1 All Definitions for Numerical Versions

$$\mathcal{C}_2 = \frac{2\cos\left(k\Delta x\right) - 2}{\Delta x^2}$$

$$\mathcal{C}_4 = \frac{-2\cos\left(2k\Delta x\right) + 32\cos\left(k\Delta x\right) - 30}{12\Delta x^2}$$

$$\mathcal{G} = \left[H - \frac{H^3}{3}\mathcal{C}\right]$$

$$\mathcal{M}_3 = \frac{24}{26 - 2\cos\left(k\Delta x\right)}$$

$$\mathcal{M}_1 = \mathcal{M}_2 = 1$$

$$\mathcal{R}_1^+ = e^{ik\Delta x} \quad , \quad \mathcal{R}_1^- = 1$$

$$\mathcal{R}_2^- = 1 + \frac{i\sin\left(k\Delta x\right)}{2}$$

$$\mathcal{R}_2^+ = e^{ik\Delta x} \left(1 - \frac{i\sin\left(k\Delta x\right)}{2}\right)$$

$$\mathcal{R}_3^- = \frac{\mathcal{M}_3}{6} \left[5 + -e^{-ik\Delta x} + 2e^{ik\Delta x}\right]$$

$$\mathcal{R}_3^+ = \frac{\mathcal{M}_3 e^{ik\Delta x}}{6} \left[5 + 2e^{-ik\Delta x} - e^{ik\Delta x}\right]$$

$$\mathcal{R}_3^u = \frac{e^{ik\Delta x} + 1}{2}$$

$$\mathcal{R}_3^u = \frac{-3e^{2ik\Delta x} + 27e^{ik\Delta x} + 27 - 3e^{-ik\Delta x}}{48}$$

$$\mathcal{F}^{h,u} = H\mathcal{R}^u$$

$$\mathcal{F}^{h,h} = -\frac{\sqrt{gH}}{2} \left[\mathcal{R}^+ - \mathcal{R}^-\right]$$

$$\mathcal{F}^{u,u} = -\frac{\sqrt{gH}}{2}\mathcal{G}\left[\mathcal{R}^{+} - \mathcal{R}^{-}\right]$$
$$\mathcal{F}^{u,h} = \frac{gH\mathcal{R}^{-} + gH\mathcal{R}^{+}}{2}$$

$$\mathcal{D} = 1 - e^{-ik\Delta x}$$

2 Taylor Expansions Of Analytic Values

We denote exact/analytic version with a subscript a

$$\mathcal{G}_{a} = H + \frac{H^{3}}{3}k^{2}$$

$$\mathcal{M}_{a} = 1 - \frac{1}{24}(k\Delta x)^{2} + \frac{1}{192}(k\Delta x)^{4} + O(x^{6})$$

$$\mathcal{R}_{a}^{+} = \mathcal{R}_{a}^{-} = \mathcal{R}_{a} = 1 + \frac{i}{2}k\Delta x - \frac{1}{8}k^{2}\Delta x^{2} - \frac{i}{48}k^{3}\Delta x^{3} + O(x^{4})$$

For the fluxes I think its best to group the $\frac{\mathcal{D}}{\Delta x \mathcal{M}} \mathcal{F}$ because its collects all the terms using spatial approximations and has a nice form. In fact these terms approximate the derivative of the flux.

$$\begin{split} \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{h,u} &= ikH \\ \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{h,h} &= 0 \\ \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{u,h} &= ikgH \\ \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{u,u} &= 0 \end{split}$$

So in particular for the mass equation

$$h_t + Hu_x = 0$$

then

$$i\omega h + \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{h,u} u_j + \frac{\mathcal{D}_a}{\Delta x \mathcal{M}_a} \mathcal{F}_a^{h,h} h_j = 0$$

Indeed it is these the values of our approximations to these terms that confirm that our methods have the correct spatial accuracy.

3 Taylor Expansions Of First Order Values

$$\mathcal{G}_1 = H + \frac{H^3}{3}k^2 - \frac{H^3k^4(\Delta x)^2}{36} + O(x^4)$$

$$\mathcal{M}_1 = 1$$

$$\mathcal{R}_1^- = 1$$

$$\mathcal{R}_{1}^{+} = 1 + ik\Delta x - \frac{1}{2}(k\Delta x)^{2} - \frac{i}{3!}(k\Delta x)^{3} + \frac{1}{4!}(k\Delta x)^{4} + O((k\Delta x)^{5})$$

$$\mathcal{R}_{1}^{u} = 1 + \frac{i}{2}k\Delta x - \frac{1}{4}(k\Delta x)^{2} - \frac{i}{12}(k\Delta x)^{3} + O(\Delta x^{4})$$

$$\mathcal{D}\mathcal{F}_{1}^{h,u} = Hik\Delta x - \frac{Hi}{6}(k\Delta x)^{3} + O(\Delta x^{4})$$

$$\mathcal{D}\mathcal{F}_{1}^{h,h} = -\frac{\sqrt{gH}}{2}k^{2}\Delta x + O(\Delta x^{3})$$

$$\mathcal{D}\mathcal{F}_{1}^{u,h} = igHk - \frac{igHk^{3}(\Delta x)^{2}}{6} + O(\Delta x^{3})$$

$$\mathcal{D}\mathcal{F}_{1}^{u,u} = \frac{i\sqrt{gH}H}{6}(H^{2} + 3)k^{2}(\Delta x) - \frac{\sqrt{gH}H}{72}(2H^{2} + 3)k^{4}(\Delta x)^{3} + O(\Delta x^{5})$$

4 Taylor Expansions Of Second Order Values

$$\mathcal{G}_{2} = H + \frac{H^{3}}{3}k^{2} - \frac{H^{3}k^{4}(\Delta x)^{2}}{36} + O(x^{4})$$

$$\mathcal{M}_{2} = 1$$

$$\mathcal{R}_{2}^{-} = 1 + \frac{i}{2}(k\Delta x) - \frac{i}{12}(k\Delta x)^{3} + O(x^{4})$$

$$\mathcal{R}_{2}^{+} = 1 + \frac{i}{2}k\Delta x + \frac{i}{6}k^{3}\Delta x^{3} + O\left(\Delta x^{4}\right)$$

$$\mathcal{R}_{2}^{u} = 1 + \frac{i}{2}k\Delta x - \frac{1}{4}(k\Delta x)^{2} - \frac{i}{12}(k\Delta x)^{3} + O(\Delta x^{4})$$

$$\mathcal{D}\mathcal{F}_{2}^{h,u} = Hik - \frac{iH}{6}k^{3}(\Delta x)^{2} + \frac{iH}{120}k^{5}(\Delta x)^{4} + O(\Delta x^{5})$$

$$\mathcal{D}\mathcal{F}_{2}^{h,h} = \frac{\sqrt{gH}}{8}k^{4}(\Delta x)^{3} - \frac{\sqrt{gH}}{48}k^{6}(\Delta x)^{5} + O(\Delta x^{7})$$

$$\mathcal{D}\mathcal{F}_{2}^{u,u} = \frac{\sqrt{gH}}{24}H(H^{2} + 3)k^{4}\Delta x^{3} - \frac{\sqrt{gH}}{96}H(H^{2} + 2)k^{6}\Delta x^{5} + O(\Delta x^{6})$$

$$\mathcal{D}\mathcal{F}_{2}^{u,h} = igHk + \frac{igH}{12}k^{3}\Delta x^{2} - \frac{13igH}{240}k^{5}\Delta x^{4} + O(\Delta x^{6})$$

5 Taylor Expansions Of Third Order Values

$$\mathcal{G}_3 = H + \frac{H^3 k^2}{3} - \frac{H^3 k^6 (\Delta x)^4}{270} + O(\Delta x^6)$$

$$\mathcal{M}_3 = 1 + \frac{1}{24} (k \Delta x)^2 - \frac{1}{288} (k \Delta x)^4 + O(x^6)$$

$$R_3^- = 1 + \frac{i}{2} k \Delta x - \frac{1}{8} (k \Delta x)^2 - \frac{5}{48} (k \Delta x)^3 + O(\Delta x^4)$$

$$R_3^+ = 1 + \frac{i}{2} k \Delta x - \frac{1}{8} (k \Delta x)^2 + \frac{1}{16} (k \Delta x)^3 + O(\Delta x^4)$$

$$R_{3}^{u} = 1 + \frac{i}{2}k\Delta x - \frac{1}{8}(k\Delta x)^{2} - \frac{i}{48}(k\Delta x)^{3} + O(\Delta x^{4})$$

$$\frac{\mathcal{D}}{\Delta x \mathcal{M}_{3}} \mathcal{F}_{3}^{h,u} = ikH - \frac{9iH}{320}k^{5}\Delta x^{4} - \frac{iH}{448}k^{7}\Delta x^{6} + O(\Delta x^{9})$$

$$\frac{\mathcal{D}}{\Delta x \mathcal{M}_{3}} \mathcal{F}_{3}^{h,h} = \frac{\sqrt{gH}}{12}k^{4}\Delta x^{3} - \frac{\sqrt{gH}}{72}k^{6}\Delta x^{5} + \frac{\sqrt{gH}}{960}k^{8}\Delta x^{7} + O(\Delta x^{9})$$

$$\frac{\mathcal{D}}{\Delta x \mathcal{M}_{3}} \mathcal{F}_{3}^{u,u} = \frac{H\sqrt{gH}}{36}(H^{2} + 3)k^{4}\Delta x^{3} - \frac{iH\sqrt{gH}}{144}(H^{2} + 3)k^{5}\Delta x^{4} + O(\Delta x^{5})$$

$$\frac{\mathcal{D}}{\Delta x \mathcal{M}_{3}} \mathcal{F}_{3}^{u,h} = igkH - \frac{igH}{30}k^{5}\Delta x^{4} + O(\Delta x^{6})$$

$$\frac{\mathcal{D}}{\Delta x \mathcal{M}_{3}} \mathcal{F}_{3}^{u,h} = igkH - \frac{igH}{30}k^{5}\Delta x^{4} + O(\Delta x^{6})$$