### EGM0004

# Sistemas Não Lineares

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## Teoremas sobre estabilidade

Teorema 6: Teorema de La Salle (princípio da invariância)

Pode ocorrer de  $V \leq 0$  e ainda assim concluir sobre estabilidade assintotica

## Definições:

Positively invariant set

A set  $\Omega$  is said to be a **positively invariant set** with respect to  $\dot{x} = f(x)$  if

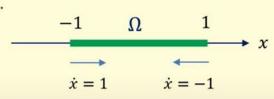
$$x(0) \in \Omega \quad \Rightarrow \quad x(t) \in \Omega, \qquad \forall t \ge 0$$



#### Example 1.

The set  $\Omega = \{x \in \mathbb{R} : |x| \le 1\}$  is a positively invariant set with respect to  $\dot{x} = -x$ .

To check if  $\Omega$  is positively invariant w.r.t the system dynamics, we *just check the boundaries*.



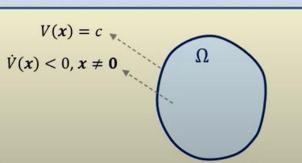
#### Example 2.

Let V be a positive definite.

and  $\dot{V}(x)$  be negative definite for  $\dot{x} = f(x)$ .

Define 
$$\Omega$$
 as  $\Omega = \{x \in \mathbb{R}^n : V(x) \le c\}$ .

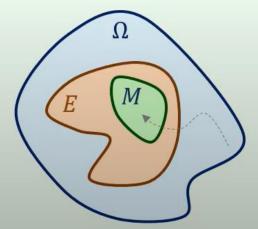
Then  $\Omega$  is a positively invariant set w.r.t.  $\dot{x} = f(x)$ . We **check the boundaries**.



## La Salle

- Let  $\Omega$  be a compact set that is **positively invariant** with respect to  $\dot{x} = f(x)$
- Let V be a continuously differentiable function on  $\Omega$  such that  $\dot{V}(x) \leq 0$  in  $\Omega$
- Let  $E \subset \Omega$  be the set of all points in  $\Omega$  such that  $\dot{V}(x) = 0$
- Let *M* be the *largest positively invariant* set in *E*

Then every solution starting in  $\Omega$  approaches M as  $t \to \infty$ 



# La Salle local e global

Teorema 6: Teorema de La Salle (princípio da invariância)

Se para um sistema  $\dot{x} = f(x), f(0) = 0$ , existe uma função V(x) tal que

- a) V(x) > 0
- b) V(x) é continuamente diferenciável
- c)  $V(x) \le 0$
- d)  $\not\exists x \neq 0$  tal que  $V(x) = 0, \forall t \geq 0$  Então é assintoticamente estável

Teorema 7: Teorema de La Salle global

Se para um sistema  $\dot{x} = f(x), f(0) = 0$ , existe uma função V(x) tal que

- a) V(x) > 0
- b) V(x) é continuamente diferenciável
- c) V(x) < 0
- d)  $\exists x \neq 0 \text{ tal que } V(x) = 0, \forall t \geq 0$
- e)  $V(x) \to \infty$  quando  $||x|| \to \infty$

Então é globalmente assintoticamente estável

O conjunto definido pela derivada de V(x)=0 não contenha

trajetórias além da trajetória identicamente nula

## La Salle local e global

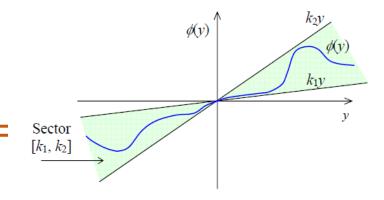
Teorema 6: Teorema de La Salle (princípio da invariância)

### Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -x_1 - x_2$$

$$V(x) = \frac{1}{2}x_1^2 + \frac{1}{2}x_2^2 > 0$$



### Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -\phi(x_1) - \mu x_2, \mu > 0$$

$$V(x) = \frac{1}{2}x_2^2 + \int_0^{x_1} \phi(t)dt > 0$$

$$\phi(x_1)x_1 > 0$$
 se  $x_1 \neq 0$  Função setorial nos quadrantes impares

$$x_1 = 0 \implies \phi(x_1) = 0$$

# La Salle local e global

Teorema 7: Teorema de La Salle global Exemplo:

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -x_1 + x_3$$

$$\dot{x}_3 = -x_2 - x_3$$

$$V(x) = \frac{1}{2}x_1^2 + \frac{1}{2}x_2^2 + \frac{1}{2}x_3^2 > 0$$

## La Salle

**Example 3 (pendulum with friction):** Show the asymptotic willity of the origin x = (0,0) using LaSalle's Theorem for

the general case where  $k \neq 0$ . For simplicity, let  $\ell = g$  and m = 1.

State equations:

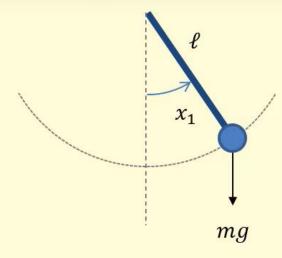
$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -\frac{g}{\ell} \sin x_1 - \frac{k}{m} x_2$$

Let 
$$V(x) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$
 
$$V(x) = mgl(1 - \cos x_1) + \frac{1}{2}ml^2x_2^2$$

$$\dot{V}(\mathbf{x}) = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = (\sin x_1)(x_2) + (x_2)(-\sin x_1 - kx_2) = -kx_2^2 \le 0$$

So  $\dot{V}(x)$  is negative semi-definite.



m: mass of the bob

 $\ell$ : length of the rod

k: friction coefficient

## La Salle

**Example 3 (pendulum with friction):** Show the asymptotic willity of the origin x = (0,0) using LaSalle's Theorem for

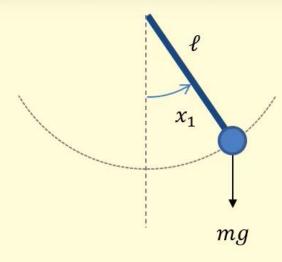
the general case where  $k \neq 0$ . For simplicity, let  $\ell = g$  and m = 1.

State equations:  $\dot{x}_1 = x_2 \\ \dot{x}_2 = -\frac{g}{\rho} \sin x_1 - \frac{k}{m} x_2$ 

Let 
$$V(x) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$
  $\dot{V}(x) = \nabla V \cdot f(x)$ 

$$\dot{V}(x) = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = (\sin x_1)(x_2) + (x_2)(-\sin x_1 - kx_2) = -kx_2^2 \le 0$$

So  $\dot{V}(x)$  is negative semi-definite.



m: mass of the bob

e length of the rod

k: friction coefficient

The set  $\Omega$  is a compact set that is **positively invariant** with respect to  $\dot{x} = f(x)$ .

The set  $\Omega$  is usually chosen as a level set of the V if V is a positive definite function.

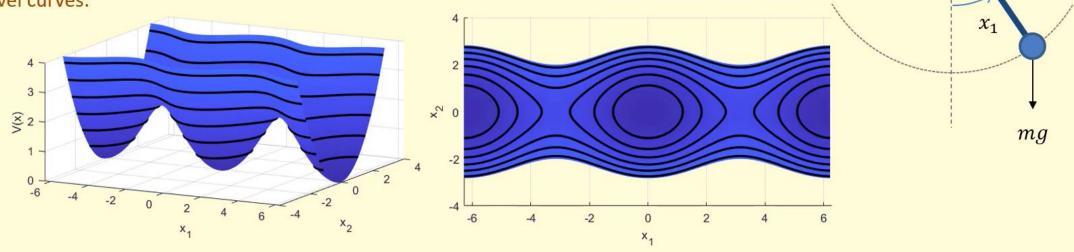
$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -\sin x_1 - kx_2$$

$$V(x) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$

$$\dot{V}(x) = -kx_2^2 \le 0$$

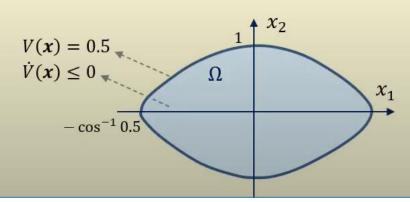
#### The level curves:



ou.. 
$$\Omega = [x_1 \ x_2]^T : -\pi < x_1 < \pi, |x_2| < k$$

### We choose the set $\Omega$ as:

$$\Omega = \{ \boldsymbol{x} \in \mathbb{R}^2 \colon V(\boldsymbol{x}) \le 0.5 \}$$



$$\dot{x}_1 = x_2$$
  
$$\dot{x}_2 = -\sin x_1 - kx_2$$

$$V(\mathbf{x}) = 1 - \cos x_1 + \frac{1}{2}x_2^2$$
  $\dot{V}(\mathbf{x}) = -kx_2^2 \le 0$ 

$$\dot{V}(x) = -kx_2^2 \le 0$$

$$E:=\{x:\dot{V}=0\}$$
 The set  $E\subset\Omega$  is the set of all points in  $\Omega$  such that  $\dot{V}(x)=0$ .

$$\dot{V}(x) = 0 \Rightarrow -kx_2^2 = 0 \Rightarrow x_2 = 0$$

$$E = \{(x_1, x_2) \in \Omega \colon x_2 = 0\}$$

The set M is the largest positively invariant set in E.

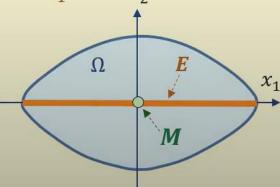
To find M, we let  $x_2 = 0$  for all  $t \ge 0$ . So  $\dot{x}_2(t) = 0$ .

Using the system's dynamical model we have:

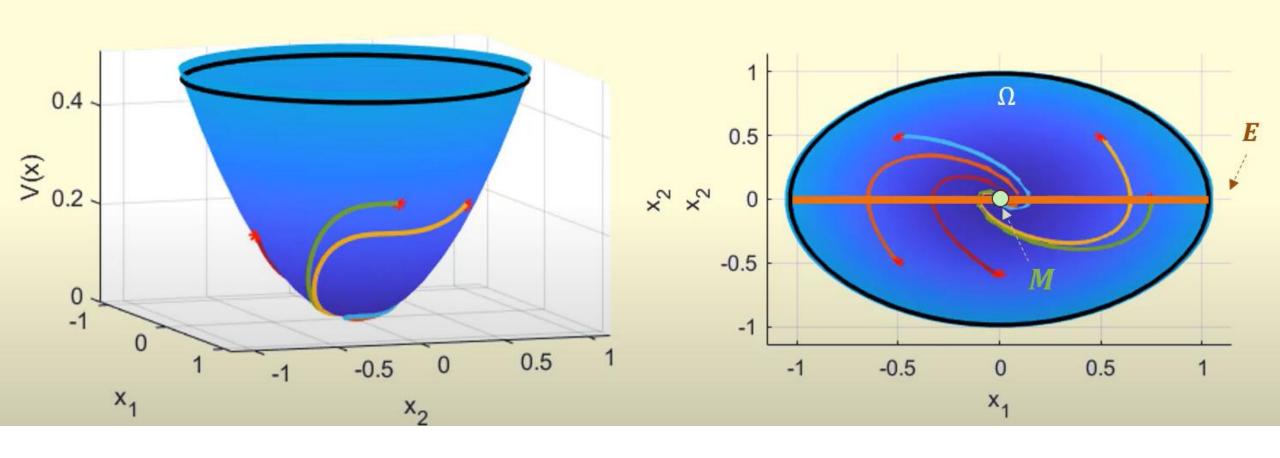
$$\dot{x}_2(t) = 0$$
  $\Rightarrow$   $-\sin x_1 - kx_2 = 0$   $\Rightarrow$   $\sin x_1 = 0$   $\Rightarrow$   $x_1 = 0$   $x_2$ 

So the set *M* is:

$$M = \{(x_1, x_2) = (0,0)\}$$



### **LaSalle's Invariance principle:** Every solution starting in $\Omega$ approaches M as $t \to \infty$



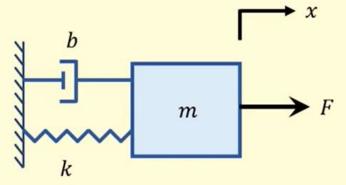
**Example 4:** Using LaSalle's Invariance Principle, show the origin is an asymptotic stability equilibrium of the mass-spring-damper system. For simplicity, let u = 0, k = b = m.

State equations: 
$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$$

Consider 
$$V(x, v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$$

$$\dot{V}(x,v) = \frac{\partial V}{\partial x}\dot{x} + \frac{\partial V}{\partial v}\dot{v} = (x)(v) + (v)(-x - v) = -v^2 \le 0$$

So  $\dot{V}(x)$  is negative semi-definite.



m: Mass

k: Spring constant

b: Damping constant

x: Displacement of the mass

 $\dot{x} = v$ : Velocity of the mass

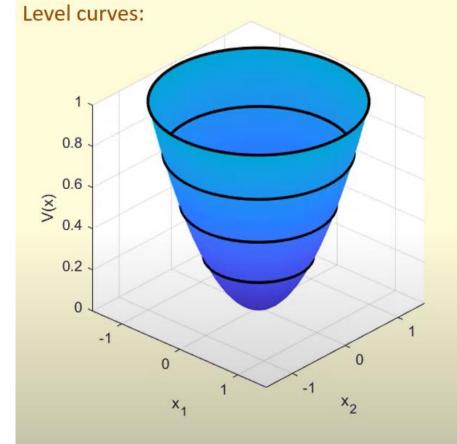
F: Input force

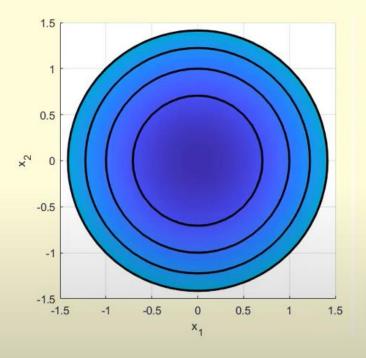
$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$$

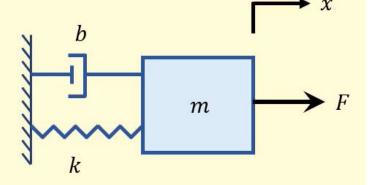
$$V(x,v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$$

$$\dot{V}(x,v) = -v^2 \le 0$$

The set  $\Omega$  is a compact set that is **positively invariant** with respect to  $\dot{x} = f(x)$ .







We choose the set  $\Omega$  as:

$$\Omega = \{(x, v) \colon V(x, v) \le 1\}$$

$$\begin{bmatrix} \dot{x} \\ \dot{v} \end{bmatrix} = \begin{bmatrix} v \\ -x - v \end{bmatrix}$$

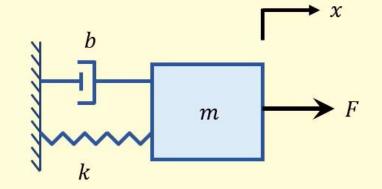
$$V(x,v) = \frac{1}{2}x^2 + \frac{1}{2}v^2$$

$$\dot{V}(x,v) = -v^2 \le 0$$

The set  $E \subset \Omega$  is the set of all points in  $\Omega$  such that  $\dot{V}(x,v) = 0$ .

$$\dot{V}(x,v) = 0 \Rightarrow -kv^2 = 0 \Rightarrow v = 0$$

$$E = \{(x, v) \in \Omega \colon v = 0\}$$



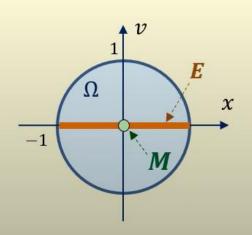
The set M is the *largest positively invariant* set in E.

In the set M we have v=0 for all t. So  $\dot{v}(t)=0$ . Using the system dynamics we have:

$$\dot{v}(t) = 0 \Rightarrow -x - v = 0 \Rightarrow x = 0$$

So the set *M* is

$$M = \{(x, v) = (0,0)\}$$



**LaSalle's Invariance Principle**: Every solution starting in  $\Omega$  approaches M as  $t \to \infty$ 

### **Lyapunov Stability Theorem**

V(x) is (continuously differentiable and) **positive** definite on  $B_r(0)$ 

 $\dot{V}(x)$  is **negative definite** on  $B_r(0)$ 

Discovered by Aleksandr Lyapunov - 1892

### LaSalle's Invariance Principle

V(x) is continuously differentiable on  $\Omega$ The set  $\Omega$  is positively invariant with respect to  $\dot{x} = f(x)$ 

 $\dot{V}\left(x\right)$  is **negative semidefinite** on  $\Omega$ 

Discovered independently by

- Nikolay Krasovsky 1959 (An extension of a result in 1952 by Barbashin & Krasovsky)
- Joseph LaSalle 1960