

# Robust Computational Models for Water Waves

Jordan Pitt, Stephen Roberts and Christopher Zoppou  
Australian National University

September 7, 2018

# Outline of the Presentation

- ▶ Motivation
- ▶ History
- ▶ Contribution
  - ▶ Method
  - ▶ Validation

# Water Waves

# Water Waves

Water wave hazards:

- ▶ Tsunamis

# Tsunamis



**Figure:** 2004 Indian Ocean Tsunami (Banda Aceh).

# Tsunamis



Figure: 2011 Tohoku Tsunami.

# Water Waves

Water wave hazards:

- ▶ Tsunamis
- ▶ Storm Surges

## Storm Surges



Figure: 2012 Hurricane Sandy Storm Surge.

## Water Waves

Water wave hazards:

- ▶ Tsunamis
- ▶ Storm Surges

Phenomena caused by water waves:

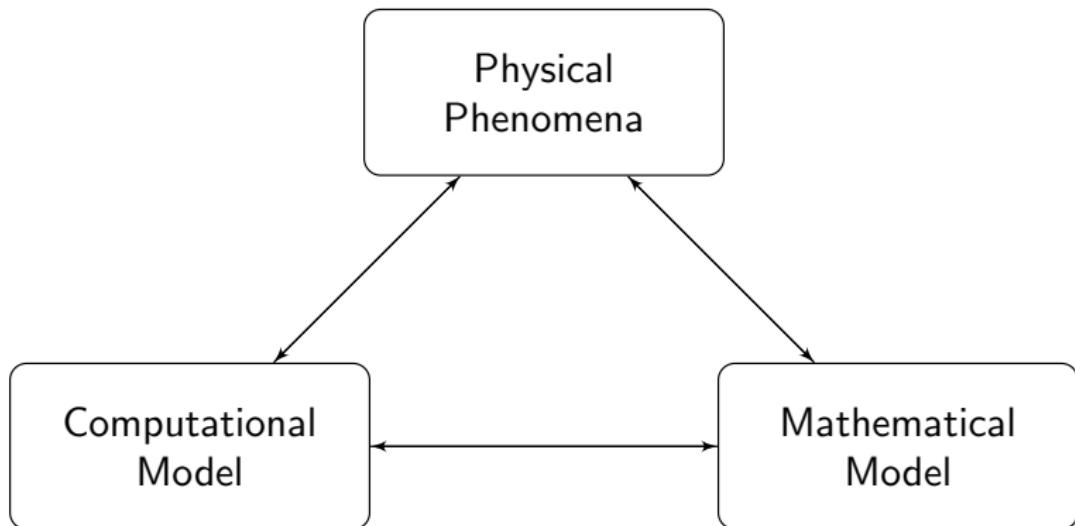
- ▶ Nutrient Transport
- ▶ Beach Erosion
- ▶ Breakup of Sea Ice

# Computational Modelling

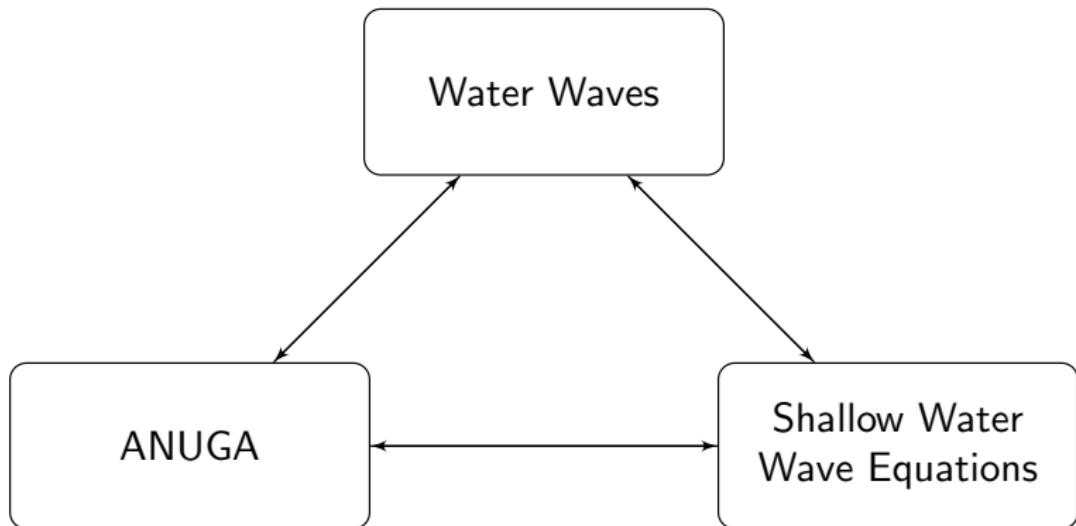
Goal: Model physics on computers.

# Computational Modelling

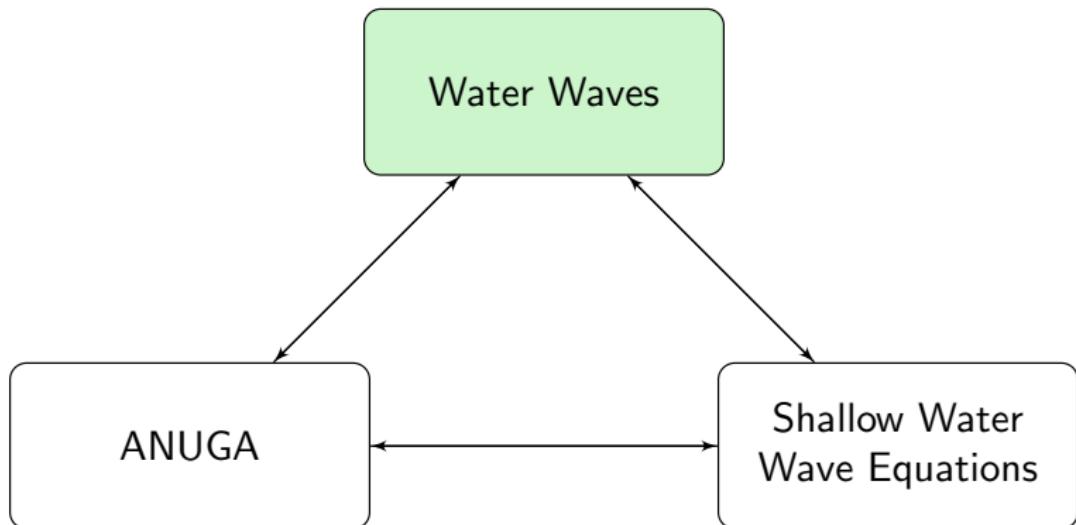
Goal: Model physics on computers.



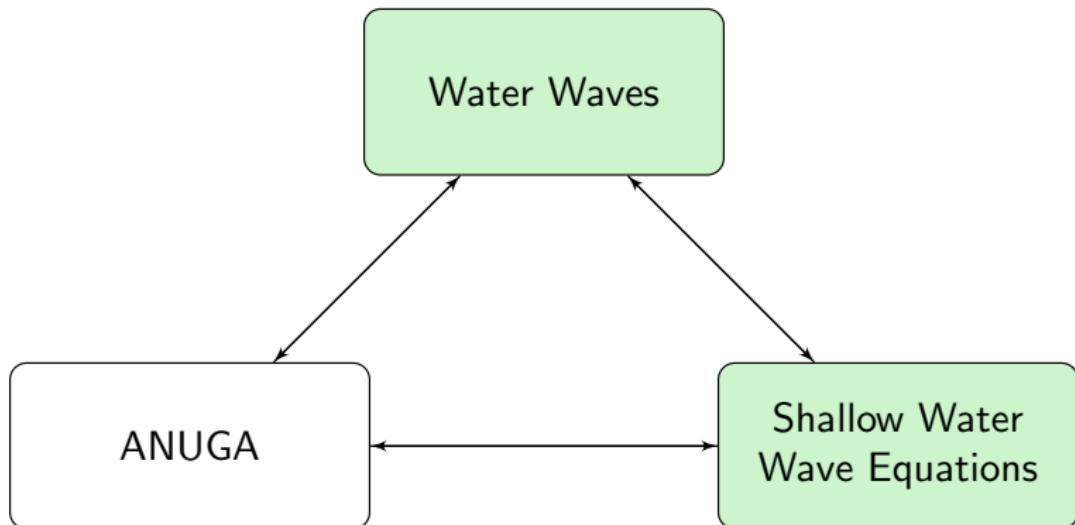
# ANUGA



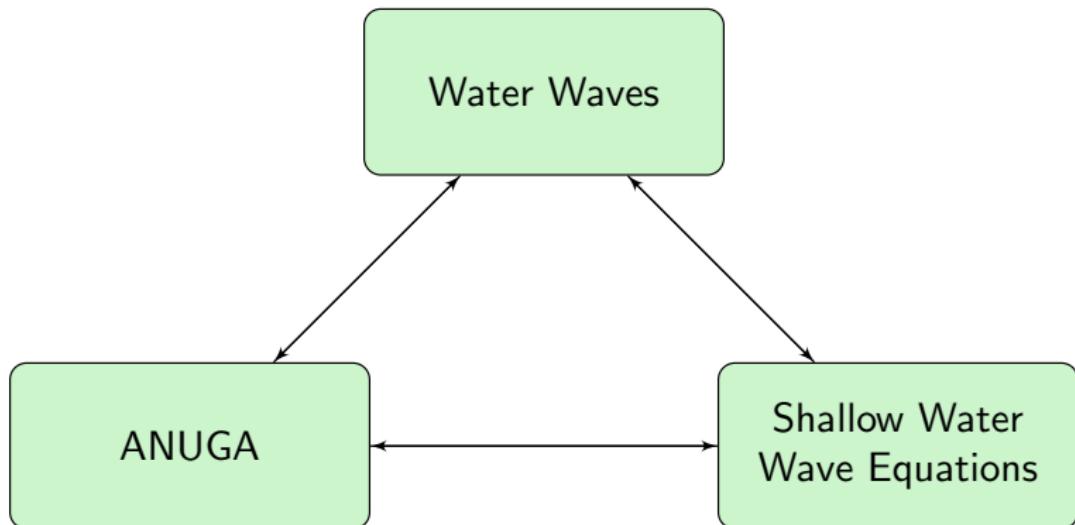
## ANUGA: Water Waves



# ANUGA: Shallow Water Wave Equations



# ANUGA



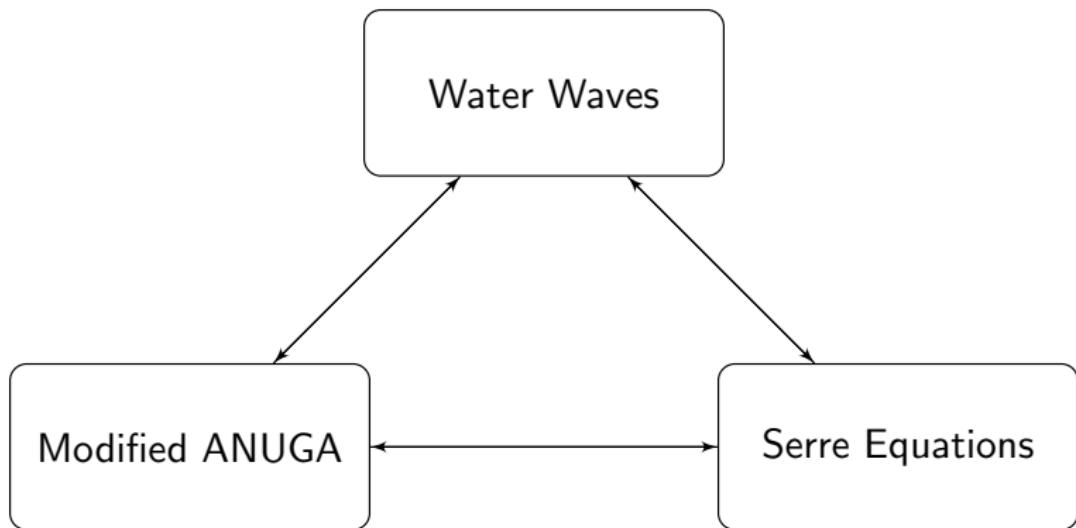
## Outcome

New project at the ANU to develop a robust computational model for the Serre equations.

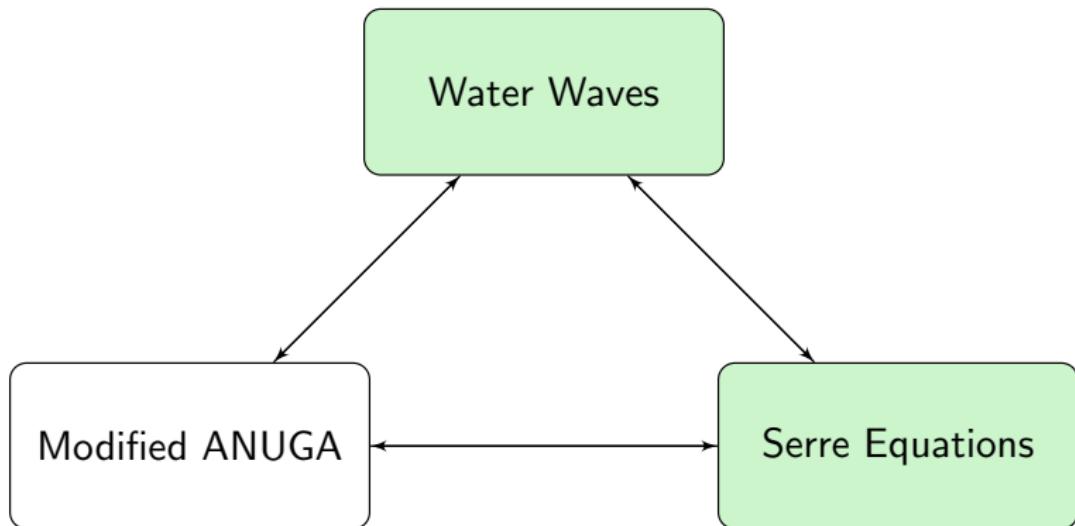
## Outcome

New project at the ANU to develop a robust computational model for the Serre equations.

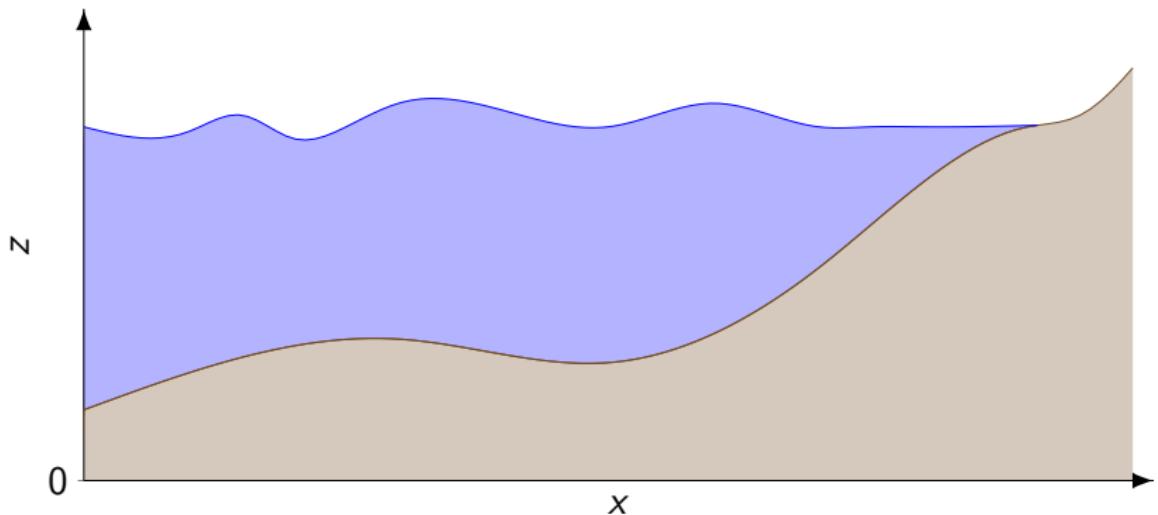
Goal:



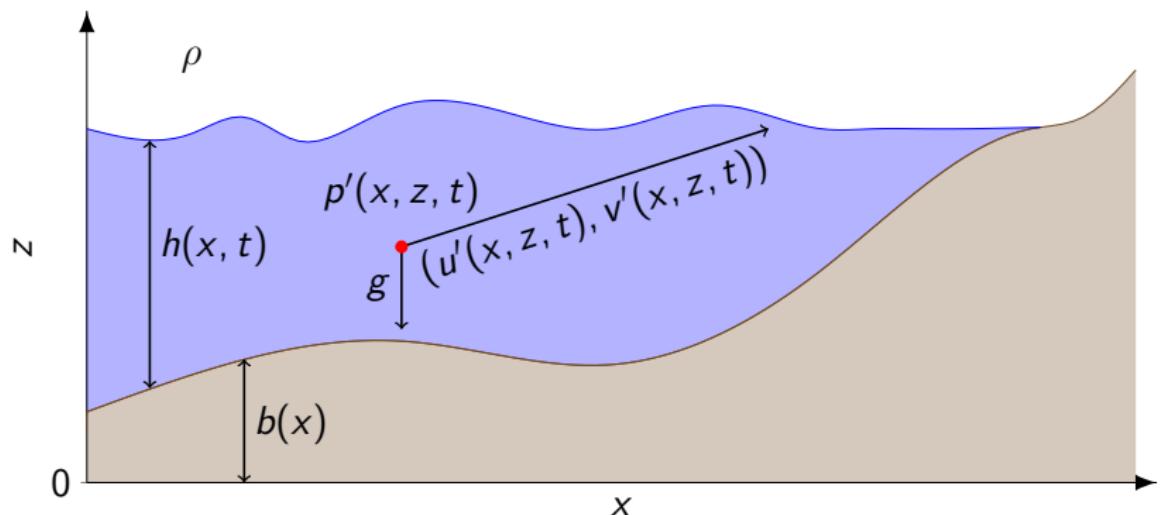
## Mathematical Model



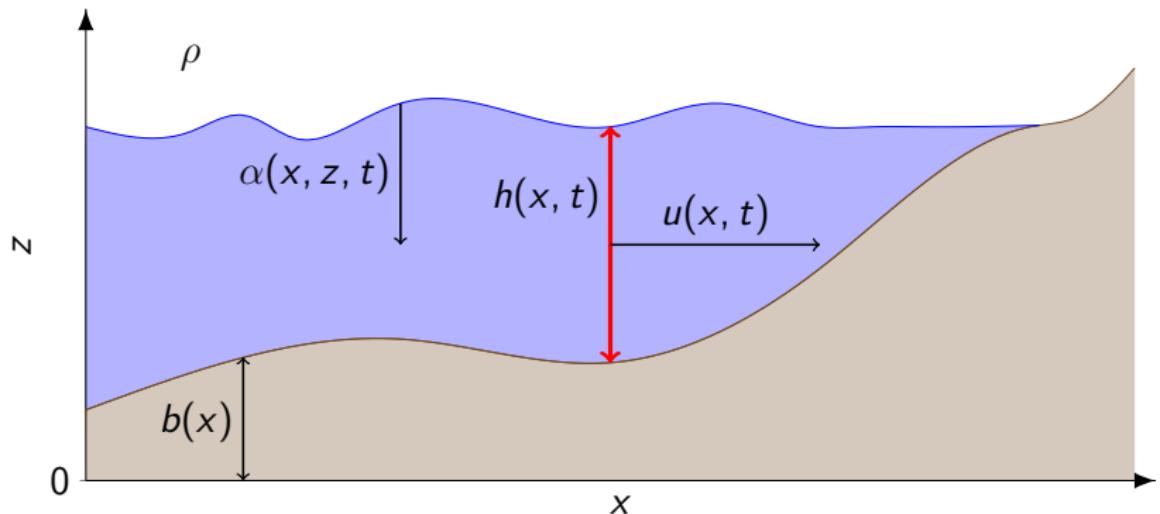
## Typical Scenario



## Navier Stokes Model



# Serre Model



## Assumptions

## Assumptions

Quantity	Shallow Water Wave Equations	Serre Equations
Particle: $v'(x, z, t)$	0	$u \frac{\partial b}{\partial x} - (h - \alpha) \frac{\partial b}{\partial x}$

## Assumptions

Quantity	Shallow Water Wave Equations	Serre Equations
Particle: $v'(x, z, t)$	0	$u \frac{\partial b}{\partial x} - (h - \alpha) \frac{\partial b}{\partial x}$
Particle: $p'(x, z, t)$	$g\rho\alpha$	$g\rho\alpha + \rho\alpha\Psi + \frac{1}{2}\rho\alpha(2h - \alpha)\Phi$

where

$$\alpha(x, z, t) = (h(x, t) + b(x)) - z$$

## Assumptions

Quantity	Shallow Water Wave Equations	Serre Equations
Particle: $v'(x, z, t)$	0	$u \frac{\partial b}{\partial x} - (h - \alpha) \frac{\partial b}{\partial x}$
Particle: $p'(x, z, t)$	$g\rho\alpha$	$g\rho\alpha + \rho\alpha \Psi + \frac{1}{2}\rho\alpha(2h - \alpha)\Phi$

where

$$\alpha(x, z, t) = (h(x, t) + b(x)) - z$$

and

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

## Equations

Mass: 
$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

Momentum: 
$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right)$$

$$+ \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0.$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

## Pros and Cons

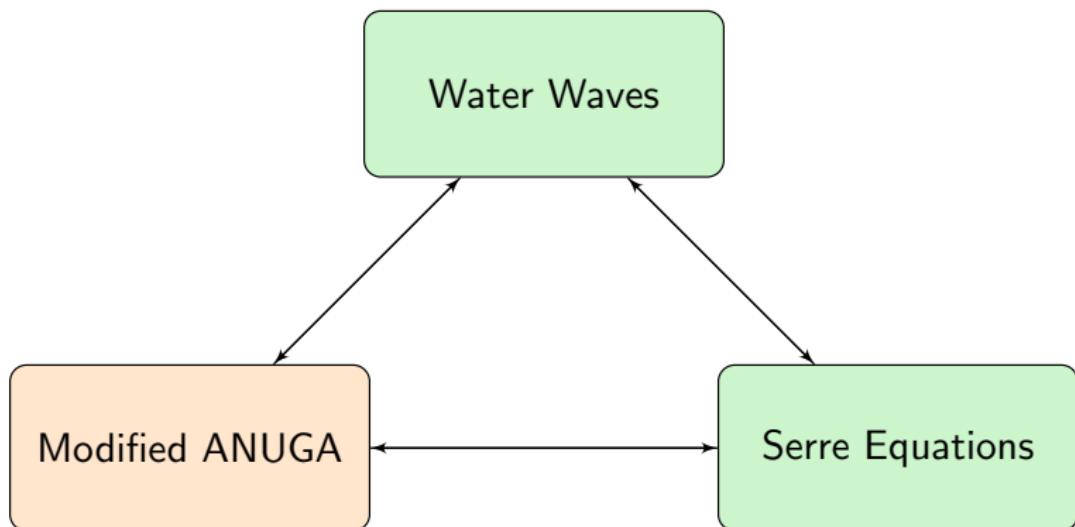
Pros:

- ▶ Includes dispersive effects
- ▶ Considered one of the best models for water waves
- ▶ Can apply techniques of ANUGA

Cons:

- ▶ More complex than the Shallow Water Wave Equations

## Computational Model



## Previous Work at the ANU

- ▶ 2014: Chris Zoppou's PhD thesis  
Developed computational model for the 1D Serre equations.
- ▶ 2014: My Honours thesis  
Independent reproduction of Chris Zoppou's computational model.

## Open Problems and Thesis Goals

**2D:** 1D method that extends well to 2D

**Robust:** Validation for steep gradients in free surface

**Robust:** Inclusion and validation of dry beds

## Open Problems and Thesis Goals

**2D:** 1D method that extends well to 2D

**Robust:** Validation for steep gradients in free surface

**Robust:** Inclusion and validation of dry beds

Technique: Develop a robust computational model from the 1D Serre equations that can be easily extended to 2D.

## Finite Volume Method

**2D:** Extends well to 2D

**Robust:** Stable in the presence of steep gradients

**Robust:** Stable in the presence of dry beds

- ▶ Maintains conservation properties of the equations
- ▶ ANUGA

## Finite Volume Method

**2D:** Extends well to 2D

**Robust:** Stable in the presence of steep gradients

**Robust:** Stable in the presence of dry beds

- ▶ Maintains conservation properties of the equations
- ▶ ANUGA

Chris Zoppou's thesis demonstrated an adaptation of the Finite Volume Method to solve the Serre Equations.

## Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right) + \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0,$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

## Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right) + \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0,$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

For a Finite Volume Method we require equations in the form

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where  $f(q)$  and  $s(q)$  do not contain temporal derivatives

## Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right) + \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0,$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

For a Finite Volume Method we require equations in the form

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where  $f(q)$  and  $s(q)$  do not contain temporal derivatives

## Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right) + \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0,$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

For a Finite Volume Method we require equations in the form

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where  $f(q)$  and  $s(q)$  do not contain temporal derivatives

## Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial}{\partial x} \left( u^2 h + \frac{gh^2}{2} + \frac{h^2}{2} \Psi + \frac{h^3}{3} \Phi \right) + \frac{\partial b}{\partial x} \left( gh + h\Psi + \frac{h^2}{2} \Phi \right) = 0,$$

$$\Psi = \frac{\partial b}{\partial x} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} \right) + u^2 \frac{\partial^2 b}{\partial x^2}, \quad \Phi = \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} - u \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial x \partial t}.$$

For a Finite Volume Method we require equations in the form

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0$$

where  $f(q)$  and  $s(q)$  do not contain temporal derivatives

## Reformulation

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\begin{aligned} \frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right) \\ + \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = 0. \end{aligned}$$

with

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2}h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3}h^3 \frac{\partial u}{\partial x} \right).$$

## Reformulation

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0,$$

$$\begin{aligned} \frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right) \\ + \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = 0. \end{aligned}$$

with

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2}h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3}h^3 \frac{\partial u}{\partial x} \right).$$

## Finite Volume Method Example

Conservation of a quantity:

$$q(x, t).$$

## Finite Volume Method Example

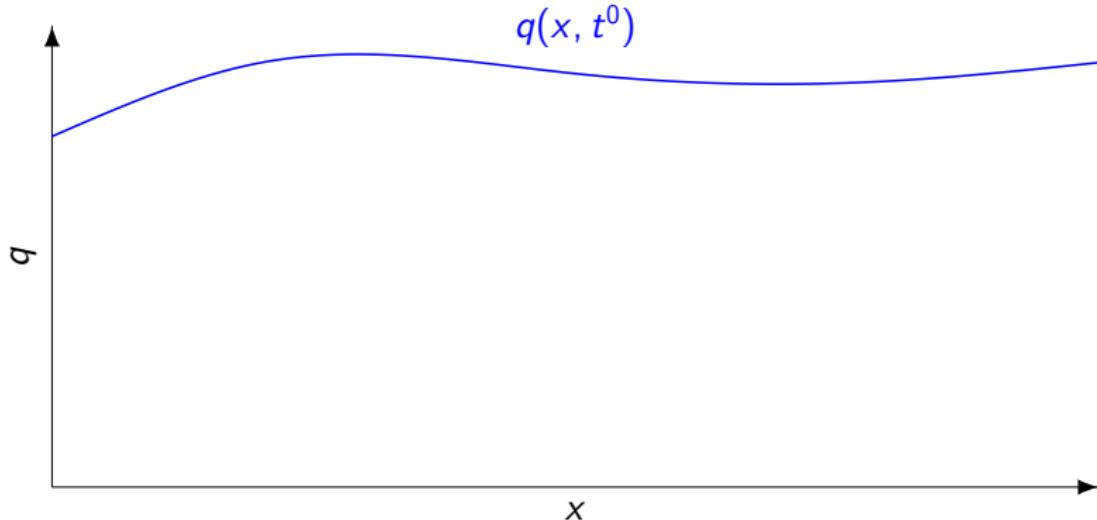
Conservation of a quantity:

$$q(x, t).$$

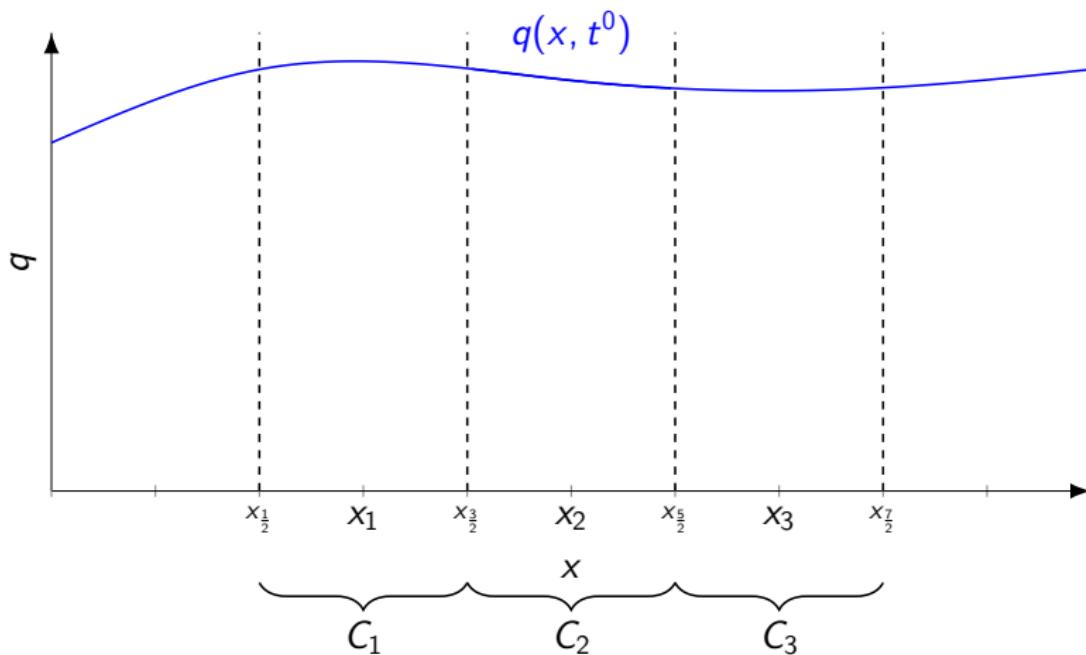
Conservation equation with a source term:

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0.$$

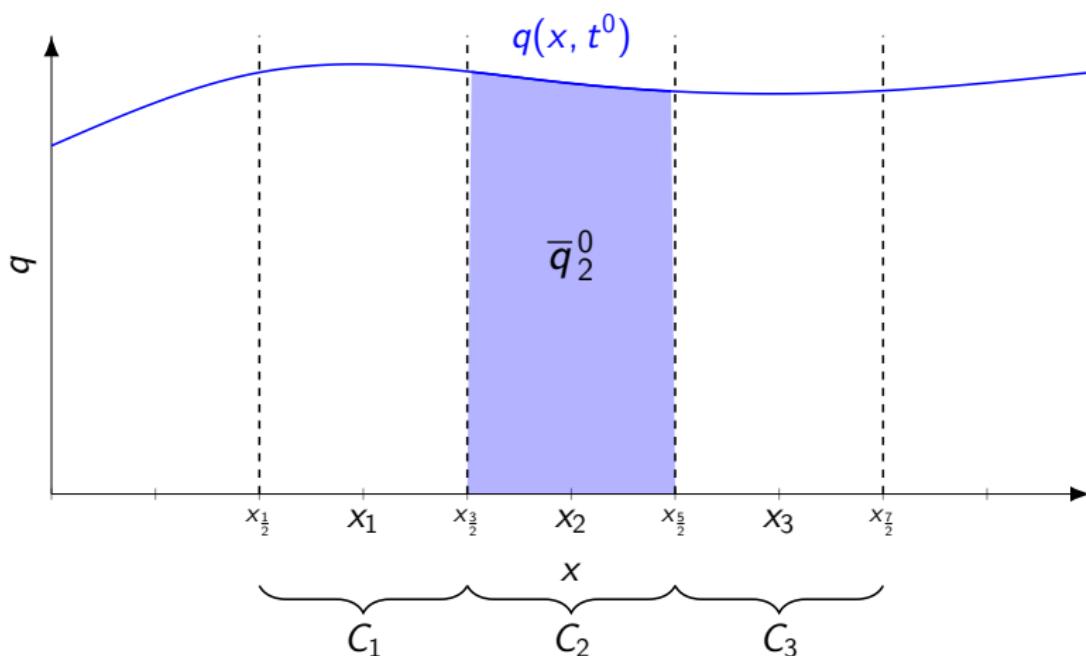
Function at  $t = t^0$



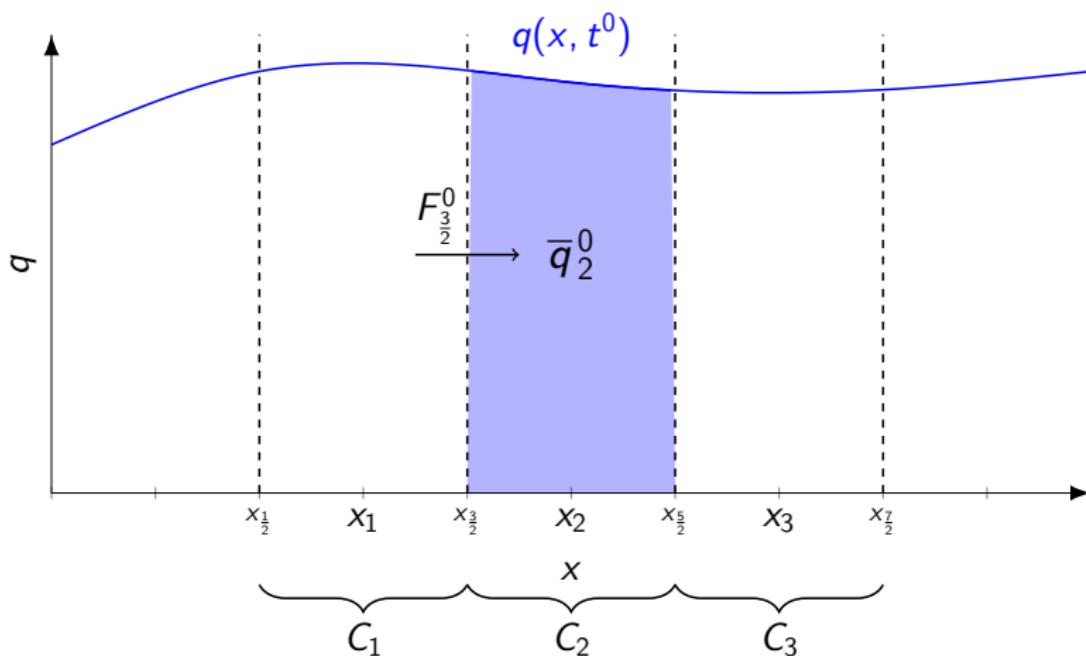
# Cell Discretisation



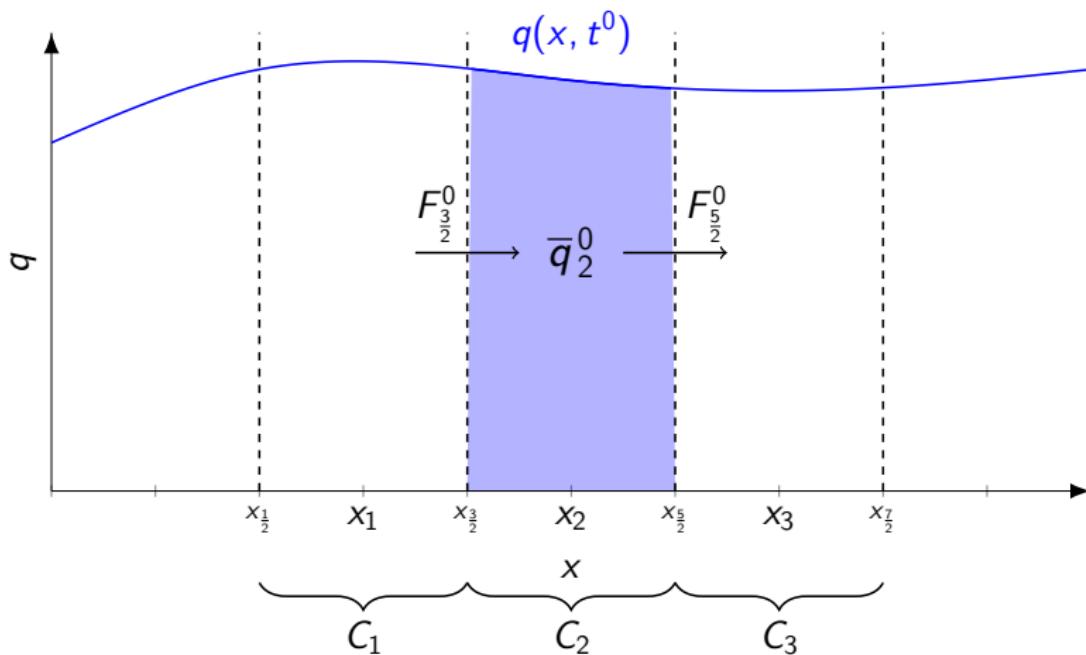
# Total Amount



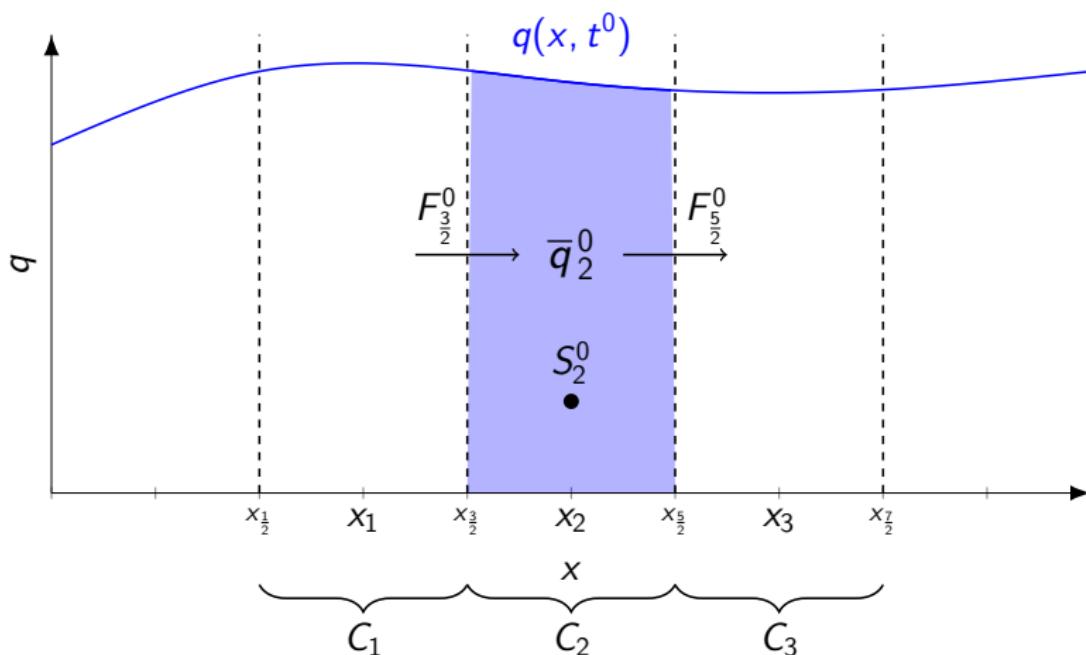
# Flux Left



# Flux Right



# Source



## Finite Volume Update

$$\bar{q}_2^1 = \bar{q}_2^0 - \left( F_{\frac{5}{2}}^0 - F_{\frac{3}{2}}^0 \right) - (S_2^0),$$

## Finite Volume Update

$$\bar{q}_2^1 = \bar{q}_2^0 - \left( F_{\frac{5}{2}}^0 - F_{\frac{3}{2}}^0 \right) - (S_2^0),$$

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0,$$

## Finite Volume Update

$$\bar{q}_2^1 = \bar{q}_2^0 - \left( F_{\frac{5}{2}}^0 - F_{\frac{3}{2}}^0 \right) - (S_2^0),$$

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0,$$

## Finite Volume Update

$$\bar{q}_2^1 = \bar{q}_2^0 - \left( F_{\frac{5}{2}}^0 - F_{\frac{3}{2}}^0 \right) - (S_2^0),$$

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} + s(q) = 0,$$

$$\overbrace{\int_{C_2} q(x, t^1) dx}^{\bar{q}_2^1} = \overbrace{\int_{C_2} q(x, t^0) dx}^{\bar{q}_2^0} - \left( \overbrace{\int_{t^0}^{t^1} f(q(x_{5/2}, t)) dt}^{F_{\frac{5}{2}}^0} \right. \\ \left. - \overbrace{\int_{t^0}^{t^1} f(q(x_{3/2}, t)) dt}^{F_{\frac{3}{2}}^0} \right) - \overbrace{\int_{t^0}^{t^1} \int_{C_2} s(q(x, t)) dt}^{S_2^0}.$$

## Update Formula for Serre Equations

$$\bar{h}_j^{n+1} = \bar{h}_j^n - [F_{j+1/2}^n - F_{j-1/2}^n],$$

$$\bar{G}_j^{n+1} = \bar{G}_j^n - [F_{j+1/2}^n - F_{j-1/2}^n] - S_j^n.$$

## Update Formula for Serre Equations

$$\bar{h}_j^{n+1} = \bar{h}_j^n - [F_{j+1/2}^n - F_{j-1/2}^n],$$

$$\bar{G}_j^{n+1} = \bar{G}_j^n - [F_{j+1/2}^n - F_{j-1/2}^n] - S_j^n.$$

- ▶ All the fluxes  $F_{j+1/2}^n$  and  $F_{j-1/2}^n$  and the source term  $S_j^n$  require  $u$  at  $t^n$

## Calculate Velocity

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

---

<sup>1</sup>Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

## Calculate Velocity

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right).$$

- ▶ Previously used a Finite Difference Method <sup>1</sup>
- ▶ Contribution: use a Finite Element Method

---

<sup>1</sup>Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

## Finite Element Method

**2D:** Extends well to 2D

**Robust:** Stable in the presence of steep gradients

**Robust:** Stable in the presence of dry beds

- ▶ Maintains conservation properties
- ▶ ANUGA

## Finite Element Method Example

Example:

$$-\frac{\partial^2 q}{\partial x^2} = f,$$

## Finite Element Method Example

Example:

$$-\frac{\partial^2 q}{\partial x^2} = f,$$

Weak Form:

$$-\int_{\Omega} \frac{\partial^2 q}{\partial x^2} v \, dx = \int_{\Omega} fv \, dx,$$

## Finite Element Method Example

Example:

$$-\frac{\partial^2 q}{\partial x^2} = f,$$

Weak Form:

$$-\int_{\Omega} \frac{\partial^2 q}{\partial x^2} v \, dx = \int_{\Omega} fv \, dx,$$

Integrate by parts

$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx = \int_{\Omega} fv \, dx.$$

## Finite Element Method

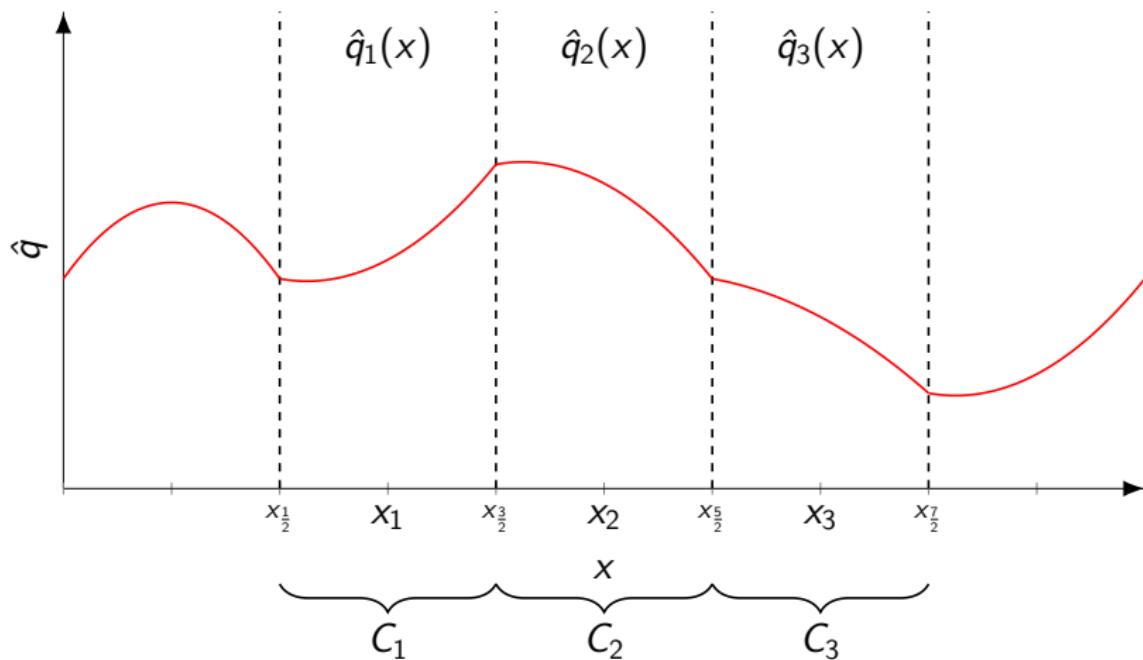
$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx = \int_{\Omega} f v \, dx,$$

## Finite Element Method

$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx = \int_{\Omega} f v \, dx,$$

$$\sum_j \left[ \int_{C_j} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx \right] = \sum_j \left[ \int_{C_j} f v \, dx \right],$$

## Piecewise Polynomial Representation



## Finite Element Method

$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx = \int_{\Omega} f v \, dx,$$

$$\sum_j \left[ \int_{C_j} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx \right] = \sum_j \left[ \int_{C_j} f v \, dx \right],$$

$$\sum_j \left[ \int_{C_j} \frac{\partial \hat{q}_j}{\partial x} \frac{\partial \hat{v}_j}{\partial x} \, dx \right] = \sum_j \left[ \int_{C_j} \hat{f}_j \hat{v}_j \, dx \right],$$

## Finite Element Method

$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx = \int_{\Omega} f v \, dx,$$

$$\sum_j \left[ \int_{C_j} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} \, dx \right] = \sum_j \left[ \int_{C_j} f v \, dx \right],$$

$$\sum_j \left[ \int_{C_j} \frac{\partial \hat{q}_j}{\partial x} \frac{\partial \hat{v}_j}{\partial x} \, dx \right] = \sum_j \left[ \int_{C_j} \hat{f}_j \hat{v}_j \, dx \right],$$

$$\mathbf{A}\vec{q} = \vec{c}.$$

## Finite Element Method

$$\int_{\Omega} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} dx = \int_{\Omega} f v dx,$$

$$\sum_j \left[ \int_{C_j} \frac{\partial q}{\partial x} \frac{\partial v}{\partial x} dx \right] = \sum_j \left[ \int_{C_j} f v dx \right],$$

$$\sum_j \left[ \int_{C_j} \frac{\partial \hat{q}_j}{\partial x} \frac{\partial \hat{v}_j}{\partial x} dx \right] = \sum_j \left[ \int_{C_j} \hat{f}_j \hat{v}_j dx \right],$$

$$\mathbf{A}\vec{q} = \vec{c}.$$

where

- ▶  $\mathbf{A}$  depends on  $\hat{v}_j$
- ▶  $\vec{q}$  determines  $\hat{q}_j$
- ▶  $\vec{c}$  depends on  $\hat{f}_j$  and  $\hat{v}_j$

## Finite Element Method for Serre Equations

$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right),$$

## Finite Element Method for Serre Equations

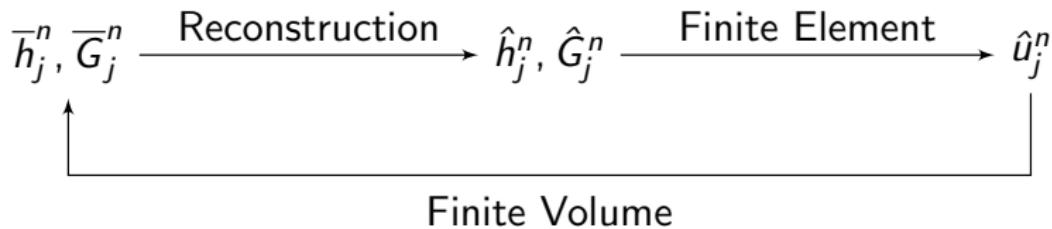
$$G = hu \left( 1 + \frac{\partial h}{\partial x} \frac{\partial b}{\partial x} + \frac{1}{2} h \frac{\partial^2 b}{\partial x^2} + \left[ \frac{\partial b}{\partial x} \right]^2 \right) - \frac{\partial}{\partial x} \left( \frac{1}{3} h^3 \frac{\partial u}{\partial x} \right),$$

$$\mathbf{A}\vec{u} = \vec{c}.$$

where

- ▶  $\mathbf{A}$  depends on the polynomial representation of  $h$ ,  $b$  and test function.
- ▶  $\vec{u}$  determines the polynomial representation of  $u$ .
- ▶  $\vec{c}$  depends on polynomial representation of  $G$  and test function.

# Method



## Progress

2D: 1D method that extends well to 2D ✓

Robust: Validation for steep gradients in free surface

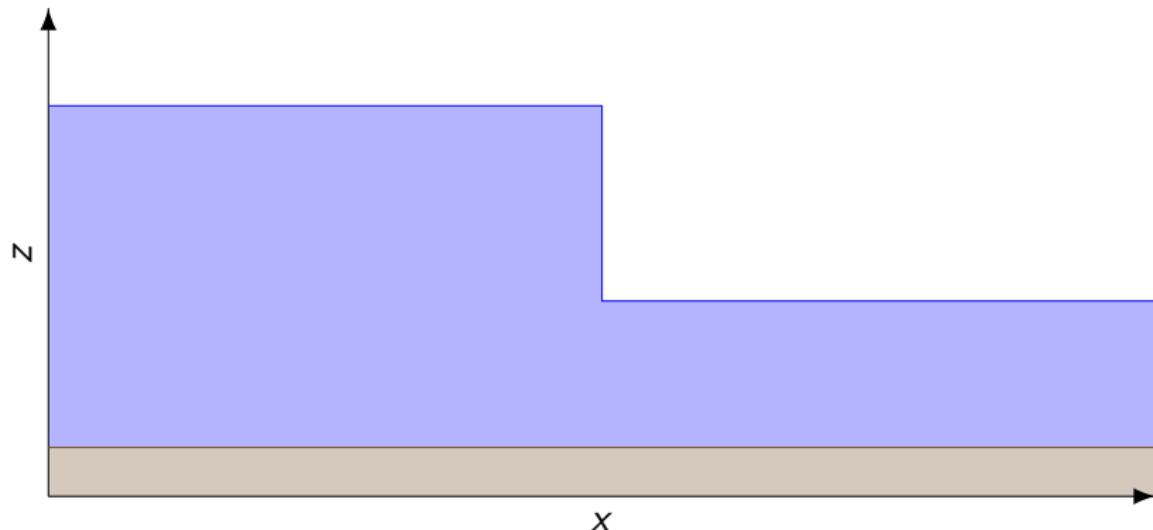
Robust: Inclusion and validation of dry beds

# Validation

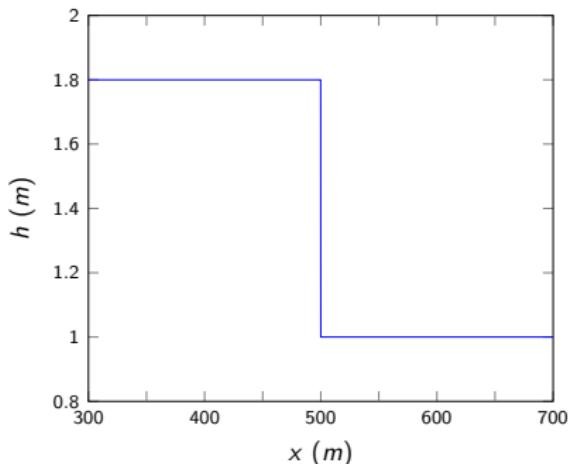
- ▶ Steep gradients in the free surface
- ▶ Dry beds

## Statement of Problem

How does this initially still body of water evolve?



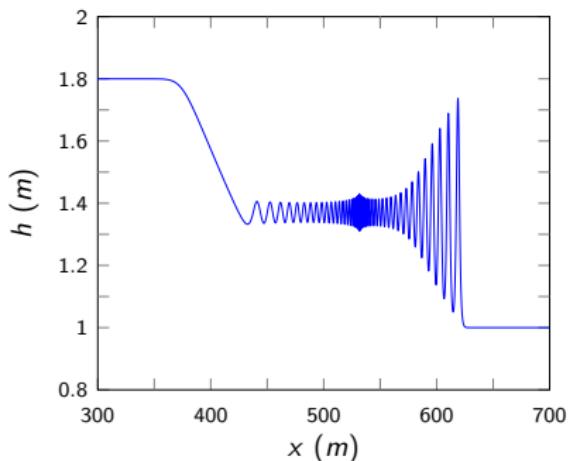
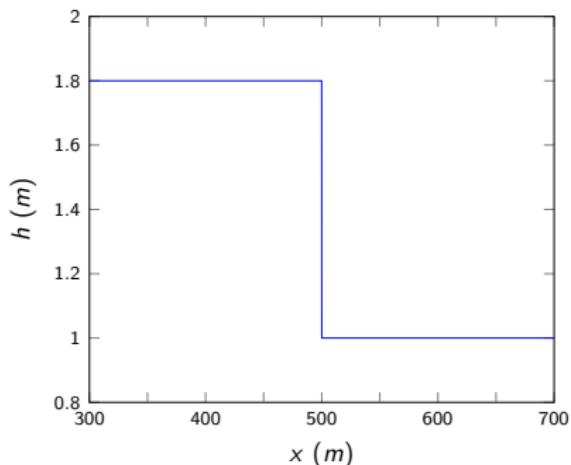
## Our New Numerical Solution



---

Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. *Wave Motion*, 76(1):6177.

## Our New Numerical Solution



---

Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. *Wave Motion*, 76(1):6177.

# Evolution

## What was known

- ▶ No analytic solutions
- ▶ Some experimental comparisons <sup>2</sup>
- ▶ Other numerical solutions from the literature

---

<sup>2</sup>Zoppou, C. (2014). Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows. PhD thesis, Australian National University.

## Contribution

- ▶ Observed this new behaviour
- ▶ Demonstrated convergence
- ▶ Comprehensive review of numerical solutions from the literature

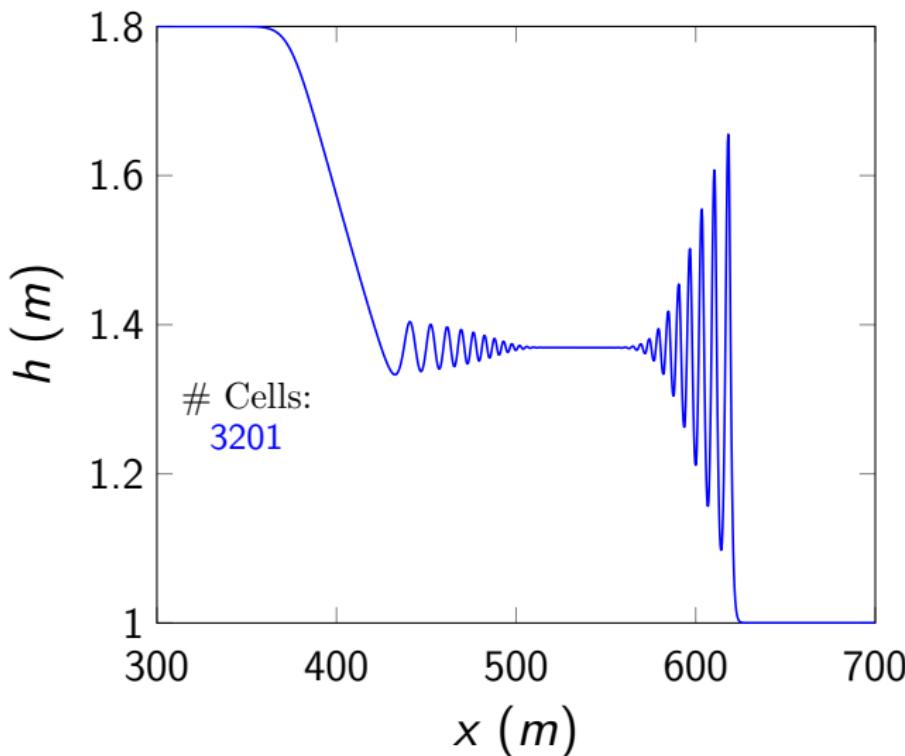
---

Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. *Wave Motion*, 76(1):6177.

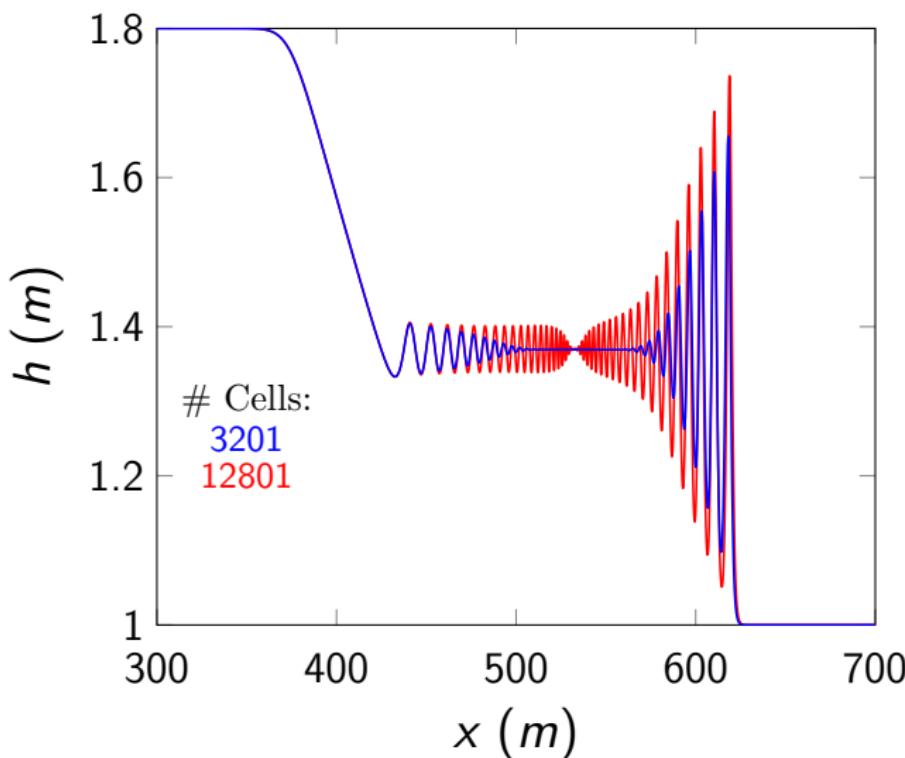
## Convergence

As we increase number of cells the numerical solutions should converge to the true solution

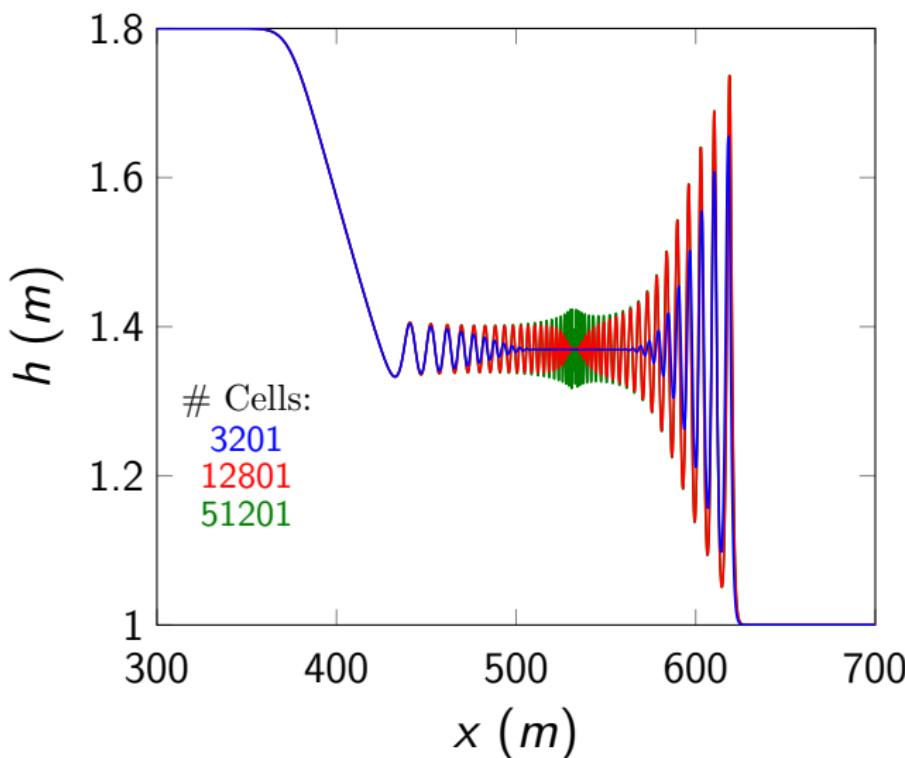
## Convergence



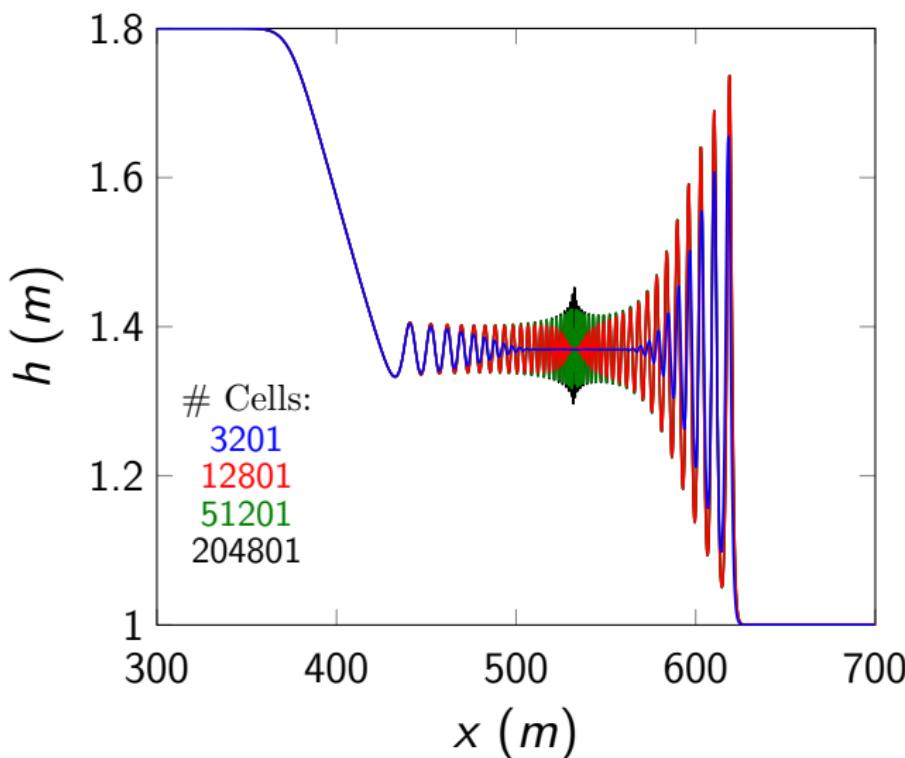
## Convergence



## Convergence



## Convergence



## Comprehensive Review

- ▶ Demonstrated consistent behaviour across many numerical methods
- ▶ Were able to explain why the behaviour had not previously been observed

## Result

Validated our computational model when steep gradients are present in the free surface.<sup>3</sup>

---

<sup>3</sup>Pitt, J., Zoppou, C., and Roberts, S. (2018). Behaviour of the Serre Equations in the Presence of Steep Gradients Revisited. Wave Motion, 76(1):6177.

## Progress

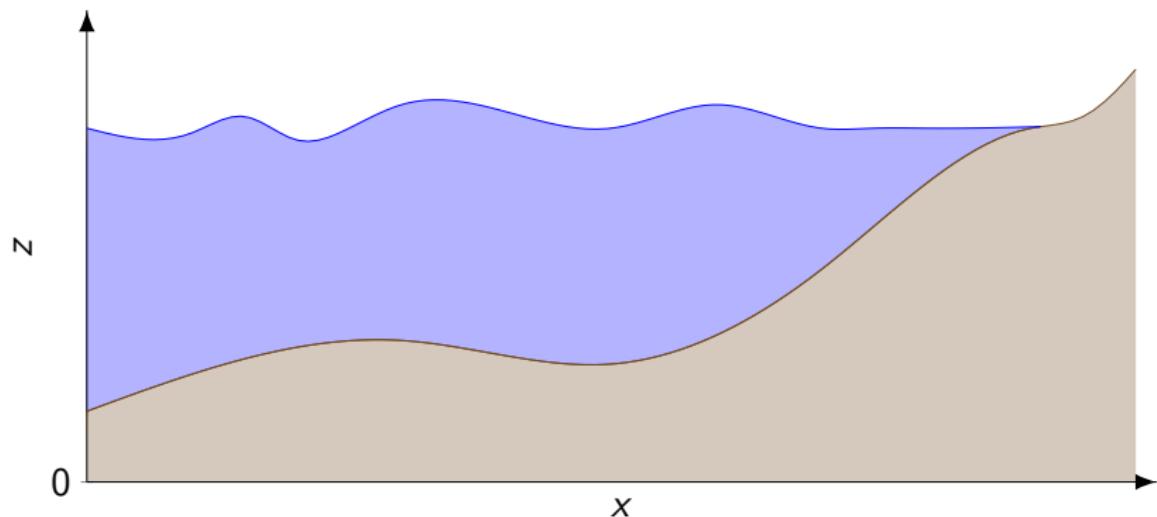
2D: 1D method that extends well to 2D ✓

Robust: Validation for steep gradients in free surface ✓

Robust: Inclusion and validation of dry beds

## Statement of Problem

Properly handle interaction of waves and the dry bed



## What was known

- ▶ No analytic solutions
- ▶ A variety of numerical techniques only compared to experimental data

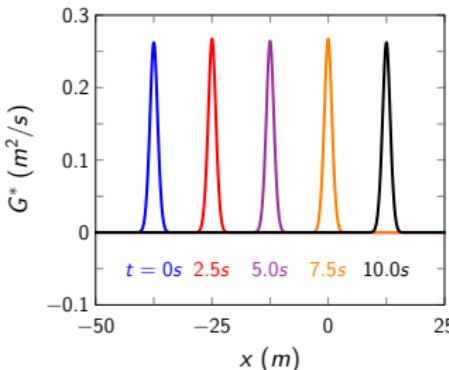
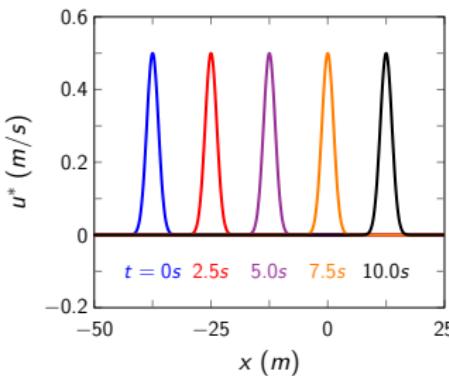
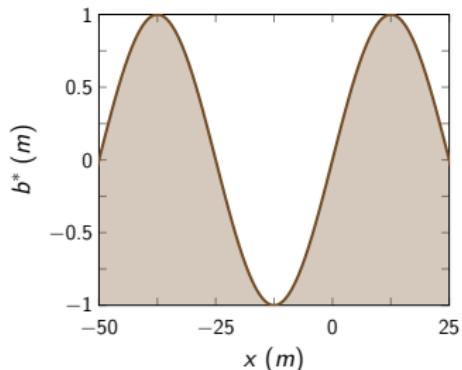
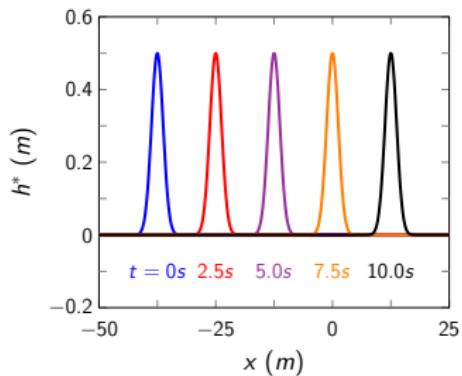
## Contribution

- ▶ Solved modified equations that did possess analytic solutions
- ▶ Compared with experimental data

## Constructing Modified Equations

- ▶ Pick functions for height, velocity and bed:  $h^*$ ,  $u^*$  and  $b^*$
- ▶ Add source terms to Serre equations that force  $h^*$ ,  $u^*$  and  $b^*$  to be solutions
- ▶ Validation tests

# Pick Functions



## Modify Equations

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = S_h^*,$$

$$\begin{aligned} \frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left( uG + \frac{gh^2}{2} - \frac{2}{3}h^3 \left[ \frac{\partial u}{\partial x} \right]^2 + h^2 u \frac{\partial u}{\partial x} \frac{\partial b}{\partial x} \right) \\ + \frac{1}{2}h^2 u \frac{\partial u}{\partial x} \frac{\partial^2 b}{\partial x^2} - hu^2 \frac{\partial b}{\partial x} \frac{\partial^2 b}{\partial x^2} + gh \frac{\partial b}{\partial x} = S_G^*. \end{aligned}$$

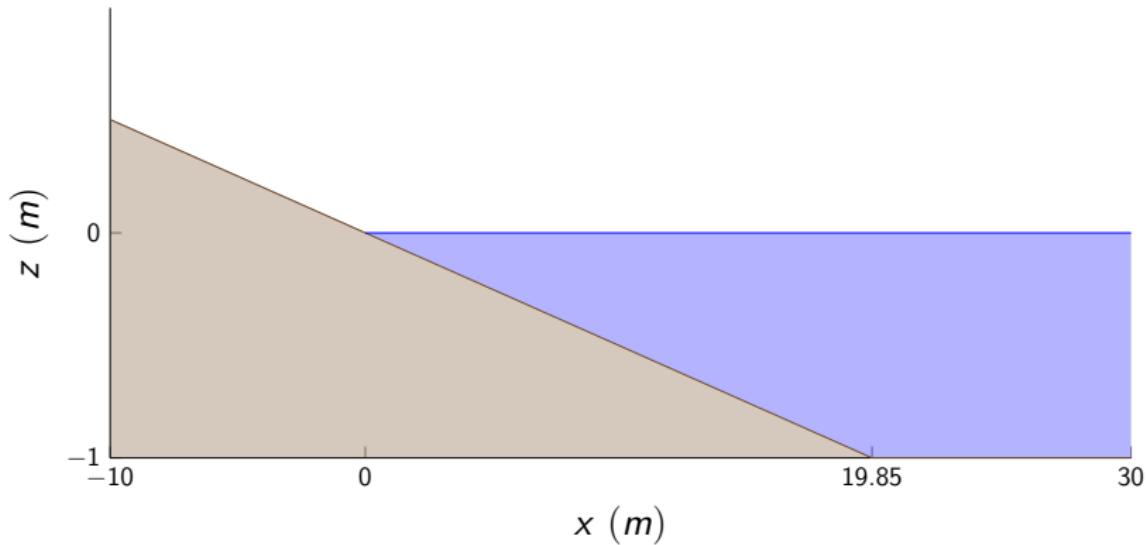
$S_h^*$  and  $S_G^*$  are just the LHS with the quantities replaced by their associated chosen function.

# Results

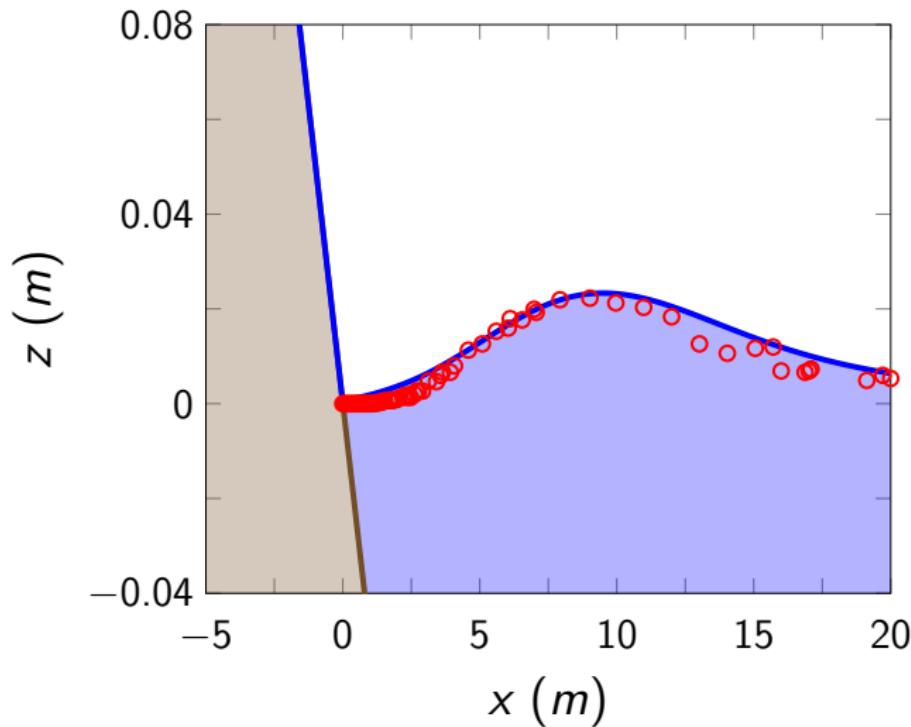
## Modified Equations Validation Conclusions

- ▶ Very strong test as all terms must be accurately approximated
- ▶ Only source of error is the method

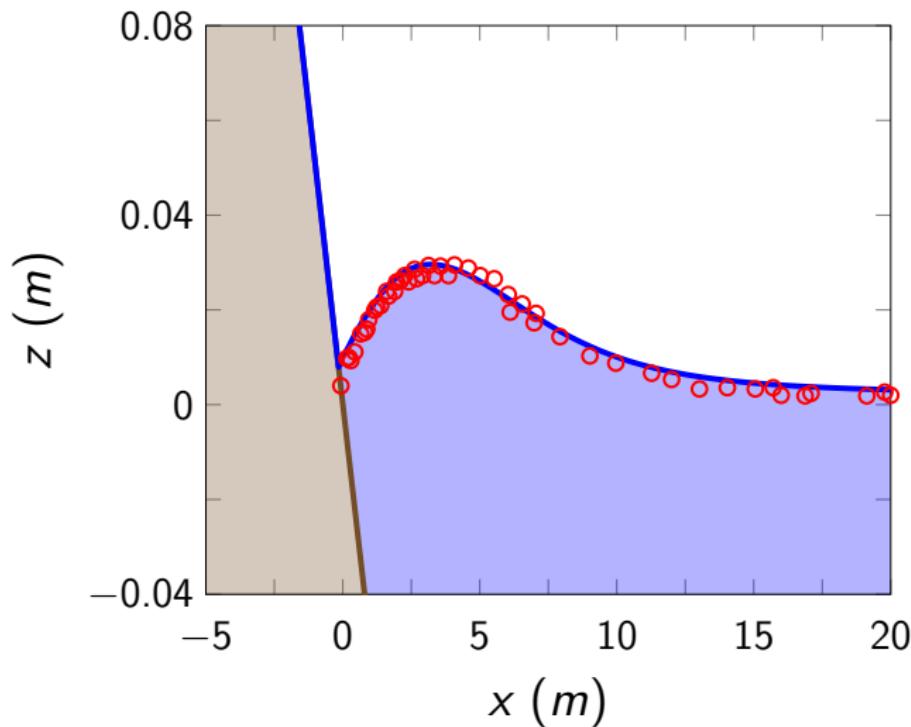
# Experimental Data



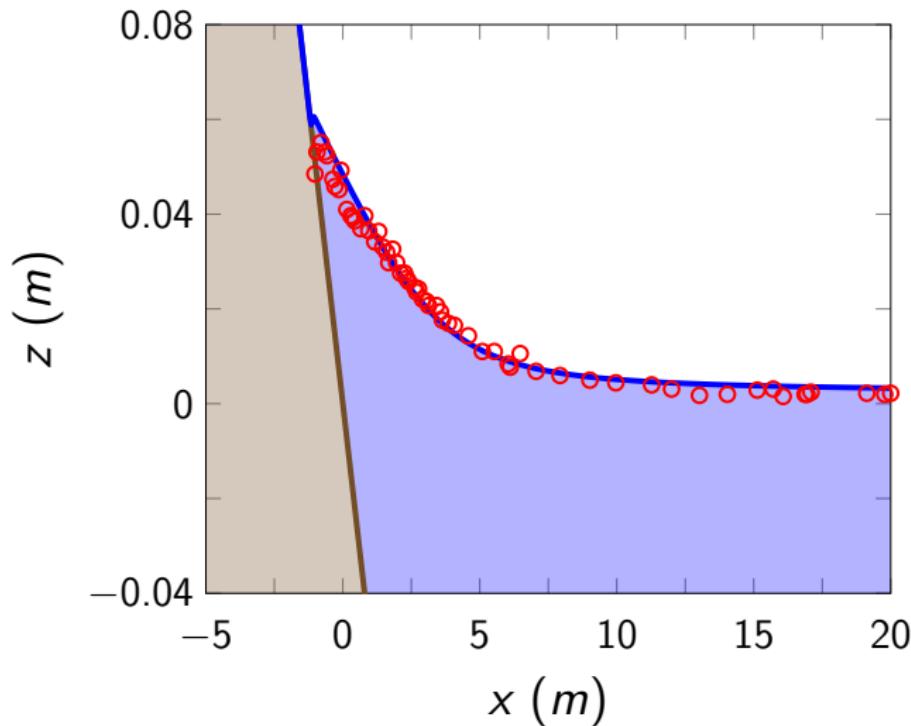
$t = 30s$



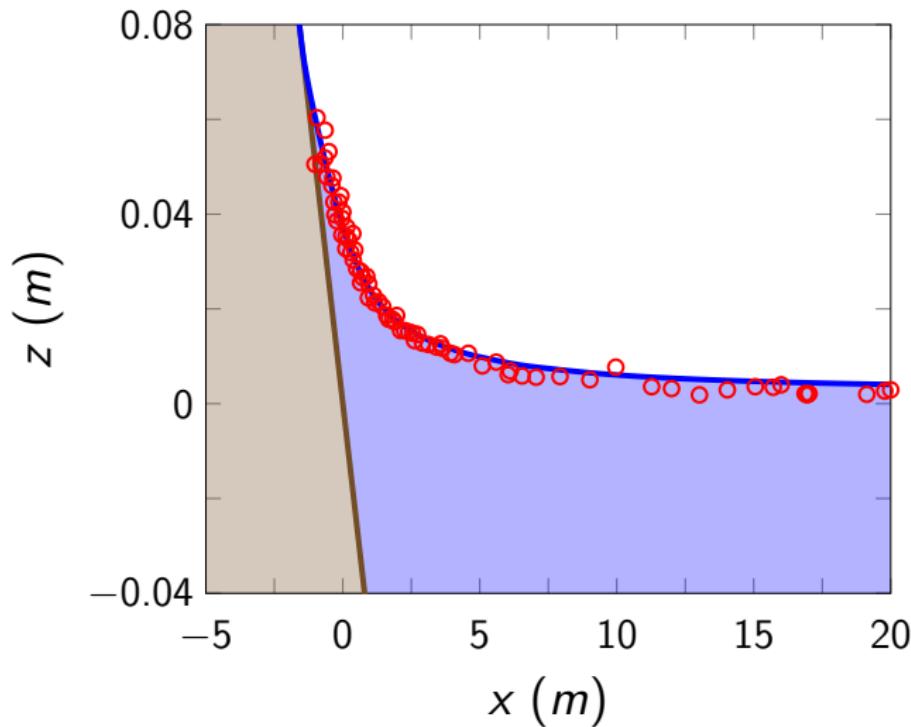
$t = 40s$



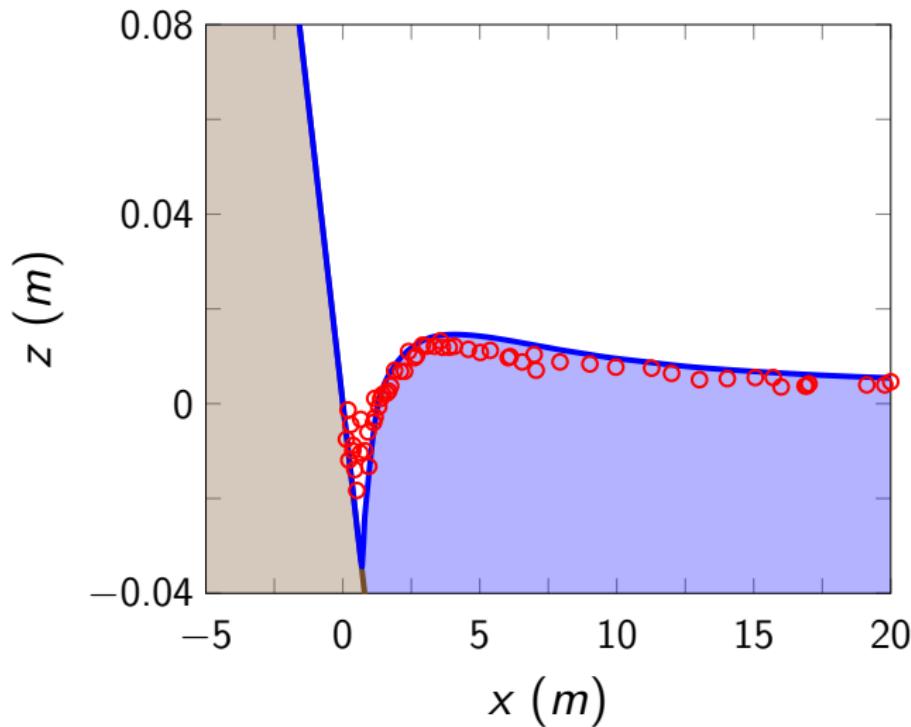
$t = 50s$



$t = 60s$



$t = 70s$



## Experimental Validation Conclusions

- ▶ Demonstrates agreement of computational model and physical process
- ▶ Many sources of errors

## Progress

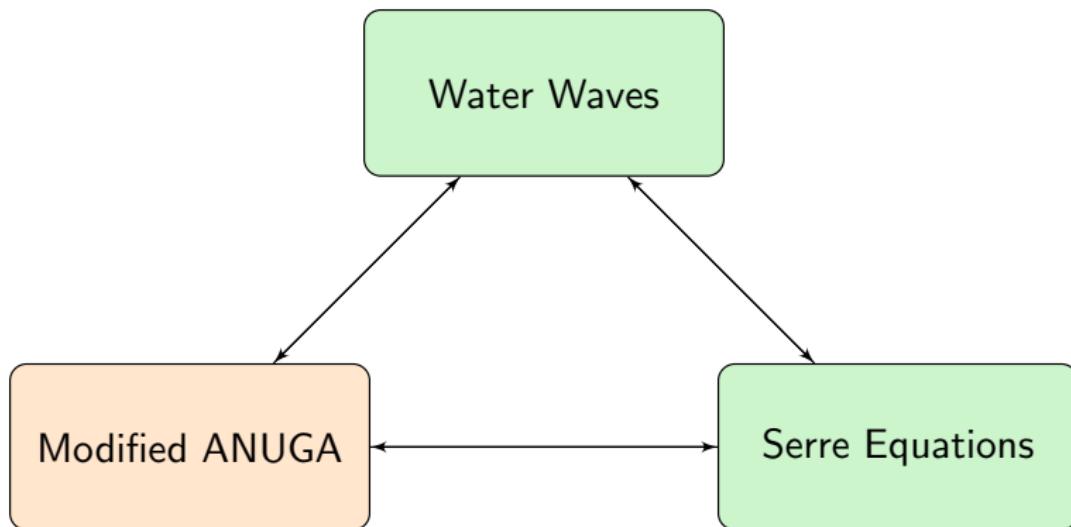
2D: 1D method that extends well to 2D ✓

Robust: Validation for steep gradients in free surface ✓

Robust: Inclusion and validation of dry beds ✓

## Conclusions

- ▶ Developed a robust computational model from the 1D Serre equations for the 2D water wave problem



## References |

Pitt, J., Zoppou, C., and Roberts, S. (2018).

Behaviour of the serre equations in the presence of steep gradients revisited.

*Wave Motion*, 76(1):61–77.

Zoppou, C. (2014).

*Numerical Solution of the One-dimensional and Cylindrical Serre Equations for Rapidly Varying Free Surface Flows.*

PhD thesis, Australian National University, Mathematical Sciences Institute, College of Physical and Mathematical Sciences, Australian National University, Canberra, ACT 2600, Australia.