

ASPECTS OF SUPERRADIANT SCATTERING OFF KERR BLACK HOLES

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- What is superradiance ?
- Kerr black hole and the Penrose process
- Newman-Penrose formalism and the Teukolsky equation
- Analytical approximated solutions
- Numerical methods
- Scattering of plane waves
- Conclusions and future work

Default

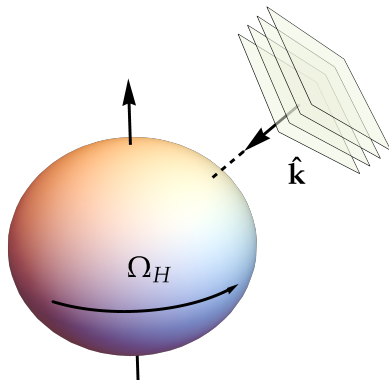
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Alert

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Example

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Boyer-Lindquist coordinates

$$g = \left(1 - \frac{2Mr}{\rho^2}\right) dt^2 + 2a \sin^2 \theta \frac{(r^2 + a^2 - \Delta)}{\rho^2} dt d\varphi$$

$$- \frac{(r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta}{\rho^2} \sin^2 \theta d\varphi^2 - \frac{\rho^2}{\Delta} dr^2 - \rho^2 d\theta^2$$

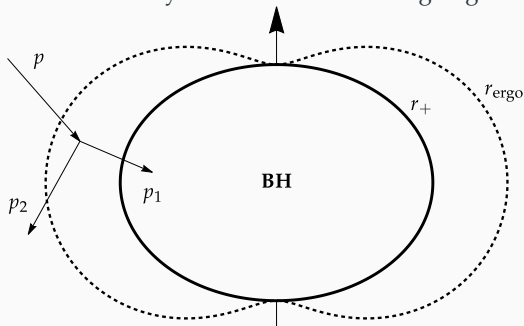
$$[\Delta = r^2 - 2Mr + a^2, \rho^2 = r^2 + a^2 \cos^2 \theta]$$

- BH angular momentum $J = aM$ [$a = 0 \Rightarrow$ Schwarzschild]
- Killing vectors $k = \partial_t$ (stationary) and $m = \partial_\varphi$ (axisymmetric)
- Horizons at $g^{rr} = 0 \Rightarrow \Delta = 0 \Rightarrow r = r_\pm \equiv M \pm \sqrt{M^2 - a^2}$
Cosmic censorship conjecture $\Rightarrow |a| \leq M$
- Infinite redshift boundary $g_{tt} = 0 \Rightarrow$ **Ergoregion** ($g_{tt} < 0$)

$$r_+ < r < r_{\text{ergo}}(\theta) \equiv M + \sqrt{M^2 - a^2 \cos^2 \theta}$$

PENROSE PROCESS

- Particle decays into two inside the ergoregion: $p = p_1 + p_2$



Killing horizon

Hypersurface $r = r_+$
with normal null vector

$$\xi = k + \Omega_H m$$

Event horizon angular
momentum:

$$\Omega_H = \frac{a}{2Mr_+}$$

- We may have $E_1 = k \cdot p_1 < 0 \Rightarrow E_2 = E + |E_1| > E$
- Local energy condition $\xi \cdot p_1 > 0$ at $r = r_+ \Rightarrow E_1 - \Omega_H L_1 > 0$

$$\psi \sim e^{-i\omega t + im\varphi} \Rightarrow \frac{\delta J}{\delta M} = \frac{\hbar m}{\hbar \omega} \Rightarrow \delta M < 0 \wedge \boxed{\omega(\omega - \Omega_H m) < 0}$$

Energy extraction !

- Electromagnetic perturbations ($A^\mu \ll 1$)

$$R_{\mu\nu} - \frac{R}{2}g_{\mu\nu} = 8\pi T_{\mu\nu} = 0 + \mathcal{O}(A^2) \Rightarrow \text{fixed background (Kerr)}$$

- Tetrad frame of complex null vectors $(e_1, e_2, e_3, e_4) = (\mathbf{l}, \mathbf{n}, \mathbf{m}, \bar{\mathbf{m}} \equiv \mathbf{m}^*)$
with $\mathbf{l} \cdot \mathbf{m} = \mathbf{n} \cdot \bar{\mathbf{m}} = 0$, $\mathbf{l} \cdot \mathbf{n} = 1$, $\mathbf{m} \cdot \bar{\mathbf{m}} = -1$

$\underbrace{\mathbf{E}, \mathbf{B}}_{\text{6 real components}}$

\longrightarrow

$$\begin{aligned}\phi_0 &= F_{\mu\nu} l^\mu m^\nu \\ \phi_1 &= \frac{1}{2} F_{\mu\nu} (l^\mu n^\nu - m^\mu \bar{m}^\nu) \\ \phi_2 &= F_{\mu\nu} \bar{m}^\mu n^\nu\end{aligned}$$

- Equations written using tetrad derivatives ($\mathbb{D}, \Delta, \delta, \bar{\delta}$) and spin connection $\gamma_{cab} = (e_c)^\mu (e_b)^\nu \nabla_\nu (e_a)_\mu \Rightarrow \lambda, \pi, \tau, \varrho, \varepsilon, \sigma, \kappa, \gamma, \mu, \nu, \alpha, \beta$

$$\begin{aligned}\nabla_\mu F^{\mu\nu} &= 0 \\ \nabla_{[\mu} F_{\nu\rho]} &= 0\end{aligned} \longrightarrow \begin{cases} \mathbb{D}\phi_2 - \bar{\delta}\phi_1 = -\lambda\phi_0 + 2\pi\phi_1 + (\varrho - 2\varepsilon)\phi_2 \\ \Delta\phi_1 - \delta\phi_2 = \nu\phi_0 - 2\mu\phi_1 + (2\beta - \tau)\phi_2 \\ \mathbb{D}\phi_1 - \bar{\delta}\phi_0 = (\pi - 2\alpha)\phi_0 + 2\varrho\phi_1 - \kappa\phi_2 \\ \Delta\phi_0 - \delta\phi_1 = (2\gamma - \mu)\phi_0 - 2\tau\phi_1 + \sigma\phi_2\end{cases}$$

Kinnersley tetrad

$$\mathbf{l} = \frac{1}{\Delta} (r^2 + a^2, \Delta, 0, a)$$

$$\mathbf{n} = \frac{1}{2\rho^2} (r^2 + a^2, -\Delta, 0, a)$$

$$\mathbf{m} = \frac{1}{\sqrt{2}\bar{\rho}} (ia \sin \theta, 0, 1, i \csc \theta)$$

$$[\bar{\rho} = r + ia \cos \theta]$$

- \mathbf{l}, \mathbf{n} are the doubly degenerate principal null of directions of the Weyl tensor

$$\mathbf{l} \sim (e_t + e_r) \quad (\text{ingoing})$$

$$\mathbf{n} \sim \frac{1}{2}(e_t - e_r) \quad (\text{outgoing})$$

- \mathbf{l}, \mathbf{n} are geodesic ($\kappa = \lambda = 0$) and shear-free ($\nu = \sigma = 0$)
 \Rightarrow equations decouple

- Reformulation of the fields: $\Phi_n = (\sqrt{2}\bar{\rho})^n \phi_n \quad (n = 0, 1, 2)$
- Spacetime symmetry $\Rightarrow \Phi_n \sim e^{-i\omega t + im\phi} \Rightarrow \partial_t \rightarrow -i\omega, \partial_\phi \rightarrow im$
- Rewrite operators: $\mathbb{D} = \mathcal{D}_0, \mathbb{\Delta} = -\frac{\Delta}{2\rho^2} \mathcal{D}_0^\dagger$

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or show `\textbf{bold}` results.

becomes

The theme provides sensible defaults to *emphasize* text, **accent** parts or show **bold** results.

Coordinate frame

- Vector (coordinate) basis

$$e_\mu \equiv \frac{\partial}{\partial x^\mu} \quad (\mu = t, r, \theta, \varphi)$$

- Spacetime metric

$$g_{\mu\nu}$$

- Tensorial components

$$A_\mu, F_{\alpha\beta}$$

- Covariant derivative

$$A_{\mu;\nu} = A_{\mu,\nu} - \Gamma_{\mu\nu}^\rho A_\rho$$

- Levi-Civita connection

$$\Gamma_{\mu\nu}^\rho = \frac{1}{2}g^{\rho\sigma}(g_{\sigma\mu,\nu} + g_{\sigma\nu,\mu} - g_{\mu\nu,\sigma})$$

Tetrad frame

- Tetrad (non-coordinate) basis

$$e_a \equiv (e_a)^\mu \frac{\partial}{\partial x^\mu} \quad (a = 1, 2, 3, 4)$$

- Raising/lowering tetrad indices

$$\eta_{ab} = e_a \cdot e_b = (e_a)^\mu (e_b)^\nu g_{\mu\nu}$$

- Tetrad (NP) scalars

$$A_a = (e_a)^\mu A_\mu, F_{ab} = (e_a)^\alpha (e_b)^\beta F_{\alpha\beta}$$

- Intrinsic derivative

$$A_{a|b} \equiv (e_a)^\mu (e_b)^\nu A_{\mu;\nu} = A_{a,b} - \gamma_{cab} A^c$$

- Spin connection

$$\gamma_{cab} = (e_c)^\mu (e_a)_\mu{}_\nu A_{\mu;\nu} = (e_b)^\nu$$

Kinnersley tetrad

$$\begin{aligned}\mathfrak{l} &= \frac{1}{\Delta} \left(r^2 + a^2, \Delta, 0, a \right) \\ \mathfrak{n} &= \frac{1}{2\rho^2} \left(r^2 + a^2, -\Delta, 0, a \right) \\ \mathfrak{m} &= \frac{1}{\sqrt{2}\bar{\rho}} \left(ia \sin \theta, 0, 1, i \csc \theta \right)\end{aligned}$$

$$\begin{aligned}\pm_1 Z_{\ell m} &\simeq -4\bar{\omega}(\bar{\omega} - m\bar{\Omega}_H) (2 - \tau)(2\bar{\omega}\tau)^{2\ell} \\ &\times \left[\frac{(\ell - 1)!(\ell + 1)!}{(2\ell)!(2\ell + 1)!} \right]^2 \prod_{n=1}^{\ell} \left(n^2 + \frac{4\omega^2}{\tau^2} \right)\end{aligned}$$

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becomes

The theme provides sensible defaults to *emphasize* text, **accent** parts or show **bold** results.

- Regular
- *Italic*
- SMALLCAPS
- **Bold**
- ***Bold Italic***
- **BOLD SMALLCAPS**
- Monospace
- *Monospace Italic*
- Monospace Bold
- *Monospace Bold Italic*

Items

- Milk
- Eggs
- Potatos

Enumerations

1. First,
2. Second and
3. Last.

Descriptions

PowerPoint Meeh.

Beamer Yeeeha.

- This is important

- This is important
- Now this

- This is important
- Now this
- And now this

- This is really important
- Now this
- And now this

Three different block environments are pre-defined and may be styled with an optional background color.

Default

Block content.

Alert

Block content.

Example

Block content.

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Example

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$$e = \lim_{n \rightarrow \infty} \left(1 + \frac{1}{n}\right)^n$$

Veni, Vidi, Vici

Metropolis defines a custom beamer template to add a text to the footer. It can be set via

```
\setbeamertemplate{frame footer}{My custom footer}
```

Conclusion

Get the source of this theme and the demo presentation from

`github.com/matze/mtheme`

The theme *itself* is licensed under a Creative Commons Attribution-ShareAlike 4.0 International License.

Questions?

LINE PLOTS

Sometimes, it is useful to add slides at the end of your presentation to refer to during audience questions.

The best way to do this is to include the `appendixnumberbeamer` package in your preamble and call `\appendix` before your backup slides.

Metropolis will automatically turn off slide numbering and progress bars for slides in the appendix [1].

References

- [1] J. G. Rosa, *Superradiance in the sky*, Phys. Rev. **D95**, 064017 (2017), arXiv:1612.01826 [gr-qc] .