Finite Volume}"]
\end{tikzcd}
$$

Progress

2D: 1D method that extends well to 2D ✓

Robust: Validation for steep gradients in free surface

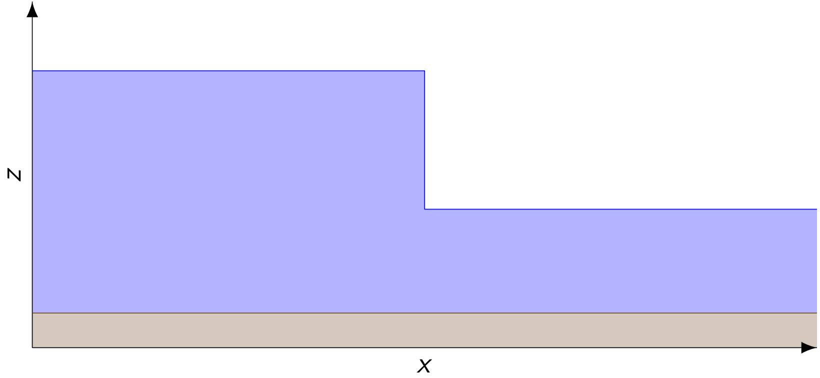
Robust: Inclusion and validation of dry beds

Validation

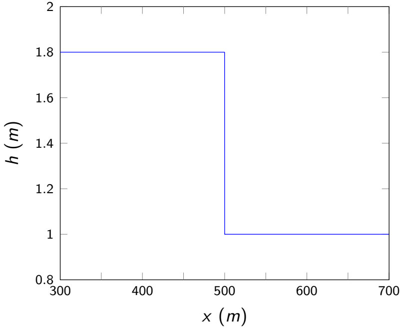
- Steep gradients in the free surface
- Dry beds

Statement of Problem

How does this initially still body of water evolve?



Our New Numerical Solution



Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

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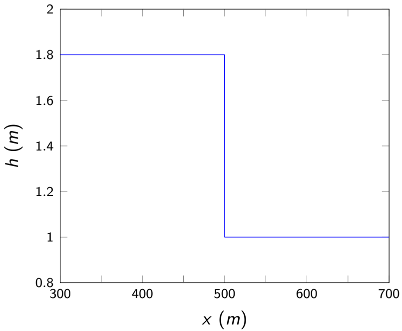
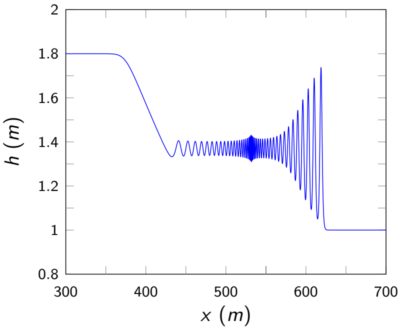




Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

Evolution

What was known

- No analytic solutions
- Some experimental comparisons [2]
- Other numerical solutions from the literature

[2] Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

Contribution

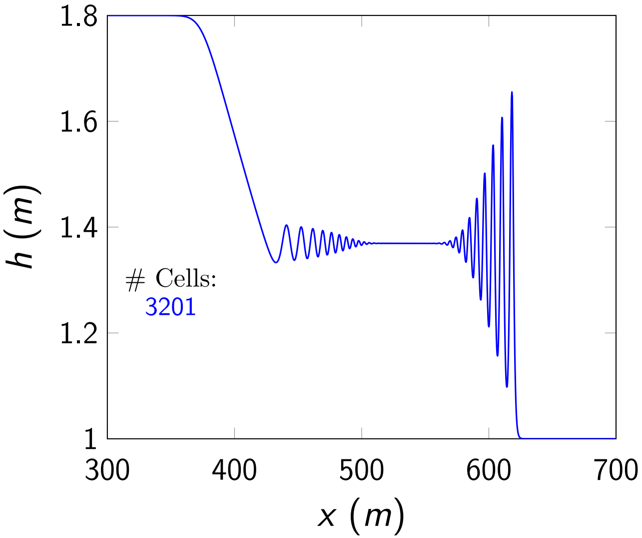
- Observed this new behaviour
- Demonstrated convergence
- Comprehensive review of numerical solutions from the literature

Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

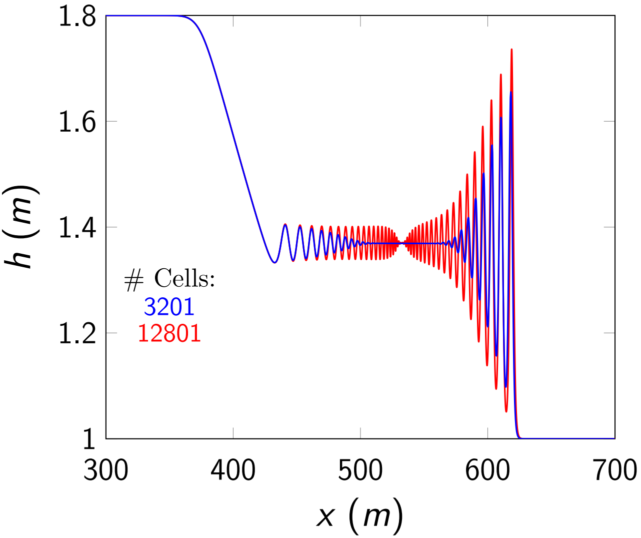
Convergence

As we increase number of cells the numerical solutions should converge to the true solution

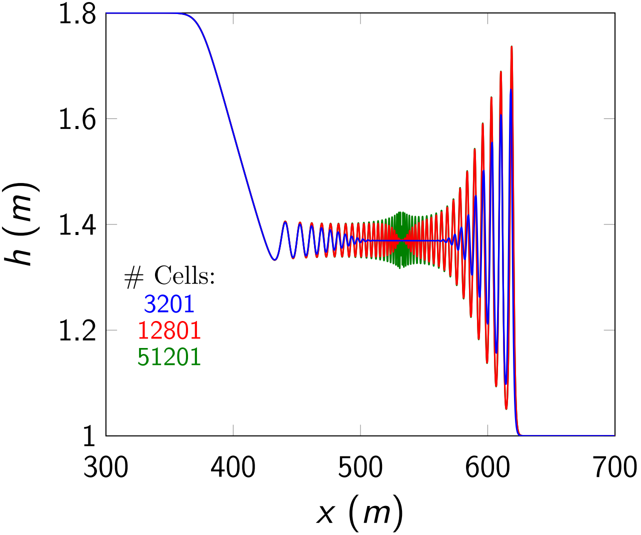
Convergence



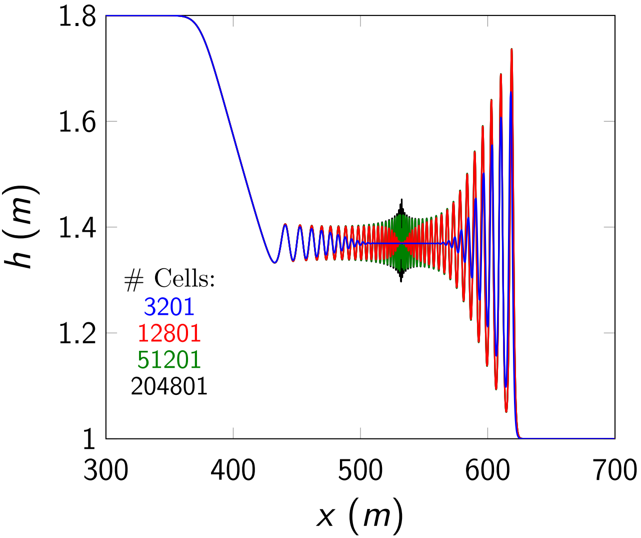
Convergence



Convergence



Convergence



Comprehensive Review

- Demonstrated consistent behaviour across many numerical methods
- Were able to explain why the behaviour had not previously been observed

Result

Validated our computational model when steep gradients are present in the free surface. [3]

[3] Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

Progress

2D: 1D method that extends well to 2D ✓

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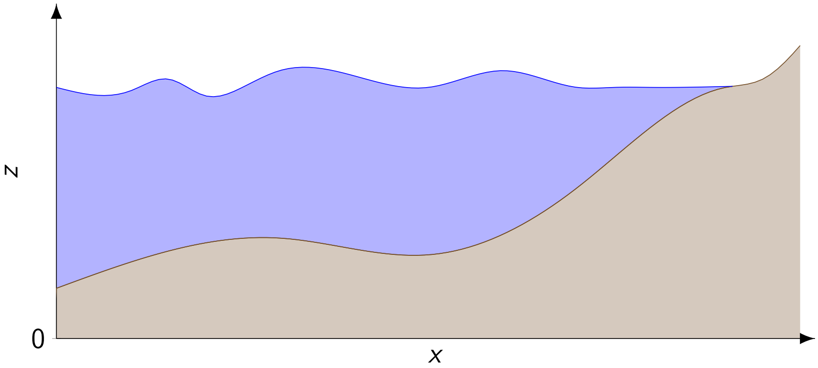
Robust: Inclusion and validation of dry beds

Statement of Problem

Properly handle interaction of waves and the dry bed



What was known

- No analytic solutions
- A variety of numerical techniques only compared to experimental data

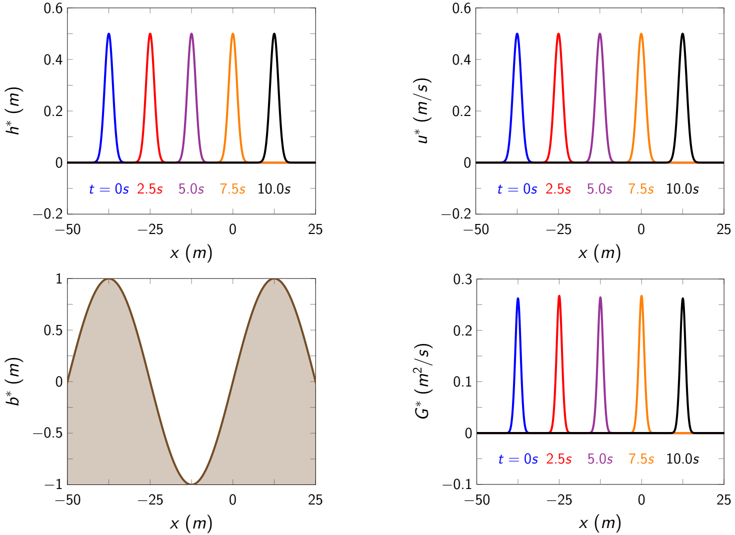
Contribution

- Solved modified equations that did possess analytic solutions
- Compared with experimental data

Constructing Modified Equations

- Pick functions for height, velocity and bed: h^*, u^* and b^*
- Add source terms to Serre equations that force h^*, u^* and b^* to be solutions
- Validation tests

Pick Functions



Modify Equations

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = S_h^*,$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x}\left(uG + \frac{gh^2}{2} - \frac{2}{3}h^3\left[\frac{\partial u}{\partial x}\right]^2 + h^2 u \frac{\partial u}{\partial x}\frac{\partial b}{\partial x}\right)$$
$$+\frac{1}{2}h^2 u \frac{\partial u}{\partial x}\frac{\partial^2 b}{\partial x^2} - hu^2\frac{\partial b}{\partial x}\frac{\partial^2 b}{\partial x^2} + gh\frac{\partial b}{\partial x} = S_G^*.$$

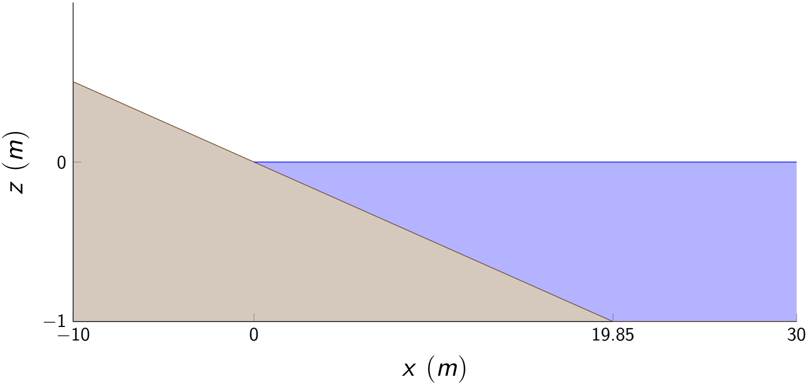
S_h^* and S_G^* are just the LHS with the quantities replaced by their associated chosen function.

Results

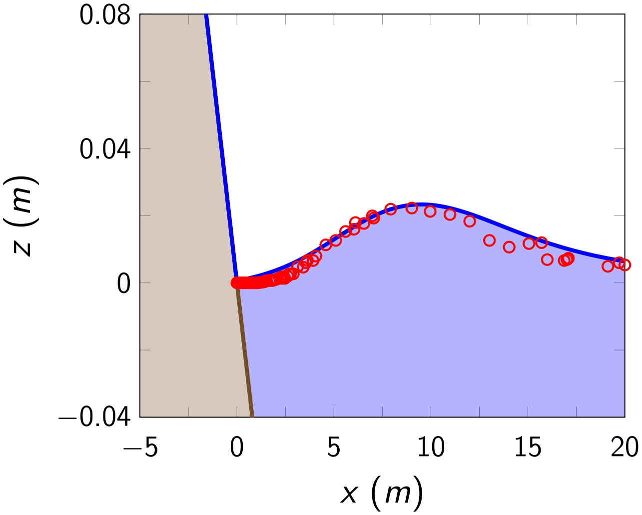
Modified Equations Validation Conclusions

- Very strong test as all terms must be accurately approximated
- Only source of error is the method

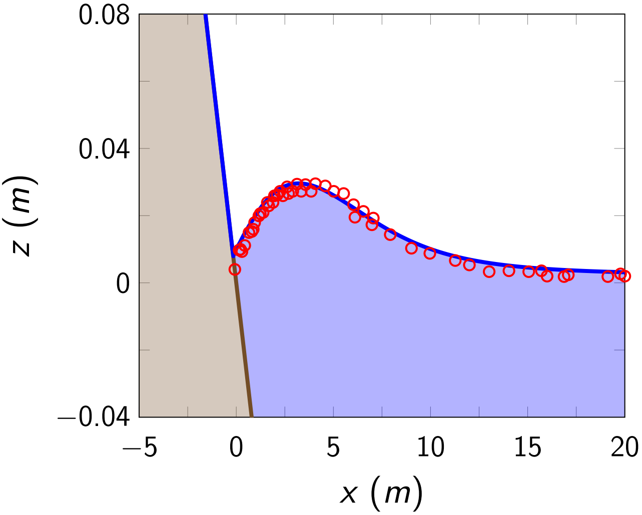
Experimental Data



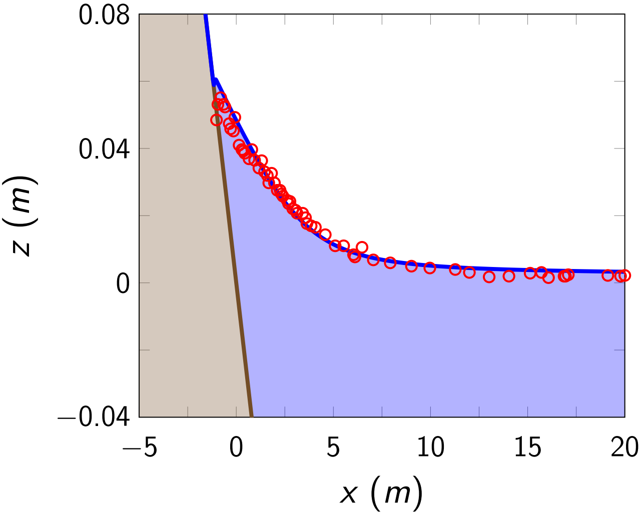
$t = 30s$



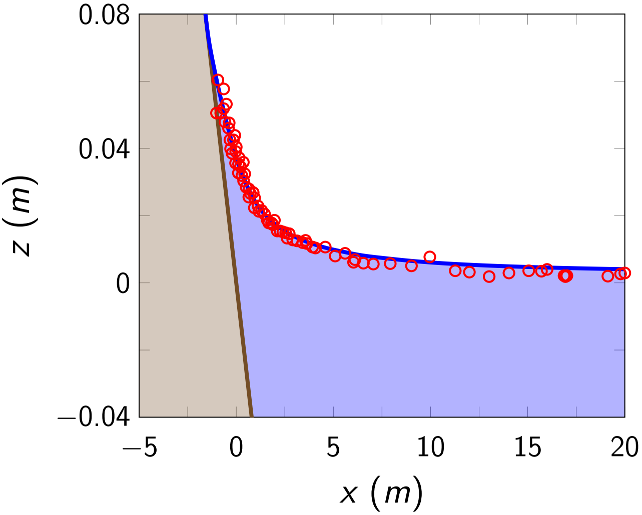
$t = 40s$



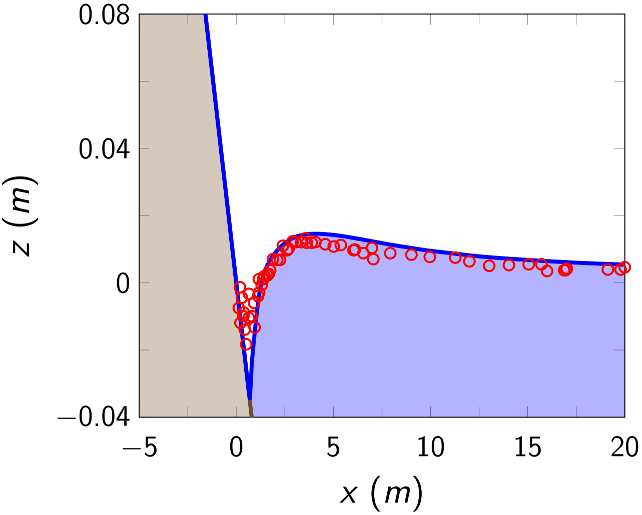
$t = 50s$



$t = 60s$



$t = 70s$



Experimental Validation Conclusions

- Demonstrates agreement of computational model and physical process
- Many sources of errors

Progress

2D: 1D method that extends well to 2D ✓

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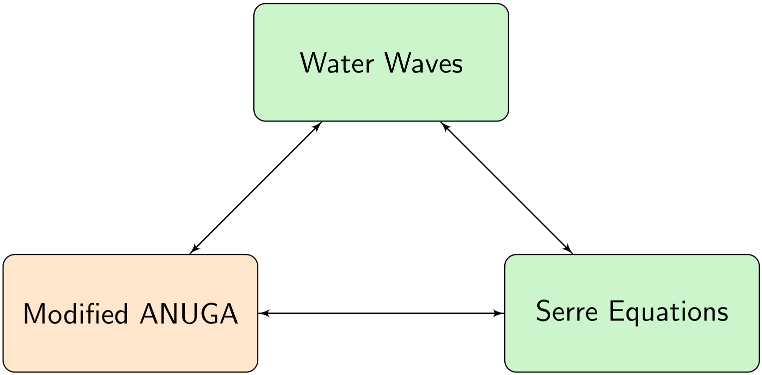
Conclusions

- Developed a robust computational model from the 1D Serre equations for the 2D water wave problem



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Wave Motion, 76(1):61–77.

Zoppou, C. (2014).
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PhD thesis, Australian National University, Mathematical Sciences Institute, College of Physical and Mathematical Sciences, Australian National University, Canberra, ACT 2600, Australia.