

EGM0004

Sistemas Não Lineares

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Programa de Pós-Graduação em Engenharia Mecatrônica

24T12 (60h) (13:00-14:40h) – 22.08.2022 : 21.12.2022

Teoremas sobre estabilidade

Teorema 6: Teorema de La Salle (princípio da invariância)

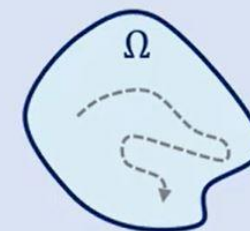
Pode ocorrer de $\dot{V} \leq 0$ e ainda assim concluir sobre estabilidade assintótica

Definições:

Positively
invariant set

A set Ω is said to be a **positively invariant set** with respect to $\dot{x} = f(x)$ if

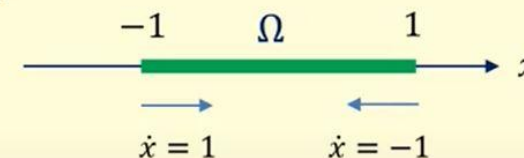
$$x(0) \in \Omega \Rightarrow x(t) \in \Omega, \quad \forall t \geq 0$$



Example 1.

The set $\Omega = \{x \in \mathbb{R} : |x| \leq 1\}$ is a positively invariant set with respect to $\dot{x} = -x$.

To check if Ω is positively invariant w.r.t the system dynamics, we **just check the boundaries**.

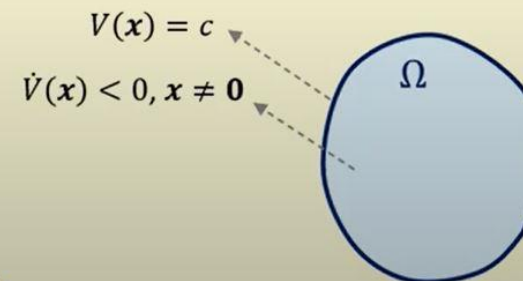


Example 2.

Let V be a positive definite,
and $\dot{V}(x)$ be negative definite for $\dot{x} = f(x)$.

Define Ω as $\Omega = \{x \in \mathbb{R}^n : V(x) \leq c\}$.

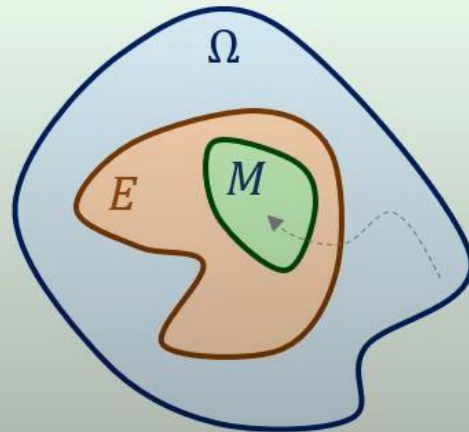
Then Ω is a positively invariant set w.r.t. $\dot{x} = f(x)$. We **check the boundaries**.



La Salle

- Let Ω be a compact set that is **positively invariant** with respect to $\dot{x} = f(x)$
- Let V be a **continuously differentiable** function on Ω such that $\dot{V}(x) \leq 0$ in Ω
- Let $E \subset \Omega$ be the set of all points in Ω such that $\dot{V}(x) = 0$
- Let M be the **largest positively invariant** set in E

Then every solution starting in Ω approaches M as $t \rightarrow \infty$



La Salle local e global

Teorema 6: Teorema de La Salle (princípio da invariância)

Se para um sistema $\dot{x} = f(x)$, $f(0) = 0$, existe uma função $V(x)$ tal que

a) $V(x) > 0$

b) $V(x)$ é continuamente diferenciável

c) $\dot{V}(x) \leq 0$

d) $\nexists x \neq 0$ tal que $\dot{V}(x) = 0, \forall t \geq 0$ Então é assintoticamente estável

Teorema 7: Teorema de La Salle global

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d) $\nexists x \neq 0$ tal que $\dot{V}(x) = 0, \forall t \geq 0$

d) $\nexists x \neq 0$ tal que $\dot{V}(x) = 0, \forall t \geq 0$

e) $V(x) \rightarrow \infty$ quando $\|x\| \rightarrow \infty$

Então é globalmente assintoticamente estável

La Salle local e global

Teorema 6: Teorema de La Salle (princípio da invariância)

Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -x_1 - x_2$$

$$V(x) = \frac{1}{2}x_1^2 + \frac{1}{2}x_2^2 > 0$$

Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -\phi(x_1) - \mu x_2, \mu > 0$$

$$V(x) = \frac{1}{2}x_2^2 + \int_0^{x_1} \phi(t)dt > 0$$

$\phi(x_1)x_1 > 0$ se $x_1 \neq 0$ Função setorial nos quadrantes ímpares

$$x_1 = 0 \implies \phi(x_1) = 0$$

La Salle local e global

Teorema 7: Teorema de La Salle global

Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -x_1 + x_3$$

$$\dot{x}_3 = -x_2 - x_3$$

$$V(x) = \frac{1}{2}x_1^2 + \frac{1}{2}x_2^2 + \frac{1}{2}x_3^2 > 0$$

La Salle

Example 3 (pendulum with friction): Show the asymptotic stability of the origin $x = (0,0)$ using LaSalle's Theorem for the general case where $k \neq 0$. For simplicity, let $\ell = g$ and $m = 1$.

State equations:

$$\dot{x}_1 = x_2$$

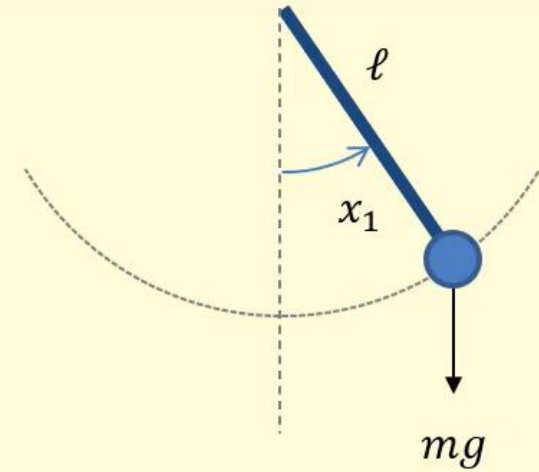
$$\dot{x}_2 = -\frac{g}{\ell} \sin x_1 - \frac{k}{m} x_2$$

Let $V(x) = 1 - \cos x_1 + \frac{1}{2}x_2^2$

$$V(x) = mgl(1 - \cos x_1) + \frac{1}{2}ml^2x_2^2$$

$$\dot{V}(x) = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = (\sin x_1)(x_2) + (x_2)(-\sin x_1 - kx_2) = -kx_2^2 \leq 0$$

So $\dot{V}(x)$ is negative semi-definite.



m : mass of the bob
 ℓ : length of the rod
 k : friction coefficient

La Salle

Example 3 (pendulum with friction): Show the asymptotic stability of the origin $x = (0,0)$ using LaSalle's Theorem for the general case where $k \neq 0$. For simplicity, let $\ell = g$ and $m = 1$.

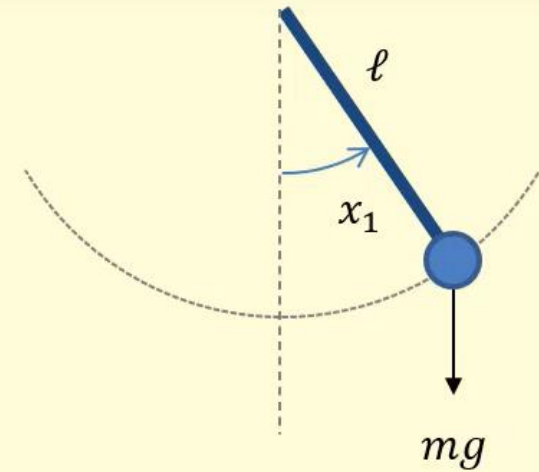
State equations:

$$\begin{aligned}\dot{x}_1 &= x_2 \\ \dot{x}_2 &= -\frac{g}{\ell} \sin x_1 - \frac{k}{m} x_2\end{aligned}$$

$$\text{Let } V(x) = 1 - \cos x_1 + \frac{1}{2} x_2^2 \quad \dot{V}(x) = \nabla V \cdot f(x)$$

$$\dot{V}(x) = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = (\sin x_1)(x_2) + (x_2)(-\sin x_1 - kx_2) = -kx_2^2 \leq 0$$

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The set Ω is a compact set that is **positively invariant** with respect to $\dot{x} = f(x)$.

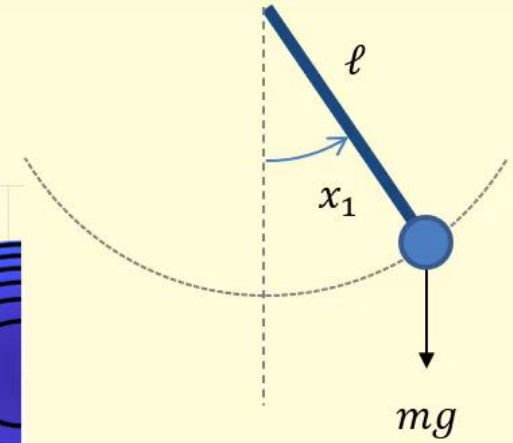
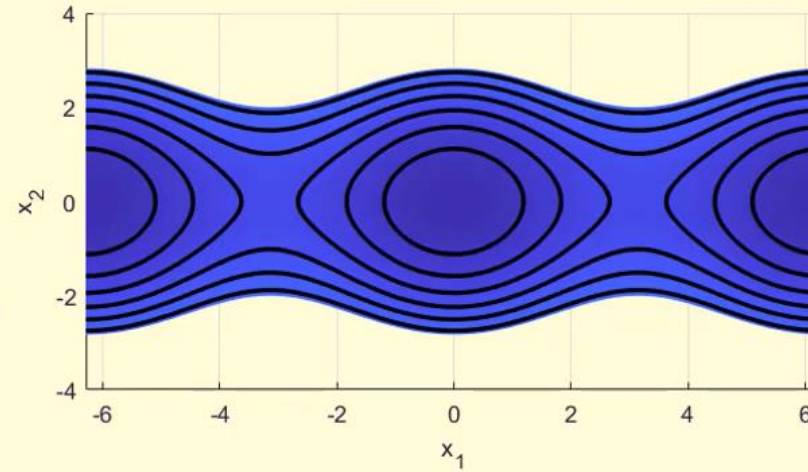
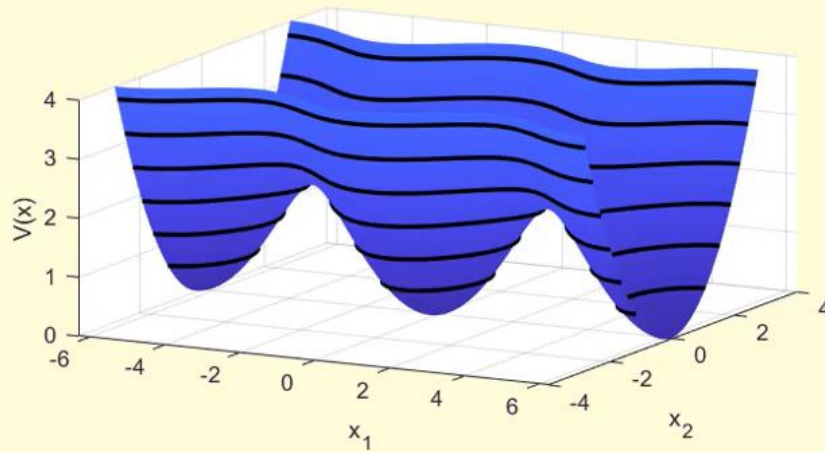
The set Ω is usually chosen as a level set of the V if V is a positive definite function.

$$\begin{aligned}\dot{x}_1 &= x_2 \\ \dot{x}_2 &= -\sin x_1 - kx_2\end{aligned}$$

$$V(\mathbf{x}) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$

$$\dot{V}(\mathbf{x}) = -kx_2^2 \leq 0$$

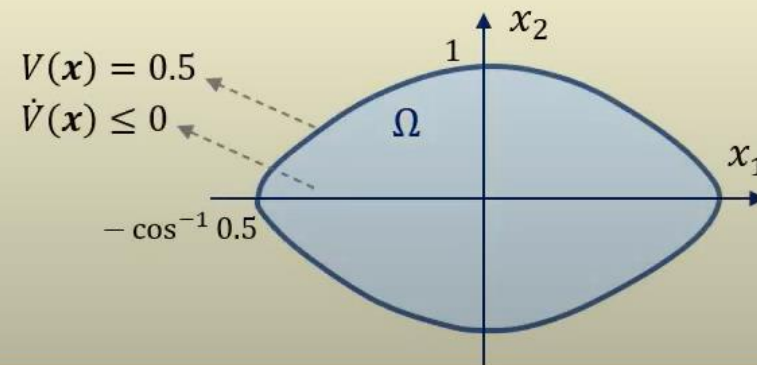
The level curves:



$$\text{ou.. } \Omega = [x_1 \quad x_2]^T : -\pi < x_1 < \pi, |x_2| < k$$

We choose the set Ω as:

$$\Omega = \{\mathbf{x} \in \mathbb{R}^2 : V(\mathbf{x}) \leq 0.5\}$$



$$\begin{aligned}\dot{x}_1 &= x_2 \\ \dot{x}_2 &= -\sin x_1 - kx_2\end{aligned}$$

$$V(\mathbf{x}) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$

$$\dot{V}(\mathbf{x}) = -kx_2^2 \leq 0$$

$$E := \{x : \dot{V} = 0\}$$

The set $E \subset \Omega$ is the set of all points in Ω such that $\dot{V}(\mathbf{x}) = 0$.

$$\dot{V}(\mathbf{x}) = 0 \Rightarrow -kx_2^2 = 0 \Rightarrow x_2 = 0$$

$$E = \{(x_1, x_2) \in \Omega : x_2 = 0\}$$

The set M is the *largest positively invariant set* in E .

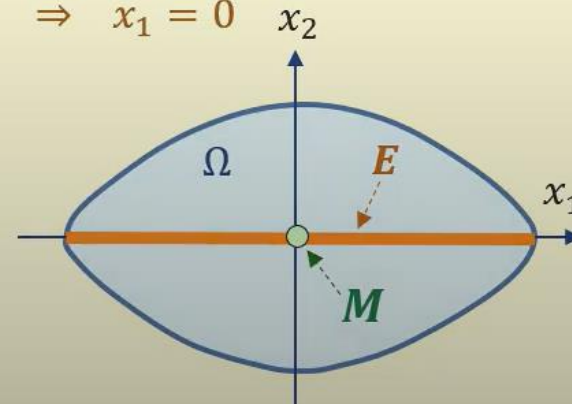
To find M , we let $x_2 = 0$ for all $t \geq 0$. So $\dot{x}_2(t) = 0$.

Using the system's dynamical model we have:

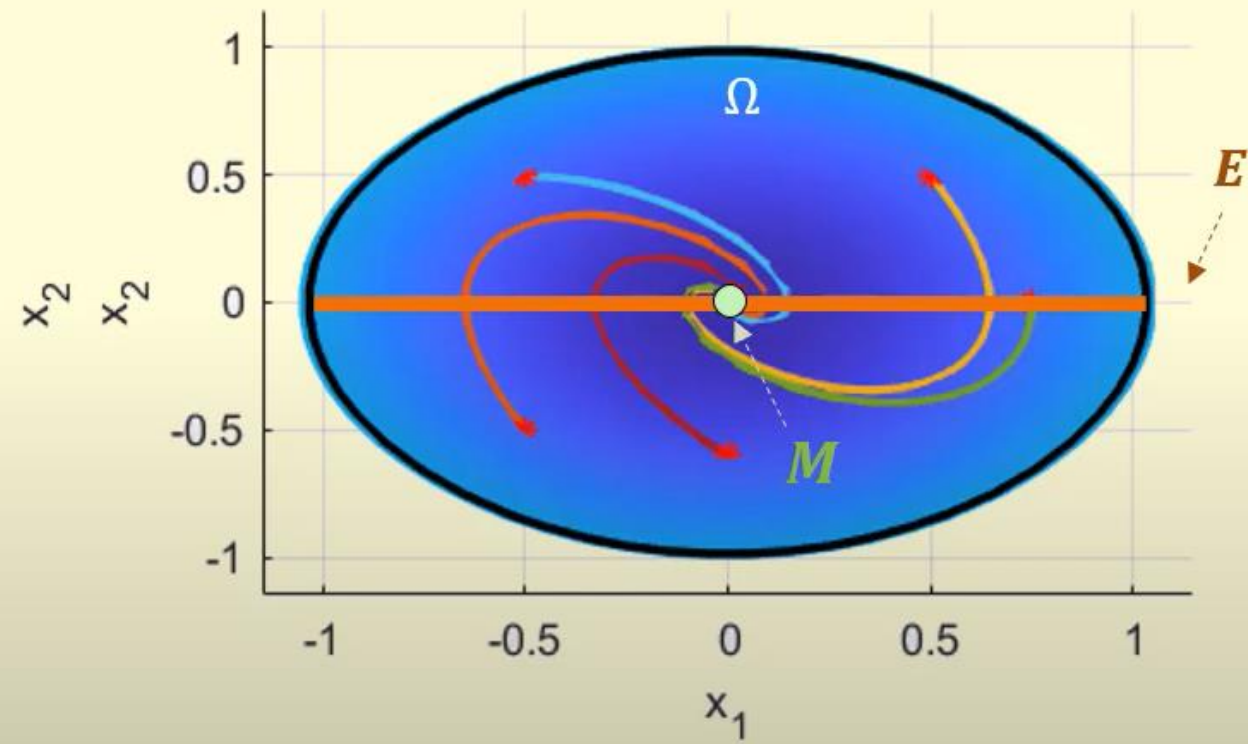
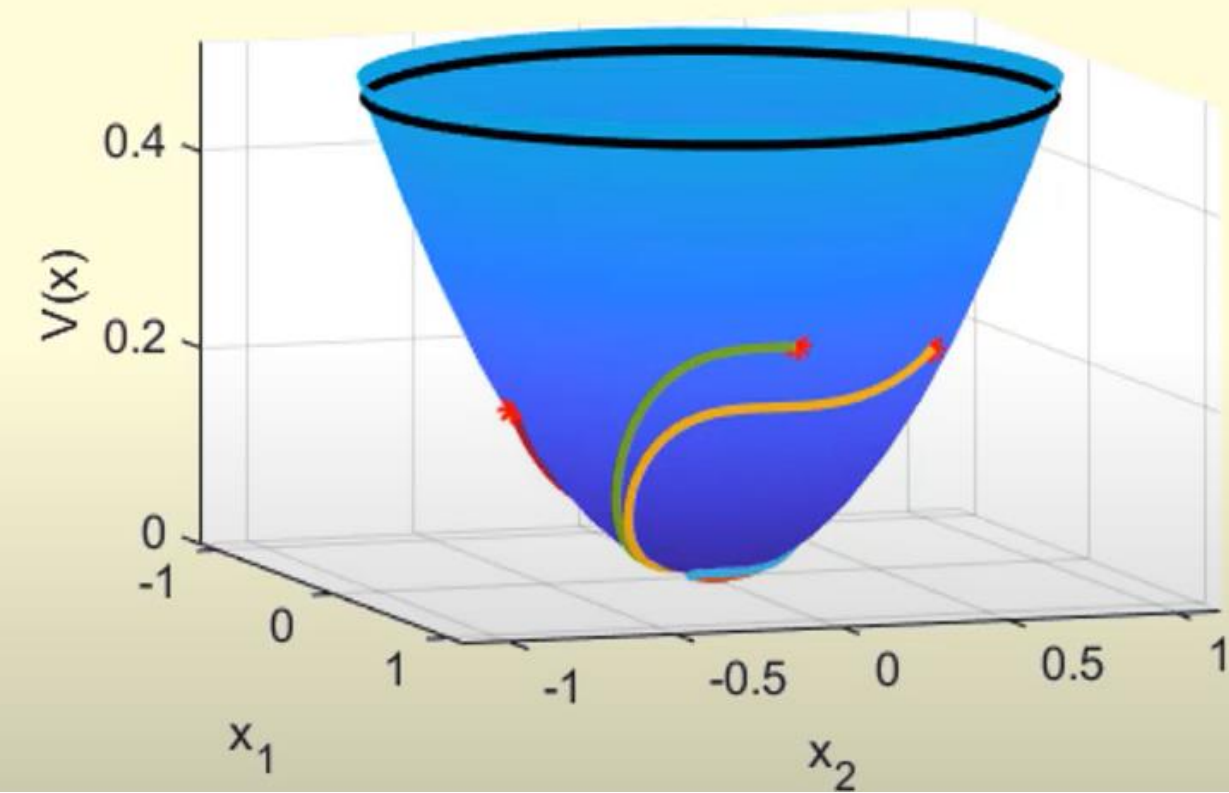
$$\dot{x}_2(t) = 0 \Rightarrow -\sin x_1 - kx_2 = 0 \Rightarrow \sin x_1 = 0 \Rightarrow x_1 = 0$$

So the set M is:

$$M = \{(x_1, x_2) = (0, 0)\}$$



LaSalle's Invariance principle: Every solution starting in Ω approaches M as $t \rightarrow \infty$



Example 4: Using LaSalle's Invariance Principle, show the origin is an asymptotic stability equilibrium of the mass-spring-damper system. For simplicity, let $u = 0$, $k = b = m$.

State space model:

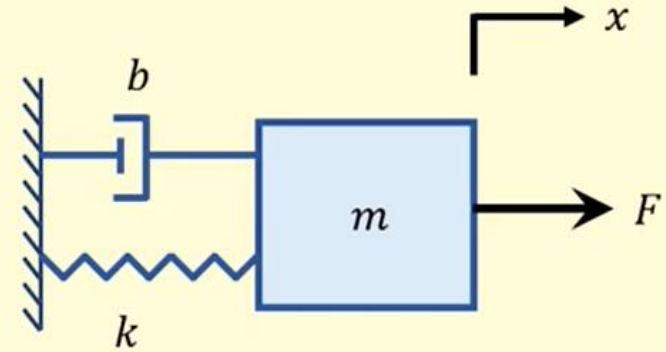
$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{k}{m} & -\frac{b}{m} \end{bmatrix} \begin{bmatrix} x \\ v \end{bmatrix} + \begin{bmatrix} 0 \\ \frac{1}{m} \end{bmatrix} u$$
$$y = [1 \quad 0] \begin{bmatrix} x \\ v \end{bmatrix}$$

State equations: $\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$

Consider $V(x, v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$

$$\dot{V}(x, v) = \frac{\partial V}{\partial x} \dot{x} + \frac{\partial V}{\partial v} \dot{v} = (x)(v) + (v)(-x - v) = -v^2 \leq 0$$

So $\dot{V}(x)$ is negative semi-definite.



m :	Mass
k :	Spring constant
b :	Damping constant
x :	Displacement of the mass
$\dot{x} = v$:	Velocity of the mass
F :	Input force

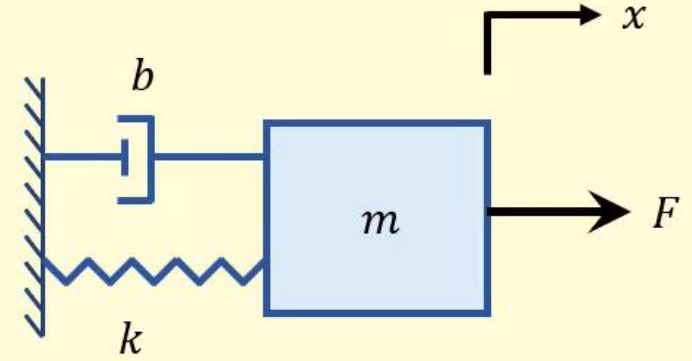
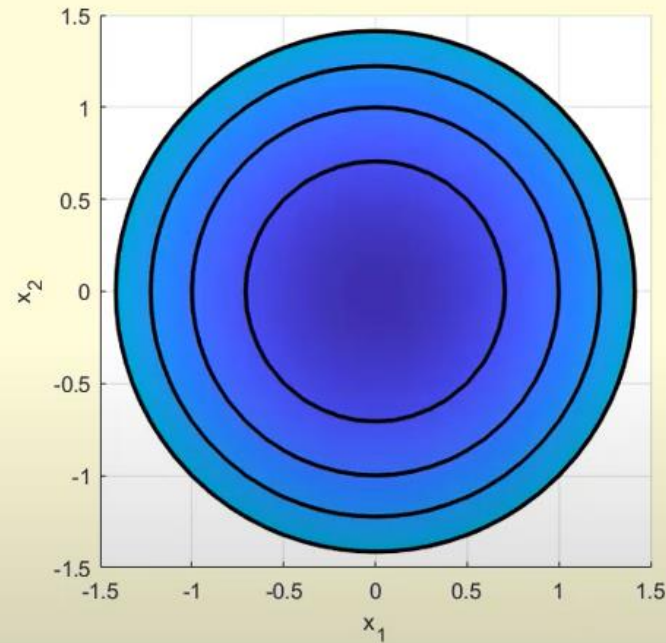
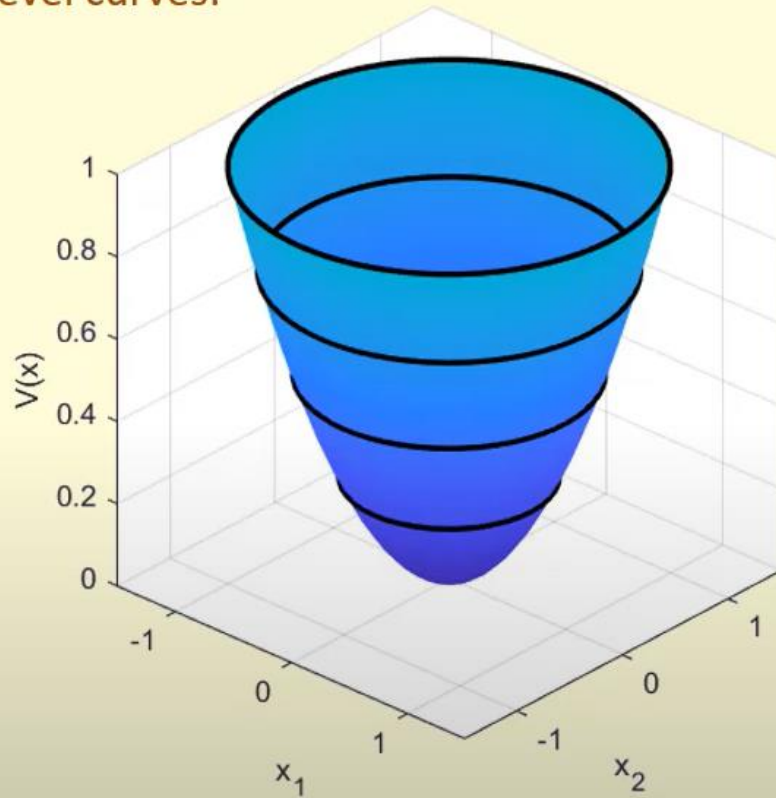
$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$$

$$V(x, v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$$

$$\dot{V}(x, v) = -v^2 \leq 0$$

The set Ω is a compact set that is **positively invariant** with respect to $\dot{x} = f(x)$.

Level curves:



We choose the set Ω as:

$$\Omega = \{(x, v): V(x, v) \leq 1\}$$

$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$$

$$V(x, v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$$

$$\dot{V}(x, v) = -v^2 \leq 0$$

The set $E \subset \Omega$ is the set of all points in Ω such that $\dot{V}(x, v) = 0$.

$$\dot{V}(x, v) = 0 \Rightarrow -kv^2 = 0 \Rightarrow v = 0$$

$$E = \{(x, v) \in \Omega: v = 0\}$$

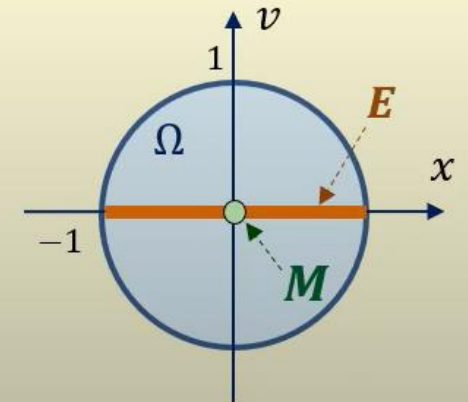
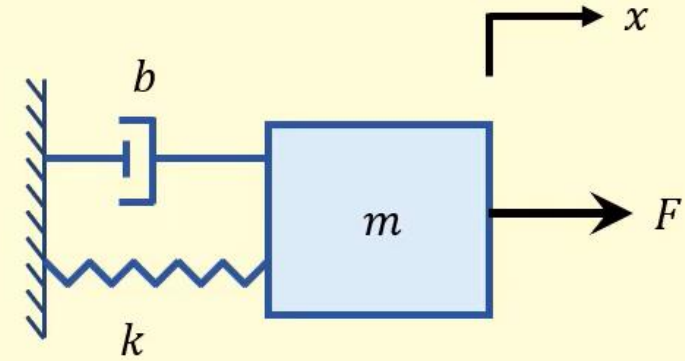
The set M is the *largest positively invariant* set in E .

In the set M we have $v = 0$ for all t . So $\dot{v}(t) = 0$. Using the system dynamics we have:

$$\dot{v}(t) = 0 \Rightarrow -x - v = 0 \Rightarrow x = 0$$

So the set M is

$$M = \{(x, v) = (0, 0)\}$$



LaSalle's Invariance Principle: Every solution starting in Ω approaches M as $t \rightarrow \infty$

Lyapunov Stability Theorem

$V(x)$ is (continuously differentiable and) **positive definite** on $B_r(0)$

$\dot{V}(x)$ is **negative definite** on $B_r(0)$

Discovered by Aleksandr Lyapunov - 1892

LaSalle's Invariance Principle

$V(x)$ is continuously differentiable on Ω
The set Ω is positively invariant with respect to $\dot{x} = f(x)$

$\dot{V}(x)$ is **negative semidefinite** on Ω

Discovered independently by

- *Nikolay Krasovsky* – 1959 (An extension of a result in 1952 by *Barbashin & Krasovsky*)
- *Joseph LaSalle* – 1960