

A comparison of different order hybrid finite difference-volume methods for solving the Serre equations in conservative law form

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ABSTRACT

In this paper a first-, second- and third-order version of a numerical method for the Serre equations are described. Such a method is described as a finite difference-volume method due to the use of a finite difference approximation and a finite volume method to solve the Serre equations in conservation law form. These methods are then validated and used to investigate a discrepancy in the literature about the results of solving the Serre equations in the presence of steep gradients.

Keywords: dispersive waves, conservation laws, Serre equation, finite volume method, finite difference method

¹ INTRODUCTION

Free surface flows occur in many important applications such as; tsunamis, storm surges and tidal bores. Because viscosity has a negligible effect on these problems the Euler equations can be used to model them. However, numerical methods for the Euler equations are not yet computationally efficient enough to deal with these problems over large domains. Thus various approximations to the Euler equations have been derived, the most crude being the shallow water wave equations which have been used to model the named applications in the past. However, the shallow water wave equations assume a hydrostatic pressure distribution in a fluid column which is not fully justified in rapidly varying flows because vertical acceleration inside the fluid becomes important. Consequently, many equations have been derived as approximations to the Euler equations in

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shallow fluids that are less restrictive in their assumptions than the shallow water wave equations. The Serre equations are one of these and are of particular interest because they do not enforce a hydrostatic pressure distribution over a fluid column allowing for fully non-linear flows (Lannes and Bonneton 2009).

The Serre equations were first derived by Serre (1953) for flat bottom topographies in one dimension. More general equations were then derived for smooth bottom topographies in one dimension (Su and Gardner 1969) and later smooth bottom topographies in two dimensions (Green and Naghdi 1976). These equations have been handled in many different ways (Mitsotakis et al. 2014; Bonneton et al. 2011; Antunes do Carmo et al. 1993; Chazel et al. 2011; Cienfuegos and Bonneton 2006; Cienfuegos and Bonneton 2007; Dutykh et al. 2011). This paper follows the decomposition of the Serre equations into conservative law form (Le Métayer et al. 2010; Li et al. 2014) and then follows the formulation of Le Métayer et al. (2010) and Zoppou (2014) to build first-, second- and third-order methods.

Zoppou and Roberts (1996) demonstrated that first- and third-order methods produce diffusive errors smearing steep gradients. While second-order methods produce dissipative errors introducing non-physical oscillations around steep gradients. Because steep gradients arise naturally in fluid flows and the Serre equations produce dispersive waves around them (El et al. 2006) it is important that for the described methods these oscillations are not significantly polluted by either diffusion or dissipation.

This paper aims to clear up the discrepancy between the results of Le Métayer et al. (2010) and El et al. (2006) by examining the behaviour of a certain dam-break problem. In particular in Le Métayer et al. (2010) it is stated that their first-order method was sufficient to capture the important behaviour of the dam-break problem, this paper will test the validity of that assertion. To accomplish this first-, second- and third-order methods to solve the Serre equations are constructed and validated by comparing the solutions to known analytical solutions and laboratory data of flows containing steep gradients. Then the results for a dam-break problem for all described methods are compared to investigate this conflict.

SERRE EQUATIONS

The Serre equations can be derived as an approximation to the full Euler equations by depth integration similar to Su and Gardner (1969). They can also be seen as an asymptotic expansion of the Euler equations (Lannes and Bonneton 2009). The former is more consistent with the perspective from which numerical methods will be developed in this paper while the latter indicates the appropriate regions in which to use these equations as a model of fluid flow.

The scenario under which the Serre approximation is made consists of a two dimensional $\mathbf{x} = (x, z)$ fluid over a bottom topography as in Figure 1, under the action of

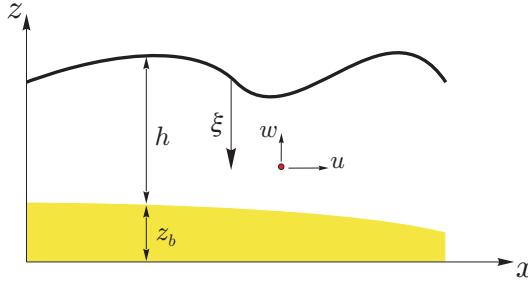


FIG. 1: The notation used for one-dimensional flow governed by the Serre equation.

50 gravity. The water depth is $h(x, t)$ and $z_b(x)$ is the bed elevation. The fluid is subject
 51 to the pressure, $p(\mathbf{x}, t)$ and gravitational acceleration, $\mathbf{g} = (0, g)^T$ and has a velocity
 52 $\mathbf{v} = (u(\mathbf{x}, t), w(\mathbf{x}, t))$, where $u(\mathbf{x}, t)$ is the velocity in the x -coordinate and $w(\mathbf{x}, t)$ is the
 53 velocity in the z -coordinate and t is time. Assuming that $z_b(x)$ is constant the Serre equa-
 54 tions read (Li et al. 2014; Zoppou 2014)

$$55 \quad \frac{\partial h}{\partial t} + \frac{\partial(\bar{u}h)}{\partial x} = 0, \quad (1a)$$

56

$$57 \quad \underbrace{\frac{\partial(\bar{u}h)}{\partial t} + \frac{\partial}{\partial x} \left(\bar{u}^2 h + \frac{gh^2}{2} \right)}_{\text{Shallow Water Wave Equations}} + \underbrace{\frac{\partial}{\partial x} \left(\frac{h^3}{3} \left[\frac{\partial \bar{u}}{\partial x} \frac{\partial \bar{u}}{\partial x} - \bar{u} \frac{\partial^2 \bar{u}}{\partial x^2} - \frac{\partial^2 \bar{u}}{\partial x \partial t} \right] \right)}_{\text{Dispersion Terms}} = 0 \quad (1b)$$

58

59 where \bar{u} is the average of u over the depth of water.

60 Alternative Conservation Law Form of the Serre Equations

61 In Le Métayer et al. (2010) and Zoppou (2014) it is demonstrated that the Serre equa-
 62 tions can be rearranged into a conservation law form, by introducing a new quantity

$$63 \quad G = uh - h^2 \frac{\partial h}{\partial x} \frac{\partial u}{\partial x} - \frac{h^3}{3} \frac{\partial^2 u}{\partial x^2}. \quad (2)$$

64 Consequently, the equations can be rewritten as

$$65 \quad \frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0 \quad (3a)$$

69 and

70
$$\frac{\partial G}{\partial t} + \frac{\partial}{\partial x} \left(Gu + \frac{gh^2}{2} - \frac{2h^3}{3} \frac{\partial u}{\partial x} \frac{\partial u}{\partial x} \right) = 0 \quad (3b)$$

71

72 where the bar over u has been dropped to simplify the notation. A hybrid method can
73 be developed for the Serre equations that solves the elliptic problem (2) for u and then
74 the conservation law (3) for h and G at some later time. This replicates the process of
75 Le Métayer et al. (2010) and Zoppou (2014).

76 **NUMERICALLY SOLVING THE SERRE EQUATIONS WRITTEN IN
77 CONSERVATION LAW FORM**

78 There are numerous ways a numerical method could be built to solve the Serre equa-
79 tions in conservation law form (3). For flows that contain steep gradients the finite volume
80 method seems the most appropriate. A finite volume method to solve (3) updates the con-
81 served quantities h and G over a single time step Δt for instance from time t^n to t^{n+1} . So
82 that

83
$$\begin{bmatrix} h^{n+1} \\ G^{n+1} \end{bmatrix} = \mathcal{L}(h^n, G^n, u^n, \Delta t) \quad (4)$$

84

85 where \mathcal{L} is some numerical solver for (3) and the superscript denotes the time at which a
86 quantity is evaluated; e.g. $u^n = u(t^n)$. The complete solution also involves solving (2) for
87 u given h and G denoted by

88
$$u^{n+1} = \mathcal{A}(h^{n+1}, G^{n+1}). \quad (5)$$

90

91 **SOLVING THE ELLIPTIC EQUATION \mathcal{A} FOR U**

92 Assuming that a discretisation in space has a fixed resolution so that $\forall i x_{i+1} - x_i = \Delta x$;
93 allows for a simple finite difference approximation to (2) as a suitable method for \mathcal{A}
94 (Le Métayer et al. 2010; Zoppou 2014). Since the goal of this paper is to develop and com-
95 pare a range of different order methods for this problem both a second- and fourth-order
96 centred finite difference approximation to (2) were used. By taking such approximations
97 to the first- and second-order spatial derivatives the second- and fourth-order analogues of
98 (2) are given by

99
$$G_i = u_i h_i - h_i^2 \left(\frac{h_{i+1} - h_{i-1}}{2\Delta x} \right) \left(\frac{u_{i+1} - u_{i-1}}{2\Delta x} \right) - \frac{h_i^3}{3} \left(\frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} \right)$$

100

101 and

$$102 \quad G_i = u_i h_i - h_i^2 \left(\frac{-h_{i+2} + 8h_{i+1} - 8h_{i-1} + h_{i-2}}{12\Delta x} \right) \left(\frac{-u_{i+2} + 8u_{i+1} - 8u_{i-1} + u_{i-2}}{12\Delta x} \right)$$

$$103 \quad - \frac{h_i^3}{3} \left(\frac{-u_{i+2} + 16u_{i+1} - 30u_i + 16u_{i-1} - u_{i-2}}{12\Delta x^2} \right)$$

104 where the subscript denotes the spatial coordinate at which the quantity is evaluated; e.g.
105 $u_i = u(x_i)$. Both of these can be rearranged into a matrix equation with the following form

$$106 \quad \begin{bmatrix} u_0 \\ \vdots \\ u_m \end{bmatrix} = A^{-1} \begin{bmatrix} G_0 \\ \vdots \\ G_m \end{bmatrix} =: \mathcal{A}(\mathbf{h}, \mathbf{G})$$

$$107$$

108 where for a second-order approximation the matrix A is tri-diagonal while for a fourth-
109 order method it is penta-diagonal.

110 SOLVING THE CONSERVATION LAW FORM OF THE SERRE EQUATIONS

111 A finite volume method of sufficient order was developed to solve (3). Unlike finite
112 difference methods which utilise nodal values of quantities, finite volume methods use the
113 cell averages, for example the average water depth over a cell which spans $[x_{i-1/2}, x_{i+1/2}]$
114 is

$$115 \quad \bar{h}_i = \frac{1}{\Delta x} \int_{x_{i-1/2}}^{x_{i+1/2}} h(x, t) dx$$

$$116$$

117 where $x_{i\pm 1/2} = x_i \pm \Delta x/2$. Finite volume methods update the cell averages by the following
118 method

$$119 \quad \bar{U}_i^{n+1} = \bar{U}_i^n - \frac{\Delta t}{\Delta x} \left(F_{i+1/2}^n - F_{i-1/2}^n \right) \quad (6)$$

$$120$$

121 where $\bar{U}_i^n = [\bar{h}_i^n \bar{G}_i^n]^T$ is an approximation of the vector of the conserved quantities aver-
122 aged over the cell at time t^n . While $F_{i\pm 1/2}^n$ is an approximation of the average flux at the
123 respective cell boundary $x_{i\pm 1/2}$ over the time interval $[t^n, t^{n+1}]$, which is obtained by solving
124 the Riemann problem at the cell boundaries.

125 Local Riemann Problem

126 Since \bar{U}_i^n is known what remains is to calculate the time averaged fluxes $F_{i\pm 1/2}$. In
127 Kurganov et al. (2002) the time averaged inter-cell flux is approximated by

$$128 \quad F_{i+1/2} = \frac{a_{i+1/2}^+ f(q_{i+1/2}^-) - a_{i+1/2}^- f(q_{i+1/2}^+)}{a_{i+1/2}^+ - a_{i+1/2}^-} + \frac{a_{i+1/2}^+ a_{i+1/2}^-}{a_{i+1/2}^+ - a_{i+1/2}^-} \left[q_{i+1/2}^+ - q_{i+1/2}^- \right] \quad (7)$$

$$129$$

130 where f is the instantaneous flux of the conserved quantity q evaluated using the recon-
 131 structed values from the cells that border the cell interface $x_{i+1/2}$. While $a_{i+1/2}^-$ and $a_{i+1/2}^+$
 132 are given by

$$133 \quad 134 \quad a_{i+\frac{1}{2}}^- = \min \left[\lambda_1 \left(q_{i+\frac{1}{2}}^- \right), \lambda_1 \left(q_{i+\frac{1}{2}}^+ \right), 0 \right]$$

135 and

$$136 \quad 137 \quad a_{i+\frac{1}{2}}^+ = \max \left[\lambda_2 \left(q_{i+\frac{1}{2}}^- \right), \lambda_2 \left(q_{i+\frac{1}{2}}^+ \right), 0 \right]$$

138 where λ_1 and λ_2 are estimates of the smallest and largest eigenvalues respectively of the
 139 Jacobian.

140 Propagation Speeds of a Local Shock

141 As demonstrated in Zoppou (2014) λ_1 and λ_2 are bounded by the phase speed of the
 142 shallow water wave equations, so that

$$143 \quad \lambda_1 := u - \sqrt{gh} \leq v_p \leq u + \sqrt{gh} =: \lambda_2$$

144 where v_p is the phase speed of the Serre equations, thus $a_{i+1/2}^-$ and $a_{i+1/2}^+$ are fully deter-
 145 mined.

146 Reconstruction

147 The quantities $q_{i+1/2}^-$ and $q_{i+1/2}^+$ in (7) are given by the two reconstructions at $x_{i+1/2}$,
 148 one from the cell to the left $[x_{i-1/2}, x_{i+1/2}]$ and one from the cell to the right $[x_{i+1/2}, x_{i+3/2}]$
 149 denoted by the superscripts $-$ and $+$ respectively. The order of the polynomials used
 150 to reconstruct the quantities inside the cells determines the spatial order of the method.
 151 Constant polynomials result in a first-order method (Godunov 1959). Similarly first- and
 152 second-degree polynomials result in second- and third-order methods respectively.

153 For a zero-degree polynomial the interpolant has the value \bar{q}_i at x_i , this is also the
 154 case for linear interpolation functions. For the zero-degree case the interpolants are fully
 155 determined i.e $q_{i-1/2}^+ = \bar{q}_i = q_{i+1/2}^-$ and monotinicity preserving. There are a variety of
 156 ways to construct higher-degree interpolants not all of which are necessarily monotonicity
 157 preserving, which can result in the introduction of numerical oscillations during the re-
 158 construction process. To suppress these non-physical oscillations in higher order methods
 159 limiting must be implemented. For the second-order method the minmod limiter was used
 160 as in Kurganov et al. (2002). While for the third-order method the Koren limiter was used
 161 (Koren 1993). This results in the following reconstruction scheme for the second-order
 162 method

$$164 \quad 165 \quad q_{i+\frac{1}{2}}^- = \bar{q}_i + a_i \frac{\Delta x}{2}$$

166 and

167

168

$$q_{i+\frac{1}{2}}^+ = \bar{q}_{i+1} - a_{i+1} \frac{\Delta x}{2}$$

169 where

170

171

$$a_i = \text{minmod} \left\{ \theta \frac{\bar{q}_{i+1} - \bar{q}_i}{\Delta x}, \frac{\bar{q}_{i+1} - \bar{q}_{i-1}}{2\Delta x}, \theta \frac{\bar{q}_i - \bar{q}_{i-1}}{\Delta x} \right\} \quad \text{for } \theta \in [1, 2].$$

172 While for the third-order method the reconstruction scheme is

173

174

$$q_{i+\frac{1}{2}}^- = \bar{q}_i + \frac{1}{2} \phi^-(r_i) (\bar{q}_i - \bar{q}_{i-1})$$

175 and

176

177

$$q_{i+\frac{1}{2}}^+ = \bar{q}_i - \frac{1}{2} \phi^+(r_i) (\bar{q}_i - \bar{q}_{i-1})$$

178 where

179

180

181

$$\phi^-(r_i) = \max \left[0, \min \left[2r_i, \frac{1+2r_i}{3}, 2 \right] \right],$$

182

183

$$\phi^+(r_i) = \max \left[0, \min \left[2r_i, \frac{2+r_i}{3}, 2 \right] \right]$$

184 and

185

186

$$r_i = \frac{\bar{q}_{i+1} - \bar{q}_i}{\bar{q}_i - \bar{q}_{i-1}}.$$

187 **Fully discrete approximations to the instantaneous flux $f(q_{i+\frac{1}{2}}^\pm)$**

188 For water depth, the fully discrete approximation to $f(h_{i+1/2}^\pm)$ is given by

189

190

$$f\left(h_{i+\frac{1}{2}}^\pm\right) = u_{i+\frac{1}{2}}^\pm h_{i+\frac{1}{2}}^\pm$$

191 which is independent of the order of accuracy of the method.

192 The flux $f(G_{i+1/2}^\pm)$ is more complicated because of the derivative and is given by

193

194

$$f\left(G_{i+\frac{1}{2}}^\pm\right) = u_{i+\frac{1}{2}}^\pm G_{i+\frac{1}{2}}^\pm + \frac{g\left(h_{i+\frac{1}{2}}^\pm\right)^2}{2} - \frac{2\left(h_{i+\frac{1}{2}}^\pm\right)^3}{3} \left[\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^\pm \right]^2.$$

195 There are multiple ways to approximate this derivative. The first- and third-order approx-
 196 imations to the derivatives can be obtained by an upwind finite difference approximation.
 197 By assuming that u is continuous, a second-order approximation that has the correct order
 198 and is simpler to implement than its corresponding upwind finite difference approximation
 199 can be used. Thus the following approximations to the derivatives were obtained, for the
 200 first-order method

201
$$\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^+ = \frac{u_{i+\frac{3}{2}}^+ - u_{i+\frac{1}{2}}^+}{\Delta x},$$

 202
 203

204
$$\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^- = \frac{u_{i+\frac{1}{2}}^- - u_{i-\frac{1}{2}}^-}{\Delta x}.$$

 205

206 For the second-order method

207
$$\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^- = \left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^+ = \frac{u_{i+1} - u_i}{\Delta x}$$

 208

209 and for the third-order method

210
$$\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^+ = \frac{-u_{i+\frac{3}{2}}^+ + 4u_{i+\frac{1}{2}}^+ - 3u_{i-\frac{1}{2}}^+}{\Delta x}$$

 211

212 and

213
$$\left(\frac{\partial u}{\partial x} \right)_{i+\frac{1}{2}}^- = \frac{3u_{i+\frac{1}{2}}^- - 4u_{i-\frac{1}{2}}^- + u_{i-\frac{3}{2}}^-}{\Delta x}.$$

 214

215 **Transforming between nodal values and cell averages**

216 The operator \mathcal{L} in (4) uses cell averages while the operator \mathcal{A} in (5) uses nodal values
 217 at the cell centres. Therefore, a transformation from the cell averages to the nodal values
 218 is required. For the first- and second-order methods this distinction is trivial since $\bar{q}_i = q_i$.
 219 However, for the third-order method this is a very important distinction and failure to
 220 handle this correctly will result in a loss of accuracy.

221 A quadratic polynomial that gives the correct cell averages for the cell centred at x_i
 222 and its two neighbours satisfies this equation

223
$$q_i = \frac{-\bar{q}_{i+1} + 26\bar{q}_i - \bar{q}_{i-1}}{24}$$

 224

225 at the cell centre x_i . This is a tri-diagonal matrix equation that transforms from cell aver-
 226 ages to nodal values with third-order accuracy and is denoted by \mathcal{M} . The inverse trans-
 227 formation \mathcal{M}^{-1} denotes the solution of the tri-diagonal matrix equation given nodal values
 228 resulting in cell averages which is also third-order accurate. This completes the solution
 229 of the Serre equations (2) and (3) with the following process denoted by \mathcal{H}

230

$$\mathcal{H}(\bar{\mathbf{U}}^n, \Delta t) = \begin{cases} \mathbf{U}^n &= \mathcal{M}(\bar{\mathbf{U}}^n) \\ \mathbf{u}^n &= \mathcal{A}(\mathbf{U}^n) \\ \bar{\mathbf{u}}^n &= \mathcal{M}^{-1}(\mathbf{u}^n) \\ \bar{\mathbf{U}}^{n+1} &= \mathcal{L}(\bar{\mathbf{U}}^n, \bar{\mathbf{u}}^n, \Delta t) \end{cases}.$$

231

232 **Strong-Stability-Preserving Runge-Kutta Scheme**

233 The process above is first-order accurate in time. There are many methods to increase
 234 the accuracy of a time step evolution. This paper will use the strong stability Runge-Kutta
 235 steps described in Gottlieb et al. (2009) to construct fully second- and third-order methods
 236 by linear combinations of \mathcal{H} . This leads to the following processes, for the first-order
 237 method

238

$$\bar{\mathbf{U}}^{n+1} = \mathcal{H}(\bar{\mathbf{U}}^n, \Delta t)$$

239

240 the second-order method

241

$$\bar{\mathbf{U}}^{(1)} = \mathcal{H}(\bar{\mathbf{U}}^n, \Delta t),$$

242

$$\bar{\mathbf{U}}^{(2)} = \mathcal{H}(\bar{\mathbf{U}}^{(1)}, \Delta t),$$

243

$$\bar{\mathbf{U}}^{n+1} = \frac{1}{2}(\bar{\mathbf{U}}^{(1)} + \bar{\mathbf{U}}^{(2)})$$

244

245 and the third-order method

246

$$\bar{\mathbf{U}}^{(1)} = \mathcal{H}(\bar{\mathbf{U}}^n, \Delta t),$$

247

$$\bar{\mathbf{U}}^{(2)} = \mathcal{H}(\bar{\mathbf{U}}^{(1)}, \Delta t),$$

248

$$\bar{\mathbf{U}}^{(3)} = \frac{3}{4}\bar{\mathbf{U}}^n + \frac{1}{4}\bar{\mathbf{U}}^{(2)},$$

249

$$\bar{\mathbf{U}}^{(4)} = \mathcal{H}(\bar{\mathbf{U}}^{(3)}, \Delta t),$$

250

$$\bar{\mathbf{U}}^{n+1} = \frac{1}{3}\bar{\mathbf{U}}^n + \frac{2}{3}\bar{\mathbf{U}}^{(4)}.$$

251

252 **Stability Constraint**

253 A necessary condition for stability of all these explicit methods based on the finite
 254 volume method is the Courant-Friedrichs-Lowy condition (Courant et al. 1928) which
 255 states that

256
$$\Delta t < \frac{\Delta x}{2 \max(|\lambda_i|)} \forall i \quad (8)$$

 257

258 where λ_i is the i th eigenvalue of the Jacobian of the flux vector.

259 **NUMERICAL SIMULATIONS**

260 The discussed methods will now be used to solve three different problems; (i) the
 261 soliton which is an analytic solution of the Serre equations; (ii) one of the experiments
 262 conducted by Hammack and Segur (1978) and (iii) a dam-break problem from El et al.
 263 (2006) and Le Métayer et al. (2010). The first two will be used to validate the models,
 264 with the soliton used to establish the order of convergence of the models. The second
 265 problem is used to validate the models using experimental data which contains flows with
 266 steep gradients. Lastly the dam-break problem will be used to compare the results of these
 267 methods with those of El et al. (2006) and Le Métayer et al. (2010). The aim of which is to
 268 verify the claim of the latter that a first-order method for the Serre equations is sufficiently
 269 accurate to capture the important behaviour of the dam-break problem.

270 **Soliton**

271 Currently cnoidal waves are the only family of analytic solutions to the Serre equa-
 272 tions (Carter and Cienfuegos 2011). Solitons are a particular instance of cnoidal waves
 273 that travel without deformation and have been used to verify the convergence rates of the
 274 proposed methods in this paper.

275 For the Serre equations the solitons have the following form

276
$$h(x, t) = a_0 + a_1 \operatorname{sech}^2(\kappa(x - ct)), \quad (9a)$$

277
$$u(x, t) = c \left(1 - \frac{a_0}{h(x, t)} \right), \quad (9b)$$

278 and
 279
$$\kappa = \frac{\sqrt{3a_1}}{2a_0 \sqrt{a_0 + a_1}} \quad (9c)$$

280
$$c = \sqrt{g(a_0 + a_1)} \quad (9d)$$

279 where a_0 and a_1 are input parameters that determine the depth of the quiescent water and
 280 the maximum height of the soliton above that respectively. In the simulation $a_0 = 10\text{m}$,
 281 $a_1 = 1\text{m}$ for $x \in [-500\text{m}, 1500\text{m}]$ and $t \in [0\text{s}, 100\text{s}]$. With $\Delta t = 0.01\Delta x$ which satisfies (8)
 282 and $\theta = 1.2$ for the second-order reconstruction. The example results for $\Delta x = 100/2^6\text{m}$
 283 can be seen in Figure 2.

284 Figure 2 demonstrates the superiority of the second- and third-order methods compared
 285 to the first-order method. With the first-order method there is significant attenuation of the
 286 wave due to its diffusive behaviour which creates a wider wave profile and some smaller
 287 trailing waves. However, the first-order method does produce the correct speed of the
 288 wave with a small phase error. While the second- and third-order methods demonstrate no
 289 noticeable deformation resolving the soliton solution well on a relatively coarse grid with
 290 less than 500 cells defining the actual wave.

291 The relative error as measured by the L_1 -norm of the method can be seen in Figure 3.
 292 For a vector \mathbf{q} and an approximation to it \mathbf{q}^* the relative error as measured by the L_1 -norm
 293 is

$$294 \quad L_1(\mathbf{q}, \mathbf{q}^*) = \frac{\sum_{i=1}^m |q_i - q_i^*|}{\sum_{i=1}^m |q_i|}.$$

$$295$$

296 Figure 3 demonstrates that the methods all have the correct order of convergence in
 297 both time and space. However, this order of convergence is not uniform over all Δx . When
 298 Δx is large the actual problem is not discretised well since the cells are too large to ade-
 299 quately resolve the problem; this causes the observed suboptimal rate of convergence in
 300 Figure 3. When Δx is sufficiently small the numerical errors become small enough that
 301 floating point errors are significant and this can also lead to suboptimal rates of conver-
 302 gence as can be seen for the third-order method in Figure 3(c). Therefore, the order of
 303 convergence for all methods is confirmed.

304 Figure 3 also demonstrates the superiority of the second- and third-order methods over
 305 the first-order method for accuracy. The third-order method is also better than the second-
 306 order method in this respect although this difference is less pronounced. These differences
 307 have a significant impact on run-time if one wishes to run a simulation up to a desired
 308 accuracy. A comparison of the methods for such a problem is presented in Table 1.

309 The first-order method has a runtime two orders of magnitude greater than the second-
 310 and third-order methods. Which is sensible because the difference in Δx is also two orders
 311 of magnitude. This is computationally restrictive as running practical problems to a rea-
 312 sonable accuracy can have run times in the order of days. In particular this demonstrates
 313 that although the second- and third- order methods are more computationally complex than
 314 the first-order method this complexity is justified if one wants to numerically solve the
 315 Serre equations up to some desired accuracy. This also demonstrates that a second-order

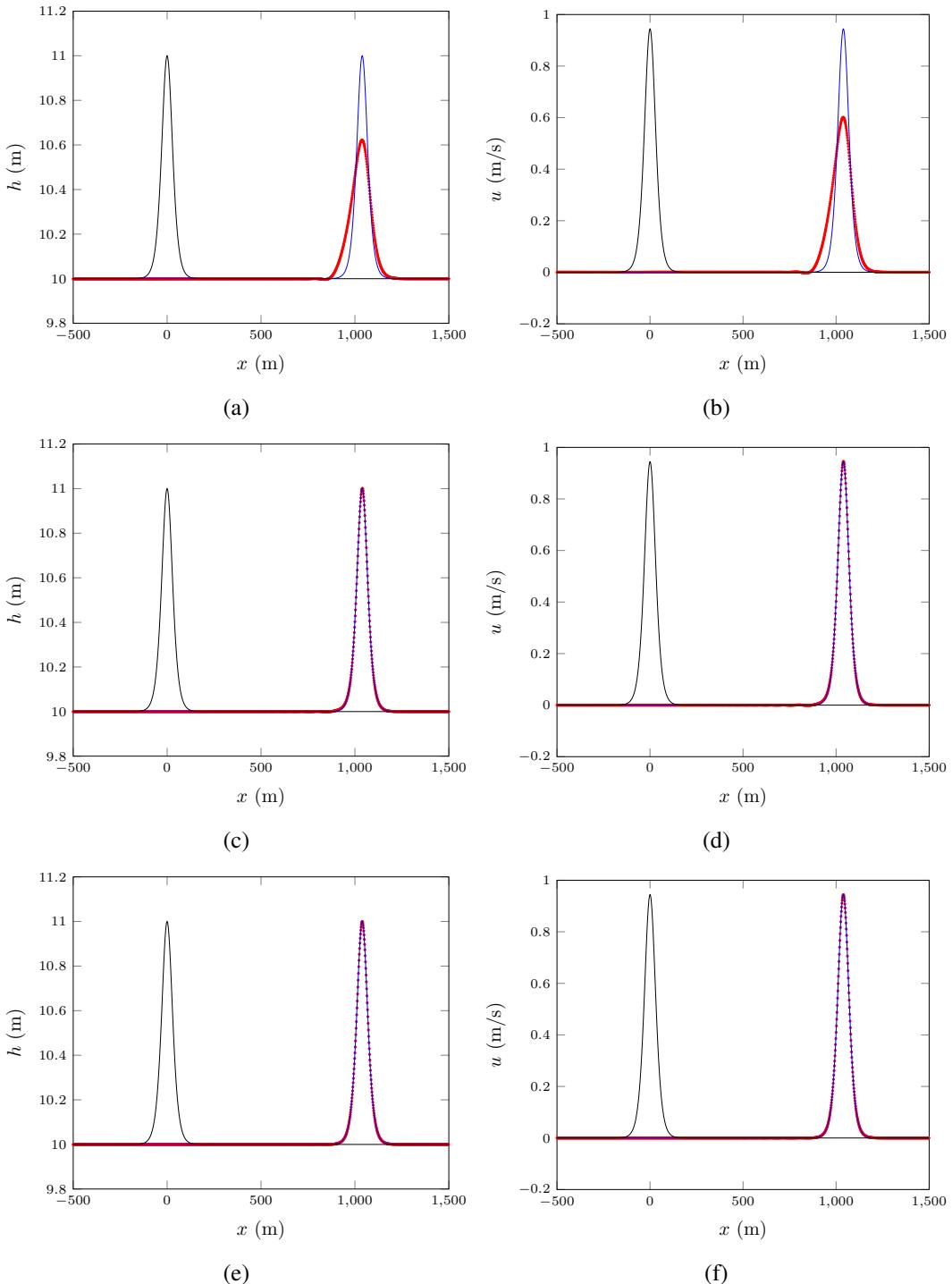


FIG. 2: The first-, second- and third-order simulation of a soliton with $\Delta x = 100/2^6$ m (\circ) plotted against the analytic solution of (9) (—) with black for $t = 0$ s and blue for $t = 100$ s.

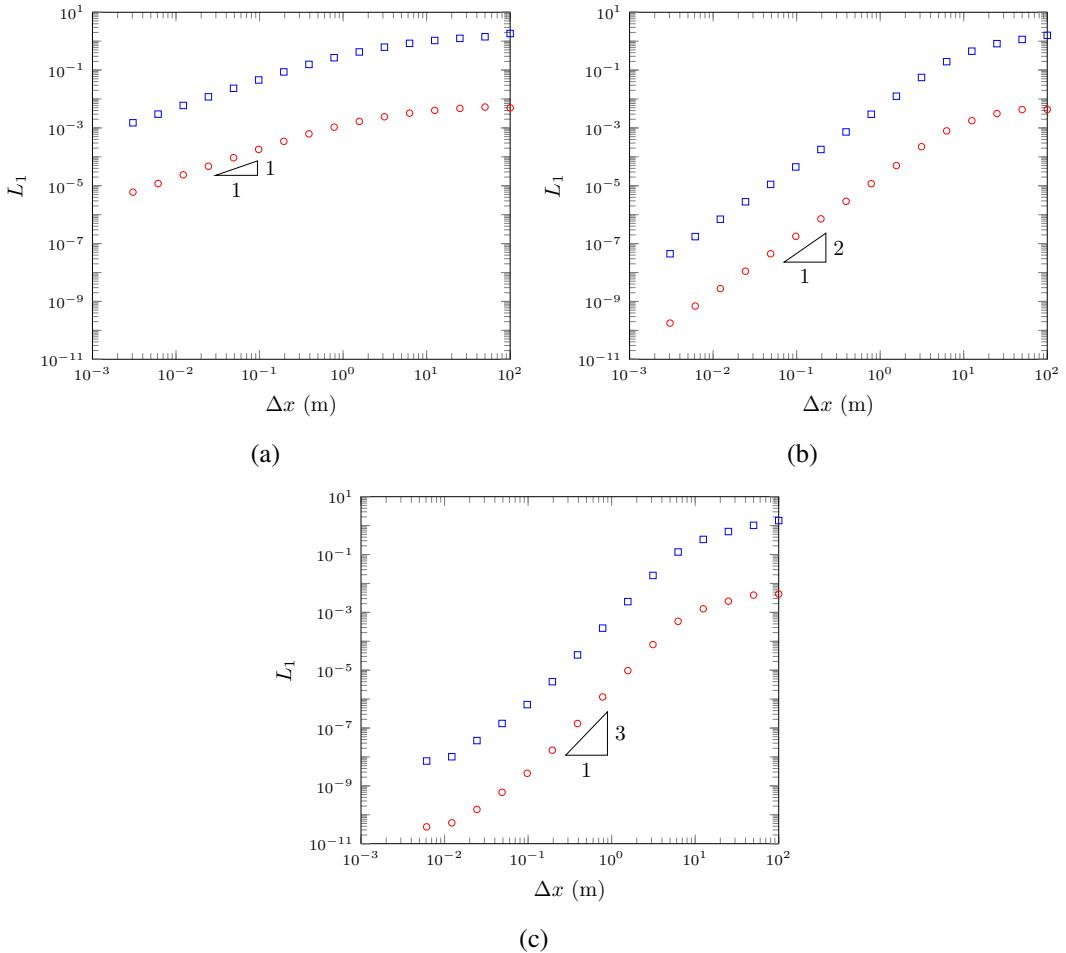


FIG. 3: Convergence of relative error using L_1 norm for analytic soliton solution for both h (\circ) and u (\diamond) for the; (a) first-, (b) second- and (c) third-order methods.

method is adequate compared to the third-order method for solving the Serre equations up to some desired accuracy.

Figure 3(b) and Figure 3(c) demonstrate that the second- and third-order methods both have similar errors and Figure 2 shows that these methods resolve the problem well. Therefore, the extra effort in running a third-order method compared to a second-order method is not justified in this case. While the effort required to go from a first-order method to a second-order method is justified since attaining a similar accuracy between them requires a restrictively small Δx for a first-order method.

Order	Δx (m)	L_1 relative error for h	L_1 relative error for u	Run Time (s)
First	$100/2^{11}$	$9.32715378271 \times 10^{-5}$	$2.34320597345 \times 10^{-2}$	481.216886997
Second	$100/2^6$	$5.00595847273 \times 10^{-5}$	$1.25175521776 \times 10^{-2}$	2.19665694237
Third	$100/2^5$	$7.681432554 \times 10^{-5}$	$1.89115440455 \times 10^{-2}$	1.3766579628

TABLE 1: Comparison of run times for different order methods to get similar relative error as measured by the L_1 norm for h

324 Segur Laboratory Experiment

325 Hammack and Segur (1978) conducted an experiment that produced rectangular waves
 326 with the stroke of a 0.61m long piston flush with the wall of a wave tank 31.6m in length.
 327 The water height was recorded at 0m, 5m, 10m, 15m and 20m from the edge of the piston
 328 furthest from the wall over time. The quiescent water height h_1 was 0.1m while the stroke
 329 of the piston caused a depression of water which was $h_0 = 0.095$ m deep. To run this as
 330 a numerical simulation the reflected problem was used. Thus the initial conditions were
 331 reflected around the origin and $h_1 - h_0$ was doubled by setting $h_0 = 0.09$ m. The domain
 332 was chosen to be from -60m to 60m and the simulation was run for 50s with $\Delta x = 0.01$ m,
 333 $\lambda = 0.2/\sqrt{gh_1}$ m/s and $\theta = 1.2$. The results of this simulation are displayed in Figures 4 -
 334 6.

335 In this experiment for the positive side of the axis the initial depression causes a right
 336 going rarefaction fan and a left going shock. The shocks from both sides then reflect at
 337 the origin and so the shock and the rarefaction fan will travel in the same direction. The
 338 leading wave in all the related figures is the rarefaction fan while the trailing dispersive
 339 waves are the result of the reflected shock.

340 From all the related figures it can be seen that all models show good agreement be-
 341 tween the arrival of the first wave and the period of all the waves. While Figure 4 shows
 342 the first-order model is too diffusive and thus under estimates the heights of the dispersive
 343 waves. Whereas the second- and third-order methods over estimate them. This overes-
 344 timation can be explained by the Serre equations not taking into account viscous effects
 345 that may diffuse the dispersive waves and so the results could be considered as an upper
 346 bound on the wave heights for fluids with viscosity. Although even without these effects
 347 these numerical methods show good agreement with the experimental data thus validat-
 348 ing them to provide a reasonable representation of rapidly-varying flows. Additionally,
 349 it demonstrates that the oscillations observed by the produced numerical solutions of the
 350 Serre equations around steep gradients are physical and not numerical. In these simula-
 351 tions the numerical oscillations that the second-order method should produce (Zoppou and
 352 Roberts 1996) do not have a significant influence on the physical oscillations.

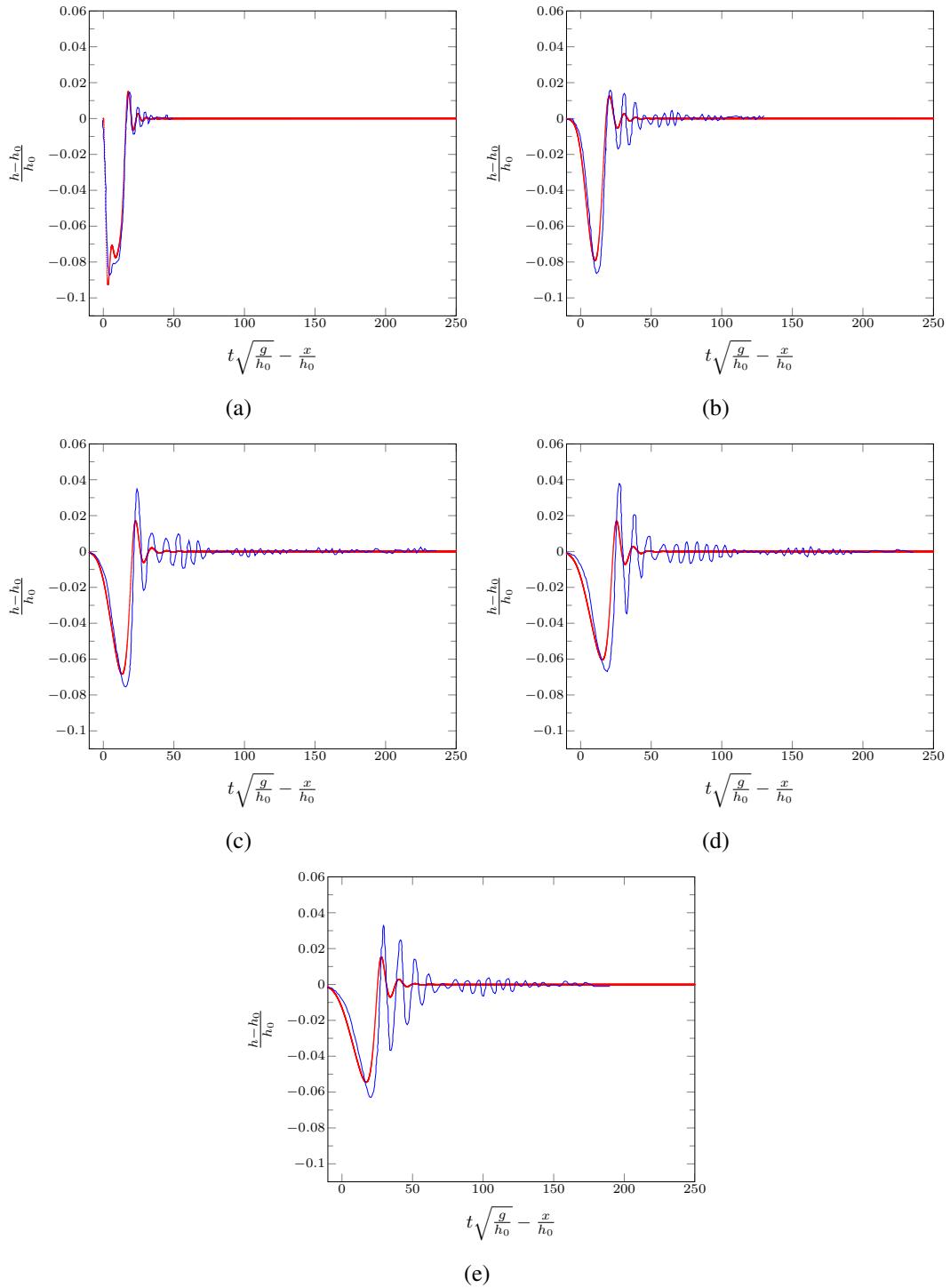


FIG. 4: Simulation of the rectangular wave experiment using the for first-order method at x/h_0 : (a) 0, (b) 50, (c) 100, (d) 150 and (e) 200.

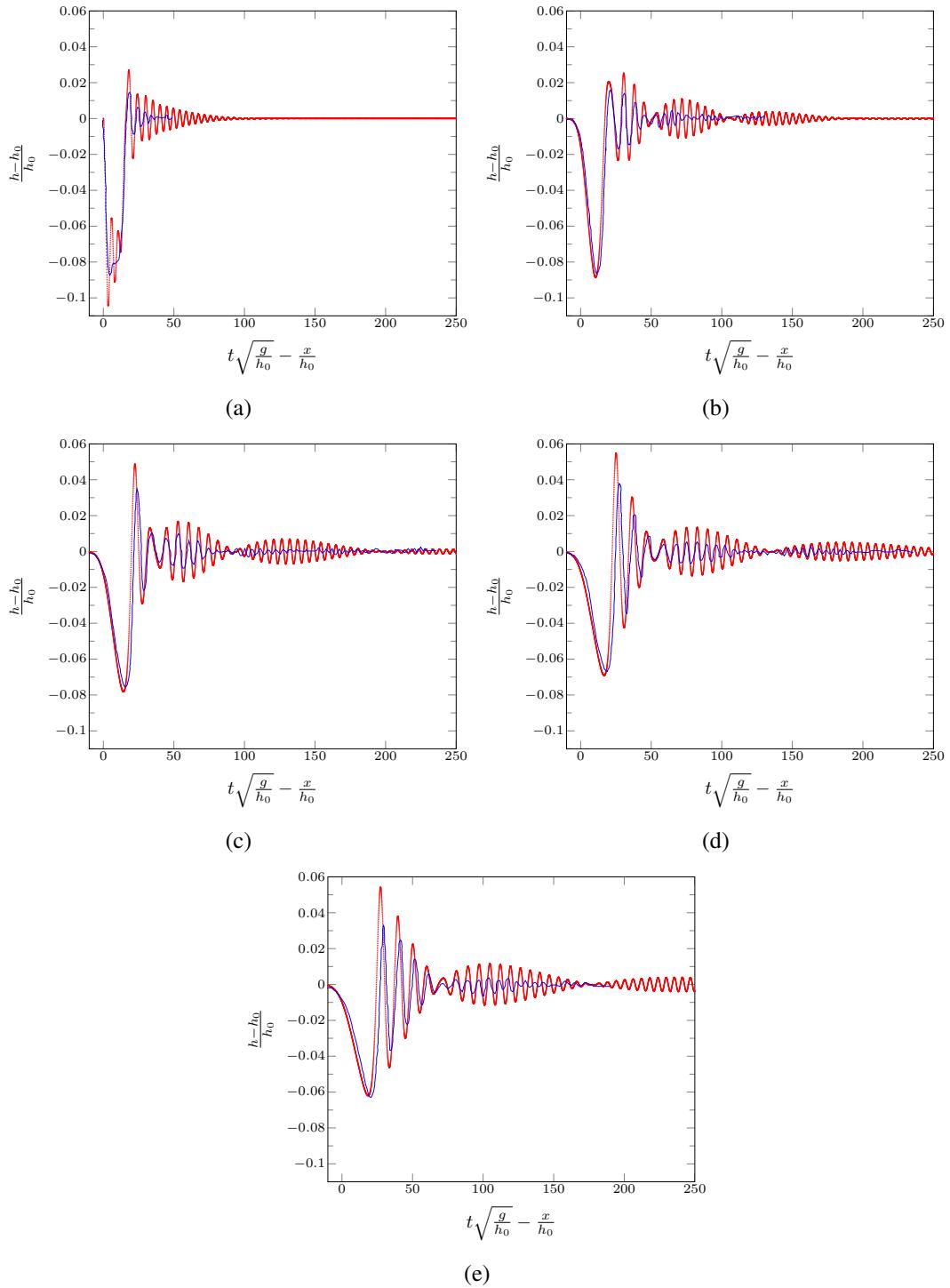


FIG. 5: Simulation of the rectangular wave experiment using the for second-order method at x/h_0 : (a) 0, (b) 50, (c) 100, (d) 150 and (e) 200.

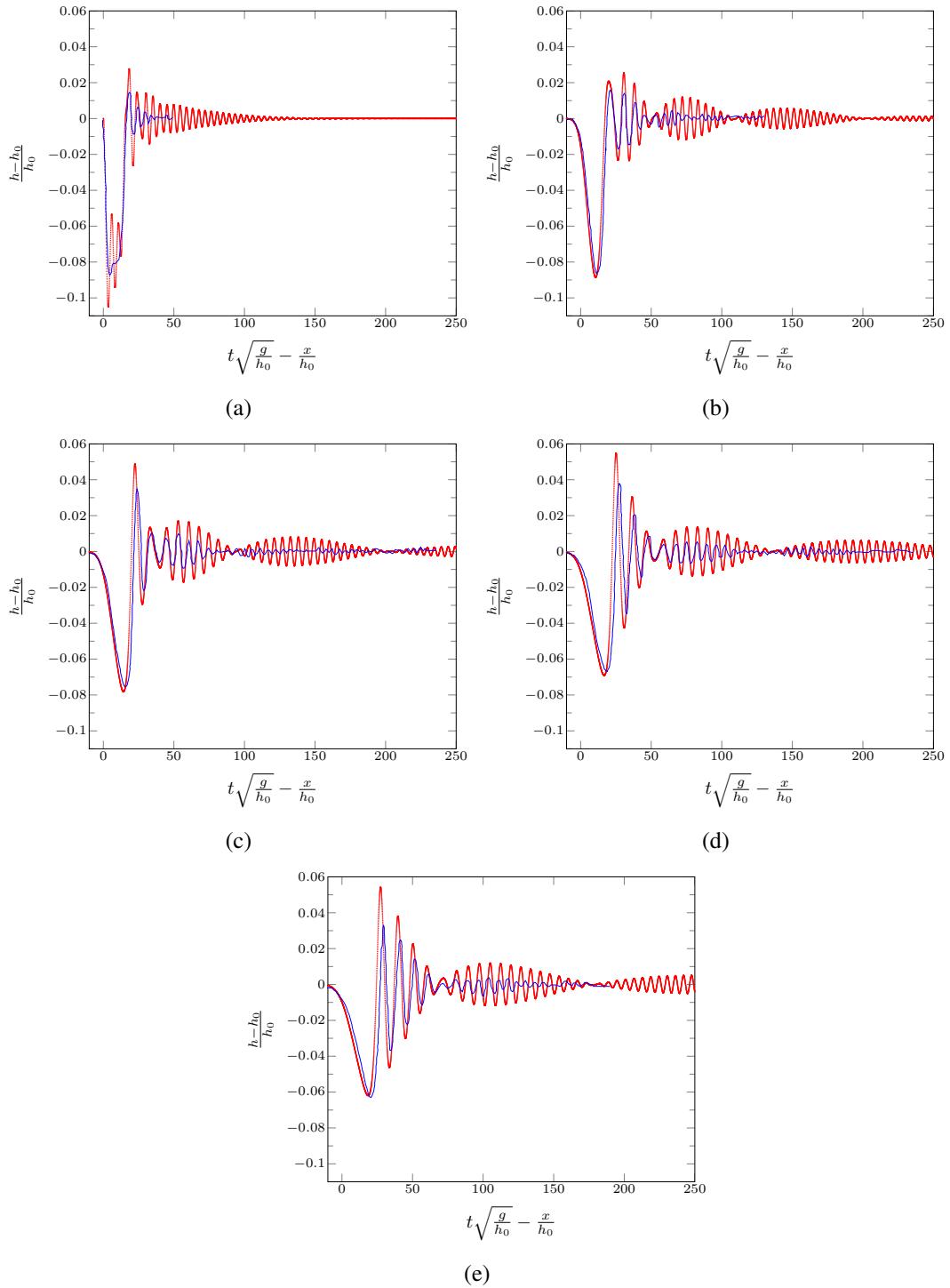


FIG. 6: Simulation of the rectangular wave experiment using the for third-order method at x/h_0 : (a) 0, (b) 50, (c) 100, (d) 150 and (e) 200.

353 **Dam-Break**

354 The dam-break problem of interest can be defined as such

$$355 \quad h(x, 0) = \begin{cases} 1.8 & x < 500 \\ 1.0 & x \geq 500 \end{cases},$$

$$356 \quad u(x, 0) = 0.0m/s.$$

357 With $x \in [0m, 1000m]$ for $t \in [0s, 30s]$. Where $\lambda = 0.01m/s$ and $\theta = 1.2$. This corresponds
 358 to sub-critical flow and was a situation demonstrated in El et al. (2006) and Le Métayer
 359 et al. (2010). An example was plotted for $\Delta x = 100/2^{10}m$ for all the methods and for
 360 $\Delta x = 100/2^{15}m$ for the first-order method in Figure 8. To determine if the oscillations that
 361 occur in the solution indeed converge to some limit as $\Delta x \rightarrow 0$ multiple Δx values were run
 362 and then the amount of variation in the solution measured. This measured how oscillatory
 363 the solution was and was used to determine the growth of the oscillations. A common way
 364 to measure this is the total variation TV (LeVeque 2002) which for \mathbf{q} is given by

$$365 \quad TV(\mathbf{q}) = \sum_{\forall i>1} |q_i - q_{i-1}|.$$

366 If the solution does indeed converge then the TV must at some point plateau, bounding
 367 the oscillations. This was indeed the findings of the experiments as can be seen by Fig-
 368 ure 7. The TV increases as Δx decreased because the models resolved more dispersive
 369 waves. As Δx decreased further the TV plateaued and so the size and number of oscilla-
 370 tions was bounded. Therefore, the method has not become unstable which supports the
 371 argument that the numerical methods do not introduce non-physical oscillations in the so-
 372 lution. Under this measure the second-order method converges rapidly to the solution of
 373 the third-order method.

374 These solutions compare very well to the findings in El et al. (2006) with both the
 375 second- and third-order methods resolving the oscillations around the “contact discontinu-
 376 ity”(El et al. 2006) between the rarefaction fan and the shock. In Le Métayer et al. (2010)
 377 it was reported that for their first-order method such oscillatory behaviour was not seen.
 378 However, for the first-order method proposed in this paper when $\Delta x = 100/2^{15}$ it was
 379 resolved as in Figure 8(d). This validates the findings in El et al. (2006). Interestingly this
 380 is a much higher resolution than one needs to represent the waves themselves. It appears
 381 that the behaviour of the dispersive waves around a steep gradient is sensitive to diffusion
 382 so that even though the first-order method in Figure 8(a) heuristically had enough cells to
 383 resolve most of the wave train, due to strong diffusion hardly any of the wave train was
 384 resolved.

385 There is a good agreement between the second- and third-order simulations of the dam-
 386 break problem as can be seen in Figures 8(b) and 8(c). Although more oscillations are

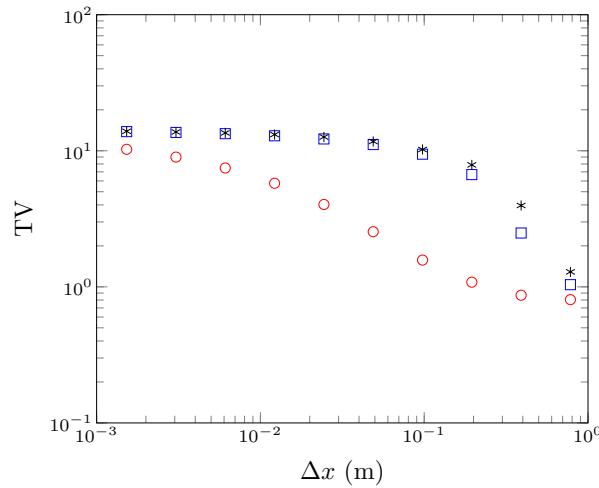


FIG. 7: The change in total variation (TV) over Δx for; (○) first-, (□) second-, and (*) third-order methods.

388 resolved by the third-order method over the second-order method, there is no significant
 389 change in the resolved behaviour of this problem between the two methods. As noted in
 390 the introduction second-order errors are dissipative; since the diffusive third-order method
 391 resolved the same oscillations it was demonstrated that none of the dissipative errors sig-
 392 nificantly polluted the wave train and so the second-order method is capable of resolving
 393 the problem well.

394 CONCLUSIONS

395 First-, second- and third-order hybrid finite difference-volume methods were devel-
 396 oped to solve the Serre equations written in conservative law form. The methods were then
 397 tested and validated. Firstly the order of the methods were all verified, secondly the meth-
 398 ods steep gradient handling capability was validated by comparison with experimental
 399 data. Thirdly the behaviour of the solutions matched previous findings in El et al. (2006).
 400 Thus it can be concluded that these methods are all valid and they properly handle steep
 401 gradients. It was also demonstrated that for these equations although second-order is not
 402 as accurate as third-order it still provides a satisfactory method for reasonable Δx unlike
 403 the first-order method which due to strong diffusion requires computationally restrictive
 404 Δx to produce satisfactory accuracy. So practical problems require at least a second-order
 405 method to solve the Serre equations.

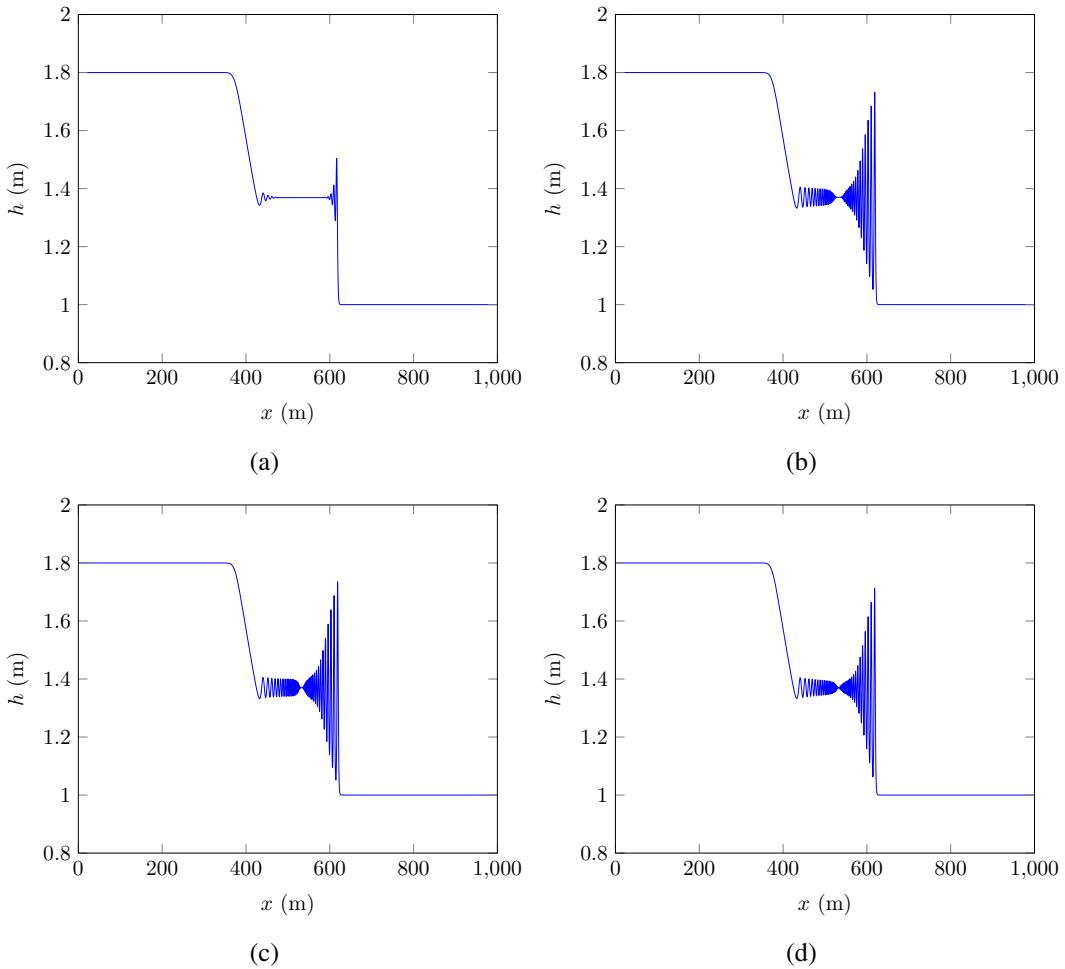


FIG. 8: Solution of the dam-break problem using the (a) first-, (b) second- and (c) third-order method with $\Delta x = 100/2^{10}\text{m}$. As well as a (d) first-order method with $\Delta x = 100/2^{15}\text{m}$.

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