

## Lecture 2 - How to Model Stars

### Equations of Stellar Structure

$$\textcircled{1} \quad \rho \frac{D\mathbf{u}}{Dt} = -\nabla p - \rho \nabla \phi \quad (\text{Euler}) \quad \mathbf{u} \rightarrow \text{velocity}$$

"ma" =  $\downarrow$  pressure       $\downarrow$  gravity       $\phi \rightarrow$  gravitational potential

$$\textcircled{2} \quad \frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{u} = 0 \quad (\text{conservation of mass})$$

$$\textcircled{3} \quad \nabla^2 \phi = 4\pi G \rho \quad (\text{gravity})$$

$$\textcircled{4} \quad \frac{DU}{Dt} - \frac{1}{\rho^2} \frac{D\rho}{Dt} = \epsilon - \frac{1}{\rho} \nabla \cdot \mathbf{F} \quad (\text{energy})$$

Some basic assumptions:

- time independence  $\rightarrow$  static  $\mathbf{u} = 0$
- spherical symmetry

$$\textcircled{1} \rightarrow \frac{dp}{dr} = -\rho \frac{d\phi}{dr}$$

$$\textcircled{2} \rightarrow 0 = 0$$

$$\textcircled{3} \rightarrow \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{d\phi}{dr} \right) = 4\pi G \rho$$

$$\textcircled{4} \rightarrow 0 = \epsilon - \frac{1}{\rho} \nabla \cdot \mathbf{F}$$

Integrate:  $r^2 \frac{d\phi}{dr} = 4\pi G \int_0^r \rho(r') r'^2 dr'$

$G M(r)$

$\hookrightarrow$  mass enclosed in radius  $r$ .

$$\frac{d\phi}{dr} = \frac{G M(r)}{r^2}$$

①  $\rightarrow$

$$\frac{dp}{dr} = -\frac{GM\rho}{r^2}$$

Hydrostatic equilibrium

Without another relationship between  $p$  and  $\rho$  we can't solve this!

### Polytropic Model

Assume

$$p = k \rho^\gamma$$

,  $k$  and  $\gamma$  are constant.

$$\textcircled{1} \rightarrow \frac{dp}{dr} = -\frac{GM\rho}{r^2} \Rightarrow \frac{r^2}{\rho} \frac{dp}{dr} = -GM$$

Take derivative of both sides:

$$\frac{d}{dr} \left( \frac{r^2}{\rho} \frac{dp}{dr} \right) = -G \frac{dM}{dr}$$

$$\text{From } \textcircled{3}, \quad \frac{dM}{dr} = 4\pi r^2 \rho$$

$$\frac{1}{r^2} \frac{d}{dr} \left( \frac{r^2}{\rho} \frac{dp}{dr} \right) = -4\pi G \rho$$

$$\text{Use } p = k \rho^\gamma \Rightarrow \frac{dp}{dr} = k \gamma \rho^{\gamma-1} \frac{d\rho}{dr}$$

$$\frac{k\gamma}{r^2} \frac{d}{dr} \left( r^2 \rho^{\gamma-2} \frac{d\rho}{dr} \right) = -4\pi G \rho$$

This is an ODE involving just density.

Clean it up:

- define a dimensionless parameter:

$$\rho(r) = \rho_c \tilde{\rho}^n(r)$$

$\rightarrow$  dimensionless density

- define  $n$  such that

$$\gamma = \frac{n+1}{n}$$

- define  $\eta = \frac{r}{\lambda_n}$ ,  $\lambda_n = \left[ \frac{(n+1)k}{4\pi G} \rho_c^{\frac{(1-n)}{n}} \right]^{1/2}$

$\Rightarrow$  The ODE becomes

$$\frac{1}{\eta^2} \frac{d}{d\eta} \left( \eta^2 \frac{d\rho}{d\eta} \right) = -\rho^n$$

Lane-Emden Equation

Boundary Conditions:

- At the surface, the density goes to zero:

$$\rho(\eta_s) = 0$$

$\hookrightarrow$  radius of surface  $\eta_s$

- At the centre of the star, the density is  $\rho_c$

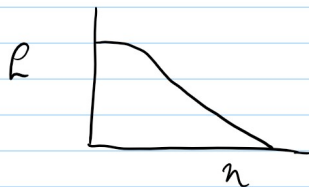
$$\Rightarrow \rho(0) = 1$$

- $\frac{d\rho}{d\eta} = 0$  at  $\eta = 0$

Analytic Solutions:

- $n = 0$ ,  $\rho(\eta) = 1 - \frac{\eta^2}{6}$

Surface is at  $\rho(\eta_s) = 0 \Rightarrow \eta_s = \sqrt{6}$ .



- $n = 1$ ,  $\rho(\eta) = \frac{\sinh \eta}{\eta}$ ,  $\eta_s = \pi$

- $n = 5$ ,  $\rho(\eta) = \frac{1}{\sqrt{1 + \frac{1}{3}\eta^2}}$ ,  $\eta_s = \infty$ . \*

## White Dwarfs

- He, O, C
    - nonrelativistic gas —
    - relativistic gas —
- } polytropes  $p \propto \rho^\gamma$

### Steps

- solve L-E eq. via RKY
- give  $\rho_c$ , calculate  $r_s$ ,  $M$
- look up the Chandrasekhar limit and typical radii of WDs —

How good are the models?