# A Finite Element-Volume Method for the Serre Equations

Jordan Pitt, Stephen Roberts and Christopher Zoppou Australian National University

November 6, 2018

## Outline

Results 00000000 0000000000

## Outline

Motivation

## Outline

- Motivation
- Method

## Outline

- Motivation
- Method
- Results



Method O 000000000000000000 0000000 00 Results 00000000 000000000

Motivation

#### Ocean Wave Hazards

Results 00000000 0000000000

Motivation

#### Ocean Wave Hazards

Tsunamis

Motivation

## Sulawesi 2018 Tsunami



Figure: Sulawesi Tsunami (Indonesia, 2018).

Method o ooooooooooooo oooooo oo Results 00000000 0000000000

Motivation

#### Ocean Wave Hazards

- ▶ Tsunamis
- Storm Surges

Motivation

# Storm Surge of Hurricane Florence and Michael

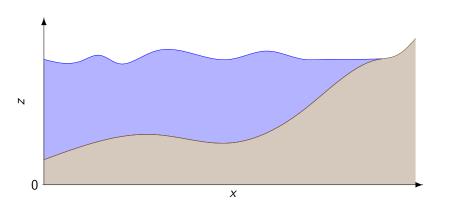


(a) Florence (U.S.A, 2018)

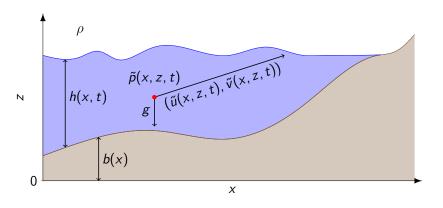


(b) Michael (U.S.A, 2018)

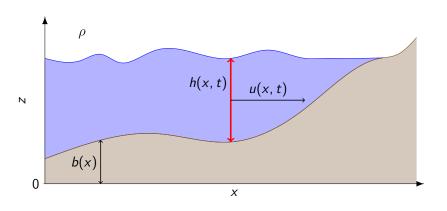
## Two Dimensional Scenario



## Navier-Stokes



# Model Simplification: Serre Equations





0000000 000000000

vvater ivious

# Assumptions

# Assumptions

Quantity	Serre Equations
Particle: $\tilde{u}(x, z, t)$	u(x,t)

# Assumptions

Quantity	Serre Equations
Particle: $\tilde{u}(x,z,t)$	u(x,t)
Particle: $\tilde{v}(x,z,t)$	$u\frac{\partial b}{\partial x}-(z-b)\frac{\partial b}{\partial x}$

# Assumptions

Quantity	Serre Equations
Particle: $\tilde{u}(x,z,t)$	u(x,t)
Particle: $\tilde{v}(x,z,t)$	$u\frac{\partial b}{\partial x}-(z-b)\frac{\partial b}{\partial x}$
Particle: $\tilde{p}(x, z, t)$	$g\rho[h+b-z]+\rho[h+b-z]\Psi$
	$+\frac{1}{2}\rho\left(h^2-[z-b]^2\right)\Phi$

where

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

## **Equations**

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

Momentum: 
$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right)$$

$$+\frac{\partial b}{\partial x}\left(gh+h\Psi+\frac{h^2}{2}\Phi\right)=0.$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

Introduction

#### Method

▶ When  $\Phi = \Psi = 0$  we have the Shallow Water Wave Equations

Introduction

#### Method

- When  $\Phi = \Psi = 0$  we have the Shallow Water Wave Equations
- Demonstrated utility of Finite Volume Methods for these equations (ANUGA)

Introduction

#### Method

- When  $\Phi = \Psi = 0$  we have the Shallow Water Wave Equations
- Demonstrated utility of Finite Volume Methods for these equations (ANUGA)

Goal: Adapt Finite Volume Methods for the Serre Equations

#### Finite Volume Method

Conservation law form

Finite Volume Method

#### Finite Volume Method

- ► Conservation law form
- ► Finite volume update

#### Finite Volume Method

Conservation law form

#### Conservation Law Form

$$\frac{\partial q}{\partial t} + \frac{\partial}{\partial x} \left[ f\left(q, \frac{\partial q}{\partial x}, \frac{\partial^2 q}{\partial x^2}, \dots, \frac{\partial^n q}{\partial x^n}\right) \right] + s\left(q, \frac{\partial q}{\partial x}, \frac{\partial^2 q}{\partial x^2}, \dots, \frac{\partial^m q}{\partial x^m}\right) = 0$$

# Equations

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

Momentum: 
$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right)$$

$$+\frac{\partial b}{\partial x}\left(gh+h\Psi+\frac{h^2}{2}\Phi\right)=0.$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

## **Equations**

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

Momentum: 
$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right)$$

$$+\frac{\partial b}{\partial x}\left(gh+h\Psi+\frac{h^2}{2}\Phi\right)=0.$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

#### Conservation Law Form

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right)$$
$$+ \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = 0.$$

#### Conservation Law Form

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right)$$
$$+ \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = 0.$$

with

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

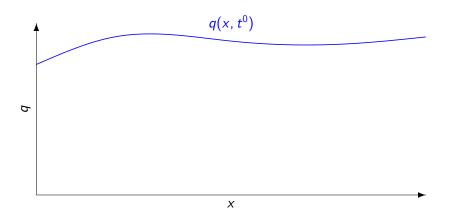
#### Finite Volume Method

- Conservation law form
- Finite volume update

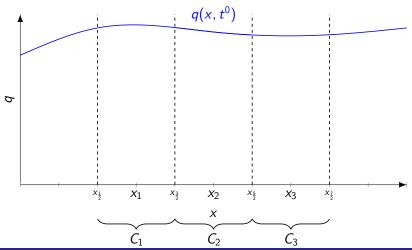
#### Conservation Law Form

$$\frac{\partial q}{\partial t} + \frac{\partial}{\partial x} \left[ f\left(q, \frac{\partial q}{\partial x}, \frac{\partial^2 q}{\partial x^2}, \dots, \frac{\partial^n q}{\partial x^n}\right) \right] + s\left(q, \frac{\partial q}{\partial x}, \frac{\partial^2 q}{\partial x^2}, \dots, \frac{\partial^m q}{\partial x^m}\right) = 0$$

## Finite Volume Method

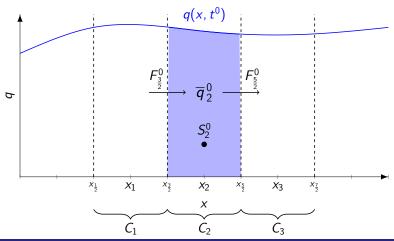


## Discretisation



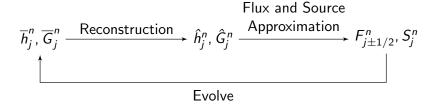
Jordan Pitt, Stephen Roberts and Christopher Zoppou Australian National University

# **Update**



Jordan Pitt, Stephen Roberts and Christopher Zoppou Australian National University

#### Finite Volume Method



 Results 00000000 0000000000

Finite Volume Method

# Require velocity to calculate flux and source

However to calculate the flux and source terms we require u

Finite Volume Method

### Require velocity to calculate flux and source

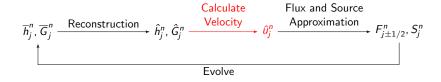
However to calculate the flux and source terms we require u

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0,$$

$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right)$$
$$+ \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = 0.$$

Finite Volume Method

#### Method



#### Reconstruction

▶ Determines spatial order of accuracy

 Results 00000000 0000000000

Reconstruction

#### Reconstruction

Determines spatial order of accuracy

Goal: Second-order accuracy

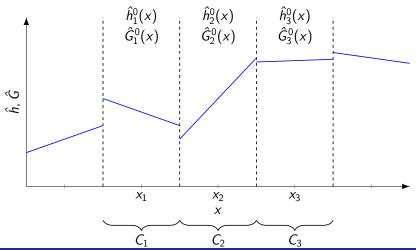
### Reconstruction Spaces

Quantity Number of		Reconstructe	
	spatial derivatives	functions	

### Reconstruction Spaces

Quantity	Number of	Reconstructed	
	spatial derivatives	functions	
h	zero	linear over cell, discontinuous at edges	
G	zero	linear over cell, discontinuous at edges	

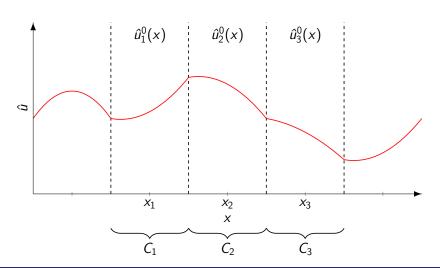
$$\hat{h}, \hat{G}$$



### Reconstruction Spaces

Quantity	Number of	Reconstructed	
	spatial derivatives	functions	
h	zero	linear over cell, discontinuous at edges	
G	zero	linear over cell, discontinuous at edges	
и	one	quadratic over cell, continuous at edges	

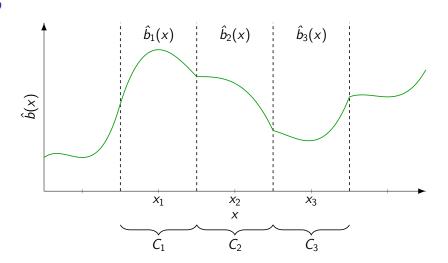




### Reconstruction Spaces

Quantity	Number of	Reconstructed	
	spatial derivatives	functions	
h	zero	linear over cell, discontinuous at edges	
G	zero	linear over cell, discontinuous at edges	
и	one	quadratic over cell, continuous at edges	
b	two	cubic over cell, continuous at edges	

ĥ



Calculation of Velocity

### Finite Element Calculation of Velocity

Finite Element Method to solve:

$$G = hu\left(1 + \frac{\partial h}{\partial x}\frac{\partial b}{\partial x} + \frac{1}{2}h\frac{\partial^2 b}{\partial x^2} + \left[\frac{\partial b}{\partial x}\right]^2\right) - \frac{\partial}{\partial x}\left(\frac{1}{3}h^3\frac{\partial u}{\partial x}\right).$$

for u given h, G and b

Calculation of Velocity

### Finite Element Calculation of Velocity

Finite Element Method to solve:

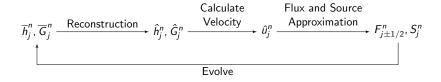
$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

for u given h, G and b

Solves the weak form replacing all quantities with their reconstructions  $\hat{h}$ ,  $\hat{G}$  and  $\hat{b}$  to get  $\hat{u}$ 

Calculation of Velocity

#### Method



#### Validation

► Analytic Solution

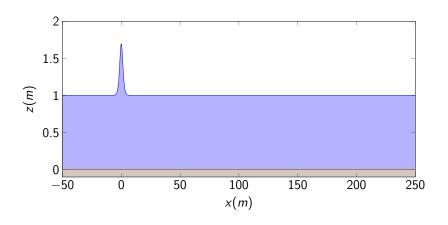
#### Validation

- ► Analytic Solution
- ► Experimental Results

#### Validation

► Analytic Solution

### Soliton Example



### Soliton Equations

$$h(x, t) = a_0 + a_1 \operatorname{sech} (\kappa (x - ct)),$$

$$u(x,t) = c\left(1 - \frac{a_0}{h(x,t)}\right),\,$$

$$b(x)=0$$

### Soliton Equations

$$h(x,t) = a_0 + a_1 \operatorname{sech} (\kappa (x - ct)),$$

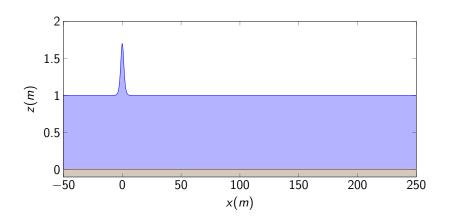
$$u(x,t)=c\left(1-\frac{a_0}{h(x,t)}\right),\,$$

$$b(x)=0$$

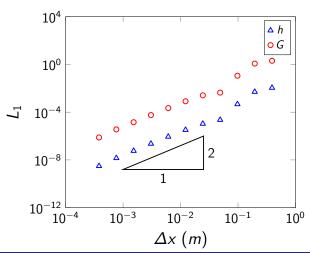
$$\kappa = \frac{\sqrt{3a_1}}{2a_0\sqrt{(a_0+a_1)}},$$

$$c=\sqrt{g(a_0+a_1)}.$$

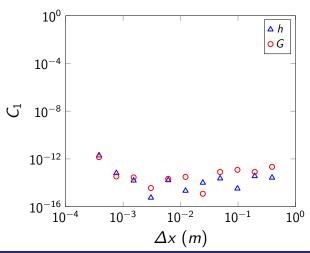
### Numerical Solution $a_0 = 1$ , $a_1 = 0.7$



### Convergence



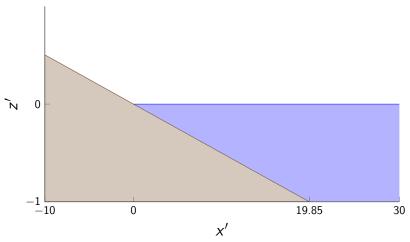
#### Conservation



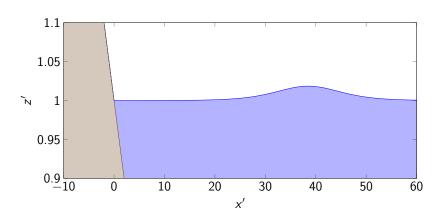
#### Validation

- ► Analytic Solution
- ► Experimental Results

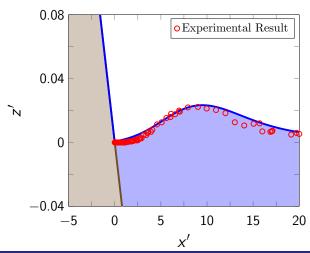
## Synolakis Experiment



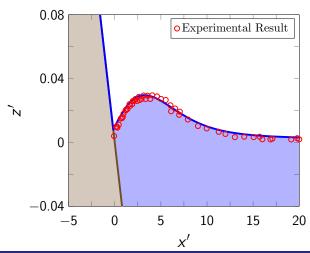
#### **Numerical Solution**



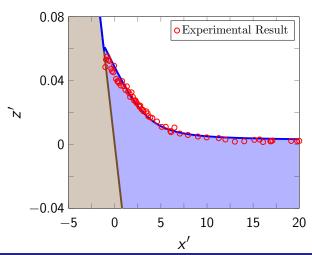
#### Comparison t = 30s



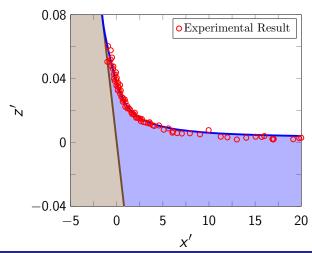
#### Comparison t = 40s



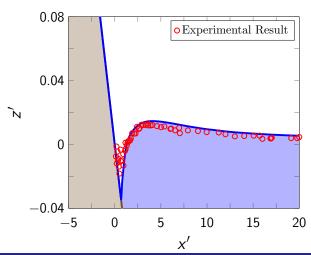
### Comparison t = 50s



#### Comparison t = 60s



### Comparison t = 70s



#### Conclusion

Finite Element Volume Method for The Serre Equations

Second-order accurate

#### Conclusion

Finite Element Volume Method for The Serre Equations

- Second-order accurate
- Conservative

#### Conclusion

Finite Element Volume Method for The Serre Equations

- Second-order accurate
- Conservative
- Reproduces analytic solutions

#### Conclusion

#### Finite Element Volume Method for The Serre Equations

- Second-order accurate
- Conservative
- Reproduces analytic solutions
- Reproduces experimental results

Results 0000000 000000000

Experimental Comparison

# Thanks!